

1 Performance of turbulence closure models for 3D RANS simulation of 2 urban flooding with exchanges between flooded streets and building 3 openings

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12 ABSTRACT

13 Urban flooding poses significant risks to densely populated areas, necessitating accurate predictive
14 models for effective risk management. While 3D computational fluid dynamics (CFD) models are
15 increasingly used to simulate flood dynamics, the selection of an appropriate turbulence closure model
16 remains critical yet underexplored in the context of urban flooding. This study evaluates the performance
17 of three widely used Reynolds-Averaged Navier-Stokes (RANS) turbulence models: standard $k-\varepsilon$, $k-\omega$,
18 and $k-\omega$ SST, in simulating urban flooding in an idealized urban district model with flow intrusion into
19 buildings. Based on laboratory experiments reported in the literature, we conducted three-dimensional
20 simulations using OpenFOAM. Nine urban configurations were considered, each with different opening
21 layouts that allow flow intrusion into buildings. The discharge distribution, surface flow patterns, and
22 velocity fields at selected cross-sections were analyzed. Results show that the simulated discharge
23 partitions match the measured values closely and lie within the $\pm 1.5\%$ uncertainty of the experimental
24 data with all three tested turbulence closures, demonstrating the reliability of the numerical approach. In
25 terms of surface flow and velocity characteristics, the $k-\varepsilon$ and $k-\omega$ SST models showed better agreement
26 with experimental observations than the $k-\omega$ model, particularly in capturing complex surface structures,
27 internal recirculation zones, and along-stream velocity variations. This study provides insights into
28 turbulence model selection for urban flood simulations, particularly those involving flow intrusion into
29 buildings. It constitutes a useful reference for future urban flood modeling.

30 **Keywords:** Urban flooding; Flow structure; Turbulence closure model; 3D modelling

31 **1 Introduction**

32 The frequency and intensity of extreme urban flooding events have increased due to rapid
33 urbanization and more frequent extreme rainfall events associated with climate change (Wang
34 et al., 2023; Tang et al., 2024; Fang, 2016; Hettiarachchi et al., 2018). Such events can cause
35 significant economic losses and even lead to a high number of fatalities (Kreibich et al., 2019),
36 which makes flood risk prediction and management a key urban management priority (Guo et
37 al., 2021; Luo et al., 2022; Rosenzweig et al., 2021). However, predicting urban flooding is a
38 major challenge because of the high variability of flow patterns resulting from the intricacy of
39 urban environments (Li et al., 2025; Saleem et al., 2025), and the multitude of features that
40 interact with the flow, such as sewer systems, buildings, underground spaces, and the flood
41 storage capacity within urban blocks (Kitsikoudis et al., 2021; Mejía-Morales et al., 2023;
42 Mignot et al., 2019; Wu et al., 2024; Yuan et al., 2024). Flood risk is generally assessed based
43 on flow characteristics such as flow depth, flow velocity, and flood duration at local scales (Xia
44 et al., 2014; Zhu et al., 2023). Therefore, advancing our understanding of urban flood processes
45 and accurately quantifying the flow characteristics within urban areas is of utmost importance
46 (Beteille, 2024; Dewals et al., 2023).

47 To thoroughly investigate the hydraulic characteristics of urban flooding, in addition to
48 collecting field data (Chanson et al., 2018), laboratory experiments designed for different urban
49 configurations have also provided valuable insights (Dong et al., 2021; Finaud-Guyot et al.,
50 2019; Li et al., 2022; Mejía-Morales et al., 2023; Mignot et al., 2019; Sañudo et al., 2024).
51 These experimental datasets offer reliable references for the development and validation of
52 numerical models, allowing urban flood hydraulics to be explored in greater depth through
53 numerical simulations (Dellinger et al., 2025; Dewals et al., 2023; Li et al., 2025; Wu et al.,
54 2024). Two-dimensional models are more commonly used to reproduce urban flood processes,
55 whereas the application of 3D models to urban flooding has only emerged in recent years (Li
56 et al., 2025; Viccione et al., 2022; Zhu et al., 2023). Two-dimensional models that solve the
57 depth-averaged shallow water equations, offer high computational efficiency and have

58 demonstrated good agreement with experimental results for flow depth and discharge
59 partitioning in urban networks (Dewals et al., 2023; Liang et al., 2024; Luo et al., 2018).
60 However, the depth-averaged formulation restricts the ability of 2D models to resolve vertical
61 flow structures and complex three-dimensional flow patterns (Haider et al., 2024; Li et al.,
62 2024). Similar limitations associated with simplified flow representations have also been
63 reported in studies comparing reduced-dimensional and three-dimensional simulations. For
64 example, Luo et al. (2018) showed in a comparative 1D–3D numerical study of open-channel
65 junction flows that simplified models tend to underestimate energy losses and secondary
66 circulations, highlighting the importance of resolving three-dimensional flow structures in
67 junction-dominated flows. Likewise, Ouro et al. (2024) compared several computational
68 hydraulics models ranging from depth-averaged shallow-water formulations to fully three-
69 dimensional approaches. They reported that three-dimensional models generally provide more
70 accurate predictions of mean velocities and flow dynamics in open-channel flows with lateral
71 cavities.

72 To address the limitations of conventional depth-averaged models in representing vertical flow
73 structures and complex three-dimensional flow patterns, several improved depth-integrated
74 models have been proposed to incorporate key three-dimensional flow effects into two-
75 dimensional frameworks. For example, Blanckaert et al. (2003) developed a nonlinear model
76 to account for the redistribution of mean flow and secondary currents in curved open-channel
77 flows. Uchida et al. (2019) further proposed a two-phase depth-integrated approach capable of
78 reproducing non-hydrostatic effects. More recently, Castro-Orgaz et al. (2025) and Lugina et
79 al. (2026) proposed enhanced depth-integrated formulations based on additional equations of
80 motion accounting for non-uniform velocity vertical profile and non-hydrostatic pressure
81 distribution, which are excluded in the conventional 2D model. These developments illustrate
82 the ongoing effort to bridge the gap between computationally efficient depth-averaged models
83 and fully three-dimensional simulations.

84 In recent years, refined three-dimensional computational fluid dynamics (CFD) models have
85 been increasingly applied to urban flood simulations. For instance, Zhu et al. (2023) analyzed
86 the influence of urban neighborhood layouts on flow intrusion into residential buildings and on
87 pedestrian safety, showing that 3D simulations can effectively capture flood exchanges between
88 buildings and streets under various urban configurations. In urban flooding, strong three-
89 dimensional flow separation, reattachment, and vortex formation occur around buildings and
90 street junctions, and 3D numerical simulations enable a more realistic representation of these
91 hydraulic structures and improve the understanding of local flow dynamics (Finaud-Guyot et
92 al., 2019; Li et al., 2025). Li et al. (2024) compared a depth-averaged $k-\varepsilon$ 2D model with a 3D
93 Reynolds-Averaged Navier–Stokes (RANS) model using the $k-\omega$ SST closure in a study of
94 urban flood modelling with different urban layouts, but no flow intrusion within buildings nor
95 urban blocks. Their results indicated that the 2D model can provide reasonable predictions of
96 overall flood dynamics and water levels, making it suitable for large-scale flood propagation
97 analyses. However, the 3D model reproduced more accurately jet trajectories and complex
98 recirculation patterns that cannot be resolved by depth-averaged approaches. Li et al. (2025)
99 further extended this modelling comparison to configurations involving exchanges between
100 flooded streets and building interiors. Their results showed that the 3D model reproduces more
101 accurately the shape and size of recirculation zones and captures vertical velocity variations
102 associated with jet flows, while the 2D depth-averaged model exhibited systematic limitations.
103 In particular, the 2D model failed to reproduce the water level fluctuations observed in the
104 experiments, showed larger errors in the prediction of discharge partitioning among outlets, and
105 produced distorted recirculation structures compared with the measurements. These results
106 highlight the limitations of depth-averaged approaches for representing complex three-
107 dimensional flow structures in urban flooding. Moreover, these results imply that while depth-
108 averaged models may be efficient for large-scale flood propagation, they can be inadequate
109 when complex three-dimensional flow structures, such as jets, recirculations, and vertical
110 velocity variations, play an important role. Viccione et al. (2022) used a 3D CFD model based
111 on RANS equations coupled with a $k-\varepsilon$ turbulence model to simulate flood–building

112 interactions, showing that local flow kinematics strongly affect flood forces acting on building
113 facades and demonstrating the capability of 3D models to assess these effects. Dellinger et al.
114 (2025) evaluated 2D models (with and without the $k-\varepsilon$ closure) and a 3D $k-\omega$ SST model
115 against experimental data for urban flood flows in a realistic street network. They showed that
116 the 3D model effectively reproduced the recirculation behind intersections. Recently, Saleem
117 et al. (2025) simulated dam-break-induced urban flooding using a 3D CFD model, revealing
118 that while grid refinement reduces numerical errors, the selection of an appropriate turbulence
119 closure is the primary factor in capturing high-inertia flow dynamics. Their sensitivity analysis
120 of RANS-based models demonstrates that the realizable $k-\varepsilon$ variant significantly outperforms
121 the standard $k-\varepsilon$ and other closures, achieving high predictive accuracy ($NSE > 0.75$) and
122 superior computational efficiency in complex urban geometries.

