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# <sup>15</sup>NH<sub>3</sub> in the atmosphere of a cool brown dwarf

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61 **Brown dwarfs serve as ideal laboratories for studying the atmospheres of giant exoplanets on**  
62 **wide orbits as the governing physical and chemical processes in them are nearly identical<sup>1,2</sup>.**  
63 **Understanding the formation of gas giant planets is challenging, often involving the endeav-**  
64 **our to link atmospheric abundance ratios, such as the carbon-to-oxygen (C/O) ratio, to for-**  
65 **mation scenarios<sup>3</sup>. However, the complexity of planet formation requires additional tracers,**  
66 **as the unambiguous interpretation of the measured C/O ratio is fraught with complexity<sup>4</sup>.**  
67 **Isotope ratios, such as deuterium-to-hydrogen and <sup>14</sup>N/<sup>15</sup>N, offer a promising avenue to gain**  
68 **further insight into this formation process, mirroring their utility within the solar system<sup>5,6,7</sup>.**  
69 **For exoplanets only a handful of constraints on <sup>12</sup>C/<sup>13</sup>C exist, pointing to the accretion of <sup>13</sup>C-**  
70 **rich ice from beyond the disks' CO iceline<sup>8,9</sup>. Here we report on the mid-infrared detection**  
71 **of the <sup>14</sup>NH<sub>3</sub> and <sup>15</sup>NH<sub>3</sub> isotopologues in the atmosphere of a cool brown dwarf with an ef-**  
72 **fective temperature of 380 K in a spectrum taken with the Mid-InfraRed Instrument of the**  
73 **James Webb Space Telescope. As expected, our results reveal a <sup>14</sup>N/<sup>15</sup>N value consistent with**  
74 **star-like formation by gravitational collapse, demonstrating that this ratio can be accurately**  
75 **constrained. Since young stars and their planets should be more strongly enriched in the**  
76 **<sup>15</sup>N isotope<sup>10</sup>, we expect that <sup>15</sup>NH<sub>3</sub> will be detectable in a number of cold, wide-separation**  
77 **exoplanets.**

78 The coldest class of brown dwarfs, so-called Y-dwarfs, span temperatures from 250 K to

79 500 K<sup>11</sup>. Their atmospheres are dominated by the absorption of water, methane, and ammonia,  
80 while water clouds likely become important for the colder Y-dwarfs<sup>12</sup>. Since their emission peaks  
81 in the mid-infrared beyond 4  $\mu\text{m}$ , Y-dwarf spectroscopic characterisation is challenging and the  
82 number of studies has been limited<sup>13,12,14</sup>. JWST is transforming the study of Y-dwarfs, by granting  
83 access to their full luminous range<sup>15</sup>. We analysed JWST/MIRI Medium Resolution Spectrometer  
84 (MRS)<sup>16</sup> observations of the Y-dwarf archetype WISEP J182831.08+265037.8 (hereafter WISE  
85 J1828), with an effective temperature of  $\sim 380$  K<sup>11</sup>. We obtain a mid-infrared spectrum at  $R \sim 3,000$   
86 to 1,500, between 4.9 and 27.9  $\mu\text{m}$ . The data reduction is described in Methods. Our observations  
87 are presented in Fig. 1, together with an exemplary best-fit model from our analysis, and reveal a  
88 spectrum rich in molecular features, namely a broad water absorption band at 5-7  $\mu\text{m}$ , methane at  
89 7.6  $\mu\text{m}$ , and ammonia at 9-13  $\mu\text{m}$ . The ammonia band is also shown in more detail in the lower  
90 panels of Fig. 1. We analyzed the atmospheric properties of WISE J1828 using several retrieval  
91 codes<sup>17,18,19</sup> and self-consistent atmosphere models in radiative-convective equilibrium<sup>20,21,22</sup>. The  
92 best-fit spectra and residuals are shown in Extended Data, Fig. 1. Since the MIRI observations  
93 mostly probe high altitudes in the atmosphere (with the contribution function peaking at  $\sim 1$  bar),  
94 we also added archival near-infrared data<sup>14</sup>, which probes deeper layers, at pressures of  $\sim 10$  bar  
95 (Extended Data Fig. 2). Due to its high luminosity, given its spectral type, WISE J1828 is suggested  
96 to be a binary system<sup>a</sup>. We thus modeled WISE J1828 as an equal mass binary system, emitting  
97 with identical atmospheres.

98 By combining the results of different retrieval approaches (see Methods), we constrain the  
99  $\log_{10}(\text{Volume Mixing Ratios})$  (VMRs) of the conspicuous absorbers  $\text{H}_2\text{O}$ ,  $\text{CH}_4$ , and  $^{14}\text{NH}_3$  to be  
100  $-3.03^{+0.18}_{-0.21}$ ,  $-3.65^{+0.21}_{-0.21}$ , and  $-4.79^{+0.15}_{-0.25}$  respectively. The surface gravity of WISE J1828 is con-  
101 strained to be  $\log(g) = 4.34^{+0.42}_{-0.88}$ , the effective temperature is  $378^{+13}_{-18}$  K, and the radius is constrained  
102 to  $1.37^{+0.26}_{-0.13} R_{\text{Jup}}$ . The uncertainties for these values are dominated by the dispersion between the  
103 various fitting approaches, and are thus larger than the actual uncertainties derived from any single  
104 analysis. A posterior plot of the various retrievals is shown in the Extended Data Fig. 3, while all  
105 retrieval results are summarized in the Extended Data Table 1. The self-consistent models constrain  
106 the atmospheric properties to be  $\log(g) = 4.5 \pm 1.0$ ,  $R = 1.27 \pm 0.21 R_{\text{Jup}}$ , and  $T_{\text{eff}} = 450 \pm 101$  K.  
107 Again, the reported uncertainties are dominated by differences between the two models. Our best-  
108 fit values of  $\log(g)$ ,  $R$ , and  $T_{\text{eff}}$ , after applying a binary correction to the inferred radii (dividing by  
109  $\sqrt{2}$ ), indicate an age of  $\sim 5$  Gyr for a  $\sim 15 M_{\text{Jup}}$  equal-mass binary system<sup>24</sup>, which is consistent  
110 with our mass constraints, see Extended Data Table 2.

111 The metallicity we derive for WISE J1828, combining the results of all approaches that  
112 included it as a free parameter (retrievals and self-consistent), is consistent with solar:  $[\text{M}/\text{H}] =$   
113  $0.02^{+0.12}_{-0.31}$ . For C/O we find  $0.22^{+0.37}_{-0.03}$  (solar is  $0.55 \pm 0.10$ <sup>25</sup>), while N/O is constrained to  $0.014^{+0.021}_{-0.002}$   
114 (solar is  $0.138 \pm 0.023$ <sup>25</sup>). These uncertainties are again dominated by the dispersion between dif-  
115 ferent approaches. The resulting posteriors for the metallicity, C/O and N/O are shown in Extended  
116 Data Fig. 3. Our findings thus indicate an atmosphere with a solar bulk metallicity, but depleted  
117 in C and N. A likely cause for this is a departure from chemical equilibrium, where gas poor in

<sup>a</sup>However, numerous studies, including recent JWST measurements, all failed at resolving its binarity, putting an upper limit of 0.5 au on the separation of its components<sup>11,23</sup>.

118 both  $\text{NH}_3$  and  $\text{CH}_4$  is mixed up from the deep interior of the object<sup>26</sup>. The resulting gas would  
119 be enriched in  $\text{N}_2$  and  $\text{CO}$ , which our observations are not sensitive to; while  $\text{N}_2$  is not spectrally  
120 active at all, our shortest wavelengths are longer than the location of the fundamental band of  $\text{CO}$   
121 at  $\sim 4.5 \mu\text{m}$ . However, it is questionable whether enough  $\text{CO}$  can be mixed up from the deep  
122 atmosphere to palpably change the inferred  $\text{C/O}$  ratio<sup>27</sup>.

123 In addition, we detect  $^{15}\text{NH}_3$  with a significance ranging from 4-6  $\sigma$ , with several lines of  
124 ammonia clearly visible in the data, see Fig. 1. We derive a VMR of  $-7.68^{+0.24}_{-0.34}$  for  $^{15}\text{NH}_3$  and  
125  $^{14}\text{N}/^{15}\text{N}=670^{+390}_{-211}$ , averaging over the results of various models. In Fig. 2, we summarize  $^{14}\text{N}/^{15}\text{N}$   
126 for a range of astrophysical objects<sup>7</sup>. Our value for WISE J1828 is consistent with solar at the  
127 1-2  $\sigma$  level. Both the Sun and WISE J1828, which we derive to have similar ages, have  $^{14}\text{N}/^{15}\text{N}$   
128 values above those observed in the interstellar medium, which has been enriched in  $^{15}\text{N}$  by galactic  
129 stellar evolution since their formation<sup>10</sup>. Our measurement thus shows that WISE J1828 most likely  
130 formed like a star, as expected<sup>28</sup>. A strong ice enrichment is unlikely, and we correspondingly rule  
131 out cometary values  $^{14}\text{N}/^{15}\text{N}<200$  by more than 3  $\sigma$ .

