

## Dihadron fragmentation functions and their relevance for transverse spin studies

This article has been downloaded from IOPscience. Please scroll down to see the full text article.

2011 J. Phys.: Conf. Ser. 295 012053

(<http://iopscience.iop.org/1742-6596/295/1/012053>)

View [the table of contents for this issue](#), or go to the [journal homepage](#) for more

Download details:

IP Address: 147.156.26.43

The article was downloaded on 11/07/2012 at 13:39

Please note that [terms and conditions apply](#).

# Dihadron fragmentation functions and their relevance for transverse spin studies

A. Courtoy<sup>1</sup>, A. Bacchetta<sup>1,2</sup> and M. Radici<sup>1</sup>

<sup>1</sup> INFN-Sezione di Pavia, 27100 Pavia, Italy.

<sup>2</sup> Dipartimento di Fisica Nucleare e Teorica, Università di Pavia, 27100 Pavia, Italy.

E-mail: aurore.courtoy@pv.infn.it

**Abstract.** Dihadron fragmentation functions describe the probability that a quark fragments into two hadrons plus other undetected hadrons. In particular, the so-called interference fragmentation functions describe the azimuthal asymmetry of the dihadron distribution when the quark is transversely polarized. They can be used as tools to probe the quark transversity distribution in the nucleon. Recent studies on unpolarized and polarized dihadron fragmentation functions are presented, and we discuss their role in giving insights into transverse spin distributions.

## 1. Introduction

Our knowledge on the hadron structure is incomplete. We know that the Parton Distribution Functions (PDFs) describe the one-dimensional structure of hadrons. At leading order, the PDFs are three: number density, helicity and transversity. However the experimental knowledge on the latter is rather poor as it is a chiral-odd quantity not accessible through fully inclusive processes. Semi-inclusive production of two hadrons [1, 2] offers an alternative way to access transversity, where the chiral-odd partner of transversity is represented by the Dihadron Fragmentation Functions (DiFF)  $H_1^{\sphericalangle}$  [3], which relates the transverse spin of the quark to the azimuthal orientation of the two-hadron plane. Since the transverse momentum of the hard parton is integrated out, the cross section can be studied in the context of collinear factorization. This peculiarity is an advantage over the  $p_T$ -factorization framework, where the cross sections involve convolutions of the relevant functions instead of simple products.

The transversely polarized DiFF has been computed only in a spectator model [4]. Recently, the HERMES collaboration has reported measurements of the asymmetry containing the product  $h_1 H_1^{\sphericalangle}$  [5]. The COMPASS collaboration has presented analogous preliminary results [6]. The BELLE collaboration has also presented preliminary measurements of the azimuthal asymmetry in  $e^+e^-$  annihilation related to the DiFF [7].

Our present goal is to extract transversity through this channel. To this end, we need an expression for the chiral-odd DiFF  $H_1^{\sphericalangle}$  obtained from  $e^+e^-$  data. This in its turn requires a knowledge of the unpolarized DiFF  $D_1$ . Hence, as a first step, we present here a parameterization of the unpolarized DiFF  $D_1$  as given from the Monte Carlo generator (MC) of the BELLE collaboration.

## 2. Two-hadron Inclusive DIS: towards Transversity

We consider the SIDIS process  $e(l) + N^\uparrow(P) \rightarrow e(l') + \pi^+(P_1) + \pi^-(P_2) + X$ , where the momentum transfer  $q = l - l'$  is space-like, with  $l, l'$ , the lepton momenta before and after the scattering. The two pions coming from the fragmenting quark have momenta  $P_1$  and  $P_2$ , respectively, and invariant mass  $M_h$ , which is considered to be much smaller than the hard scale of the process. We introduce the vectors  $P_h = P_1 + P_2$  and  $R = (P_1 - P_2)/2$ . We describe a 4-vector  $a$  as  $[a^-, a^+, a^x, a^y]$ , i.e. in terms of its light-cone components  $a^\pm = (a^0 \pm a^3)/\sqrt{2}$  and its transverse spatial components. We introduce the light-cone fraction  $z = P_h^-/k^-$ .  $P$  is the momentum of the nucleon target with mass  $M$ . We refer to Refs. [3, 4] for details and kinematics.

