



Multiple Timescale Spectral Analysis in Wind Engineering

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Multiple Timescale Spectral Analysis

$$u''(t) + \varepsilon f(u, u'(t)) + u(t) = 0 \quad \varepsilon \ll 1$$

$$\varepsilon = 0 \rightarrow u(t) = R \cos(t + \varphi_0)$$

Perturbation methods

Slow timescale : $T = \varepsilon t$ Fast timescale : t

$$d_t \rightarrow \partial_t + \varepsilon \partial_T$$

$$u(t) \rightarrow u_0(t, T) + \varepsilon u_1(t, T) + \dots$$

Leading order solution :

$$u_0(t, T) = R(T) \cos(t + \varphi(T))$$

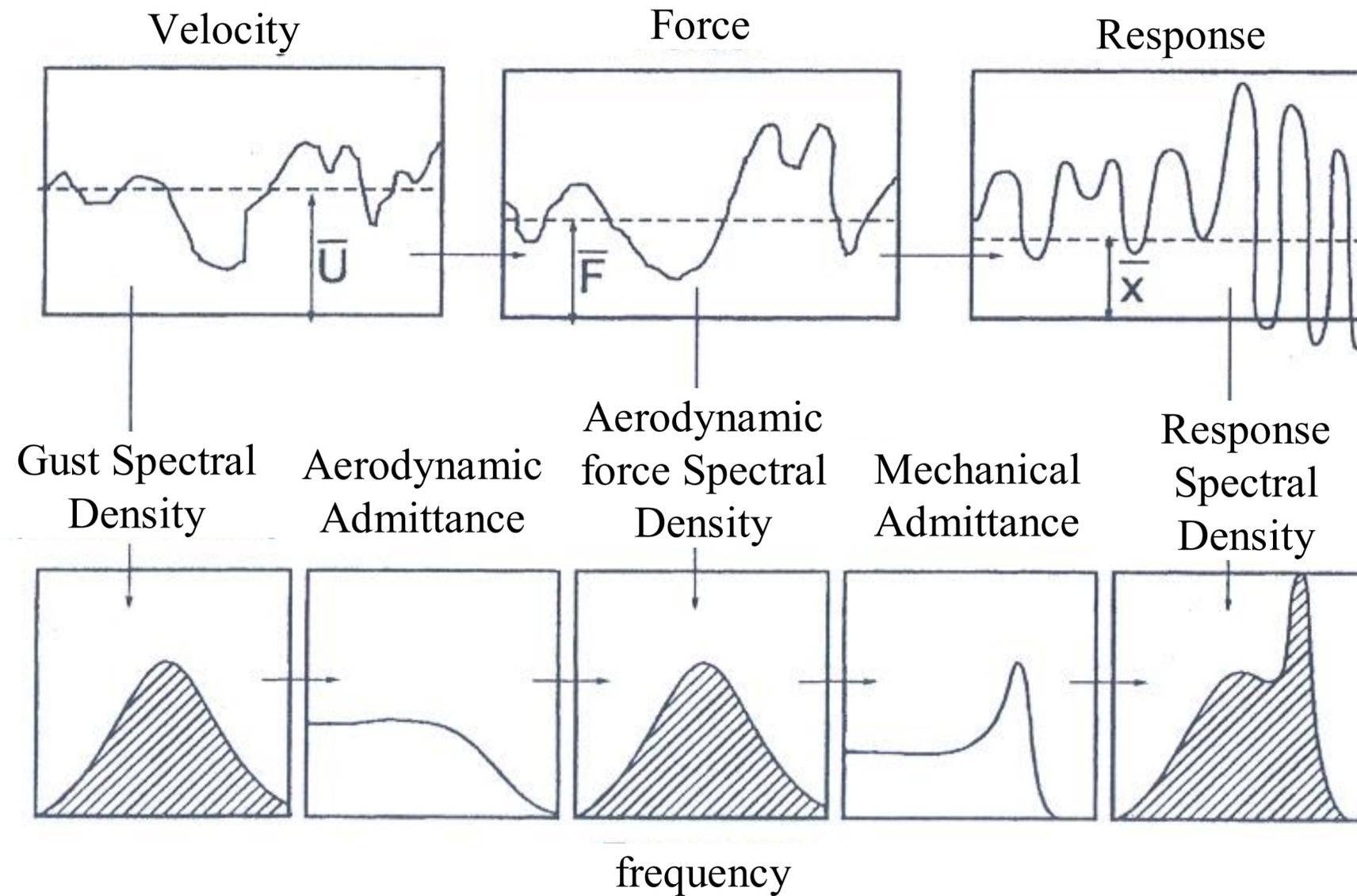
Secularity conditions :

$$R'(T) = \dots$$
$$\varphi'(T) = \dots$$

- ▶ Multiple timescale, in time domain
- ▶ Specific features :
 - ▶ Drop derivatives by one order
 - ▶ Slow is momentarily considered as constant for fast dynamics

Multiple Timescale Spectral Analysis

- ▶ Spectral analysis = integration of a spectrum over frequency

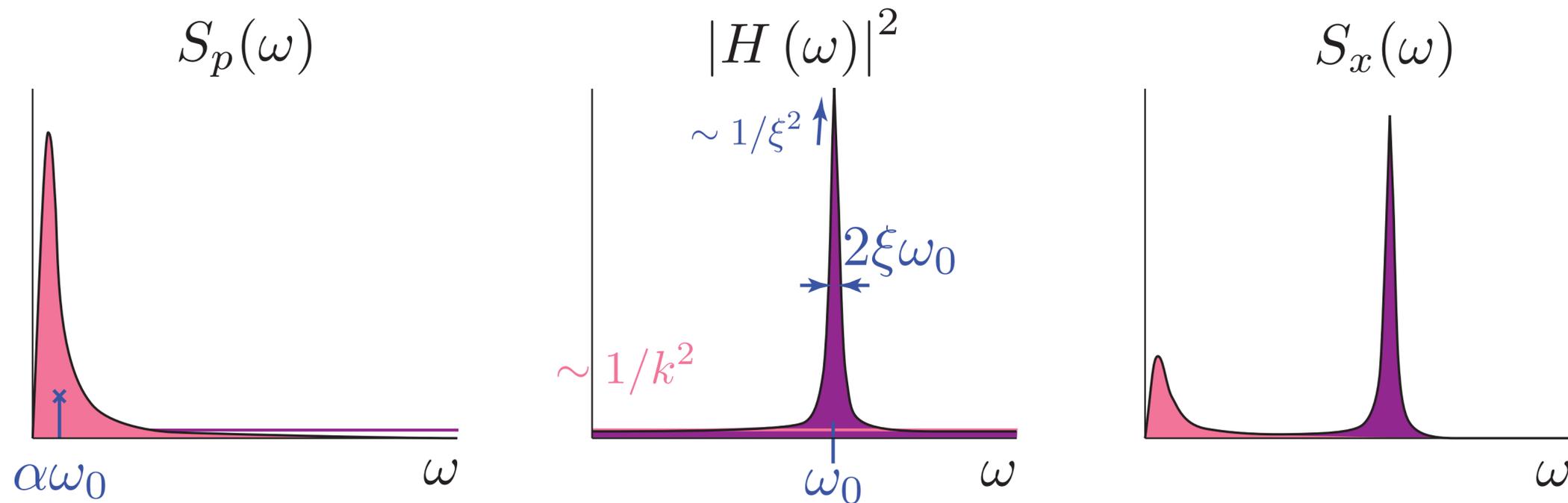


Multiple Timescale Spectral Analysis

- ▶ Spectral analysis = integration of a spectrum over frequency
- ▶ Spectrum (of response) = product of 2 functions
- ▶ Notice existence of two timescales for fast and slow dynamics
- ▶ Notice two small parameters in integrand : α and ξ