132 Constraints on  $^{14}\text{N}/^{15}\text{N}$  can serve as a formation tracer. For example, comets in the solar  
133 system, fundamental planetary building blocks, are enriched in  $^{15}\text{N}$  by a factor of 2-3 when com-  
134 pared to solar, due to  $^{15}\text{N}$ -rich  $\text{NH}_3$  and  $\text{HCN}$  ice. In contrast,  $\text{N}_2$  gas in the solar accretion disk is  
135 thought to have been depleted in  $^{15}\text{N}$ <sup>7</sup>. In the solar system both Jupiter and Saturn are enriched in  
136 bulk nitrogen, but show  $^{14}\text{N}/^{15}\text{N}$  values consistent with the Sun. This may mean that they accreted  
137 ice cold enough to contain even the volatile  $\text{N}_2$  ice, which requires temperatures  $<30 \text{ K}$ , and cor-  
138 responds to orbital distances  $>25 \text{ au}$ <sup>29,30</sup>. An enriched nitrogen content through accretion of ice,  
139 but at solar  $^{14}\text{N}/^{15}\text{N}$ , may therefore require accretion close to where even the highly volatile  $\text{N}_2$  can  
140 condense.

141 The understanding of how nitrogen fractionation actually occurs is incomplete<sup>7</sup>, but we sum-  
142 marize some processes during different stages of the stellar and planetary evolution in Fig. 3. In  
143 the denser parts of the clouds an increase of  $^{15}\text{N}$  in  $\text{NH}_3$  is inferred, with a candidate for this frac-  
144 tionation being isotope-selective photo-dissociation<sup>31</sup>. Subsequently,  $\text{NH}_3$  and  $\text{HCN}$  ices condense  
145 in the colder clumps, potentially producing  $^{15}\text{N}$ -rich ice. Once a protostar is formed, there is some  
146 evidence that  $^{14}\text{N}/^{15}\text{N}$  decreases further<sup>32,33</sup>. The final consequence is thought to be a  $^{15}\text{N}$  increase  
147 in the less volatile nitrogen-carriers, leading to the observed increase in  $^{15}\text{N}$  in  $\text{HCN}$  for protoplan-  
148 etary disks and in  $\text{NH}_3$  and  $\text{HCN}$  for solar system comets. We note that models predict  $^{15}\text{N}$ -poor  
149  $\text{NH}_3$  gas in protoplanetary disks<sup>33</sup>.

150 To further assess  $^{14}\text{N}/^{15}\text{N}$  as a formation tracer, we used a simplified planet formation model<sup>4</sup>.  
151 We tracked  $^{14}\text{N}/^{15}\text{N}$  as a function of the mass accreted as icy and rocky material, for a planet  
152 located between the  $\text{N}_2$  iceline of the disk, at about 20-80 au, and the  $\text{NH}_3$  iceline, ten times  
153 closer in<sup>30,34</sup>. The ices were therefore likely enriched in  $^{15}\text{N}$ . We find that for Saturn-like metal  
154 enrichment ( $\sim 6 \times \text{solar}$ <sup>35</sup>) the  $^{14}\text{N}/^{15}\text{N}$  decreases by 30-40% when compared to solar (see Extended  
155 Data Fig. 4), indicating that  $^{14}\text{N}/^{15}\text{N}$  can significantly vary when compared to the stellar value for  
156 a planet forming between the  $\text{N}_2$  and  $\text{NH}_3$  icelines.

157 With JWST MIRI, the formation-sensitive isotopologues  $^{14}\text{NH}_3$  and  $^{15}\text{NH}_3$  become accessi-  
158 ble for objects with low effective temperatures. In the mid-infrared,  $\text{NH}_3$  is a dominating absorber  
159 from  $T_{\text{eff}} = 1000 \text{ K}$ <sup>36</sup>, down to at least 380 K, the effective temperature of WISE J1828. As  
160 demonstrated above,  $^{14}\text{N}/^{15}\text{N}$  can constrain formation locations with respect to the  $\text{NH}_3$  and  $\text{N}_2$   
161 icelines of the disk. This is in addition to constraints on N/O, that the detection of  $\text{H}_2\text{O}$  and  $\text{NH}_3$   
162 enables, which has been suggested as another useful formation tracer<sup>37,38</sup>, but which needs careful  
163 interpretation due to chemical disequilibrium processes<sup>26</sup>. Simultaneous constraints on C/O, N/O  
164 and  $^{14}\text{N}/^{15}\text{N}$ , based on  $\text{CH}_4$ ,  $\text{CO}$ ,  $\text{H}_2\text{O}$  and  $\text{NH}_3$ , can be obtained for cold directly imaged exoplan-  
165 ets, further elucidating their formation history. These planets are found in orbits ranging from ten  
166 to hundreds of au, challenging the core-accretion paradigm for planetary formation<sup>39</sup>. They either  
167 formed at their detected locations via a star-like gravitational instability, or originated closer to  
168 their star via core accretion and subsequent outward migration<sup>40</sup>.

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## Figure Legends

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Legend Figure 1: **MIRI/MRS spectrum and exemplary best-fit model (here: pRT-free) of the Y-dwarf WISE J1828.** Panel a: full MIRI wavelength range considered in our models at retrieval resolution ( $\lambda/\Delta\lambda = 1000$ ). All other panels show the data at the original, higher MRS resolution; models have been post-processed to the same resolution. Panel b: like a, but zoomed in to show the  $\text{NH}_3$  absorption band at  $10 \mu\text{m}$  in more detail. Panels c to f: individual  $^{15}\text{NH}_3$  lines in the data, including a best-fit retrieval model with and without accounting for the opacity of  $^{15}\text{NH}_3$ . The error bars shown for the observations correspond to  $1-\sigma$  confidence levels. Panel f contains two overlapping MIRI MRS sub-channels. Colored lines denote the theoretical positions of the absorption lines of  $\text{H}_2\text{O}$ ,  $\text{CH}_4$ ,  $\text{NH}_3$  and  $^{15}\text{NH}_3$ .

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Legend Figure 2: **Comparison of the  $^{14}\text{N}/^{15}\text{N}$  ratio in the solar neighbourhood.** The values are based on either ammonia isotopologues (orange circles) or other molecules (black symbols), in different environments, subdivided by classes<sup>7</sup>. Our estimate for WISE J1828 appears as the brown dwarf class and is consistent with solar values (dashed horizontal yellow line) at the  $1-2 \sigma$  level. Lower values indicate enrichment in  $^{15}\text{N}$ , while the dashed blue line represents the current value of the interstellar medium (ISM).

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Legend Figure 3: **Different phases of the star and planetary formation.** We also show the relationship with the ammonia fractionation and the evolution of the  $^{14}\text{N}/^{15}\text{N}$  ratio at these stages: inside a molecular clouds with prestellar cores, during the formation of a protostar and in a circumstellar disk around a young star.

## Methods

### JWST/MIRI observations and data processing

The MIRI/MRS targeted WISE J1828 on July 28th 2022 without the use of target acquisition as part of JWST Guaranteed Time Observing program with PID 1189. All three dichroic/grating settings (SHORT, MEDIUM, LONG) were obtained in order to cover the full wavelength range, from 4.9 - 27.9  $\mu\text{m}$ <sup>16</sup>. The observations were executed with the point source optimized two-point dither (negative direction), and the detector set up with 180 frames per integration. No dedicated background observations were obtained.