123 These studies demonstrate the value of 3D turbulence models for urban flood simulations.
124 However, the reliability of such simulations depends strongly on the turbulence closure used,
125 as previous studies have shown that different RANS turbulence models can lead to significant
126 differences in predicted flow fields. For example, Farhadi et al. (2018) compared the
127 performance of the standard $k-\varepsilon$, the classical $k-\omega$ and an improved $k-\omega$ model incorporating a
128 cross-diffusion term and a modified stress limiter in three-dimensional racetrack flume
129 experiments, showing that these turbulence closure methods did not exhibit significant
130 differences in capturing the overall flow structures. Shaheed et al. (2019) compared the
131 performance of the standard $k-\varepsilon$ model and the realizable $k-\varepsilon$ model in curved and confluent
132 channels, showing that the standard $k-\varepsilon$ model captures better velocity profiles in curved
133 channels, while the realizable $k-\varepsilon$ model was more accurate in confluent channels. This
134 indicates that the performance of turbulence closures can depend strongly on the flow geometry.

135 Table 1 provides a summary of previous numerical studies employing turbulence closure
136 models to investigate the detailed flow characteristics of urban environments and selected
137 examples of open-channel flows. The table lists the type of numerical model (2D or 3D), the
138 turbulence closure adopted, the model configuration, and whether the study considered urban

139 flooding or included a comparative analysis. This overview highlights the limited number of
140 comparative studies focusing on urban flooding, particularly regarding flow intrusion into
141 buildings or urban blocks.

142 Overall, several studies based on 3D RANS turbulence models have produced insightful results
143 that have advanced our understanding of urban flooding processes. Likewise, comparative
144 analyses have been conducted to evaluate different turbulence closures for open-channel flow
145 conditions. However, evaluations focusing specifically on turbulence closures for urban
146 flooding, particularly those involving flow exchanges between building interiors and the
147 surrounding streets, remain limited. Thus, information that can help selecting turbulence
148 models in such configurations is scarce.

149 The present study aims to evaluate the performance of turbulence closure models in 3D RANS
150 simulations of urban flooding within an idealized district with flow intrusion into buildings. By
151 systematically comparing the standard k - ϵ , k - ω , and k - ω SST models, the study seeks to clarify
152 their relative strengths and limitations in predicting urban flow variables. Special attention is
153 paid to the representation of near-wall effects and flow mixing at bifurcations and junctions,
154 which are critical to capturing flow patterns in urban environments and assessing the flood
155 hazard. The work of Mejía-Morales et al. (2021) offers a valuable reference, as their
156 experiments on an idealized urban district with different openings around an urban block
157 provide representative configurations for model validation. Building upon recent research
158 (Dewals et al., 2023; Li et al., 2025), this study further analyses the ability of three turbulence
159 closures to reproduce complex flow patterns resulting from street-building interactions during
160 urban flooding, across nine geometric configurations.

161 Section 2 briefly introduces the experimental setup of Mejía-Morales et al. (2021), the tested
162 urban configurations, and the computational framework with the selected turbulence models.
163 Sections 3 and 4 present the simulation results and discussion, respectively, while Section 5
164 summarizes the main conclusions.

165 Table 1 Previous numerical research with turbulence closures on detailed flow characteristics in open
 166 channels/urban area.

Reference	Model (2D or 3D)	Turbulence Closure	Model Configurations	Urban Flooding Scenarios	Comparative Studies of Turbulence Closure
Luo et al. (2018)	3D	$k-\omega$ SST	Junction	No	No
Farhadi et al. (2018)	3D	$k-\varepsilon$ and $k-\omega$	Racetrack flumes	No	Yes
Shaheed et al. (2019)	3D	$k-\varepsilon$ and $k-\omega$	Curved and confluent channels	No	Yes
Viccione et al. (2022)	3D	$k-\varepsilon$	Urban basin	Yes	No
Zhu et al. (2023)	3D	$k-\varepsilon$	Grid of building blocks	Yes	No
Saleem et al. (2025)	3D	$k-\varepsilon$	Grid of building blocks	Yes	No
Li et al. (2024)	3D	$k-\omega$ SST	District	Yes	No
Li et al. (2025)	3D	$k-\omega$ SST	District	Yes	No
Dellinger et al. (2025)	2D and 3D	$k-\omega$ SST	District	Yes	No
Current study	3D	$k-\varepsilon$, $k-\omega$ and $k-\omega$ SST	District	Yes	Yes

167 2 Methodology

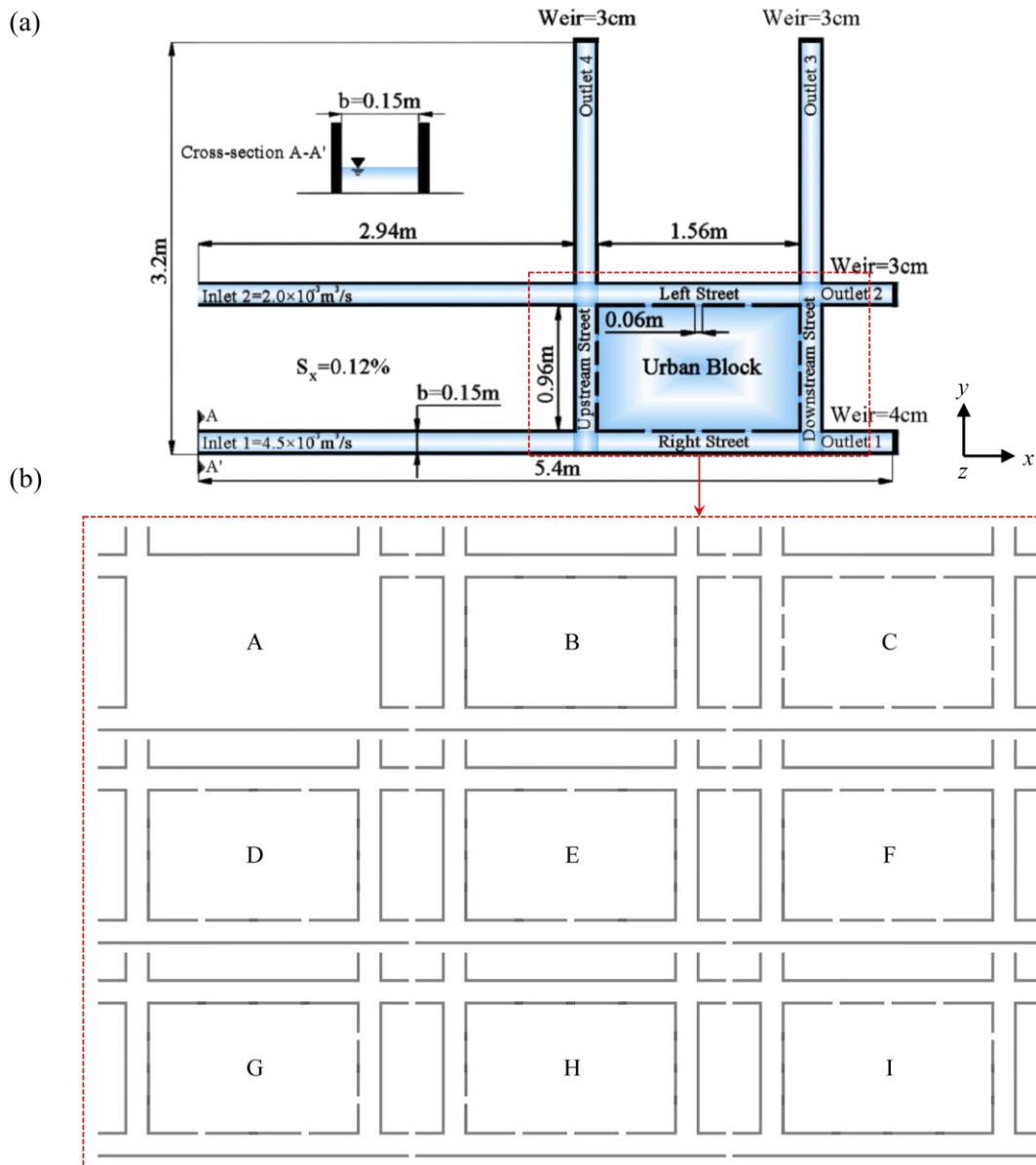
168 2.1 Experimental setup

169 This study is based on experimental research on urban flooding at a city block scale conducted
 170 by Mejía-Morales et al. (2021). The physical model represents a rectangular urban block
 171 surrounded by four streets as depicted in Figure 1(a). Openings in the central block enable flow
 172 exchanges between the streets and the interior of the block, with the aim of representing the
 173 intrusion of water into buildings through open or broken doors during an urban flood. All streets
 174 had a uniform rectangular cross-section with a width of 0.15 m. The x -direction had a slope of

175 0.12%, while the y -direction remained flat. The dimensions of the central urban block were
176 $1.56 \text{ m} \times 0.96 \text{ m}$. The central block was delimited by walls with a thickness of 2 cm and a height
177 of 15 cm.