$$S_x(\omega) = |H(\omega)|^2 S_p(\omega)$$

$$\sigma_x^2 = \int_{-\infty}^{+\infty} S_x(\omega) d\omega$$



$$B = \left(\frac{\sigma_p}{k}\right)^2 \quad \text{Background}$$

$$R = \int_{-\infty}^{+\infty} |H(\omega)|^2 S_p(\omega) d\omega$$

$$R = \frac{\pi\omega_0}{2\xi} \frac{S_p(\omega_0)}{k^2} \quad \text{Resonant}$$

Another example of spectral analysis: covariance

- **covariances** of modal responses obtained from cross-PSDs

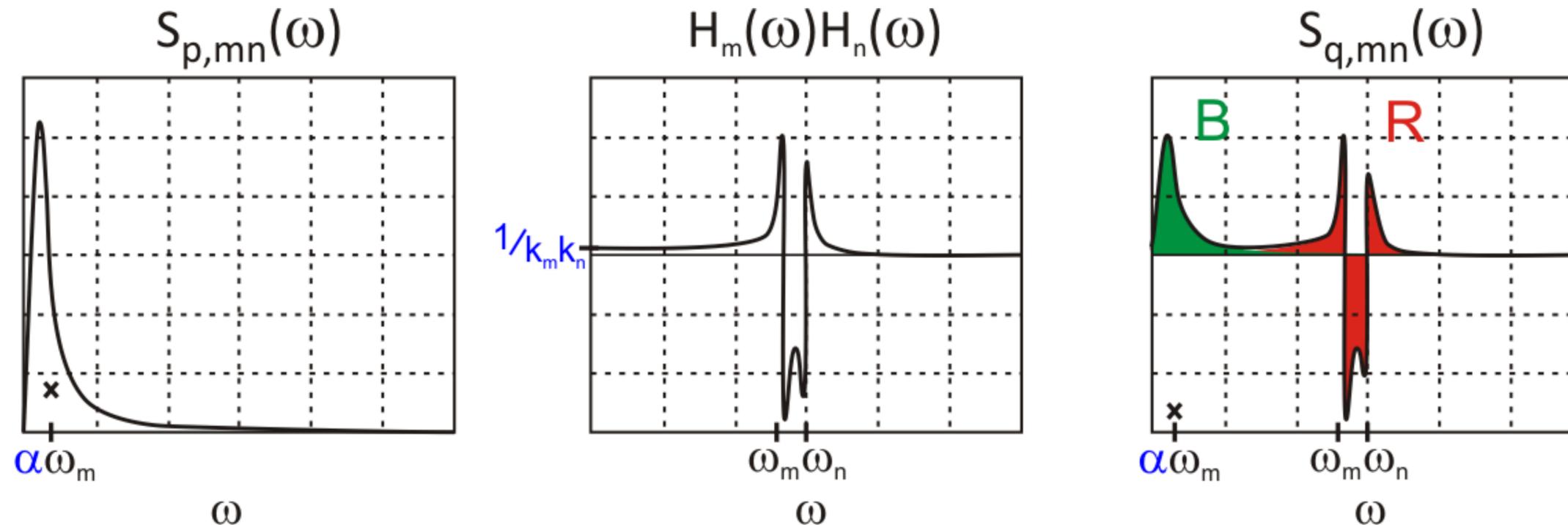
$$\sigma_{ij} = \int_{-\infty}^{+\infty} H_i(\omega) H_j^*(\omega) S_{p_{ij}}(\omega) d\omega$$

Background : replace $H_i(\omega) H_j^*(\omega)$ by $1/k_i k_j$

Resonant : replace $S_{p_{ij}}(\omega)$ by $S_{p_{ij}}(\bar{\omega})$

Variance

Background : replace $|H(\omega)|^2$ by $1/k^2$
 Resonant : replace $S_p(\omega)$ by $S_p(\omega_0)$



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Resonant component requires :

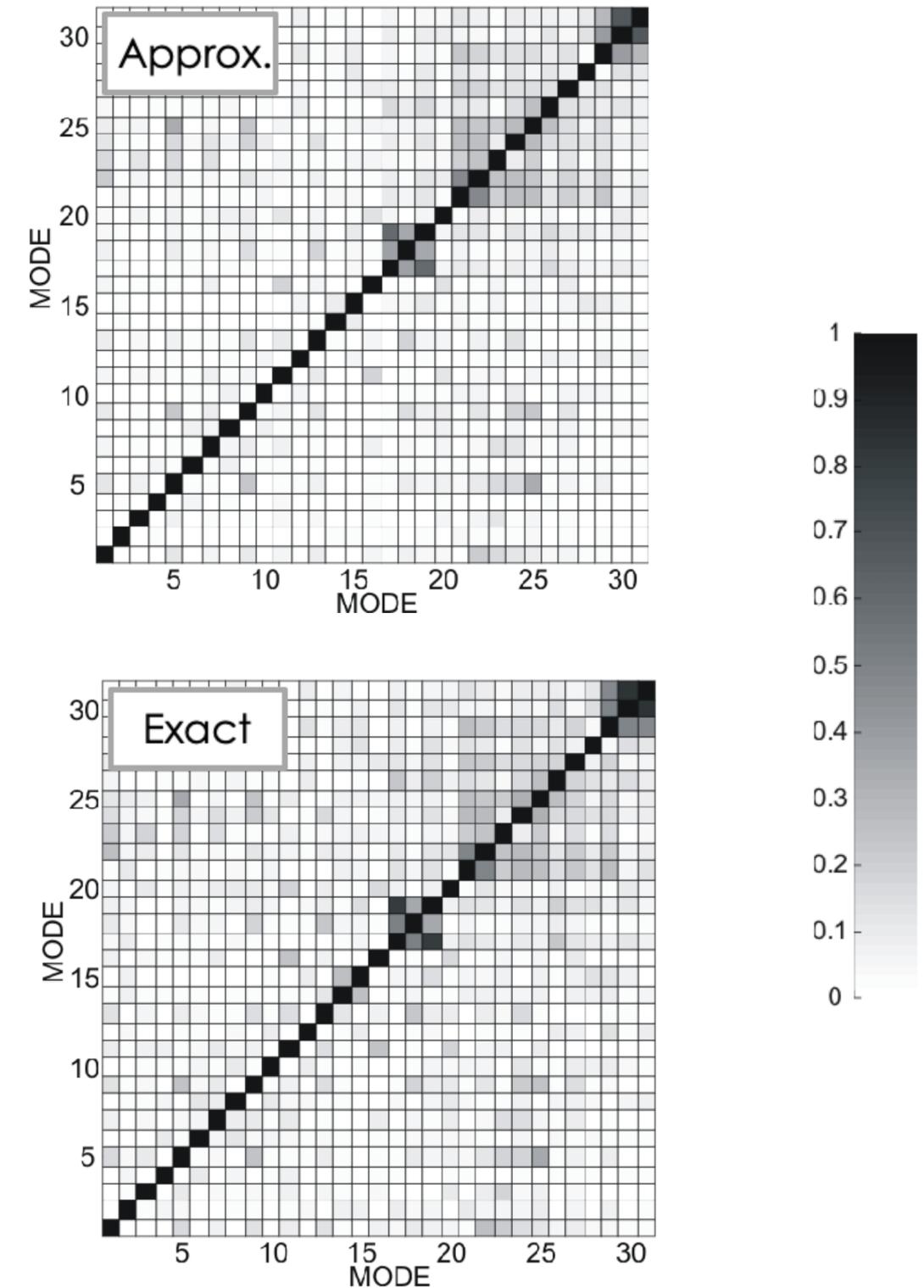
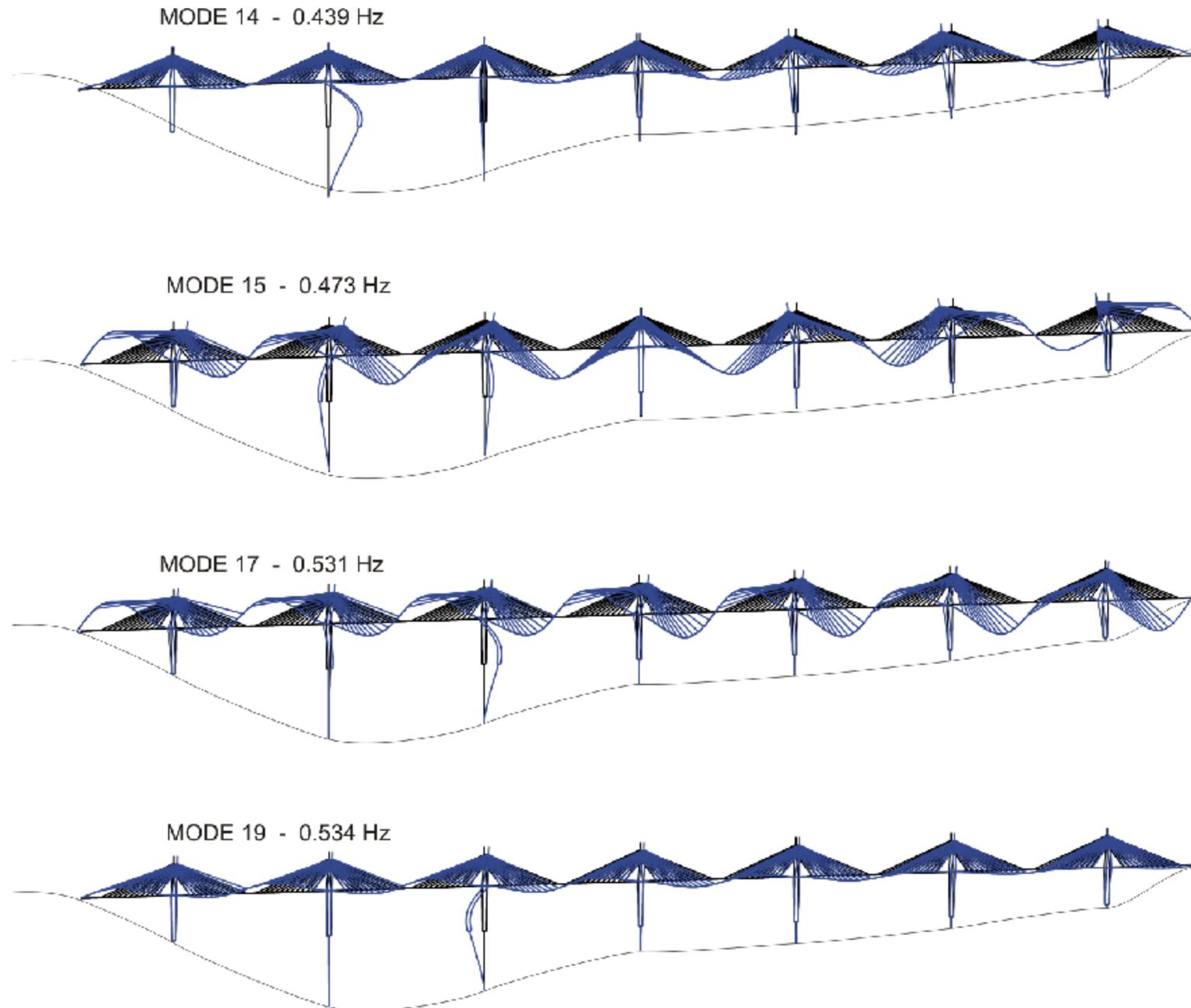
$$\int_{-\infty}^{+\infty} H_i(\omega) H_j^*(\omega) d\omega = \frac{1}{k_i k_j} \frac{4\pi\omega_i^2 \omega_j^2 (\xi_i \omega_i + \xi_j \omega_j)}{2\omega_i^2 \omega_j^2 (2\xi_i^2 + 2\xi_j^2 - 1) + 4\xi_i \xi_j \omega_i^3 \omega_j + 4\xi_i \xi_j \omega_i \omega_j^3 + \omega_i^4 + \omega_j^4}$$