The jwst pipeline was used (pipeline version: 1.9, CRDS version: 11.16.20, CRDS context: jwst\_1045.pmap) to process the data. The raw (level 1B) files were processed with the detector level pipeline (CALWEBB\_DETECTOR1) to produce calibrated rate files. This pipeline applies corrections for the nonlinearity of the ramps, dark current, detects jumps in the ramp due to cosmic rays, and finally fits the ramp signal to obtain slope values (rate.fits). The detector images were inspected and no sign of cosmic ray showers were found<sup>16</sup>. Since the WISE J1828 detector images were clean, we could use the fact that the point source itself is faint to perform a nod subtraction between the two dither points. First, we had to make sure that for every pair of detector images the flux levels were on the same level before subtracting them. A small time dependent difference between the individual integrations has been seen in MRS observations, with the first integration of the visit having a brighter flux level originating from the detector idling prior to the start of the exposures<sup>16</sup>. We used the region between the MRS channels on the detectors that do not contain any astrophysical signal to estimate a single value using the median, and subtract it from the whole detector. Next, the two dithers were subtracted from each other to remove the thermal background contribution. For Channel 1 of the MRS, a detector artifact that manifests as vertical stripes was still present after the nod subtraction. These stripes were around 10% of the science signal, and since the dispersion direction also closely follows the detector columns<sup>16</sup> the stripes affect the continuum of the extracted spectrum. We chose a region of the detector where the source signal was almost zero, and estimated the stripe contribution as the median of ten rows, which was then subtracted from every row of the detector.

With the detector images clean from detector artifacts and the thermal background, the spectroscopy pipeline (CALWEBB\_SPEC2) was run in order to obtain calibrated detector images (cal.fits), assigning the astrometric and wavelength information, correcting for the scattered light and detector fringing, and applying the photometric calibration. Finally, with CALWEBB\_SPEC3 we built the dither combined cubes<sup>41</sup>, with outlier rejection enabled. With background subtracted cubes, the spectral extraction from the cube is done by performing aperture photometry for each wavelength slice of the cube. We first determined the center of the point source by fitting a 2D Gaussian in the wavelength collapsed cube, then placed an

315 aperture of 1 Full Width Half Max (FWHM) centered on the source, applying an aperture  
316 correction for each wavelength to account for the missing flux of the point spread function  
317 outside the aperture. Some outliers remain in the extracted spectrum, which were removed for  
318 plotting the spectrum in Fig. 1 but not while fitting it. We traced back these outliers to the  
319 detector where a few cosmic rays overlap with the spectral trace, but are not bright enough for  
320 the outlier algorithm to detect. We therefore clip these values manually for each spectral band  
321 by setting a threshold for the flux. The outliers affected in total 0.4% of the spectrum.

## 322 Retrieval analysis

323 We carried out independent analyses of the MIRI/MRS spectra using a diversity of models,  
324 namely with the radiative-convective equilibrium codes ATMO and ARCiS+ML, and with the  
325 retrieval codes ARCiS, Brewster, and petitRADTRANS<sup>20,19,17,18,21,22</sup>. For computational fea-  
326 sibility, the retrievals with ARCiS, Brewster, and petitRADTRANS were run at a resolution of  
327  $\lambda/\Delta\lambda = 1000$ ; the MRS data was binned down correspondingly. For the retrievals we decided  
328 on a setup that assumed vertically constant absorber abundances that were retrieved freely,  
329 and a flexible parameterization of the pressure-temperature structure, which varied slightly be-  
330 tween the setups, see below. For the retrievals presented here we assumed WISE J1828 to be  
331 a single object, but allowed for a radius prior wide enough to account for an equal-brightness  
332 binary scenario. Clouds were neglected, which appears to be justified from a population-wide  
333 Y-dwarf retrieval analysis on Hubble Space Telescope (HST) data<sup>42</sup>. We note, however, that  
334 the impact of clouds should increase towards longer wavelengths, and for colder Y-dwarfs<sup>43,44</sup>.  
335 As shown in Extended Data Fig. 2, our inferred  $P$ - $T$  profiles cross the water saturation vapor  
336 pressure curve at the top of the photosphere probed by HST and MIRI, so a cloud could have  
337 some moderate impact on our results and the effect of its inclusion should be assessed in future  
338 studies.

339 We also observed that the reported uncertainties of the JWST reduction can be much smaller  
340 than the differences observed in the overlapping regions of MRS subchannels. In addition,  
341 all best-fit models had a  $\chi^2$  considerably larger than the number of wavelength channels. We  
342 thus opted for retrieving the magnitude of the uncertainties via the  $10^b$  treatment<sup>45</sup>, where the  
343 error bars  $\sigma$  considered during the retrieval are calculated from the uncertainties reported from  
344 the reduction  $\sigma_{\text{red}}$  as follows:  $\sigma = \sqrt{\sigma_{\text{red}}^2 + 10^b}$ . Separate  $b$ s were retrieved for MIRI and  
345 HST data, respectively. The retrieved pressure-temperature structures of the retrievals would  
346 also exhibit kinks sometimes, which are challenging to reconcile with radiative-convective  
347 equilibrium models. In this case we implemented a regularization of the pressure-temperature  
348 structure<sup>45</sup>. With this modification, the ARCiS and petitRADTRANS retrievals optionally put  
349 a penalty on  $d^2\log T/d\log P^2$ , which strives towards a constant power-law dependence between  
350 pressure and temperature, since  $d\log T/d\log P = \text{cst}$  implies  $T \propto P^\alpha$ , with  $\alpha$  being the  
351 constant power law coefficient. This setup may therefore also reproduce the relation between  
352 pressure and temperature in the deep atmosphere, which is expected to be convective. The  
353 individual models we used for the analysis are described below.

354 The results of all model inferences for WISE J1828 are found in Extended Data Table 1,  
355 and the combined result of all retrievals is presented in Extended Data Table 2. Elemental  
356 abundance ratios (C/O, N/O,  $^{14}\text{N}/^{15}\text{N}$ ) were computed from the retrievals' VMR constraints  
357 for the atmospheric absorbers. They may thus miss additional atoms locked up in clouds (in  
358 the case of oxygen) or affected by quenching in species which are spectrally inactive ( $\text{N}_2$  in the  
359 case of N) or have features outside the HST and MIRI wavelength range (CO in the case of C  
360 and O). The one-dimensional projection of the posteriors, for key forward model parameters,  
361 is shown in Extended Data Fig. 3 for all individual models. In Extended Data Fig. 2 we show  
362 the associated pressure-temperature uncertainties derived from all model analyses.

## 363 ATMO

364 We briefly describe the main properties of the ATMO<sup>20</sup> models and the grids that have been  
365 used in our study. These grids are publicly available at [https://opendata.erc-atmo.](https://opendata.erc-atmo.eu)  
366 [eu](https://opendata.erc-atmo.eu). The ATMO models assume that clouds are not needed to reproduce the shape of the  
367 SED of brown dwarfs (apart from the 10- $\mu\text{m}$  silicate feature). The authors have proposed  
368 that diabatic convective processes<sup>46</sup> induced by out-of-equilibrium chemistry of CO/CH<sub>4</sub> and  
369 N<sub>2</sub>/NH<sub>3</sub> can reduce the temperature gradient in the atmosphere and reproduce the reddening  
370 previously thought to occur by clouds. The grids assume a modification of the temperature  
371 gradient with an effective adiabatic index. The levels modified are between 2 and 200 bars  
372 at  $\log(g) = 5.0$  and are scaled by  $10^{\log(g)-5}$  at other surface gravities. Out-of-equilibrium  
373 chemistry is used with  $K_{zz} = 10^5 \text{ cm}^2 \text{ s}^{-1}$  at  $\log(g) = 5.0$  and is scaled by  $10^{2[5-\log(g)]}$  at  
374 other surface gravities. The mixing length is assumed to be two scale heights at 20 bars and  
375 higher pressures at  $\log(g) = 5.0$ , and is scaled down by the ratio between the local pressure  
376 and the pressure at 20 bars for lower pressures. The 20 bars limit is scaled by  $10^{\log(g)-5}$  at  
377 other surface gravities. The chemistry includes 277 species and out-of-equilibrium chemistry  
378 has been performed using a relaxation model<sup>47</sup>. Rainout is assumed to occur for species that  
379 are not included in the out-of-equilibrium model. Opacity sources include H<sub>2</sub>-H<sub>2</sub>, H<sub>2</sub>-He, H<sub>2</sub>O,  
380 CO<sub>2</sub>, CO, CH<sub>4</sub>, NH<sub>3</sub>, Na, K, Li, Rb, Cs, TiO, VO, FeH, PH<sub>3</sub>, H<sub>2</sub>S, HCN, C<sub>2</sub>H<sub>2</sub>, SO<sub>2</sub>, Fe, H-,  
381 and the Rayleigh scattering opacities for H<sub>2</sub>, He, CO, N<sub>2</sub>, CH<sub>4</sub>, NH<sub>3</sub>, H<sub>2</sub>O, CO<sub>2</sub>, H<sub>2</sub>S, SO<sub>2</sub>.  
382 The grids explore the following parameters: effective temperatures between 250 and 1200K;  
383  $\log(g)$  between 2.5 and 5.5 (step 0.5); effective adiabatic index (reddening) at a value of 1.25.  
384 A standard  $\chi^2$ -minimization procedure was used to find the best fitting model.