178 In the experimental model, two inlets, Inlet 1 and Inlet 2, were positioned at the upstream end
179 of the x -direction. Inlet 1 had a consistent inflow rate of $4.5 \times 10^{-3} \text{ m}^3/\text{s}$, while the Inlet 2 flow
180 rate was set at $2 \times 10^{-3} \text{ m}^3/\text{s}$. Four outlets were placed at the ends of the four street sections,
181 with vertical weirs enabling adjustment of the water depth. Outlet 1 featured a weir measuring
182 4 cm in height, whereas the remaining three outlets were equipped with weirs measuring 3 cm
183 in height. The measurement methods are explained in detail in Mejía-Morales et al. (2021).

184 Experimental measurements of key hydraulic variables, including water depth, flow velocity,
185 and discharge, were conducted for a series of urban-block configurations with different
186 numbers and positions of openings Mejía-Morales (2022). From these configurations, nine
187 representative cases were selected for the present study (Figure 1b). For clarity and simplicity,
188 these configurations are renamed sequentially as configurations A, B, C, D, E, F, G, H, and I.



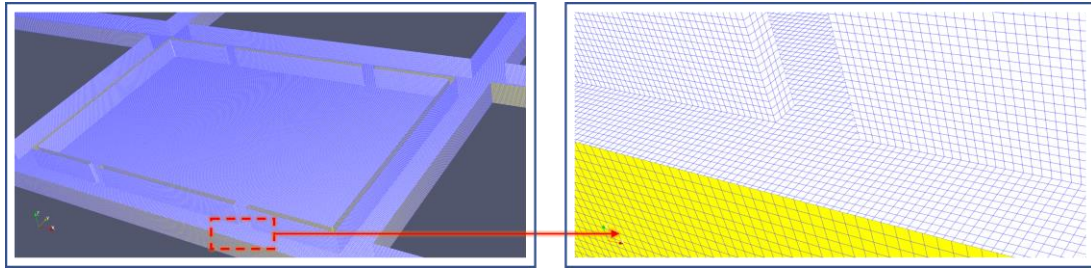
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190 Figure 1 Details of the modeling approach. (a) Top view of the experimental setup of Mejía-Morales
 191 et al. (2021). (b) The nine block configurations selected for numerical modeling, obtained by closing
 192 different combinations of openings.

193 2.2 3D Numerical modelling

194 Model description

195 The experimental model was reproduced by solving numerically the 3D incompressible RANS
 196 equations in OpenFOAM (v9). The adopted solver is interFoam, which is primarily used for
 197 two-phase incompressible and turbulent flows. A structured mesh is chosen to build the
 198 computational domain (Figure 2).



199

200 Figure 2 Model grid structure (using Configuration D as an example).

201 The free surface was captured with the volume of fluid (VOF) method introduced by Hirt et al.
 202 (1981). This method is computationally efficient, it is one of the most used methods to predict
 203 free surface flows and has been extensively validated (Torres et al., 2022). Cells containing
 204 only water or only air are assigned values of $\alpha_{\text{water}} = 1$ and $\alpha_{\text{water}} = 0$, respectively. The free
 205 surface is identified at locations where $\alpha_{\text{water}} = 0.5$. A Courant-Friedrichs-Lewy (i.e., CFL)
 206 restriction of 1 was imposed, which led to an adaptive time step of the order of 0.002 s.

207 In this study, three two-equation Reynolds-averaged turbulence closures available in
 208 OpenFOAM, namely the standard $k-\varepsilon$ (Farhadi et al., 2018), $k-\omega$ (Wilcox, 2008), and $k-\omega$ SST
 209 models (Li et al., 2025), were selected for the simulations. These models were considered
 210 appropriate because they are widely used in hydraulic engineering and free-surface turbulent
 211 flow simulations, while remaining computationally affordable for three-dimensional
 212 computations involving multiple configurations. Compared with higher-fidelity approaches
 213 such as large-eddy simulation (LES), they offer a more practical balance between
 214 computational cost and predictive capability. In addition, these three turbulence closures are
 215 among the most used models in engineering CFD and therefore provide a representative basis
 216 for assessing the influence of turbulence-model selection on simulated urban flood structures.
 217 The Wilcox $k-\omega$ formulation and its later developments have also played an important role in
 218 the evolution of widely used RANS turbulence closures. In the following, these models are
 219 referred to as $k-\varepsilon$, $k-\omega$, and $k-\omega$ SST, respectively.

220 *Boundary conditions*

221 The models require consideration of three distinct boundary condition types: inlet, outlet, and
222 solid wall boundaries (Figure 1a). At the beginning of each simulation, the water inside the
223 domain was initialized as a stagnant layer with a uniform depth of 5 cm.

224 a) Inlet

225 The inlet boundary conditions are prescribed to reproduce the experimental inflow conditions.
226 The inflow section corresponds to the experimental inlet configuration, with an opening height
227 of 4 cm. The velocity at the inlet is specified to match the experimental discharge, ensuring that
228 the inflow rate in the simulations is consistent with the experimental conditions. This
229 configuration reproduces the subcritical flow regime observed in the experiments. Turbulence
230 quantities at the inlet are prescribed as constant values following standard RANS practice,
231 including the turbulent kinetic energy k and the second turbulence variable required by each
232 turbulence model, i.e., the turbulent kinetic energy dissipation rate, ε , or the specific dissipation
233 rate, ω .

234 b) Outlet

235 A free-outflow boundary condition is prescribed at the outlet. It includes zero-gradient
236 treatment for the water phase fraction and constant atmospheric pressure for the free surface.

237 c) Solid wall boundaries

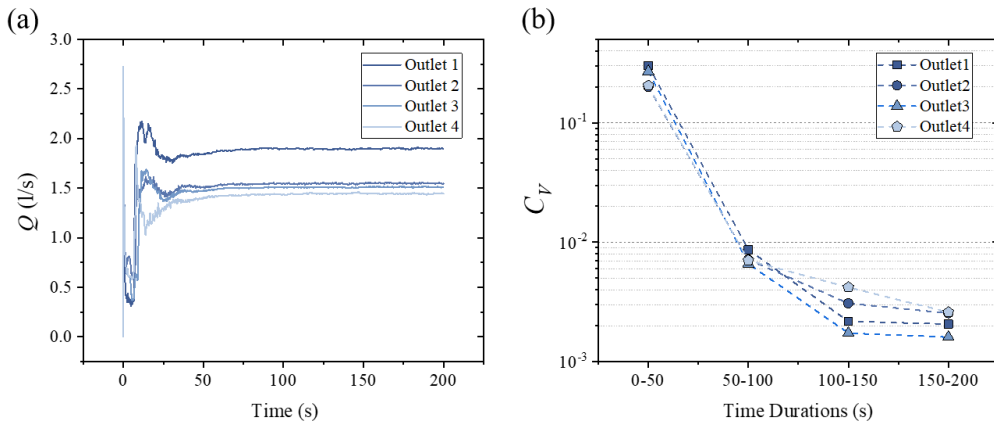
238 At solid boundaries, all three turbulence models use wall-function approaches to treat turbulent
239 kinetic energy, turbulent viscosity, and dissipation rates. These treatments provide constraints
240 on turbulence production and dissipation at walls, which are essential for accurately
241 representing near-wall turbulence in both low- and high-Reynolds number flows. Because the
242 laboratory experiment was built with smooth plexiglasses, the walls are treated as hydraulically
243 smooth in the simulations.