(Not bad but quite long — Same as der Kiureghian's formula for correlation in earthquake engineering, 1993)

Variance
$R = \int_{-\infty}^{+\infty} H(\omega) ^2 S_p(\omega_0) d\omega$
$R = \frac{\pi\omega_0}{2\xi} \frac{S_p(\omega_0)}{k^2} \quad \text{Resonant}$

Another example of spectral analysis: covariance

- **covariances** of modal responses obtained from cross-PSDs



A counter-example : bispectral analysis

- ▶ Non-Gaussian stochastic analysis : **skewness** obtained from integration of bispectrum

$$B_x(\omega_1, \omega_2) = H(\omega_1) H(\omega_2) H^*(\omega_1 + \omega_2) B_p(\omega_1, \omega_2)$$

$$m_3 = \iint_{-\infty}^{+\infty} B_x(\omega_1, \omega_2) d\omega_1 d\omega_2$$

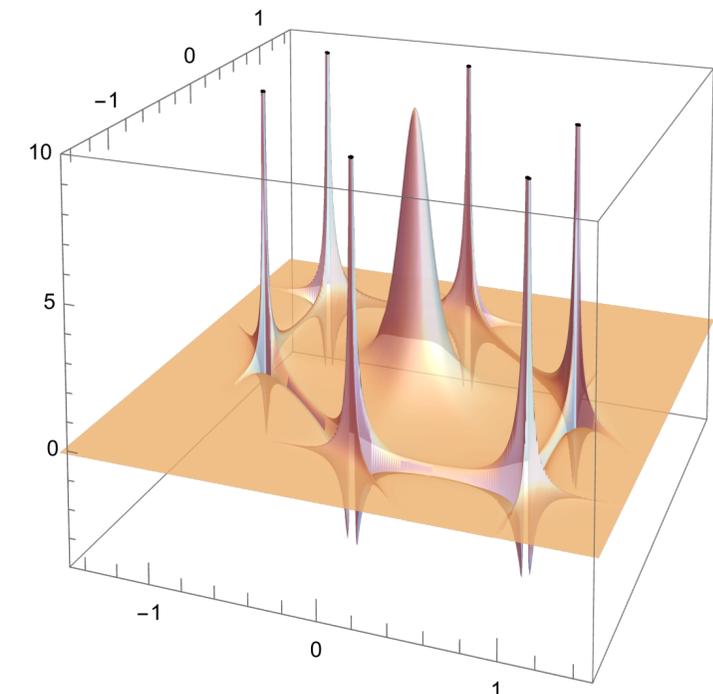
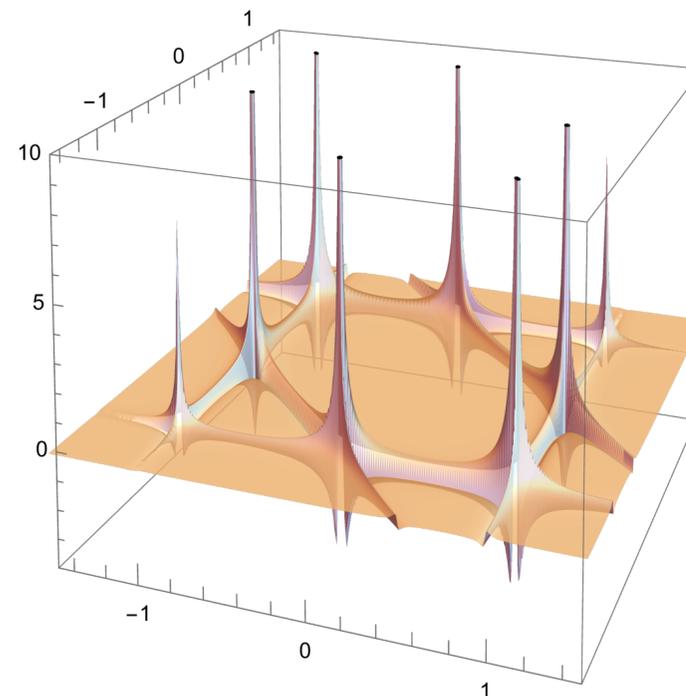
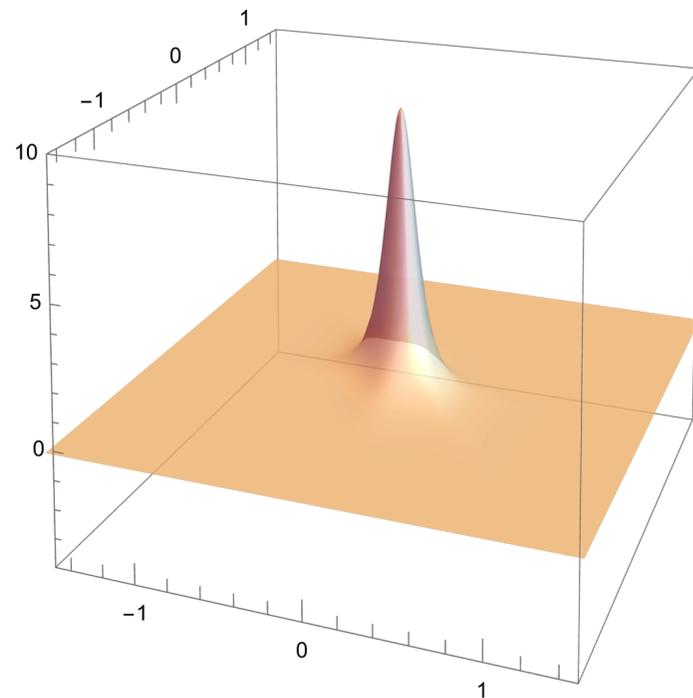
Background : replace $H(\omega_1) H(\omega_2) H^*(\omega_1 + \omega_2)$ by $1/k^3$