## 385 petitRADTRANS

386 petitRADTRANS, or pRT<sup>18</sup> (available from [https://petitradtrans.](https://petitradtrans.readthedocs.io)  
387 [readthedocs.io](https://petitradtrans.readthedocs.io)), is a Python package for the spectral synthesis of exoplanets and  
388 allows users to calculate transmission, emission or reflectance spectra. It offers a wide  
389 selection of opacities (gas line and continuum, and cloud opacities). Spectra can be calculated  
390 at any resolution, up to a wavelength spacing of  $\lambda/\Delta\lambda = 10^6$ . Coupled to a Bayesian

391 inference method such as PyMultiNest<sup>48,49</sup>, which pRT provides as a pre-implemented  
 392 retrieval package, posterior distributions for atmospheric parameters can be derived, given  
 393 an observation. For WISE J1828 we assumed a forward model setup as described for  
 394 the retrievals above, with the priors and forward model details set up as described in the  
 395 following. The prior on  $\log(g)$  was uniform from 2.5 to 6, while the radius prior ranged  
 396 from 0.5 to 3 Jupiter radii. In addition, the pressure-temperature profile was parameterized  
 397 by retrieving temperature values at 10 locations, equidistantly spaced in log-space between  
 398  $10^{-6}$  and 1000 bar, and then quadratically interpolating the temperature between these nodes  
 399 in  $\log(\text{pressure})$ . The priors were set up such that the temperature at 1000 bar was uniformly  
 400 sampled between 100 and 9000 K, and the temperatures at lower-pressure nodes was allowed  
 401 to be between 0.2-1.0 the temperature of the neighboring deeper atmospheric node. Since  
 402 kinks in the P-T profiles could be observed in the standard setup, the P-T regularization  
 403 described above was optionally turned on when deriving the atmospheric model posteriors,  
 404 fitting both MIRI/MRS and the archival HST WFC3 data. Our constraints on  $^{14}\text{N}/^{15}\text{N}$  are  
 405 not affected by the regularization. However, we observed a trend that a regularized P-T leads  
 406 to a higher atmospheric metal enrichment, higher gravity  $\log(g)$ , smaller radii, and higher  
 407 effective temperatures, see Extended Data Fig. 2. For the detection of  $^{15}\text{NH}_3$  we only used  
 408 the MRS data. We also turned the regularization off, to allow for maximum model flexibility.  
 409 This leads to a conservative estimate of the detection significance. The following opacity  
 410 species were included in the retrievals:  $\text{H}_2\text{O}$ <sup>50</sup>,  $\text{CH}_4$ <sup>51</sup>,  $\text{CO}$ <sup>52</sup>,  $\text{CO}_2$ <sup>53</sup>,  $^{15}\text{NH}_3$ <sup>54</sup>,  $\text{H}_2\text{S}$ <sup>55</sup>,  $\text{NH}_3$ <sup>56</sup>,  
 411 and  $\text{PH}_3$ <sup>57</sup>. The abundances of said molecules were retrieved using log-uniform priors from  
 412  $10^{-10}$  to 1 on their mass fractions. For the  $^{15}\text{NH}_3$  detection we used PyMultiNest, with 2000  
 413 live points, `constant_sampling_efficiency` set to False, and a sampling efficiency  
 414 of 0.3. The detection significance of  $^{15}\text{NH}_3$  was determined using a standard method<sup>58</sup>. We  
 415 report a detection significance of  $4.2 \sigma$  for  $^{15}\text{NH}_3$ . For the retrievals constraining the properties  
 416 of WISE J1828, which included HST in addition to MRS data, we ran in constant sampling  
 417 efficiency mode, with the efficiency set to 5 %. This needed to be done because otherwise  
 418 retrievals ran for  $10^8$  models, but did not finish. The partially filled weighted posterior files  
 419 of the  $10^8$ -model retrievals were consistent with the results using the constant sampling  
 420 efficiency mode. With petitRADTRANS we found  $^{14}\text{N}/^{15}\text{N} = 560_{-115}^{+165}$  for the flexible P-T  
 421 and  $^{14}\text{N}/^{15}\text{N} = 642_{-192}^{+365}$  for the regularized P-T model.

## 422 ARCIS

423 The ARTfull Modeling Code for exoplanet Science (ARCIS) is a forward modelling and re-  
 424 trieval code that can be used to analyse and simulate exoplanet atmosphere spectra. It con-  
 425 tains many physical and chemical processes including cloud formation<sup>59</sup> and disequilibrium  
 426 chemistry<sup>60</sup>. The free P-T structure used in this work is parameterized by the slope at several  
 427 pressure points in the atmosphere. We parameterize  $d\log T/d\log P$  with a prior range between  
 428  $-4/7$  to  $+4/7$ . The adiabatic gradient expected for a diatomic gas is  $d\log T/d\log P = 2/7$  so  
 429 this prior range gives a very generous range. We fix the absolute value of the temperature  
 430 structure at a pressure of  $P = 0.1$  bar. We retrieve the value of the slope at 12 pressure

431 points equally distributed over the atmosphere in  $\log -P$  space. For the detection of  $^{15}\text{NH}_3$   
432 we follow this procedure, allowing full flexibility and thereby constructing a conservative de-  
433 tection significance. For the final fits deriving the isotopic ratio and the planet parameters  
434 presented, we restrict the gradient of the P-T structure to be positive as expected for a non-  
435 irradiated atmosphere. Following a standard approach<sup>58</sup>, we find evidence that  $^{15}\text{NH}_3$  is present  
436 in the atmosphere of WISE J1828 at  $6.3 \sigma$ . We observe the same trend between metal enrich-  
437 ment,  $\log(g)$ , radii, and effective temperature as petitRADTRANS when regularizing the P-T  
438 structure, see Fig. 2. However, our derived  $^{14}\text{N}/^{15}\text{N}$  values are somewhat less stable when  
439 turning on regularization. Our regularized values are consistent with the petitRADTRANS  
440 (regularized and flexible P-T), namely  $^{14}\text{N}/^{15}\text{N} = 591^{+432}_{-190}$ . The value derived for the flexi-  
441 ble P-T setup is higher ( $^{14}\text{N}/^{15}\text{N} = 949^{+322}_{-208}$ ). More information on ARCiS can be found at  
442 <http://www.exoclouds.com>.

### 443 ARCiS+ML

444 For the self-consistent retrieval with ARCiS we assume a one-dimensional atmosphere in  
445 radiative-convective equilibrium. The retrievals were run on the MIRI MRS data only (i.e.,  
446 not including the HST data). The atmospheric composition is parameterized using [M/H],  
447 C/O and N/O, and we account for disequilibrium chemistry of  $\text{CH}_4$ - $\text{H}_2\text{O}$ -CO and  $\text{NH}_3$ - $\text{N}_2$  due  
448 to vertical mixing<sup>60</sup>, where the vertical eddy diffusion coefficient  $K_{zz}$  is another free param-  
449 eter. For WISE J1828 observed that, although  $\text{NH}_3$ - $\text{N}_2$  quenching was active in our models,  
450 it quenched from the lowest layer in the atmosphere where  $\text{NH}_3$  was still the dominating ab-  
451 sorber. Likely quenching actually occurs outside our simulated pressure domain, deeper inside  
452 the atmosphere. Therefore the ARCiS+ML constraints may be too low for N/O, similar to the  
453 constraints from the retrievals, which are insensitive to the spectrally inactive  $\text{N}_2$ . The radiative  
454 transfer module was benchmarked<sup>21</sup>. Due to the high computational load of the self-consistent  
455 models we cannot run nested sampling retrievals, which require millions of model evaluations  
456 to converge. Instead, we use a machine learning method based on SNPE (sequential neural  
457 posterior estimation<sup>61</sup>) that allows us to perform the retrieval using on the order of  $10^4$  models.  
458 The details of this retrieval method will be presented in an upcoming publication<sup>22</sup>.

### 459 Brewster

460 *Brewster*<sup>62,63</sup>, is a retrieval code that has mainly been employed in the context of explor-  
461 ing clouds in brown dwarfs. However, here, we utilised a simple cloud-free retrieval recipe.  
462 *Brewster* uses the two stream radiative transfer architecture<sup>64</sup>. We use the default 64 layer  
463 atmosphere with intervals of 0.1 dex across the pressures range  $\log P(\text{bar})$  of -4 to 2.3. The  
464 temperature is set using P-T parameterisation<sup>65</sup> linked to three atmospheric zones via expo-  
465 nential gradients. We included the molecules  $\text{H}_2\text{O}$ ,  $\text{CH}_4$ ,  $\text{NH}_3$ , CO,  $\text{PH}_3$  and  $\text{H}_2\text{S}$  which origi-  
466 nate from a compendium<sup>66,67</sup> and updated opacities<sup>62</sup>. The abundances of these molecules are  
467 modeled assuming vertically constant mixing ratios, which are retrieved as free parameters.