244 *Simulation convergence*

245 To evaluate the convergence of the numerical models, a trial simulation was conducted using
246 the $k-\omega$ SST model with a mesh size of 5 mm for Configuration H. The hydrographs at the four
247 outlets and Coefficient of variation (C_V) of outflow discharge at different time periods are
248 shown in Figure 3 (b). C_V here was calculated as the ratio of the standard deviation to the mean
249 value of the each outlet discharge. It is used to describe the variability in time of the flow
250 variables. After about 50 s, the discharge exhibited quasi-steady behavior, with the value of
251 mean C_V dropping from 2.44×10^{-1} to 7.36×10^{-3} .

252 During the simulation period of 50 to 100 seconds, the C_V of discharge exhibits a steady small
253 value compared to the period of 0 to 50 seconds. Based on these observations, the time-averaged
254 values after 50 seconds (from 50 s to 100 s) are used in this study as the reference for steady-
255 state analysis.

256 Simulations ran in parallel using the public domain OpenMPI implementation of standard MPI
257 on the high-performance computing platform of Guangxi University, with the Sugon 7000A
258 supercomputer as the core.



259
260 Figure 3 (a) Outflow discharge hydrograph of numerical simulation. (b) Coefficient of variation (C_V) of
261 outflow discharge at different time periods (taking Configuration H with the $k-\omega$ SST turbulence model
262 and a mesh resolution of 5 mm as an example)

263 3 Results

264 3.1 Grid independence test

265 Mesh sensitivity tests are conducted for Configuration H using uniform hexahedral meshes with
266 three cell sizes of $d = 4$ mm, 5 mm, and 10 mm. All the other modelling parameters were kept
267 the same. The total number of mesh cells is 6.5×10^6 (4mm), 3.3×10^6 (5mm) and 0.4×10^6
268 (10mm).

269 Mesh independence is assessed by comparing both the discharge partitioning and the flow
270 patterns across the domain. The volumetric flow rate through the outlet patches is calculated
271 based on the VOF method as:

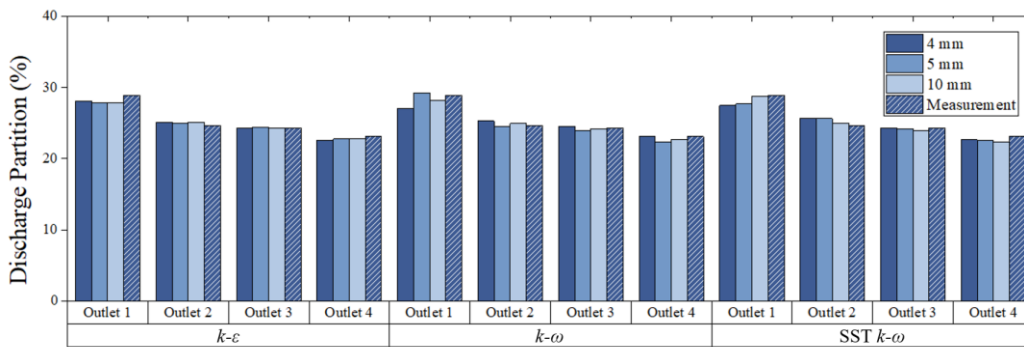
$$272 \quad \dot{V}_{water} = \sum_{f \in surface} \phi_f \cdot \alpha_{water,f} \quad (1)$$

273 where \dot{V}_{water} is the volumetric flow rate of water (m^3/s), ϕ_f is the volumetric flux across face
274 f (m^3/s), and $\alpha_{water,f}$ is the volume fraction of water at face f .

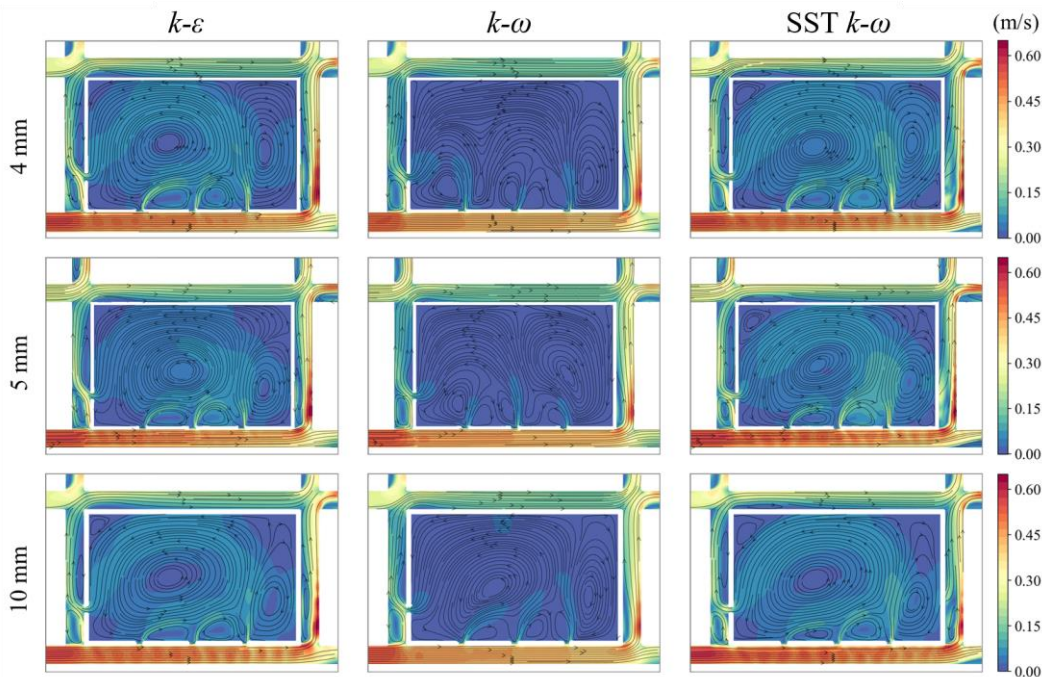
275 Figure 4 presents the discharge partition at the four outlets and under the three mesh resolutions
276 for each turbulence model. The results indicate good mesh convergence for all three turbulence
277 models. Variations in discharge distribution caused by changes in mesh resolution are generally
278 small, with average differences below 1.5%. The largest difference occurs at Outlet 1 between
279 the 4 mm and 5 mm meshes using the $k-\omega$ model (2.16%).

280 The influence of mesh resolution on simulated surface flow patterns is shown in Figure 5. For
281 the $k-\varepsilon$ and $k-\omega$ SST models, surface flow characteristics such as velocity field, flow
282 contractions and separations, and recirculation locations appear not significantly affected by
283 mesh resolution. For the $k-\omega$ model, flow velocity and overall patterns are also relatively
284 insensitive to mesh resolution in most regions. However, the flow pattern inside the block varies
285 substantially. Velocity magnitudes predicted by the $k-\omega$ model under all three mesh resolutions
286 are lower than those from the other two models. Moreover, with a 4 mm mesh, large
287 recirculations inside the building barely form, whereas in the 5 mm and 10 mm cases,
288 recirculations develop clearly. This indicates that the $k-\omega$ model's prediction of internal

289 building flow is more sensitive to mesh resolution than the other two models. This behavior is
 290 consistent with the known characteristics of the classical $k-\omega$ formulation, which can be more
 291 responsive to changes away from solid boundaries (Menter, 1994) and may therefore contribute
 292 to the variations observed here. Considering the mesh sensitivity results and computational
 293 efficiency, a mesh size of 5 mm is chosen to be used for all subsequent simulations.



294
 295 Figure 4 Outflow discharge distribution at the four outlets for measurements and numerical modeling
 296 with different mesh resolutions (4 mm, 5 mm, and 10 mm) for three turbulence closure models. The
 297 values represent the time-averaged proportion of each outlet discharge relative to the total outflow, based
 298 on the converged results from each simulation.



299
 300 Figure 5 Surface flow velocity patterns from three turbulence models at different mesh resolutions. The
 301 figures show the surface layer data, with streamlines overlaid on velocity contour maps, with Figure 9
 302 providing the measured surface flow pattern.