Resonant : replace $B_p(\omega_1, \omega_2)$ by $B_p(\omega_0, \omega_0)$

Variance

Background : replace $|H(\omega)|^2$ by $1/k^2$

Resonant : replace $S_p(\omega)$ by $S_p(\omega_0)$



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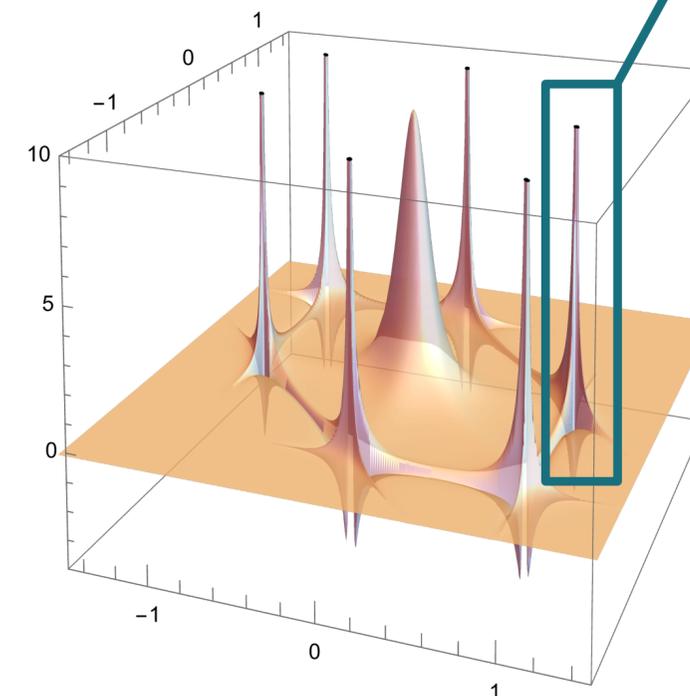
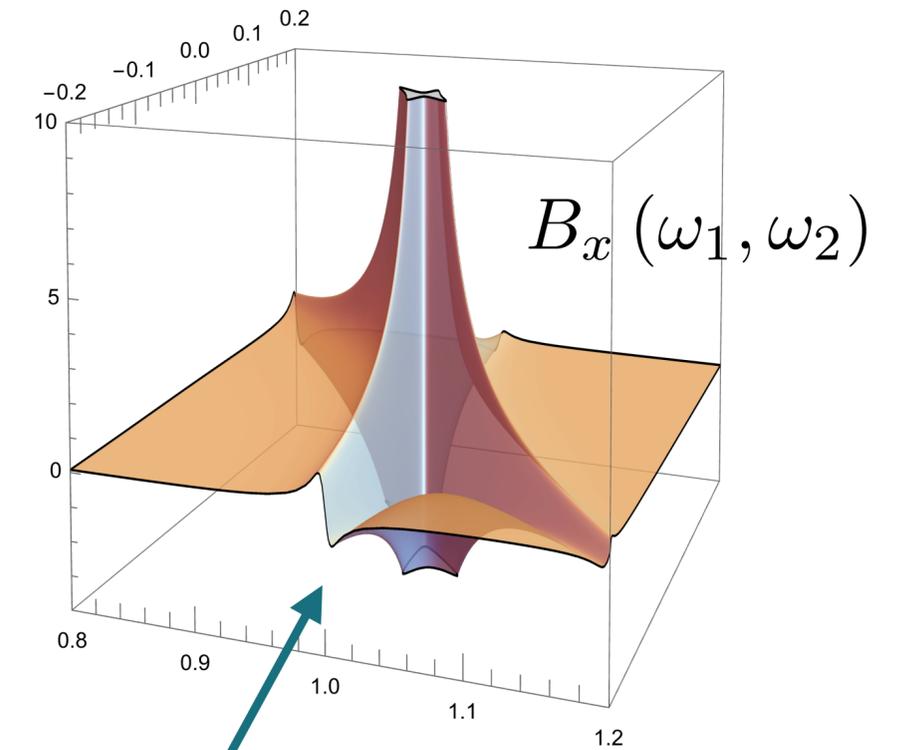
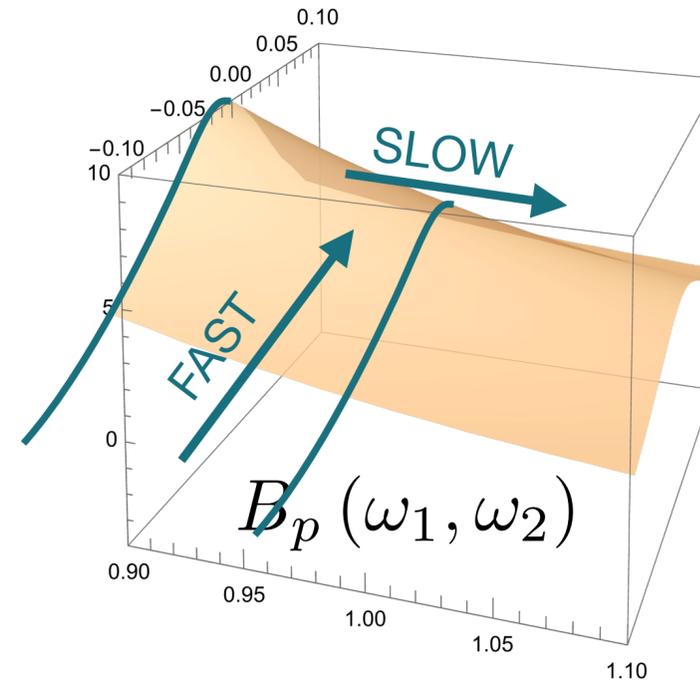
Resonant component requires :

$$\iint_{-\infty}^{+\infty} H_i(\omega_1) H_j(\omega_2) H_k^*(\omega_1 + \omega_2) d\omega_1 d\omega_2 = \dots \quad (\text{Much longer than covariance, but still ok})$$

This approximation unfortunately fails. One order of magnitude error.

A counter-example : bispectral analysis

- ▶ **zoom** on a peak
- ▶ The bispectrum of the loading changes:
 - ▶ slowly along an axis
 - ▶ fast in another direction
- ▶ The replacement by a constant in one direction is fine, but not both



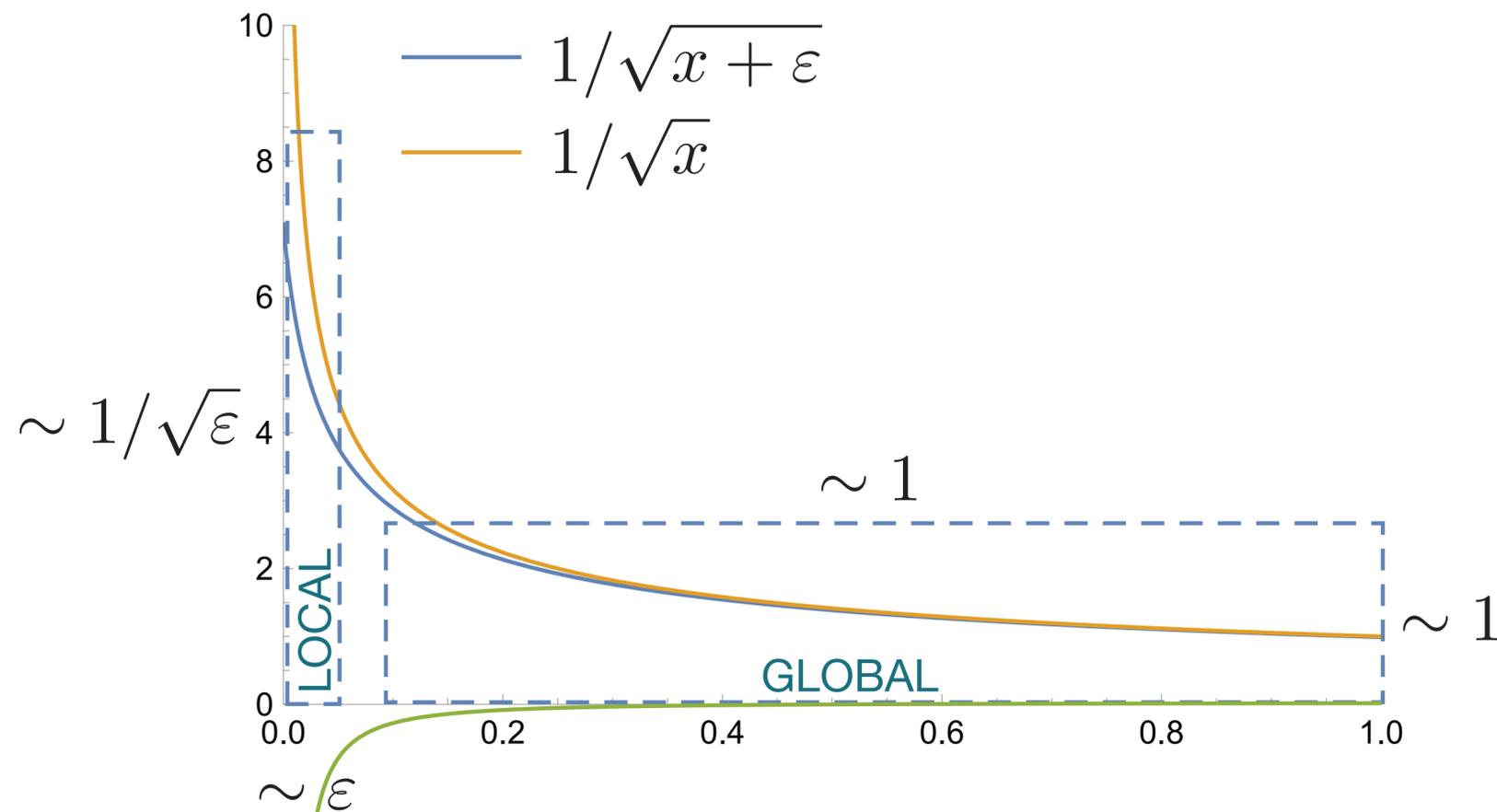
1. White noise approximation : the replacement by a constant is not universally correct.
2. We can drop the order of integration by 1, not 2.

Integral with small parameters $\varepsilon \ll 1$

Method of ranked contributions

Evaluate :

$$\int_0^1 \frac{1}{\sqrt{x + \varepsilon}} dx = 2\sqrt{1 + \varepsilon} - 2\sqrt{\varepsilon} \sim \underline{2} - 2\sqrt{\varepsilon} + \varepsilon + \dots$$



Two contributions to the integral :

- ▶ **Global** (background) : ~ 1
- ▶ **Local** (resonant) : $\sim \sqrt{\varepsilon}$

Global contribution : $f \sim \frac{1}{\sqrt{x}}$

$$\int_0^1 \frac{1}{\sqrt{x + \varepsilon}} dx \sim \int_0^1 \frac{1}{\sqrt{x}} dx = \underline{2}$$

Create **residual**, ...