468 The opacities are ingested at a resolution of  $R=10,000$ , putting the native model resolution  
469 an order of magnitude above the data being fit. Continuum opacity's for  $H_2$ - $H_2$  and  $H_2$ -He  
470 collisionally induced absorption, Rayleigh scattering due to  $H_2$ , He and  $CH_4$ , and continuum  
471 opacities due to bound-free and free-free absorption by H- and free-free absorption by  $H_2$  are  
472 also included. We apply an error inflation "tolerance" framework<sup>45</sup> and used in all subsequent  
473 published works using *Brewster*. Only the JWST/MIRI MRS observations were retrieved  
474 with *Brewster*, so the HST data was not included. Here we used the EMCEE<sup>68</sup> as our sam-  
475 pling algorithm. As we used the standard retrieval recipe for this code, an extensive list of  
476 parameter sampling priors has been used<sup>63</sup>.

## 477 Combining model results

478 Whenever we report values for the properties of WISE J1828 in the main body of the text,  
479 these have been obtained from combining the posteriors of the retrievals at equal weight,  
480 and calculating the corresponding median and 16-84 %percentile values, corresponding to  
481 the  $1 \sigma$  credible interval in case the resulting distribution is approximately Gaussian. The  
482 combined value of the self-consistent codes was obtained by taking the average of their best-fit  
483 (ATMO) and median (ARciS+ML) values, while the uncertainty is obtained from calculating  
484 the difference  $d$  between these two values, and then calculating  $\sqrt{d^2 + a^2}$ , where  $a$  is the "1  $\sigma$ "  
485 uncertainty obtained from the ARciS+ML posterior.

## 486 Exploring $^{14}N/^{15}N$ as a formation tracer

487 To approximate the impact of volatile ice accretion between the  $NH_3$  and  $N_2$  icelines in a proto-  
488 planetary disk, we generated a simplified planet formation model for calculating  $^{14}N/^{15}N$   
489 as a function of the total solid mass (rock and ices) a planet incorporated during formation.  
490 For this we used a specific framework<sup>4</sup>, which is available at <https://gitlab.com/mauricemolli/formation-inversion>. In short, we used a solar disk composition<sup>4</sup>,  
491 their Table 2. This means that the mass ratio between  $NH_3$  and  $N_2$  is 1:7 in the disk (combining  
492 gas and ice reservoirs). We then assumed, conservatively, that  $^{14}N/^{15}N$  is reduced by a factor  
493 of 2 in  $NH_3$ , when compared to the total  $^{14}N/^{15}N$  value, while the total  $^{14}N/^{15}N$  (summing  
494 over all species and phases) is conserved, which we assumed to be 300, and call ISM value  
495 in the in Extended Data Fig. 4. This figure shows the ratio of the planetary and ISM values of  
496  $^{14}N/^{15}N$  as a function of solids accreted by a planet between the  $NH_3$  and  $N_2$  icelines. Since  
497 the solids are rich in  $NH_3$ , and  $N_2$  is only in the gas phase, a higher accreted solid mass results  
498 in a lower planetary  $^{14}N/^{15}N$ . We note that the picture could be more complicated, since the  
499 disk gas could be enriched in  $^{15}N$ -poor  $N_2$  gas that evaporated off pebbles that drift in from  
500 outside the  $N_2$  iceline<sup>69</sup>.

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570 **Data availability** The JWST MIRI data presented in this paper are part of the JWST MIRI GTO pro-  
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572 Barbara A. Mikulski Archive for Space Telescopes (MAST; <https://archive.stsci.edu/>) after  
573 July 28th, 2023, and can be found either using the program identifier or using the Data Object Identifier  
574 (DOI): <https://doi.org/10.17909/as3s-x893>. The HST WFC3 spectrum is available from:  
575 <https://cdsarc.cds.unistra.fr/viz-bin/cat/J/ApJ/920/20#/article>.

576 **Code availability** The codes used in this publication to extract, reduce, and analyse the data are as fol-  
577 lows; The data reduction pipeline `jwst` can be found at <https://jwst-pipeline.readthedocs.io/en/latest/>. The atmospheric model codes used to fit the data can be found at <https://www.exoclouds.com/>  
578 for the ARCIS-code<sup>19</sup> and at <https://petitradtrans.readthedocs.io/en/latest/>  
579 for the petitRADTRANS-code<sup>18</sup>. The simplified planet formation model<sup>4</sup> used to study  
580  $^{14}\text{N}/^{15}\text{N}$  as a function of accreted ice mass can be found at <https://gitlab.com/mauricemolli/formation-inversion>. The detailed setups of the open source tools for the analyses presented here  
581 are described in the methods section of this paper, and can be made available to interested parties upon  
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583  
584

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652 **Competing Interests** The authors declare that they have no competing financial interests.

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## 656 **Extended Data**

### 657 **Figure Legends extended data**

658 **Legend Figure 1: The spectrum of WISE J1828 (black solid lines) and best-fit model.** We show  
659 the MIRI/MRS spectrum of WISE J1828 (black solid lines) and best-fit model of the regularized-P-  
660 T retrieval of petitRADTRANS (red line). Residuals (models – observed spectrum) are displayed  
661 at the bottom panel. pRT-reg and ARCiS-reg stand for the regularized P-T retrieval of petitRAD-  
662 TRANS and ARCiS, respectively.

663 **Legend Figure 2: Model inferences on the various P-T profiles derived for WISE J1828.** The  
664 individual panels always highlight the constraint from one given model, while the results of the  
665 other models are shown in the background. pRT-reg and ARCiS-reg stand for the regularized P-T  
666 retrievals, while pRT-free and ARCiS-free stand for the unregularized P-T retrievals of petitRAD-  
667 TRANS and ARCiS, respectively. The contribution functions of the HST and MIRI observations,  
668 constrained from the best-fit pRT-reg model, are shown as dotted and dashed lines, respectively.  
669 The condensation curve for water (at solar metallicity) is shown as a blue dashdotted curve, indi-  
670 cating that, while neglected in our models, water clouds could impact the spectrum in a modest  
671 away.

672 **Legend Figure 3: One-dimensional projection of the posterior distributions of the WISE J1828**  
673 **retrievals.** Values correspond to key atmospheric quantities shown in Extended Data Table 1. pRT-  
674 reg and ARCiS-reg stand for the regularized P-T retrievals, while pRT-free and ARCiS-free stand  
675 for the unregularized P-T retrievals of petitRADTRANS and ARCiS, respectively.

676 **Legend Figure 4: Evolution of the planetary  $^{14}\text{N}/^{15}\text{N}$  as a function of the mass accreted in**  
677 **solids (rock and ice).** This computation assumes a planet that forms outside the  $\text{NH}_3$  ice line but  
678 inside the  $\text{N}_2$  ice line. The black dotted line denotes the value expected for pure  $\text{NH}_3$  ice.