303 3.2 Discharge partition

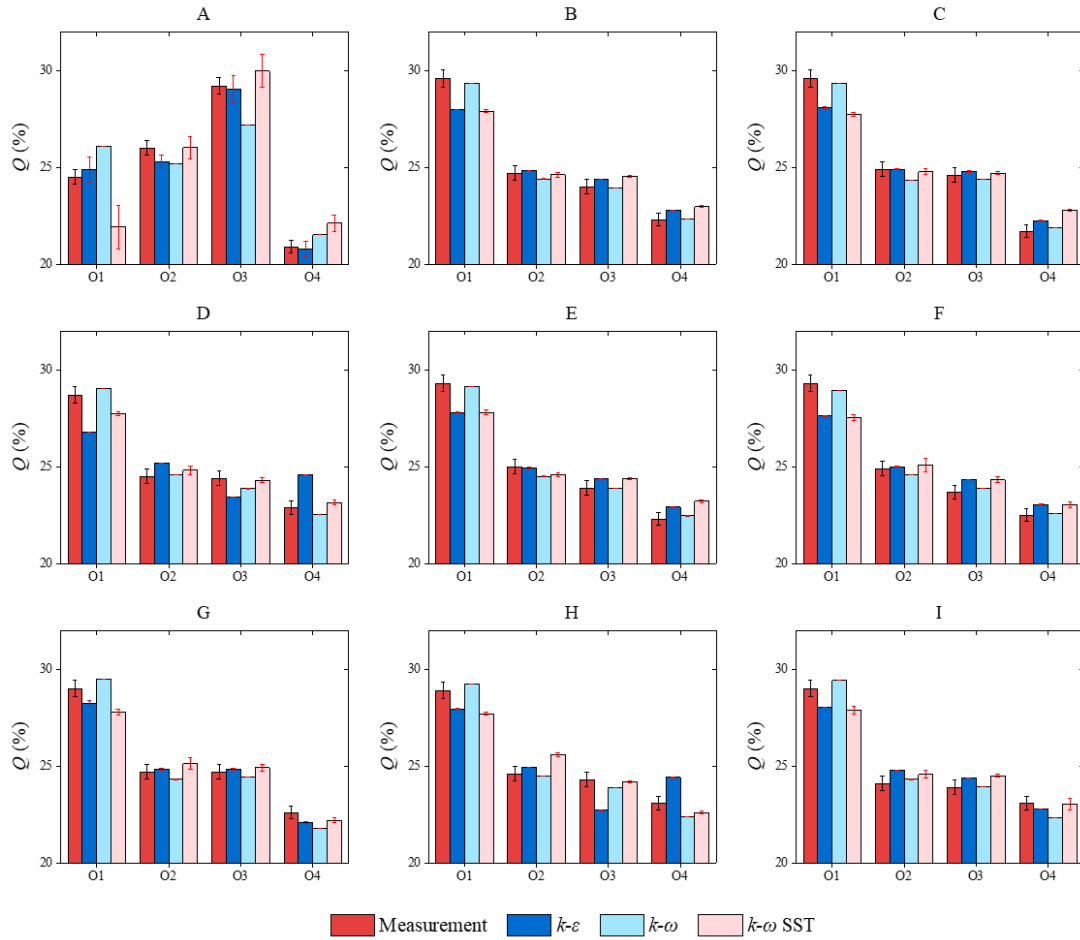
304 Figure 6 shows the discharge partition in the experiments and simulations. The error bars
305 represent the experimental measurement uncertainty (1.5% according to Mejía-Morales et al.,
306 2021) and the standard deviation of the simulated discharge ratios after convergence.

307 In Configurations B–I, Outlet 1 consistently receives the largest share of the total outflow,
308 Outlets 2 and 3 have similar proportions, and Outlet 4 receives the least. All simulations capture
309 these trends reasonably well, although the $k-\varepsilon$ model shows discrepancies in Configurations D
310 and H. Indeed, as shown in Figure 7, most simulated discharge fractions fall within the $\pm 1.5\%$
311 experimental uncertainty when compared to the respective measurements. For Configurations
312 B–I, the $k-\omega$ model matches the experimental measurements most closely, while the $k-\varepsilon$ and $k-$
313 ω SST models consistently underestimate the discharge in Outlet 1.

314 Configuration A, by contrast, behaves differently due to the absence of guiding streets. In this
315 configuration, the $k-\omega$ model does not reproduce the correct ranking, and both the $k-\varepsilon$ and $k-$
316 ω SST models show a larger standard deviation than $k-\omega$, with particularly pronounced
317 variations at Outlet 1 for the $k-\omega$ SST model. In general, the $k-\varepsilon$ and $k-\omega$ SST models exhibit
318 larger temporal variations than the $k-\omega$ model across Configurations B–I, and these differences
319 become more pronounced in Configuration A.

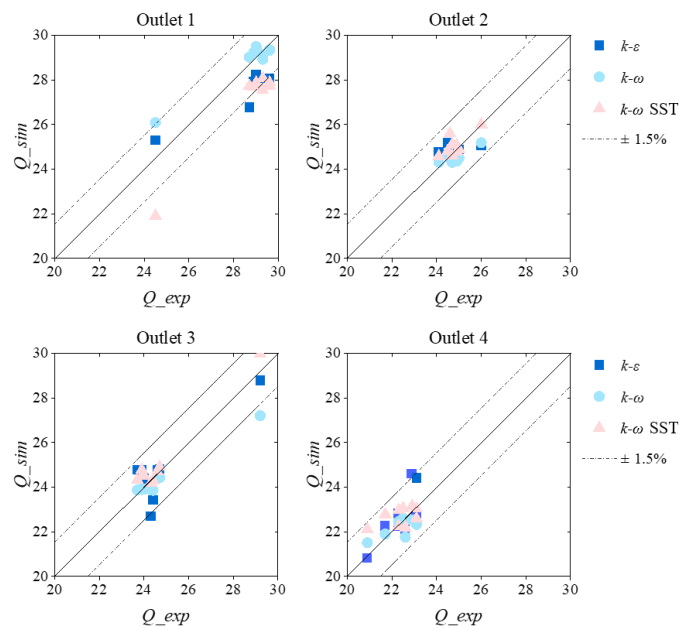
320 After convergence of the simulations, the standard deviation of the time series of computed
321 discharge fractions ranges from 0.001% to 1.149%, with the $k-\omega$ SST model generally
322 associated with the highest level of temporal variability among the tested closures.

323 In several configurations, the root means square errors (RMSE) (Figure 8) indicate that the $k-$
324 ω model generally performs particularly well, with RMSE values below 0.5 in all cases except
325 Configuration A. In comparison, the $k-\varepsilon$ and $k-\omega$ SST models remain valuable alternatives,
326 with the $k-\varepsilon$ model performing better than $k-\omega$ SST in an open geometric configuration such
327 as Configuration A.



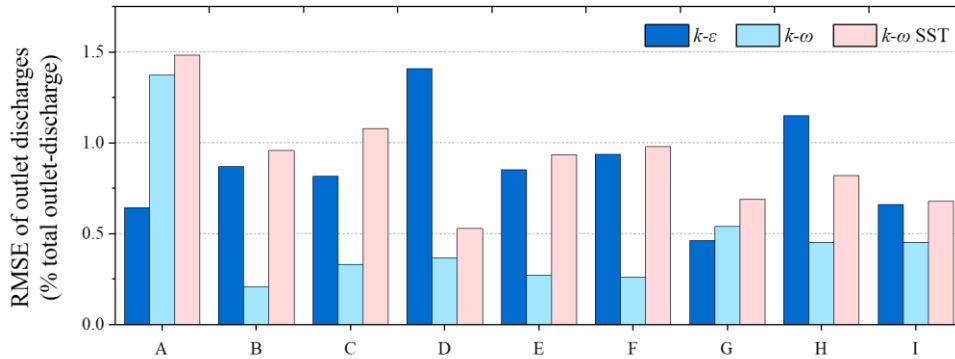
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329 Figure 6 Measured and modelled discharge partitions by the three turbulence models at the four outlets
 330 for all configurations (A - I). O1–O4 denote Outlets 1 to 4. The error bars represent the measurement
 331 uncertainty associated with the experimental instrumentation (Mejía Morales, 2022) and sample standard
 332 deviation for simulations.



333

334 Figure 7 Comparison of discharge partition at the four outlets between physical experiments (x-axis) and
 335 3D CFD simulations (y-axis) for the three tested turbulence closures and all the studied cases (Figure
 336 1b). Each subplot corresponds to one outlet, showing the percentage of total outflow through that outlet.