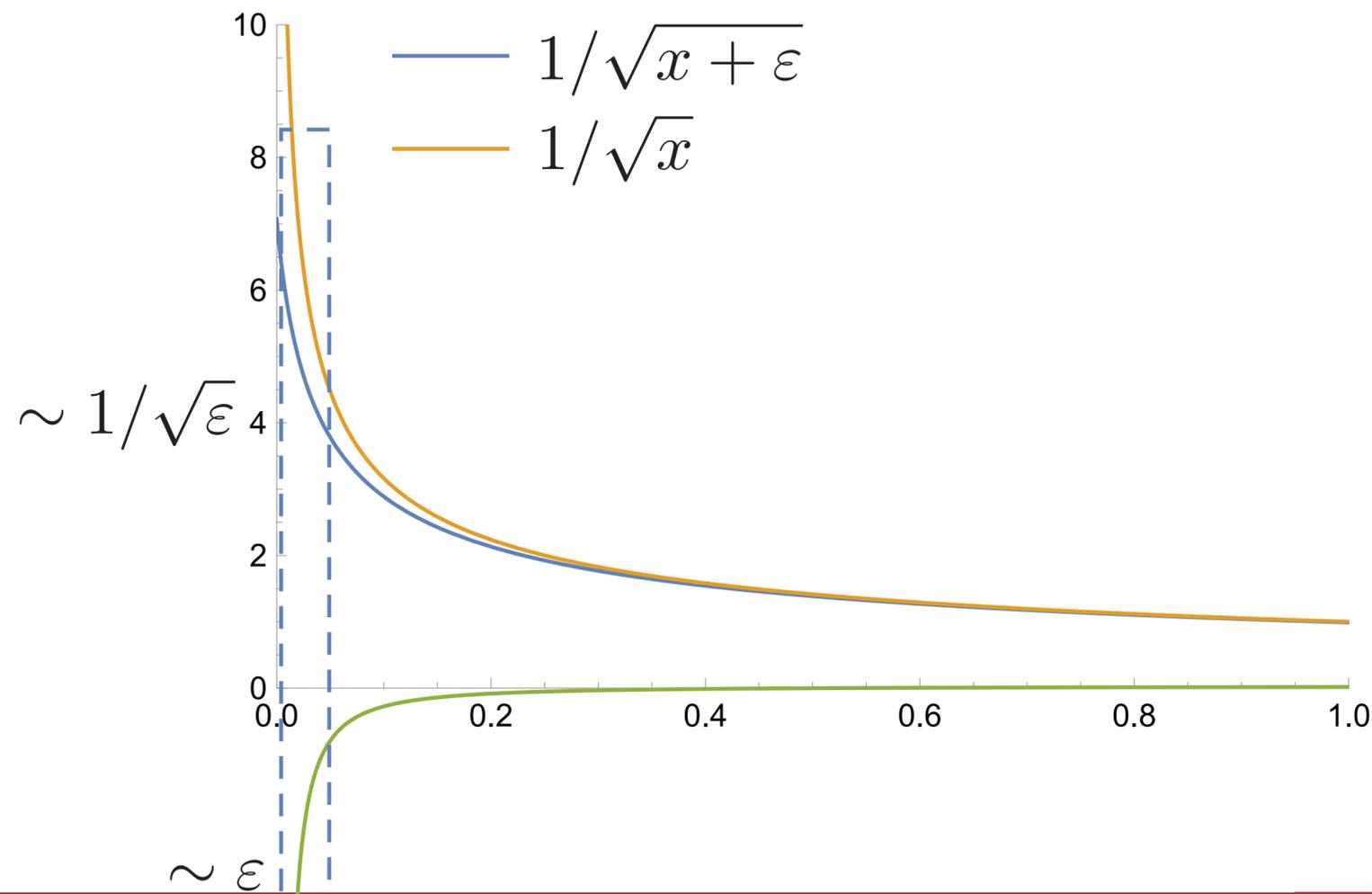
... focus on following contribution.

The Multiple Timescale Spectral Analysis

Method of ranked contributions

Evaluate :

$$\int_0^1 \frac{1}{\sqrt{x + \varepsilon}} dx = 2\sqrt{1 + \varepsilon} - 2\sqrt{\varepsilon} \sim \underline{2} - \underline{2\sqrt{\varepsilon}} + \varepsilon + \dots$$



Two contributions to the integral :

- ▶ **Global** (background) : ~ 1
- ▶ **Local** (resonant) : $\sim \sqrt{\varepsilon}$

Local contribution : ZOOM with stretched coord.

$$x = \varepsilon u$$

$$f - f_0 \sim \frac{1}{\sqrt{\varepsilon}} \left(\frac{1}{\sqrt{1 + u}} - \frac{1}{\sqrt{u}} \right)$$

$$\int_0^{1/\varepsilon \rightarrow +\infty} \frac{1}{\sqrt{\varepsilon}} \left(\frac{1}{\sqrt{1 + u}} - \frac{1}{\sqrt{u}} \right) du = -\frac{2}{\sqrt{\varepsilon}}$$

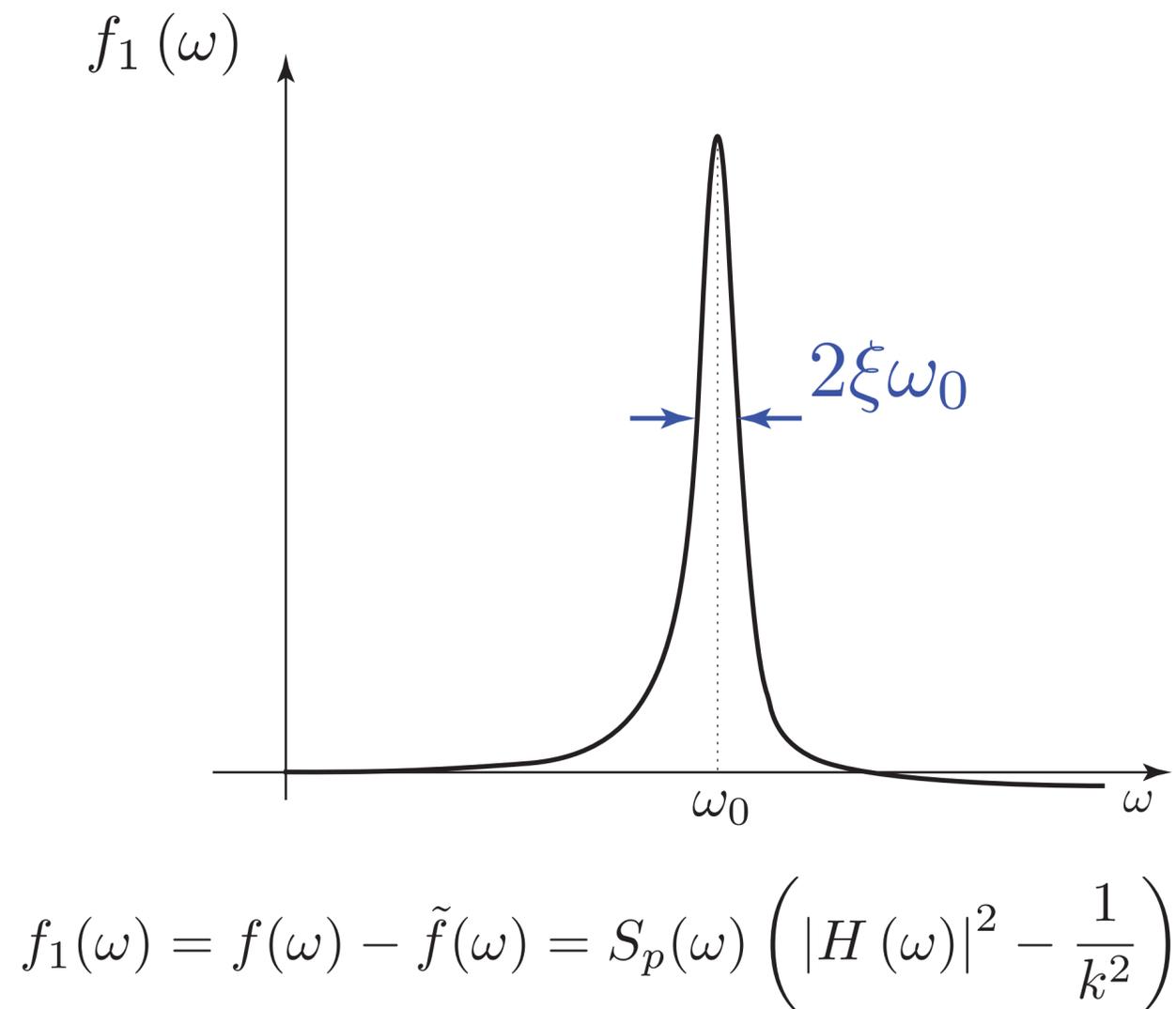
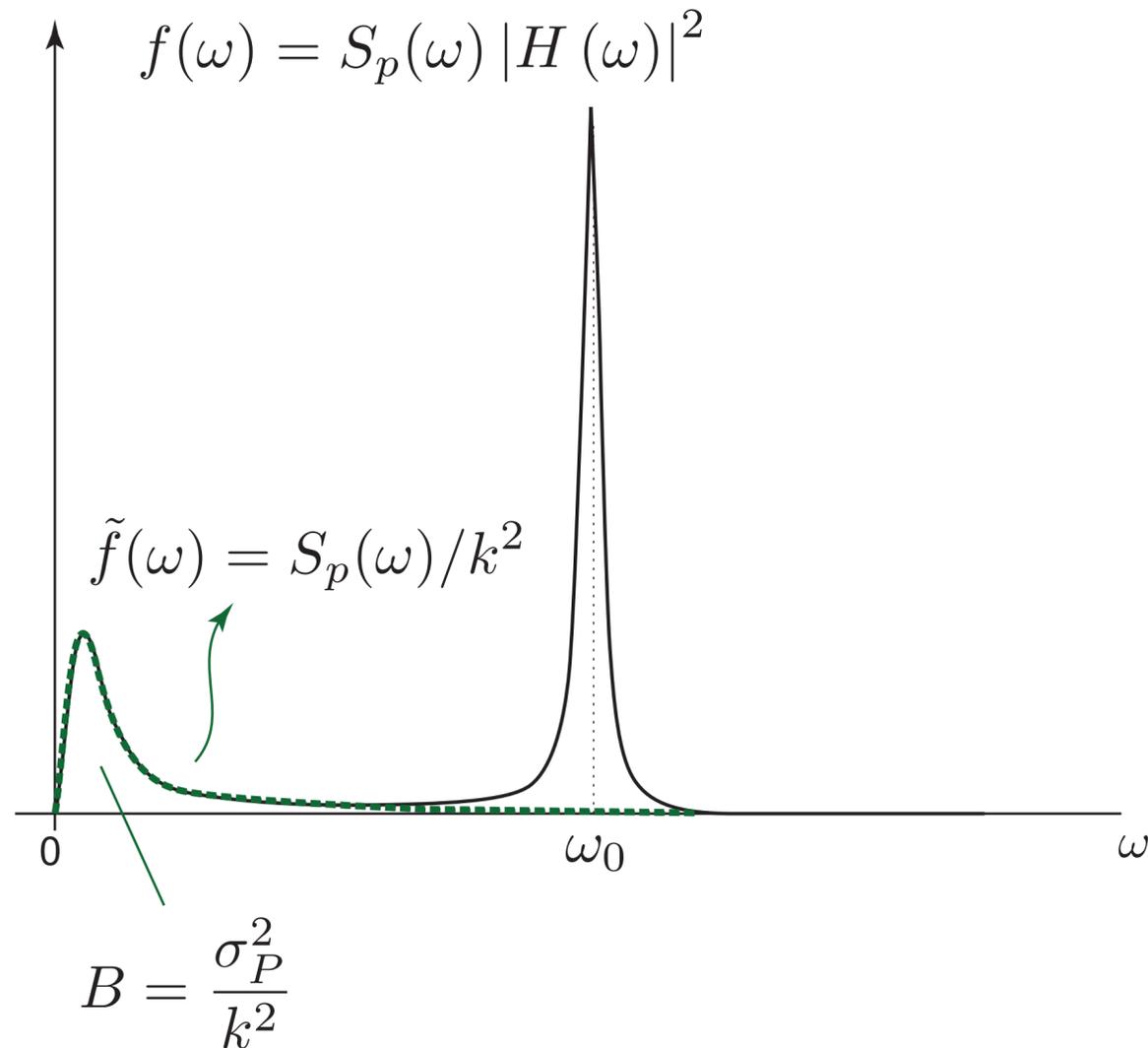
$$\text{Second contribution : } -\frac{2}{\sqrt{\varepsilon}} \cdot \varepsilon = \underline{-2\sqrt{\varepsilon}}$$