The spin asymmetry  $A_{UT}^{\sin(\phi_R + \phi_S) \sin\theta}(x, y, z, M_h^2)$  is related to an asymmetric modulation of pion pairs in the angles  $\phi_S$  and  $\phi_R$ , which represent the azimuthal orientation with respect to the scattering plane of the target transverse polarization and of the plane containing the pion pair momenta, respectively. The polar angle  $\theta$  describes the orientation of  $P_1$ , in the center-of-mass frame of the two pions, with respect to the direction of  $P_h$  in the lab frame. The asymmetry is expressed as

$$A_{UT}^{\sin(\phi_R + \phi_S) \sin\theta}(x, y, z, M_h^2) \propto -\frac{|\mathbf{R}|}{M_h} \frac{\sum_q e_q^2 h_1^q(x) H_{1,q}^{\leftarrow sp}(z, M_h^2)}{\sum_q e_q^2 f_1^q(x) D_{1,q}^{ss+pp}(z, M_h^2)}, \quad (1)$$

where the  $x$ -dependence is given by the PDFs only. The  $z$  and  $M_h$  dependence are governed by the DiFFs whose functional form we need to determine. The procedure allowing us to give the required parameterizations for the DiFFs is detailed in the following sections.

## 3. The Artru-Collins Asymmetry

We further consider the process  $e^+(l)e^-(l') \rightarrow (\pi^+\pi^-)_{\text{jet1}}(\pi^+\pi^-)_{\text{jet2}}X$ , with (time-like) momentum transfer  $q = l + l'$ . Here, we have two pairs of pions, one originating from a fragmenting parton and the other one from the related antiparton.<sup>1</sup>

The differential cross sections also depend on the invariant  $y = P_h \cdot l / P_h \cdot q$  which is related, in the lepton center-of-mass frame, to the angle  $\theta_2 = \arccos(\mathbf{l}_{e^+} \cdot \mathbf{P}_h / (|\mathbf{l}_{e^+}| |\mathbf{P}_h|))$ , with  $\mathbf{l}_{e^+}$  the momentum of the positron, by  $y = (1 + \cos\theta_2)/2$ .

The dihadron Fragmentation Functions are involved in the description of the fragmentation process  $q \rightarrow \pi^+\pi^-X$ , where the quark has momentum  $k$ . They are extracted from the correlation function [8]

$$\Delta^q(z, \cos\theta, M_h^2, \phi_R) = \frac{z|\vec{R}|}{16 M_h} \int d^2\vec{k}_T dk^+ \Delta^q(k; P_h, R) \Big|_{k^- = P_h^-/z}, \quad (2)$$

where

$$\Delta^q(k, P_h, R)_{ij} = \sum_X \int \frac{4\xi}{(2\pi)^4} e^{+ik \cdot \xi} \langle 0 | \mathcal{U}_{(-\infty, \xi)}^{n+} \psi_i^q(\xi) | P_h, R; X \rangle \langle P_h, R; X | \bar{\psi}_j^q(0) \mathcal{U}_{(0, -\infty)}^{n+} | 0 \rangle (3)$$

Since we are going to perform the integration over the transverse momentum  $\vec{k}_T$ , the Wilson lines  $\mathcal{U}$  can be reduced to unity using a light-cone gauge. The only fragmentation functions surviving after integration over the azimuthal angle defining the position of the lepton plane w.r.t. the laboratory plane [9].

$$D_1^q(z, \cos\theta, M_h^2) = 4\pi \text{Tr}[\Delta^q(z, \cos\theta, M_h^2, \phi_R) \gamma^-], \quad (4)$$

$$\frac{\epsilon_T^{ij} R_{Tj}}{M_h} H_1^{\leftarrow q}(z, \cos\theta, M_h^2) = 4\pi \text{Tr}[\Delta^q(z, \cos\theta, M_h^2, \phi_R) i \sigma^{i-} \gamma_5]. \quad (5)$$

<sup>1</sup> Variables with an extra ‘‘bar’’ refer to the pair coming from the antiquark.

We perform an expansion in terms of Legendre functions of  $\cos\theta$  (and  $\cos\bar{\theta}$ ) and keep only the  $s$ - and  $p$ -wave components of the relative partial waves of the pion pair. By further integrating upon  $d\cos\theta$  and  $d\cos\bar{\theta}$ , we isolate only the specific contributions of  $s$  and  $p$  partial waves to the respective DiFFs.