### 679 **Table Legends extended data**

#### 680 **Legend Table 1: Physical constraints on WISE J1828**

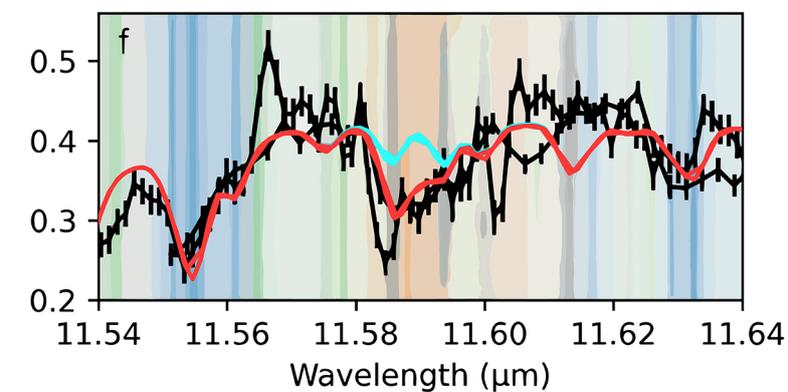
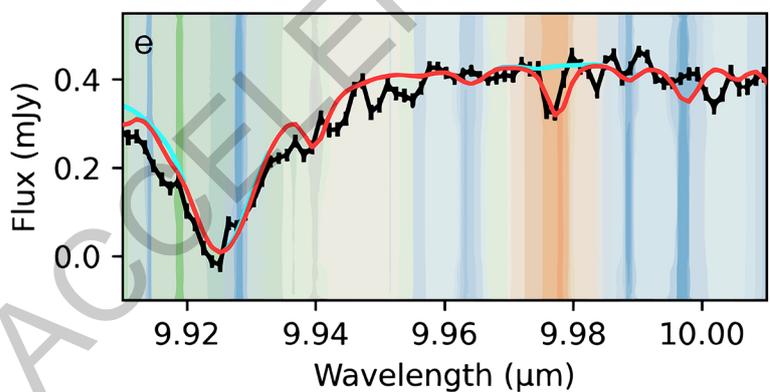
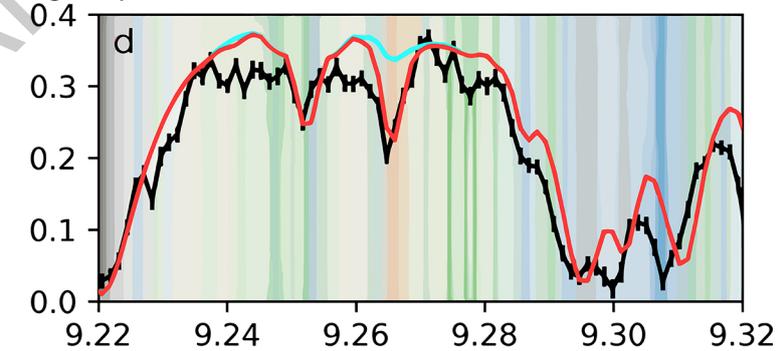
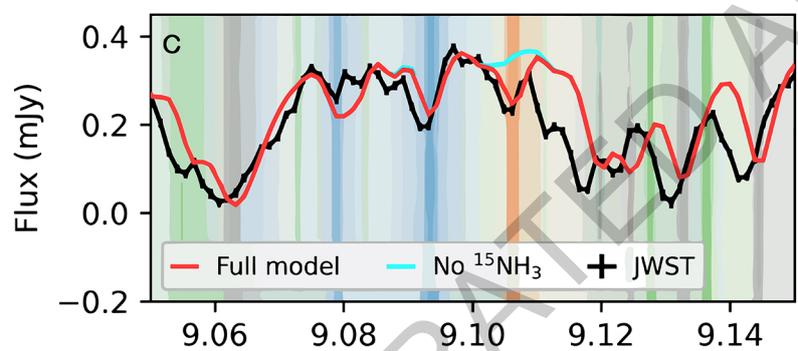
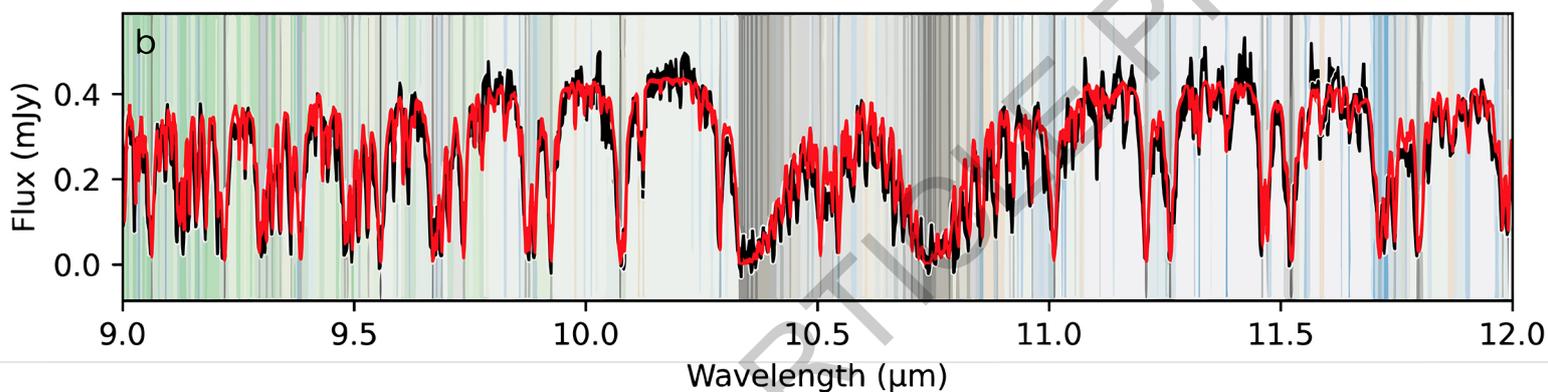
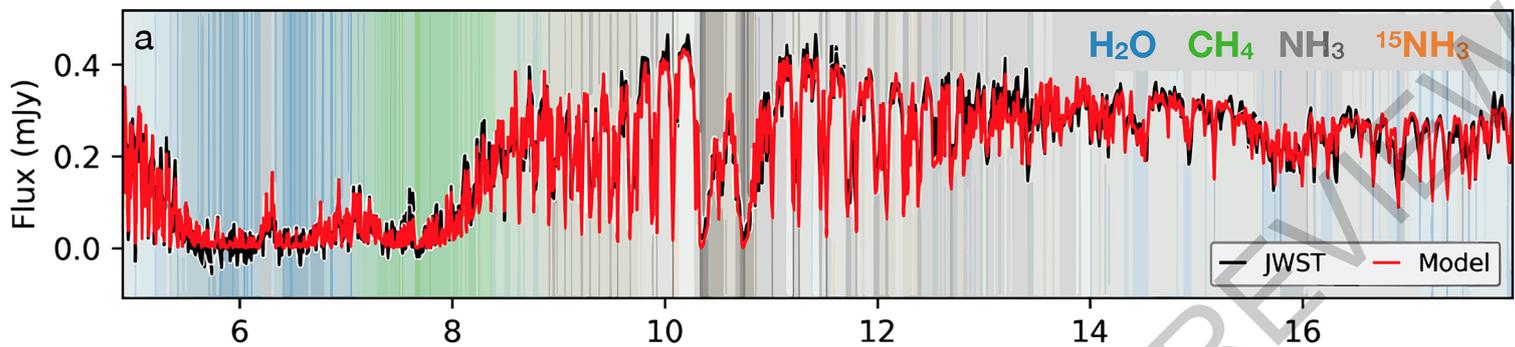
681 “free” in the code name means that the P-T structure was not regularized, while this was done  
682 for the “reg” cases.  $R_{\text{bin}}$  model radii have been multiplied by  $1/\sqrt{2}$ , assuming WISE J1828 is an  
683 equal-property binary.  $R_{\text{bin}}$  was used for calculating  $M_{\text{bin}}$  from the inferred gravity. The units of  
684  $K_{zz}$  are  $\text{cm}^2 \text{s}^{-1}$ . “chem” means that absorber abundances have been determined from a chemical  
685 model. “solar” means that the parameter was not varied, and that a solar composition was assumed  
686 instead.