337 Figure 8 RMSE between three turbulence models simulations and measured outlet discharges at the four
 338 outlets. The RMSE is calculated based on the discharge portion in each outlet relative to the total outflow.
 339

340 3.3 Surface velocity

341 Flow patterns within the streets

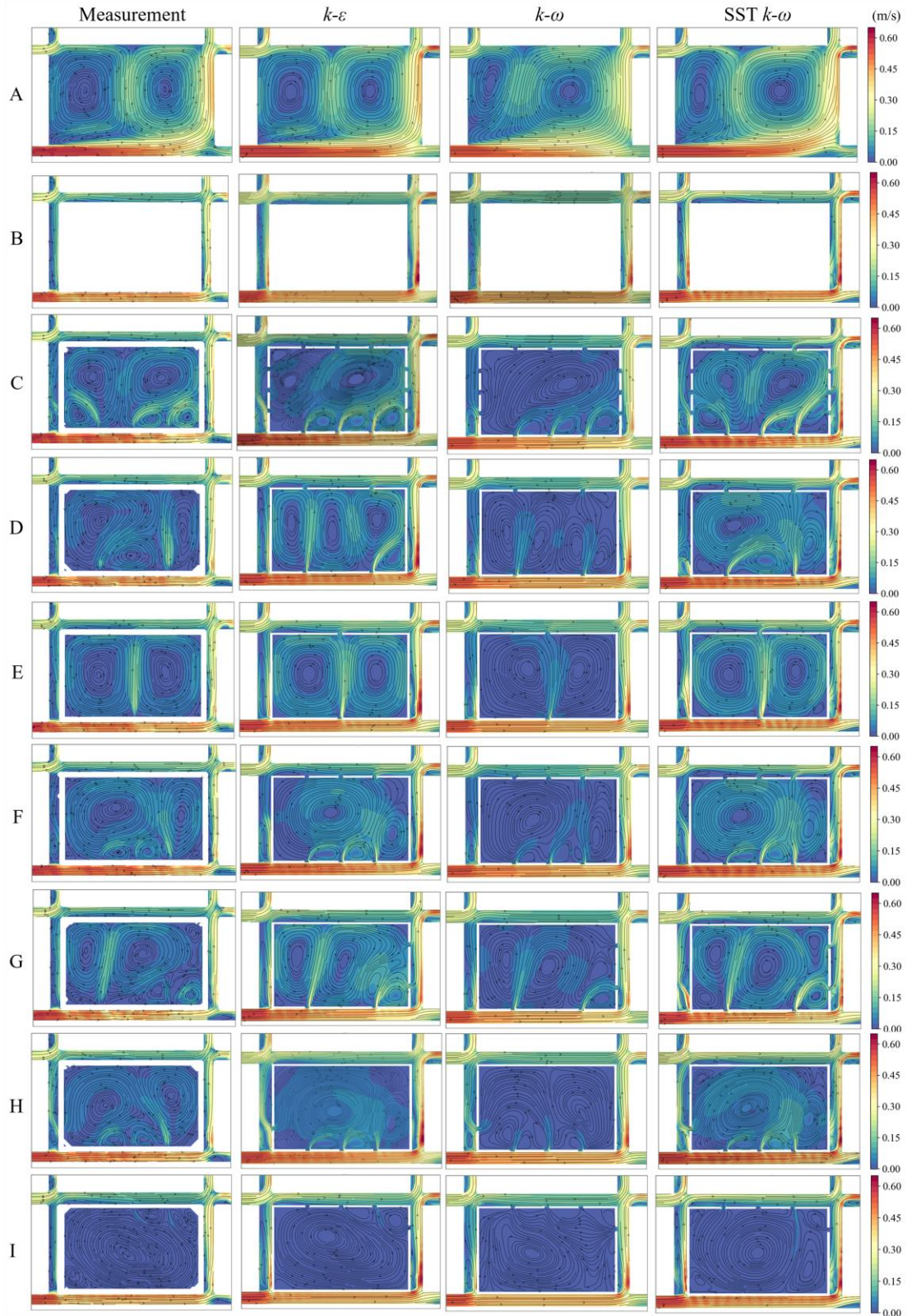
342 Figure 9 shows the time-averaged surface flow patterns. Results are presented for the
 343 measurements and for the three turbulence models. In the Right street of the physical model,
 344 all configurations exhibit a general decrease in flow velocity from upstream to downstream.
 345 This trend is also reproduced in numerical simulations. The $k-\epsilon$ and $k-\omega$ SST models
 346 successfully capture the superimposed along-stream fluctuations observed in the measurements,
 347 whereas the $k-\omega$ model tends to under-represent these fluctuations, showing instead a more
 348 uniform, monotonic velocity decline. These differences are further illustrated in Figure 10
 349 which presents the time-averaged surface velocity profiles along the centerline of the Right
 350 street for both the measurements and the three turbulence models. Configuration A is not
 351 included, as it does not contain the Right street. The $k-\epsilon$ and $k-\omega$ SST models reproduce
 352 satisfactorily the general trend of the measured data. The magnitude of the fluctuations is
 353 generally higher in the $k-\omega$ SST results than in the $k-\epsilon$ results and is closer to the fluctuation
 354 level measured in the experiments, as seen in Configurations B, C, F, H, and I. In some cases,
 355 such as Configuration D, the two models show nearly identical fluctuation levels, while in

356 Configuration E, it is difficult to determine which model aligns more closely with the
357 measurements.

358 Although the $k-\omega$ model fails to reproduce these local fluctuations in the Right street in most
359 configurations, its overall trend in flow velocity variation remains reasonably consistent with
360 the measurements.

361 While Figure 10 highlights the model performance along the Right street, Figure 9 further
362 illustrates differences in the simulated surface flow patterns at street intersections, where all
363 three models successfully reproduce the characteristic surface flow patterns formed by flow
364 confluences and bifurcations. The recirculation zones along streets show distinct differences
365 between the models. For instance, in Configuration D, the experimental flow pattern inside the
366 Upstream street indicates that the recirculation generated from the inflow of Right street spans
367 near half the street length. The $k-\varepsilon$ model overestimates this zone, extending it to nearly the
368 entire street, while the $k-\omega$ model provides the most accurate approximation. In contrast, the $k-$
369 ω SST model generates a unique "8"-shaped recirculation zone, deviating significantly from
370 both the experimental data and the other models.

371 Similar discrepancies in the $k-\omega$ SST model are observed in upstream streets at configurations
372 D, E, F, G, and H. These differences extend to Downstream streets, where the flow pattern in
373 the confluence simulated by the $k-\omega$ SST model diverges consistently from both the
374 experimental measurements and the outputs of the other two models. A detailed analysis of
375 these discrepancies is presented in the subsequent section on vertical flow structures.