The azimuthal Artru-Collins asymmetry  $A(\cos\theta_2, z, \bar{z}, M_h^2, \bar{M}_h^2)$  [10] corresponds to a  $\cos(\phi_R + \phi_{\bar{R}})$  modulation in the cross section for the process under consideration. It can be written in terms of DiFF in the following way,

$$A(\cos\theta_2, z, M_h^2, \bar{z}, \bar{M}_h^2) = \frac{\sin^2\theta_2}{1 + \cos^2\theta_2} \frac{\pi^2}{32} \frac{|\mathbf{R}| |\bar{\mathbf{R}}|}{M_h \bar{M}_h} \frac{\sum_q e_q^2 H_{1,q}^{\langle sp \rangle}(z, M_h^2) \bar{H}_{1,q}^{\langle sp \rangle}(\bar{z}, \bar{M}_h^2)}{\sum_q e_q^2 D_{1,q}^{ss+pp}(z, M_h^2) \bar{D}_{1,q}^{ss+pp}(\bar{z}, \bar{M}_h^2)}, \quad (6)$$

with  $|\mathbf{R}| = \frac{M_h}{2} \sqrt{1 - 4m_\pi^2/M_h^2}$ . To extract a parameterization of the function  $H_1^{\langle \rangle}$ , we need to know the function  $D_1$ .

#### 4. Electron-Positron Annihilation: The Unpolarized Cross-Section from BELLE

A model independent parameterization of a function means a huge freedom on the functional form one will choose. First, one can guess the causes of the shape of the data from physical arguments. One can get inspired in comparing the model results with the data: here, we take into account the results of Ref. [4] —including a critical eye on its shortcomings— in defining the shape of the MC histograms for the unpolarized cross section.

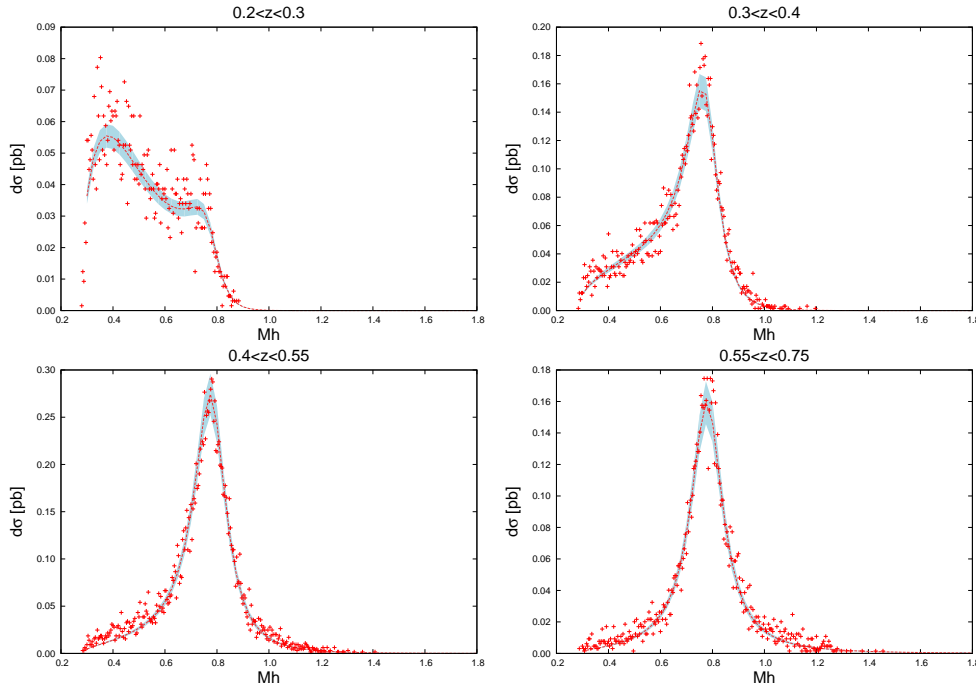
In the process  $q \rightarrow \pi^+\pi^-X$ , the prominent channels for an invariant mass of the pion pair ranging  $2m_\pi < M_h \lesssim 1.5$  GeV are, basically :

- the fragmentation into a  $\rho$  resonance decaying into  $\pi^+\pi^-$ , responsible for a peak at  $M_h \sim 770$  MeV ;
- the fragmentation into a  $\omega$  resonance decaying into  $\pi^+\pi^-$ , responsible for a small peak at  $M_h \sim 782$  MeV plus the fragmentation into a  $\omega$  resonance decaying into  $\pi^+\pi^-\pi^0$  ( $\pi^0$  unobserved), responsible for a broad peak around  $M_h \sim 500$  MeV ;
- the continuum, i.e. the fragmentation into an “incoherent”  $\pi^+\pi^-$  pair, is probably the most important channel. It is also the most difficult channel to describe with purely model-based physical arguments.