#### 688 **Legend Table 2: Physical constraints on WISE J1828, combining different codes**

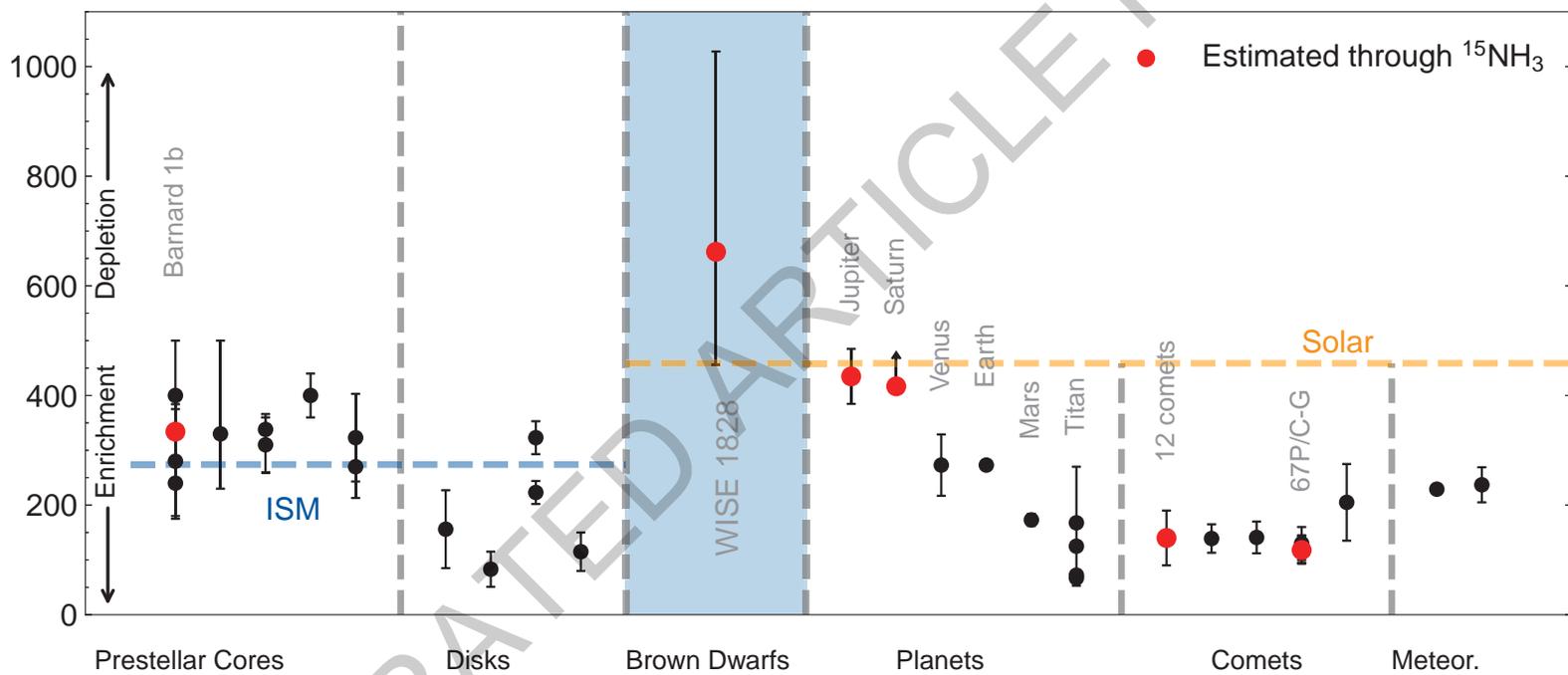
689  $R_{\text{bin}}$  model radii have been multiplied by  $1/\sqrt{2}$ , assuming WISE J1828 is an equal-property bi-  
690 nary.  $R_{\text{bin}}$  was used for calculating  $M_{\text{bin}}$  from the inferred gravity. The units of  $K_{zz}$  are  $\text{cm}^2 \text{s}^{-1}$ .  
691 “chem” means that absorber abundances have been determined from a chemical model. “no comb”  
692 means that the respective parameter was only varied in one of the two codes, please see Extended  
693 Data Table 1 for the inferred values. \* for the inferred average mass denotes that the distance be-  
694 tween the two codes’ solutions was larger than the average value, so no uncertainties, derived as  
695 explained in the method section, are given here.