The Multiple Timescale Spectral Analysis

- ▶ By means of local/global approximations we can provide the asymptotic expansion of integral
- ▶ Background = global — Resonant = local
- ▶ Important to **zoom** in the neighborhood of local contributions, then only, approximate locally

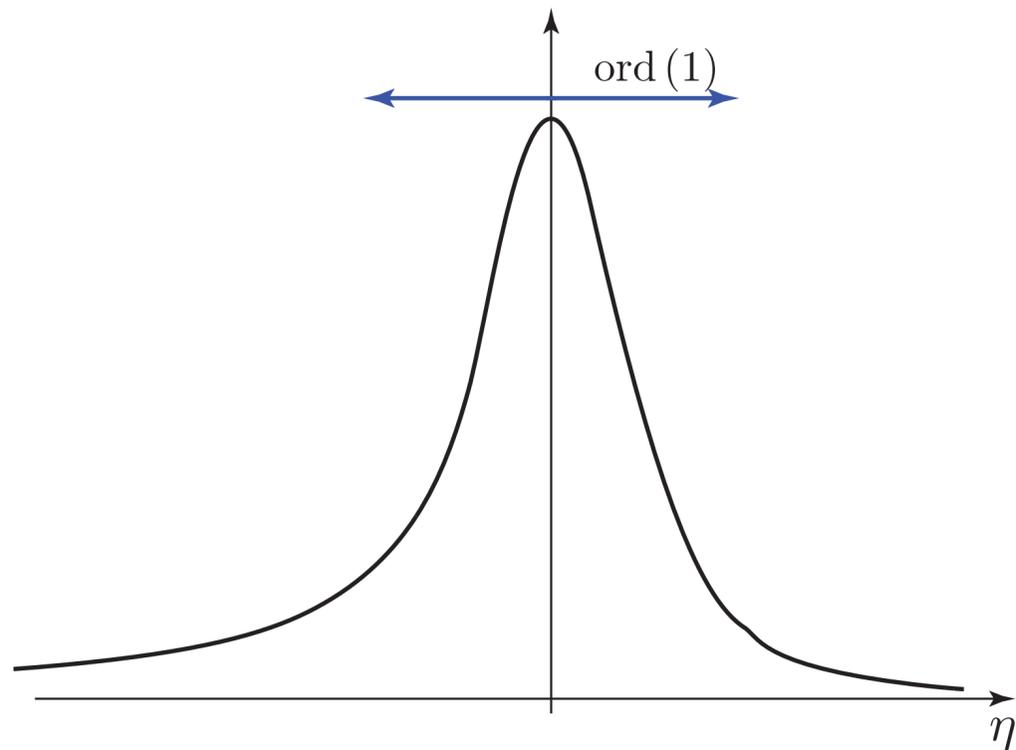
Revisiting the estimation of variance

1. Identify the two components to the integral
2. Choose the **Background** to start, determine first contribution
3. Subtract off approximation and construct first residual

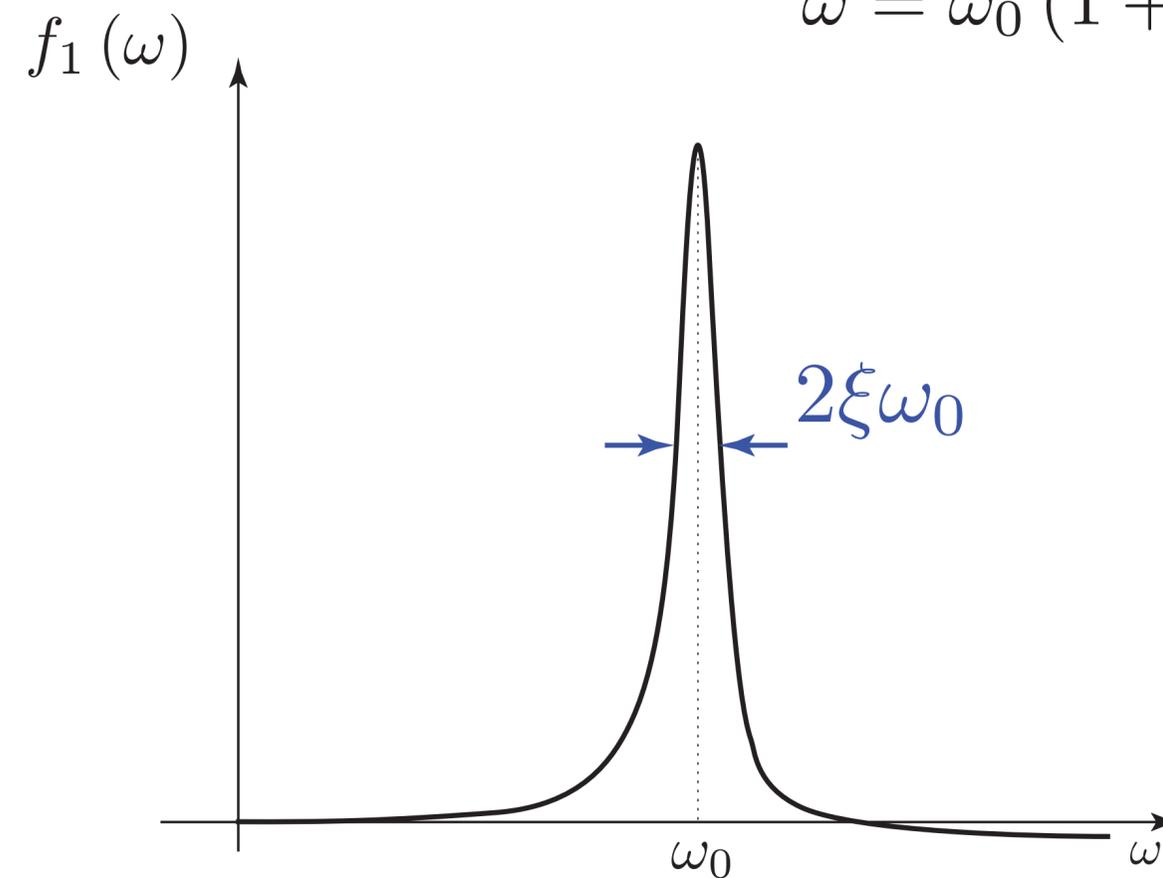


Revisiting the estimation of variance

4. Now focus on the Resonant component
5. Zoom-in with the appropriate rescaling > create a domain of order 1



Use **stretched coordinate**
 $\omega = \omega_0 (1 + \xi\eta)$



$$f_1(\omega) = f(\omega) - \tilde{f}(\omega) = S_p(\omega) \left(|H(\omega)|^2 - \frac{1}{k^2} \right)$$