In addition to the channel decomposition, one has to take into account the flavor decomposition of the cross section. This further decomposition is particularly important if one wants to be able to use the resulting parametrization in another context, e.g., SIDIS. For the time being, the MC data provided by the BELLE collaboration are additionally separated into flavors, i.e.,  $uds$  contributions and  $c$  contributions. The experimental analyses conclude that the charm contribution to the unpolarized cross section is non-negligible at BELLE’s energy.<sup>2</sup>

The main considerations one can do, before fitting the data, are the following. First, the most important contribution from the charm is in the continuum and cannot be neglected. The determination of a functional form for  $D_1$  consists then in four parallel steps, i.e. the 2-dimensional parameterization of the  $\rho$  and  $\omega$  channels and of the continuum for  $uds$  and only of the continuum for the charm. Second, it can be deduced that both the  $\rho$  and  $\omega$  channels play a role at high  $z$  values, while it seems that the  $\rho$  is less important at lower  $z$  values, as it can be seen in Fig. 1. On the other hand, the continuum decreases with  $z$ , and this behavior is different for the  $uds$  and the  $c$  flavors. Also, it can be observed, e.g. in Fig. 1, that the behavior in  $M_h$  changes from  $z$ -bin to  $z$ -bin. Those are signs that the dependence on  $z$  and  $M_h$  cannot be factorized.

<sup>2</sup> R. Seidl’s talk at TMD workshop, ECT\*, June 2010.



**Figure 1.** Number of events  $N$  for the unpolarized  $e^+e^-$  annihilation into 2 pions in a jet (plus anything else) at BELLE, normalized by the integrated luminosity  $647.26 \text{ pb}^{-1}$ . We show only the resonant channel for the  $\rho$  production. The data are represented by the dots. The error on the data (not plotted here) is assumed to be  $\sqrt{N}$ . The dashed lines represent the parameterization, and the band its errorband.  $M_h$  in GeV.

Following Eq. (6), the unpolarized cross section that we are considering here is differential in  $(\cos \theta_2, z, M_h^2, \bar{z}, \bar{M}_h^2)$ . The  $\theta_2$ -dependence is provided by the BELLE collaboration, and the set of variables  $(\bar{z}, \bar{M}_h)$  is integrated out within the experimental bounds. The methodology is as follows. The unpolarized cross section, differential in  $M_h$  and  $z$ , is

$$\frac{d\sigma^U}{2M_h dM_h dz} = \sum_{a,\bar{a}} e_a^2 \frac{1}{3} \frac{6\alpha^2}{Q^2} \langle 1 + \cos^2 \theta_2 \rangle z^2 D_1^a(z, M_h^2) \int_0^1 d\bar{z} \int_{2m_\pi}^{M_h^{\max}} 2\bar{M}_h d\bar{M}_h \bar{z}^2 \bar{D}_1^a(\bar{z}, \bar{M}_h^2) \quad ,$$

with the integration limits to be modified according to the experiment, and where  $D_1 = D_1^{ss+pp}$ . For both the  $uds$  and  $c$  flavors, the fitted function takes the form

$$FF_U = \frac{1}{3} \frac{6\alpha^2}{Q^2} \langle 1 + \cos^2 \theta_2 \rangle \sum_a e_a^2 \int_{z_{bin}} dz f_{D_1}^a(z, M_h) \int_{0.2}^1 d\bar{z} \int_{2m_\pi}^{1.5\text{GeV}} d\bar{M}_h f_{D_1}^{\bar{a}}(\bar{z}, \bar{M}_h) \quad , (7)$$

where  $\int_{z_{bin}} dz$  means that we average over the  $z$ -dependence each  $z$ -bin, and with our functional form appearing in  $f_{D_1}^a(z, M_h) = 2M_h z^2 D_1^a(z, M_h^2)$ . The upper integration limit in Eq. (7) is chosen to be in agreement with the condition  $M_h \ll Q$ . The DiFFs for quarks and antiquarks are related through the charge conjugation rules described in Ref. [4].