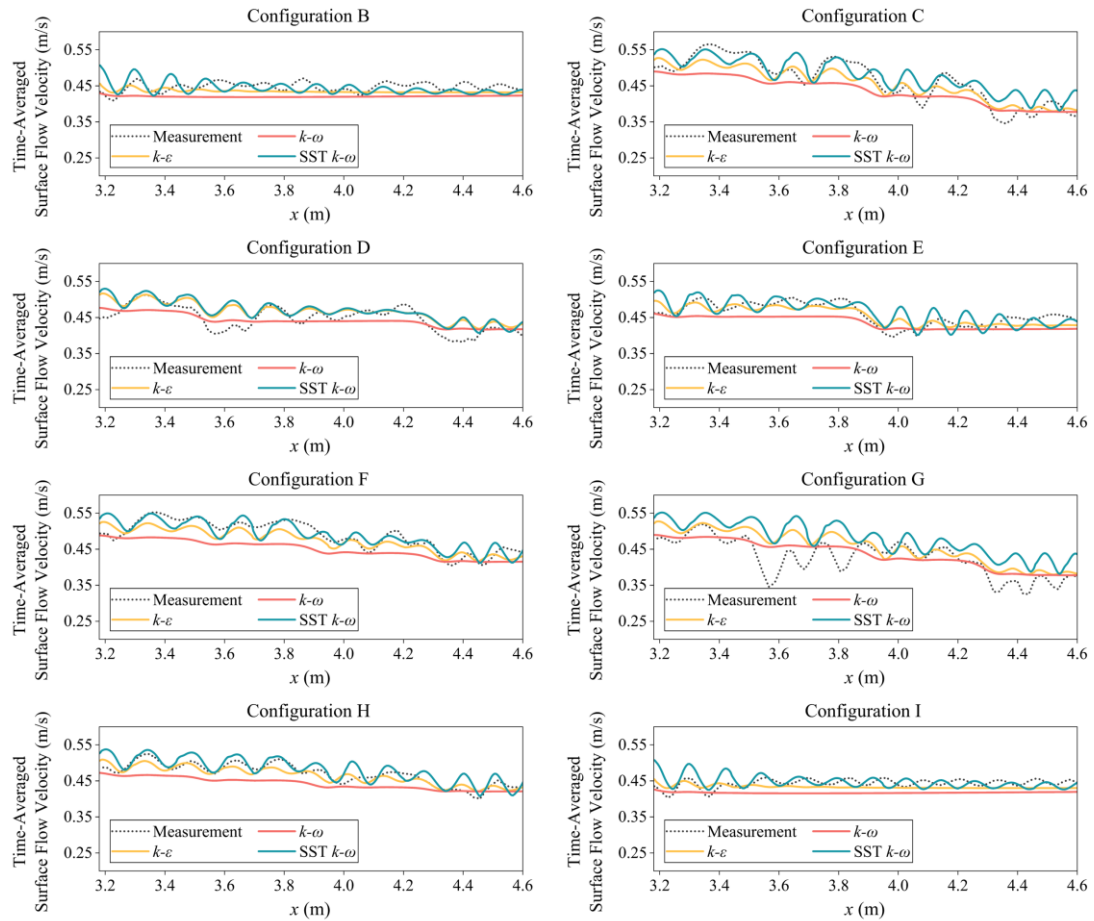
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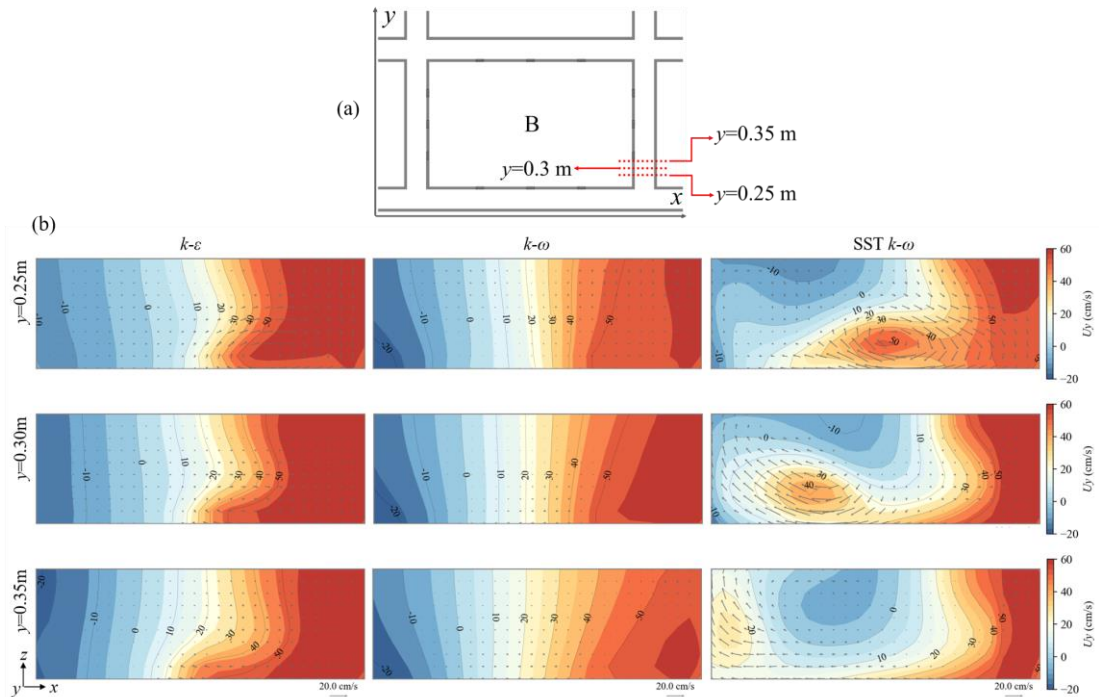
Figure 9 Time-averaged surface flow velocity patterns from measurements and those computed with the turbulence models $k-\epsilon$, $k-\omega$ and $k-\omega$ SST. The color scale represents the horizontal velocity magnitude (m/s).



380

381 Figure 10 Time-averaged surface flow velocity profiles from measurement and three turbulence models

382 along the centerline of Right street.



383

384 Figure 11 (a) The red dashed lines show cross-sections of $y = 0.25$ m, 0.30 m, and 0.35 m in the
385 Downstream street of Configuration B. (b) Contours of U_y overlaid with quiver plots of U_x and U_z at
386 cross-sections in three turbulence models.

387 Taking Configuration B as an example, Figure 11 compares the flow patterns predicted by the
388 three turbulence models at three cross-sections ($y = 0.25$ m, 0.30 m, and 0.35 m) within the
389 recirculation zone in the Downstream street formed by flow separation at the junction with the
390 Right street. It should be noted that the experimental flow field data for these cross-sections
391 were not recorded. In all three numerical models, each cross-section shows a flow pattern with
392 positive U_y near the downstream side, where peak velocities exceed 60 cm/s, and negative U_y
393 appear near the upstream side, indicating a recirculating flow structure in the vertical flow
394 pattern. However, the detailed velocity structures vary considerably among the models. The
395 flow field predicted by the $k-\omega$ model is structurally the simplest, featuring a clear division with
396 positive U_y near the right side and negative U_y near the left side of the sections, and no obvious
397 velocity layering along the z direction, indicating relatively uniform velocity distribution from
398 the surface to the bottom. In addition, the velocity gradient is more pronounced near the bed
399 and becomes weaker toward the free surface.

400 In terms of U_y contours, the $k-\varepsilon$ model exhibits a moderately complex flow structure. In the
401 upper half near the free surface, U_y and its gradient are similar to those predicted by the $k-\omega$
402 model. In the lower half near the bed, the high-velocity region adjacent to the downstream side
403 shows a clear tendency to extend toward the upstream side, resulting in velocity variations in
404 the vertical direction. This difference becomes even more pronounced in the simulation results
405 of the $k-\omega$ SST model. At the cross-section located at $y = 0.25$ m, the downstream region
406 displays a high-velocity positive U_y zone. Near the bottom center of the section, an elliptical
407 high-velocity region accompanied by a recirculation zone (as indicated by the U_x and U_z quiver
408 plot) is observed. As the flow progresses to the cross-sections at $y = 0.30$ m and 0.35 m, this
409 high-speed zone shifts upstream, its magnitude gradually decreases, and the recirculation
410 becomes weaker until it approaches the wall. Meanwhile, the area of reverse flow also shifts,
411 moving from the upstream wall toward the free surface at the center of the section.

412 All three models exhibit significant spatial variability in the vertical distribution of flow
413 velocity. The structure of the velocity variation differs markedly among the models, which is
414 closely related to the distinct approaches each turbulence model employs in resolving near-wall
415 flows. As shown in the velocity profiles, the $k-\varepsilon$ model shows a wider near-wall high-speed
416 region than the $k-\omega$ and $k-\omega$ SST models. In terms of the velocity gradient along the x -direction,
417 however, the $k-\varepsilon$ and $k-\omega$ models show a more comparable trend, both exhibiting a nearly linear
418 variation.

419 *Flow patterns within the urban block and open spaces*

420 In addition to the flow features observed in the streets, Figure 9 also shows the flow patterns
421 inside the model buildings. Table 2 summarizes the number and relative size of recirculation
422 zones predicted by the different turbulence models. This qualitative comparison helps identify
423 systematic differences in the flow structures reproduced by the turbulence closures in the
424 idealized domain. In real cases, where the scale of the problem is different and the urban layout
425 is considerably more complex, additional uncertainties may occur that are related to
426 parameterizations for roughness or drag coefficients.

427 Configuration A features an open-space layout, where the simulated velocity distributions are
428 generally close to the experimental observations. The $k-\varepsilon$ model reproduced two large and
429 nearly symmetric recirculations that closely matched the measurements. In contrast, the $k-\omega$
430 and $k-\omega$ SST models produced two asymmetric recirculations, with noticeable deviations in
431 their shape and center position compared with the experiments.

432 The $k-\varepsilon$ and $k-\omega$ SST models produced more accurate results, showing better agreement with
433 the measured data for the size and location of large recirculations in most configurations.
434 However, in some configurations, one model performed better than the other. For example, in
435 Configuration D, the $k-\varepsilon$ model produced significant errors, while the $k-\omega$ SST model performed
436 well. Conversely, in Configuration I, the $k-\varepsilon$ model yielded a better simulation, while the $k-\omega$
437 SST model generated slightly less accurate results.

438 Further differences emerged when simulating cases with smaller recirculations and local flow
439 velocities. The $k-\omega$ SST model provided a more accurate representation of recirculation velocity
440 in Configuration C, reproduced the right-side recirculation in Configuration F more accurately
441 both in terms of size and extent, and captured the position of the small recirculation in
442 Configuration G more effectively than the $k-\varepsilon$. In Configuration H, the results of the $k-\varepsilon$ and
443 the $k-\omega$ SST were almost identical, although both deviated substantially from the experimental
444 observations.