Revisiting the estimation of variance

6. Find a local approximation
7. Integrate and determine Resonant contribution

We recovered Davenport B/R's approximation

Major difference :
the use of a stretched coordinate makes it neat and avoids « replace the PSD by constant value (equivalent white noise) »

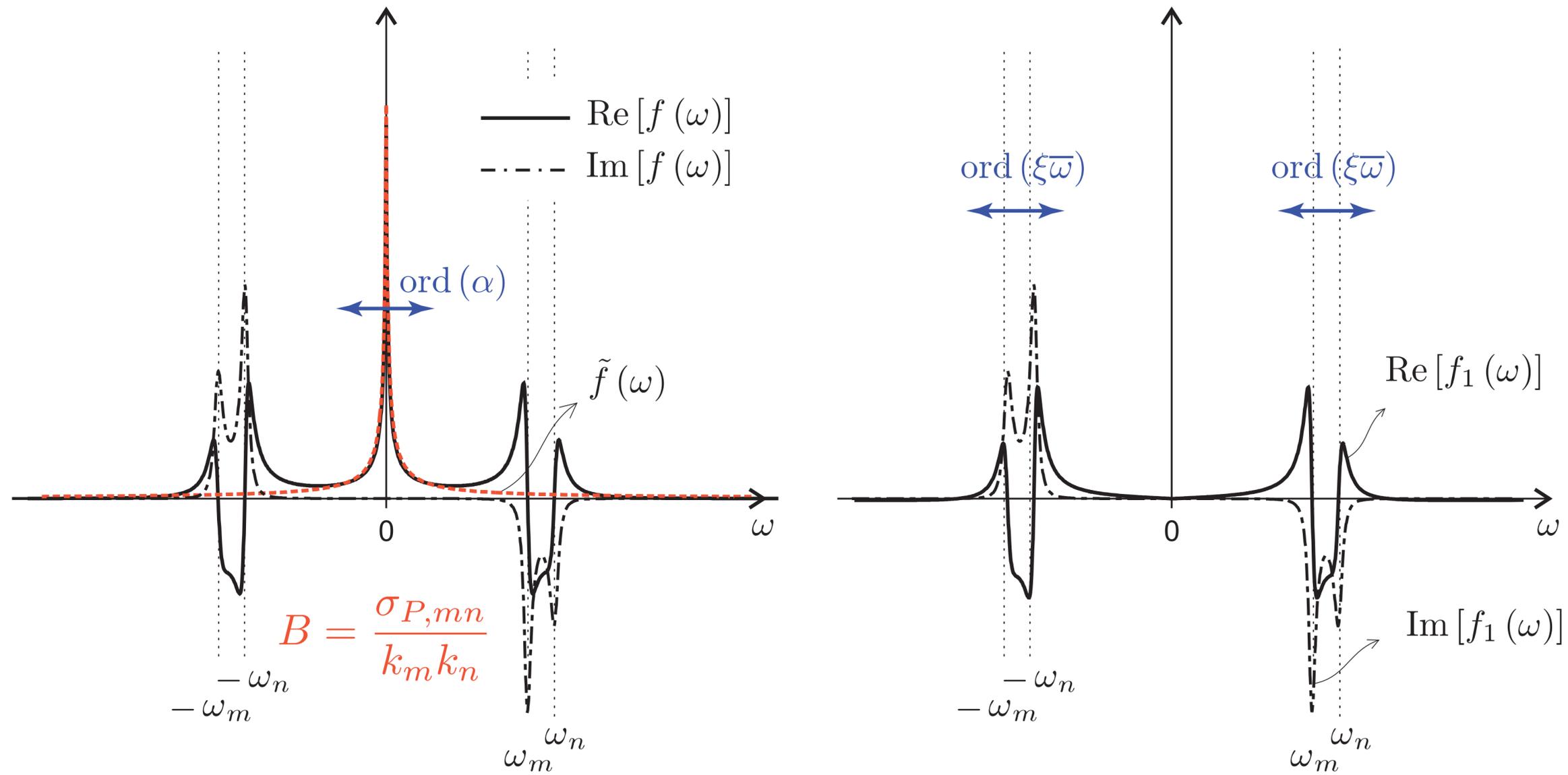
$$f_1(\omega) = S_p(\omega) \left(\underbrace{|H(\omega)|^2}_{\frac{1}{4\xi^2 k^2} \frac{1}{1+\eta^2} + \text{ord}(\xi^{-1})} - \frac{1}{k^2} \right)$$

$$\tilde{f}_1(\omega) = S_p(\omega_0) \frac{1}{4k^2} \frac{1}{\xi^2 + \left(\frac{\omega}{\omega_0} - 1\right)^2}$$

$$\rightarrow \int \tilde{f}_1(\omega) d\omega \rightarrow R = \frac{\pi\omega_0}{2\xi} \frac{S_p(\omega_0)}{k^2}$$

Revisiting the estimation of covariance

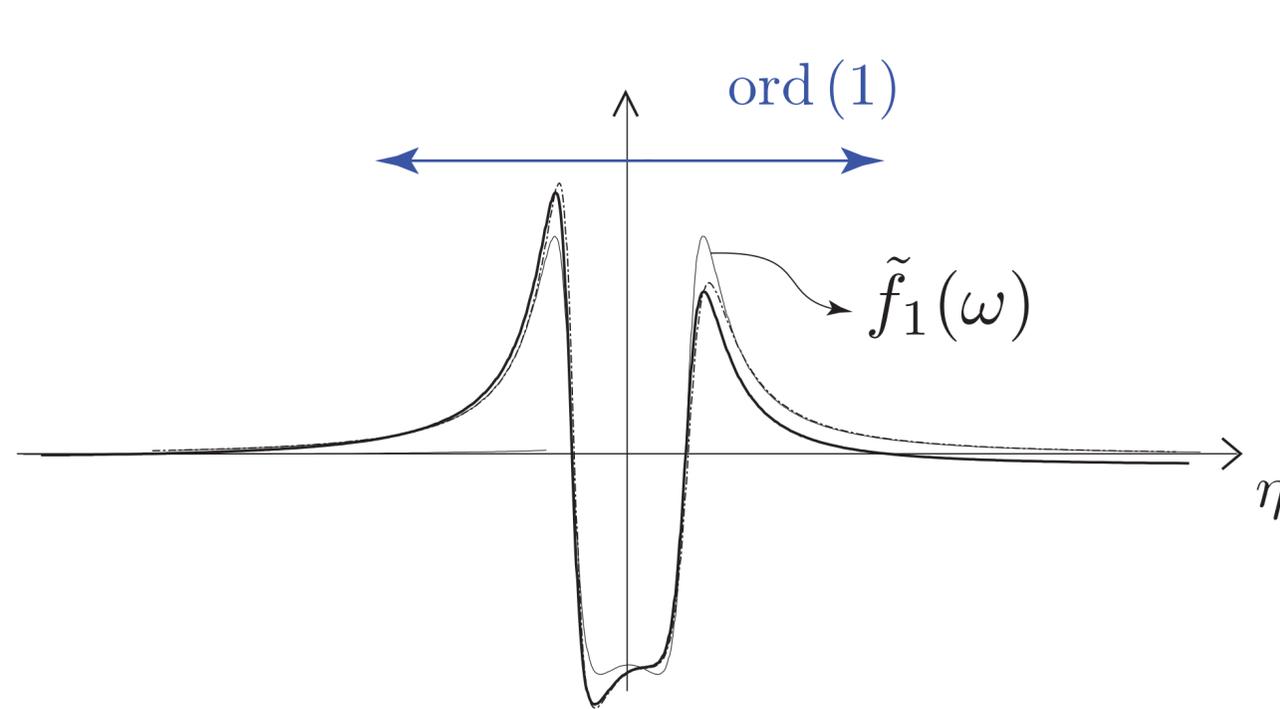
Develop same approach for the covariance



Revisiting the estimation of covariance

$$\omega = \frac{\omega_m + \omega_n}{2} + \eta(\omega_n - \omega_m)$$

Relative separation of natural frequencies : $\varepsilon = \frac{\omega_n - \omega_m}{\omega_n + \omega_m}$



$$\int \tilde{f}_1(\omega) d\omega$$



$$R = \frac{1}{k_m k_n} \operatorname{Re} \left[\frac{S_p(\omega_m) + S_p(\omega_n)}{2} \frac{\omega_m + \omega_n}{2} \frac{\pi(\xi - i\varepsilon)}{2(\varepsilon^2 + \xi^2)} \right]$$