The determination of a functional form  $f_{D_1}$  is done by fitting, by means of a  $\chi^2$  goodness-of-fit test, the MC histograms (4  $z$ -bins and about 300  $M_h$ -bins) for each channel. The best-fit functional forms lead to interesting results. The most important point is that there is no way the  $z$  and the  $M_h$  dependence can be factorized. Moreover, we have realized that no acceptable

fit would be reached with a trivial functional form for the continua. In Fig. 1 we show, as an example, the MC of the  $\rho$  production together with its parametrization. In the depicted case, the joint  $\chi^2/\text{d.o.f.}$  is  $\sim 1.25$  [11]. We quote the  $\chi^2$  values for the other channels:  $\omega$ -production ( $\chi^2/\text{d.o.f.} \sim 1.3$ ) ;  $uds$ -background ( $\chi^2/\text{d.o.f.} \sim 1.4$ ) ;  $c$ -background ( $\chi^2/\text{d.o.f.} \sim 1.55$ ) [11]. The propagation of errors gives rise to the  $1\text{-}\sigma$  error band shown in light blue.

### 5. Towards an extraction of $H_1^\triangleleft$

The DiFF  $H_1^\triangleleft(z, M_h)$  can be extracted from the Artru-Collins asymmetry. The preliminary data from the BELLE collaboration [7] will be our starting point. Those data are binned in  $(z, \bar{z})$  and  $(M_h, \bar{M}_h)$ . While we have stated in the previous section that no factorization of the  $(z, M_h)$  variables is possible for  $D_1$ , the data do not allow us to make a similar statement for  $H_1^\triangleleft$ .

The next step consists in the determination of a functional form, e.g.,

$$f_{H_1^\triangleleft}(z, M_h, \bar{z}, \bar{M}_h) \propto f(z)f(\bar{z})g(M_h)g(\bar{M}_h) \quad . \quad (8)$$

Even if we expect the  $H_1^\triangleleft$  to arise from an  $sp$ -wave interference, we presently have no guidance on the interplay of the  $(z, M_h)$  variables in the asymmetry. We opt for the simpler functional form (8) instead. Given the large uncertainties —we sum statistical and systematic errors in quadrature— on the asymmetry as well as the shape of the  $(z, \bar{z})$  dependence, it is easily realized that more than one functional form could fit the data. We are currently working in improving our fitting procedure in order to get as much information as we can from the data.

Once we will have determined the  $z$  as well as the  $M_h$ -dependence of the  $H_1^\triangleleft$  DiFF, we will have to face the flavor decomposition problem. This step will crucially influence the extraction of transversity, see Eq. (1).

We conclude by highlighting the importance of DiFFs in the extraction of transversity. We are eagerly looking forward to analyzing the published data on  $e^+e^-$  from the BELLE collaboration and to going through the described methodology.

### Acknowledgments

We are thankful to the BELLE collaboration for useful information on the data.

### References

- [1] Collins J C, Heppelmann S F and Ladinsky G A 1994 *Nucl. Phys.* **B420** 565–582 (*Preprint hep-ph/9305309*)
- [2] Jaffe R L, Jin X m and Tang J 1998 *Phys. Rev. Lett.* **80** 1166–1169 (*Preprint hep-ph/9709322*)
- [3] Radici M, Jakob R and Bianconi A 2002 *Phys. Rev.* **D65** 074031 (*Preprint hep-ph/0110252*)
- [4] Bacchetta A and Radici M 2006 *Phys. Rev.* **D74** 114007 (*Preprint hep-ph/0608037*)
- [5] Airapetian A *et al.* (HERMES) 2008 *JHEP* **06** 017 (*Preprint 0803.2367*)
- [6] Martin A (COMPASS) 2006 *Czech. J. Phys.* **56** F33–F52 (*Preprint hep-ex/0702002*)
- [7] Vossen A *et al.* 2009 (*Preprint 0912.0353*)
- [8] Bacchetta A and Radici M 2003 *Phys. Rev.* **D67** 094002 (*Preprint hep-ph/0212300*)
- [9] Boer D, Jakob R and Radici M 2003 *Phys. Rev.* **D67** 094003 (*Preprint hep-ph/0302232*)
- [10] Artru X and Collins J C 1996 *Z. Phys.* **C69** 277–286 (*Preprint hep-ph/9504220*)
- [11] Bacchetta A, Courtoy A and Radici M in preparation