445 The flow patterns predicted by the $k-\omega$ turbulence model showed noticeable discrepancies
446 compared to the measurements, primarily due to its systematic underestimation of flow velocity.
447 This underprediction not only reduced the strength of recirculation but also caused deviations
448 in their location and size. In Configurations C, D, H, and I, the $k-\omega$ model results exhibited
449 clear differences from the measurements in both velocity magnitude and flow structures,
450 particularly in the position and extent of the recirculation centers. In Configurations E, F, and
451 G, the $k-\omega$ model simulations reproduced the overall positioning of the recirculation reasonably
452 well, although the predicted velocities inside the buildings remained more than 50% lower than
453 the measured values on average.

454 Table 2 Qualitative appraisal of the agreement between observed and computed flow patterns within the
 455 block interior (Configurations B - H) and in a large open space (Configuration A). Notation “L” stands
 456 for “large recirculation(s)”, “S” for “small recirculation(s)” and “s” for “very small recirculation(s)”.

Configuration	Measurement	$k-\varepsilon$	$k-\omega$	$k-\omega$ SST
A	2L, symmetrically	2L, symmetrically	2L, asymmetrically	2L, asymmetrically
B	-	-	-	-
C	2L+3S	2L+3L+1s, inaccurately	1L+3S, inaccurately	2L+3S
D	3L	3L+1S, inaccurately	2L+2S, inaccurately	3L
E	2L, symmetrically	2L, symmetrically	2L, symmetrically	2L, symmetrically
F	2L+2S	2L+2S	2L+2S	2L+2S
G	2L+2S+3s	2L+2L	2L+2S	2L+3S+1s
H	2L+2S	2L+2S, incorrect relative sizes	2L+3S	2L+2S+1s, incorrect relative sizes
I	1L+4S	1L+2S	1L	1L+3S

457 **4 Discussion**

458 Evaluating the performance of turbulence models in urban flood simulations requires balancing
 459 modeling accuracy and computational efficiency, with flow patterns serving as a key indicator
 460 of model fidelity. Under identical computational settings (grid resolution of 5 mm, about 3.3
 461 million cells, and parallel computation with 24 cores on an Intel i9-13900 processor with DDR5
 462 memory), the $k-\omega$, $k-\omega$ SST, and $k-\varepsilon$ models completed the 200-second simulation in
 463 approximately 212, 249, and 295 hours, respectively. Taking the average runtime of 252 hours
 464 as a reference, the $k-\omega$ model was about 16% faster, while the $k-\varepsilon$ model required about 17%
 465 longer computational time.

466 Beyond runtime, numerical oscillation also influences model efficiency. An analysis of the
 467 coefficient of variation of outlet discharges revealed some systematic differences among the
 468 turbulence models, with the $k-\omega$ model showing the smallest fluctuations, followed by the $k-\varepsilon$

469 model, and the $k-\omega$ SST model exhibiting slightly higher variability in most cases (eight out of
470 nine cases). In Case A, the $k-\omega$ SST model presented a relatively larger variation, with a
471 standard deviation of approximately 1.14% in the discharge partition.

472 It is worth noting that these numerical fluctuations cannot be directly compared with the
473 experimental results, because the experimental data used here represent time-averaged
474 conditions rather than instantaneous flow variations. Therefore, the observed differences in
475 numerical oscillations should be interpreted as indicators of computational robustness, not as
476 measures of physical unsteadiness. Nevertheless, video references provided by the
477 experimental dataset (Mejía Morales, 2022) reveal that the flow sometimes exhibits noticeable
478 temporal variability within the flow domain despite the fact that incoming flow is steady. In
479 particular, the flow pattern evolves over time, especially in the jet formed as the flow enters the
480 building through the opening (see Supplementary Figure S1). This jet-induced fluctuation was
481 also reproduced in the simulations: both the $k-\varepsilon$ and $k-\omega$ SST models captured similar
482 oscillatory behavior, whereas the $k-\omega$ model, due to its notably lower velocities inside the
483 building, failed to reproduce this feature clearly. Although this temporal variability has not been
484 quantitatively analyzed within the scope of the present study, its qualitative characteristics
485 provide a valuable reference, suggesting that the experimentally observed fluctuations and the
486 turbulence model's ability to capture such flow variability merit further investigation.

487 This analysis highlights that higher numerical resolution and sensitivity to transient flow
488 features may increase computational time and introduce more variability in simulated discharge.
489 The selection of an appropriate turbulence model should be guided by the objectives of the
490 study, considering whether the priority lies in computational efficiency or in capturing detailed
491 transient behaviors. Table 3 summarizes the main characteristics of each turbulence closure
492 model.

493 Table 3 Summary of the main strengths and limitations of the three turbulence models applied in this
 494 study.

Turbulence model	Computational time (hours)	Strengths	Limitations	Performance in narrow streets with surface velocity fluctuations
$k-\varepsilon$	295	Best for open, unobstructed flows.	Long runtime, less accurate in surface flow patterns at intersections	Captures trend, less pronounced fluctuations.
$k-\omega$	212	About 16% faster, accurate discharge	Underestimates interior flows	Smooth velocity decay, minimal fluctuations.
$k-\omega$ SST	249	Captures recirculation well	May generate spurious vortices	Consistent trend, closest match to measured fluctuations.

495 **5 Conclusions**

496 This study systematically evaluated the performance of three turbulence closure models: $k-\varepsilon$, $k-$
 497 ω , and $k-\omega$ SST, in simulating urban flooding through an idealized district with flow exchanges
 498 between streets and building openings.

499 Mesh sensitivity tests confirmed that the simulation results are generally mesh independent for
 500 all three turbulence models. Variations in discharge distribution across different mesh
 501 resolutions remained below 1.5%, indicating good numerical convergence. Surface flow
 502 structures, such as velocity fields and recirculation zones, were largely unaffected by mesh
 503 refinement, except for minor differences in recirculation formation within the building when
 504 using the $k-\omega$ model.

505 When compared to the respective measurements, most simulated discharge partitions at the
 506 different outlets fall within the $\pm 1.5\%$ experimental uncertainty. The $k-\varepsilon$ and $k-\omega$ SST models
 507 outperformed the $k-\omega$ model at capturing complex flow features within building interiors and
 508 street networks, including recirculation zones, as well as contraction and expansion of flow at
 509 street intersections. In narrow street segments where surface velocity fluctuations were
 510 observed, all models reproduced the general decreasing trend in flow velocity from upstream
 511 to downstream. However, the $k-\varepsilon$ and $k-\omega$ SST models more accurately represented

512 superimposed along-stream fluctuations. The $k-\omega$ SST model showed the closest agreement
513 with the measurements, while the $k-\omega$ model tends to smooth out these variations. This makes
514 the $k-\varepsilon$ and $k-\omega$ SST models more suitable for assessing flood risk in streets with complicated
515 flow patterns, and in large open spaces, such as squares and car parks.

516 Of the three models tested, the differences in computational efficiency remained at around 15%.
517 The $k-\omega$ model required the least computational time, followed by $k-\omega$ SST and standard $k-\varepsilon$.
518 No single model excelled in all aspects. The choice of turbulence model should be guided by
519 the specific objectives of the simulation, whether prioritizing discharge accuracy, flow velocity
520 details, or computational speed.

521 Future research should extend this work by incorporating more complex urban drainage
522 systems to explore the interactions between surface flooding and sewer networks, particularly
523 during surcharge conditions. Incorporating pollutant transport and contaminant intrusion into
524 buildings would also enhance the applicability of the model for urban risk assessment (Han et
525 al., 2025). Moreover, simulating the movement and accumulation of floating debris or drifters
526 could provide new insights into flood-induced hazards, such as the impacts on building facades
527 or other structural elements (Albano et al., 2016; Bayón et al., 2024). In addition, coupling
528 hydrodynamic simulations with pedestrian stability and exposure analyses would allow for
529 assessing risks faced by pedestrians during urban flood events (Arrighi et al., 2017; Zhu et al.,
530 2023). These extensions would contribute to a more comprehensive understanding of urban
531 flood dynamics and improve predictive capabilities for flood resilience planning.

532 **Acknowledgements**

533 This work was funded by the National Natural Science Foundation of China (No. 52309015)
534 and Guangxi Specific Research Project for Research Bases and Talents (No. AD23026176).
535 This is also funded by Guangxi Young Elite Scientist Sponsorship Program
536 (GXYESS2025022). The support from the French National Research Agency (ANR) for the
537 project DEUFI (ANR-18-CE01-0020) is also acknowledged

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