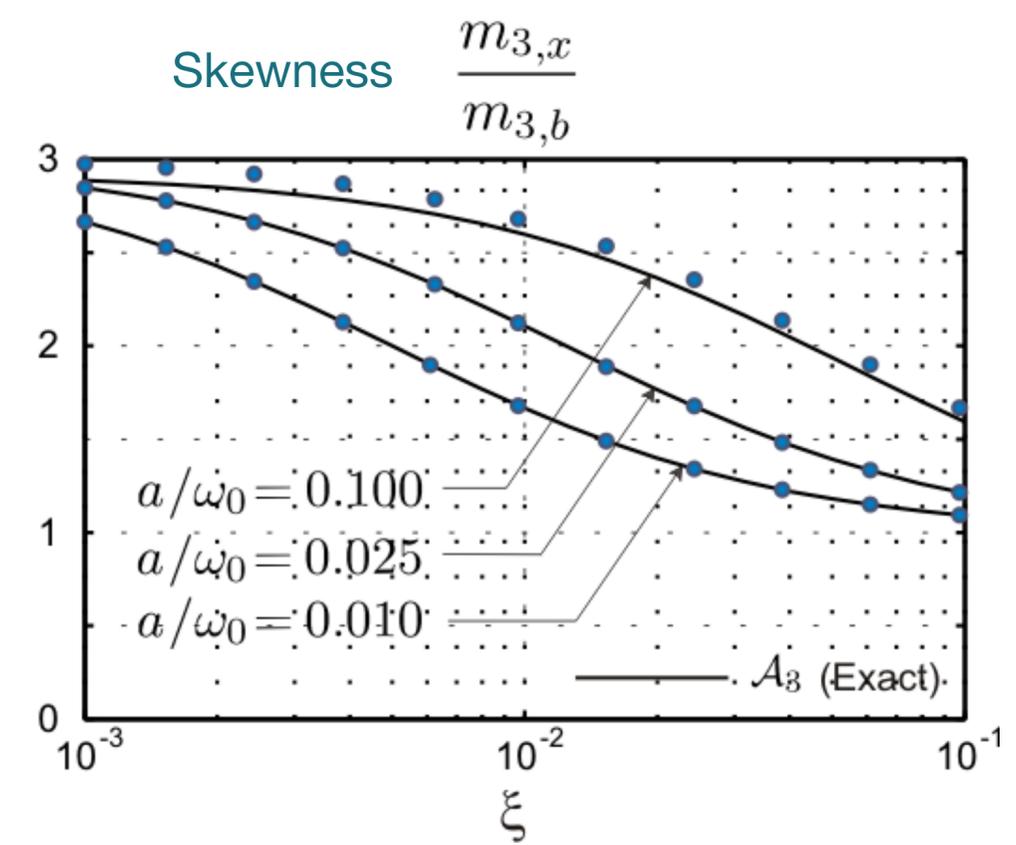
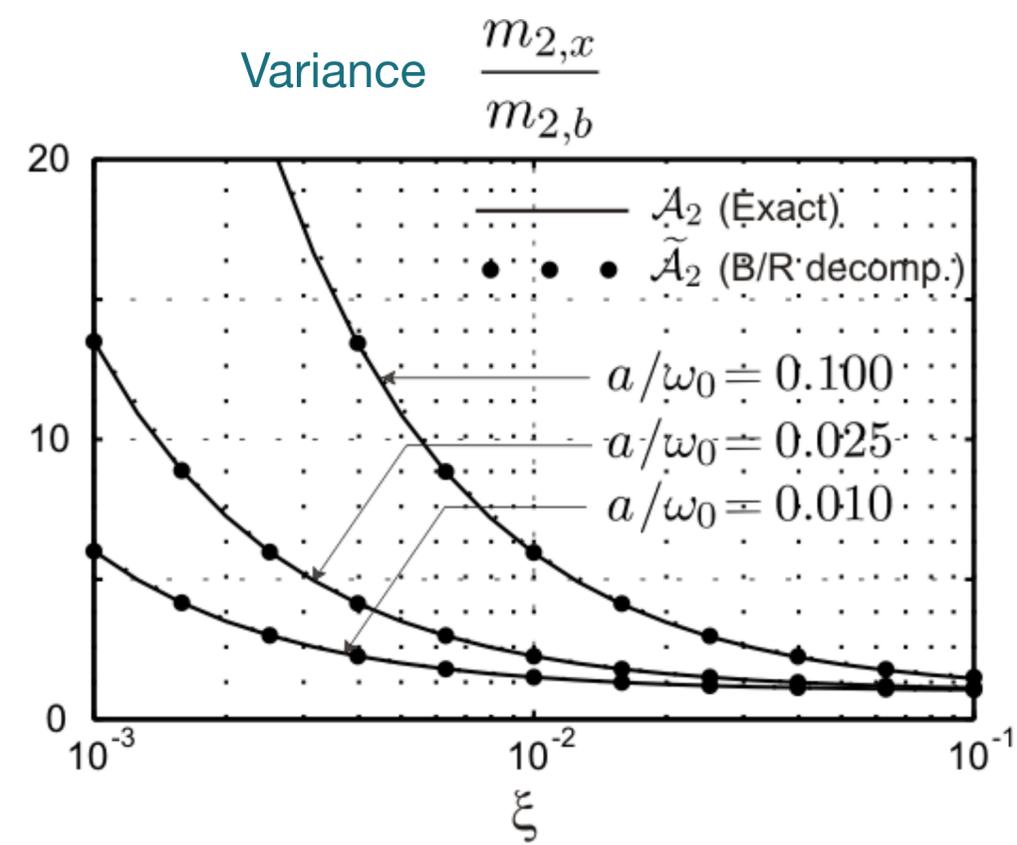
“White noise approximation”

$$\int_{-\infty}^{+\infty} H_i(\omega) H_j^*(\omega) d\omega = \frac{1}{k_i k_j} \frac{4\pi\omega_i^2 \omega_j^2 (\xi_i \omega_i + \xi_j \omega_j)}{2\omega_i^2 \omega_j^2 (2\xi_i^2 + 2\xi_j^2 - 1) + 4\xi_i \xi_j \omega_i^3 \omega_j + 4\xi_i \xi_j \omega_i \omega_j^3 + \omega_i^4 + \omega_j^4}$$

Revisiting the estimation of bispectrum

Develop same approach for bispectrum

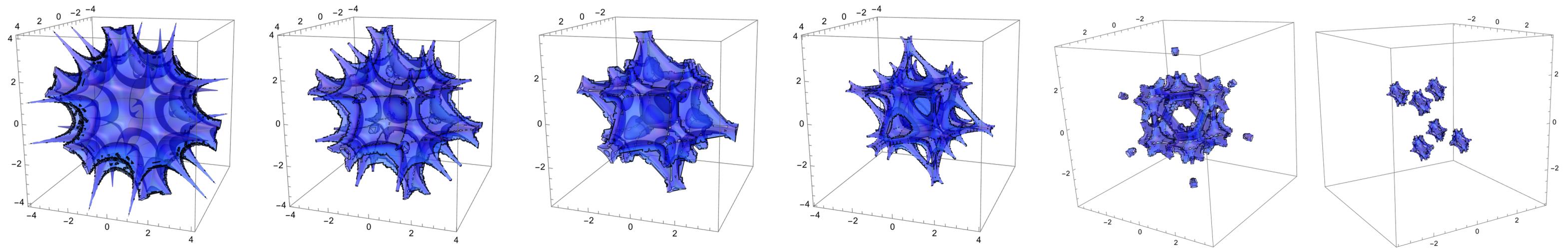
$$m_3 = \iint_{-\infty}^{+\infty} B_x(\omega_1, \omega_2) d\omega_1 d\omega_2 \quad \longrightarrow \quad m_3 = B + bR = \frac{m_{3,f}}{k^3} + 6\pi \frac{\xi \omega_0^3}{k^3} \int_{-\infty}^{+\infty} \frac{B_f(\omega_0, \omega_2)}{(2\xi \omega_0)^2 + \omega_2^2} d\omega_2$$



MTSA integration of trispectrum

Develop same approach for trispectrum (4th order non-Gaussian analysis, kurtosis)

$$m_4 = \iiint_{-\infty}^{+\infty} T_x(\omega_1, \omega_2, \omega_3) d\omega_1 d\omega_2 d\omega_3 \quad \longrightarrow \quad m_4 = B + tR = \frac{m_{4,f}}{k^4} + \iiint_{-\infty}^{+\infty} \dots d\omega_1 d\omega_2$$



Contours of fourth kernel

The other applications of the MTSA

- ▶ 1963 : Variance 1-DOF linear system [Davenport]
- ▶ 2005 : Covariance, modal responses
- ▶ 2011 : Skewness 1-DOF linear system (NG loading)
- ▶ 2012 : Kurtosis 1-DOF linear system (NG loading)
- ▶ 2015 : **Multiple Timescale Spectral Analysis**
- ▶ 2015 : Nonlinear 1-DOF system, Volterra series
- ▶ 2018 : Variance 1-DOF with fractional derivatives
- ▶ 2019-20 : Wind and waves : B-R-I decomposition
- ▶ 2022-23 : Floating structures
- ▶ 2022 : SDOF Flutter (frequency dependent systems)
- ▶ 2023 : Non-Gaussian M-DOF linear system
- ▶ 2025 : MDOF Flutter (long-span bridges)



Takehome messages

1. The (naive) replacement of PSD by a constant does not always work
2. MTSA = very generic can be applied to many contexts (integral with small numbers)
3. Drop by one order of integration \rightarrow comes with significant speedup, therefore nice for parametric studies, wrapping up in probabilistic framework or sensitivity analysis
4. Still room for applications in many other contexts
 - 3rd order Non-Gaussian MDOF to real structures (real finite element code)
 - 4th order Non-Gaussian MDOF (to be completely developed)
 - Transient loads (e.g. evolutionary spectral analysis)
 - Time-varying systems
5. Implementation in your own finite element code



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