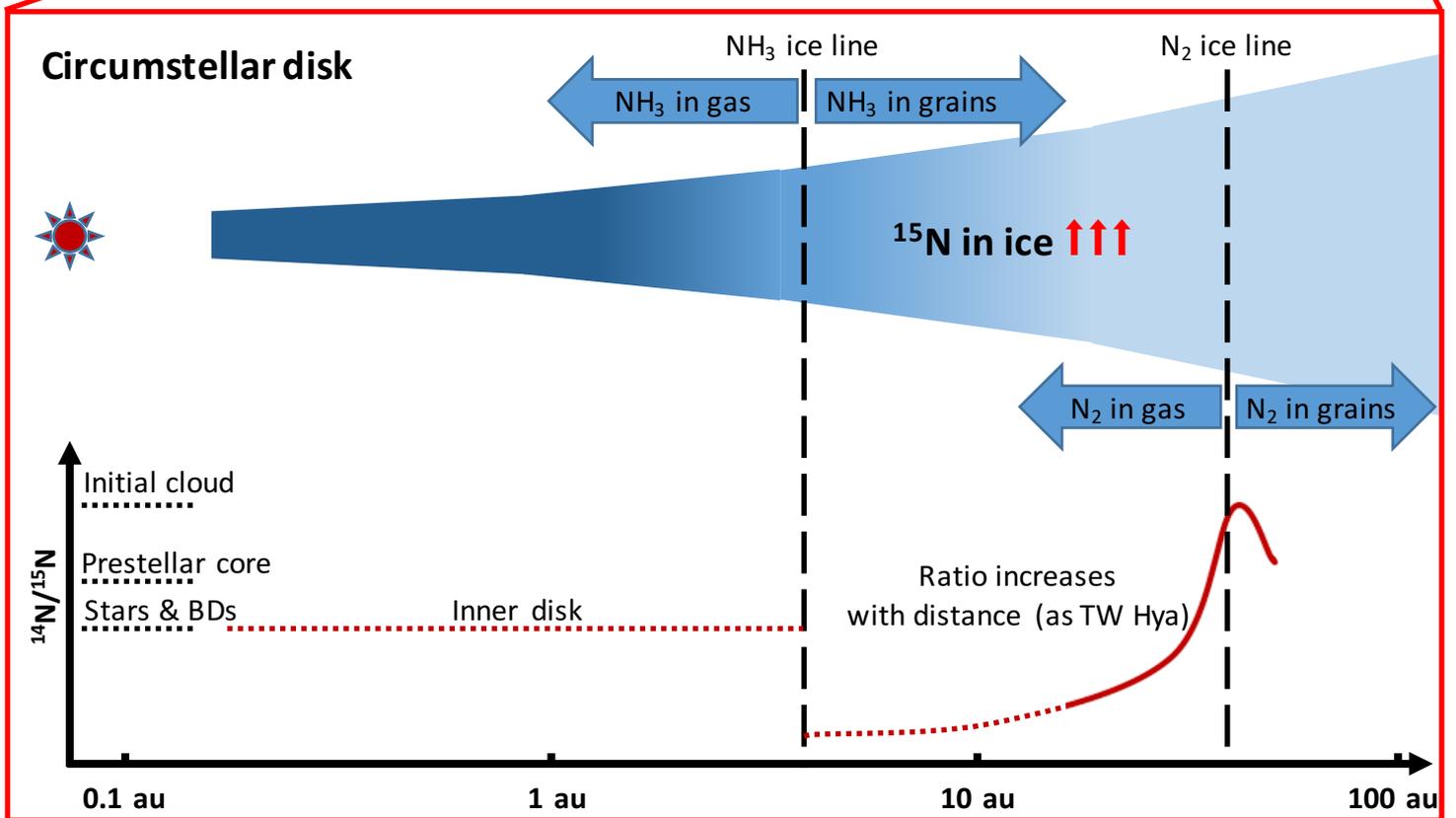
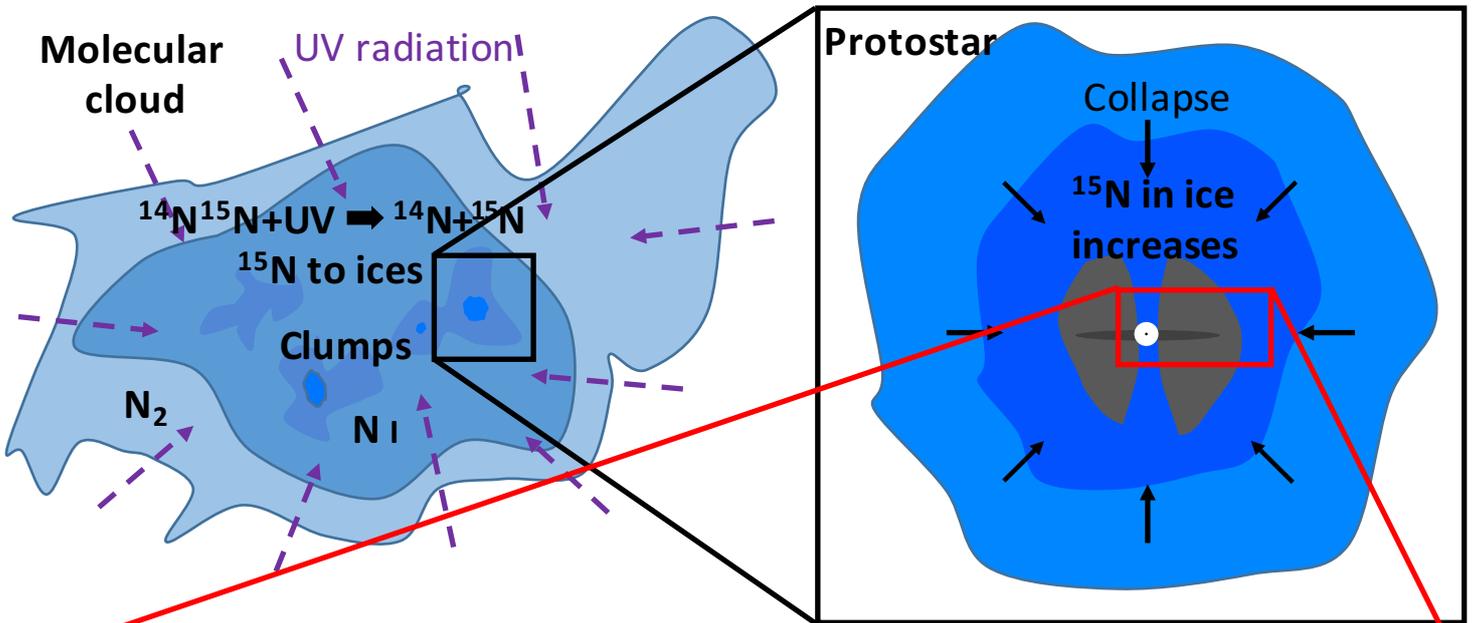
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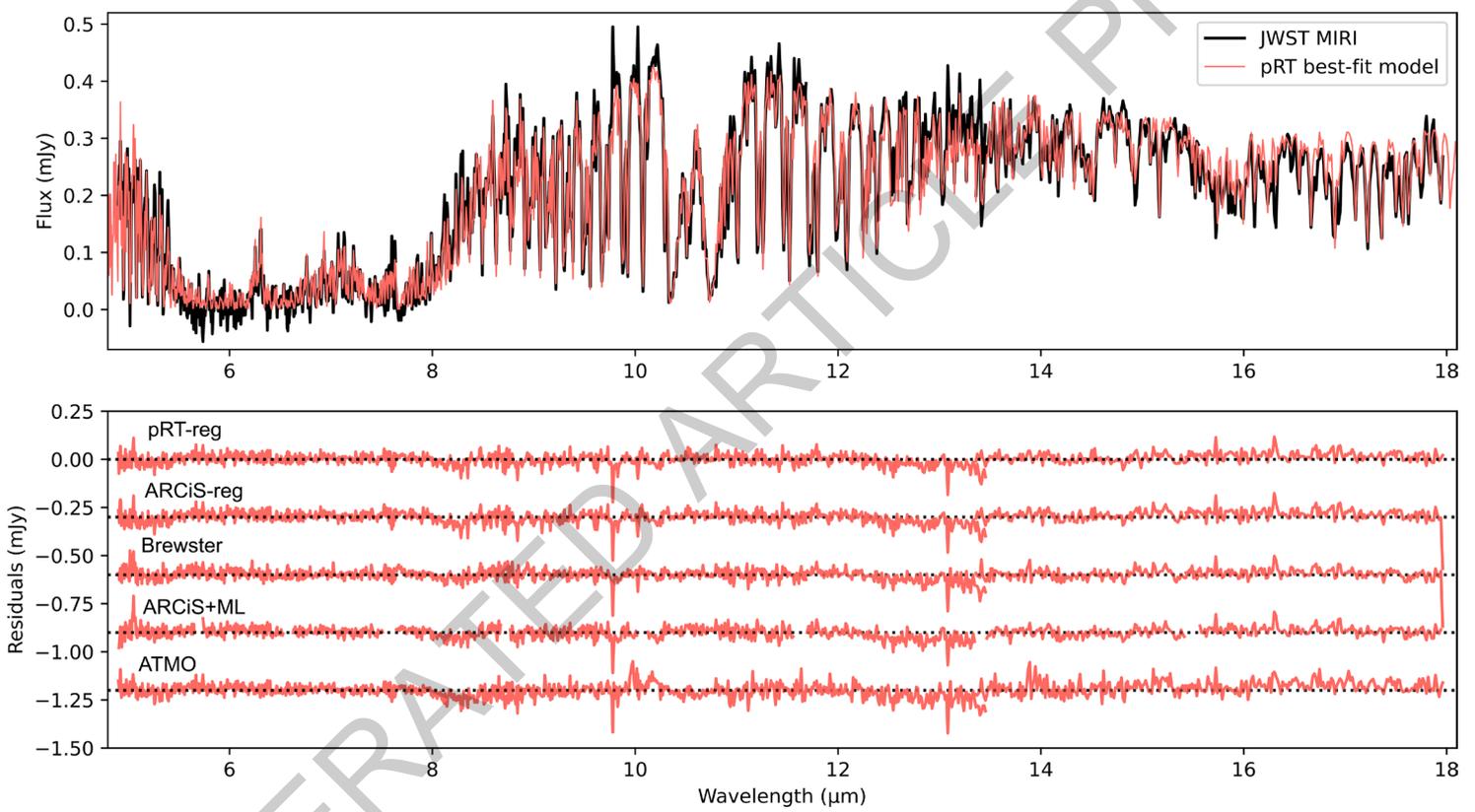
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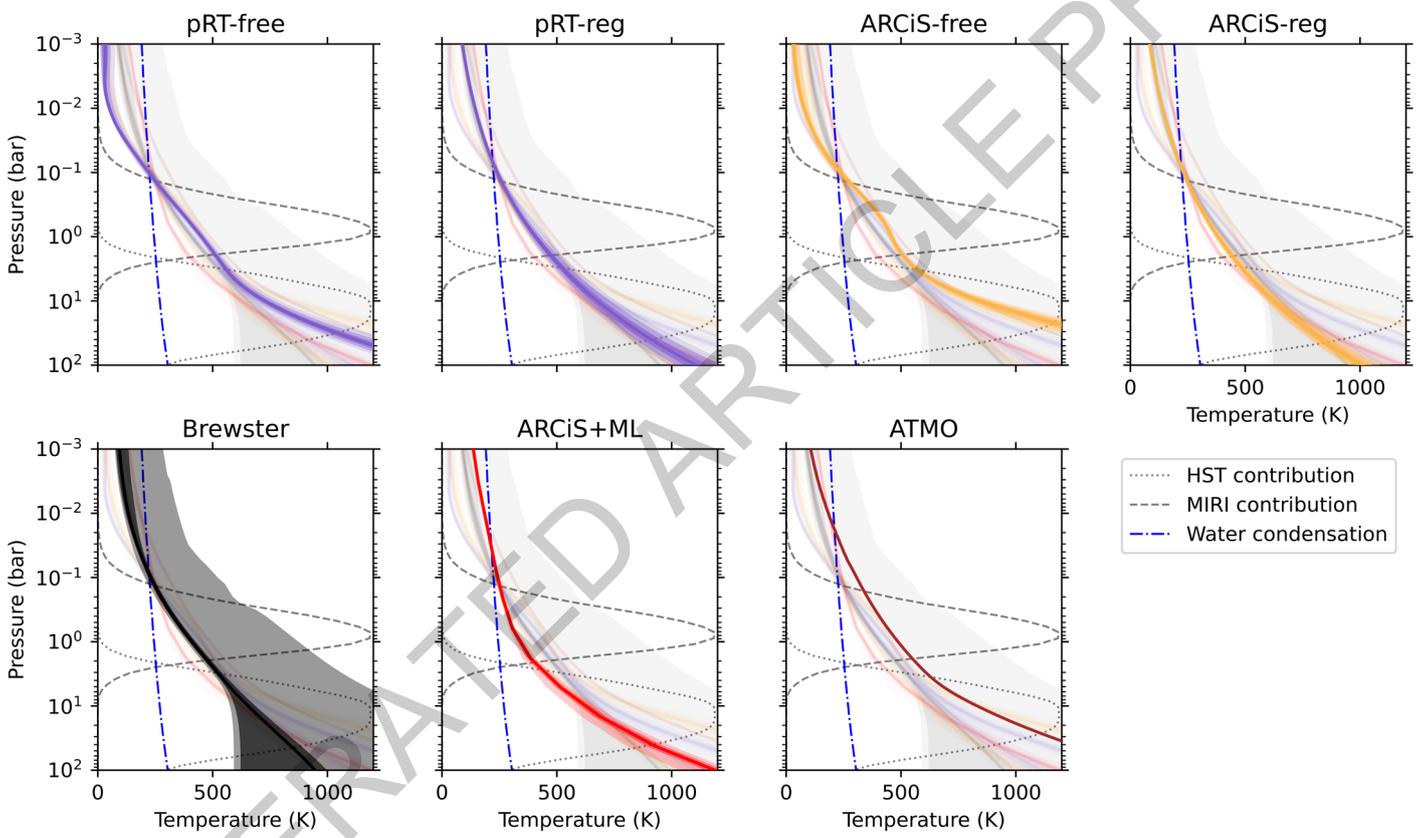
$^{14}\text{N}/^{15}\text{N}$



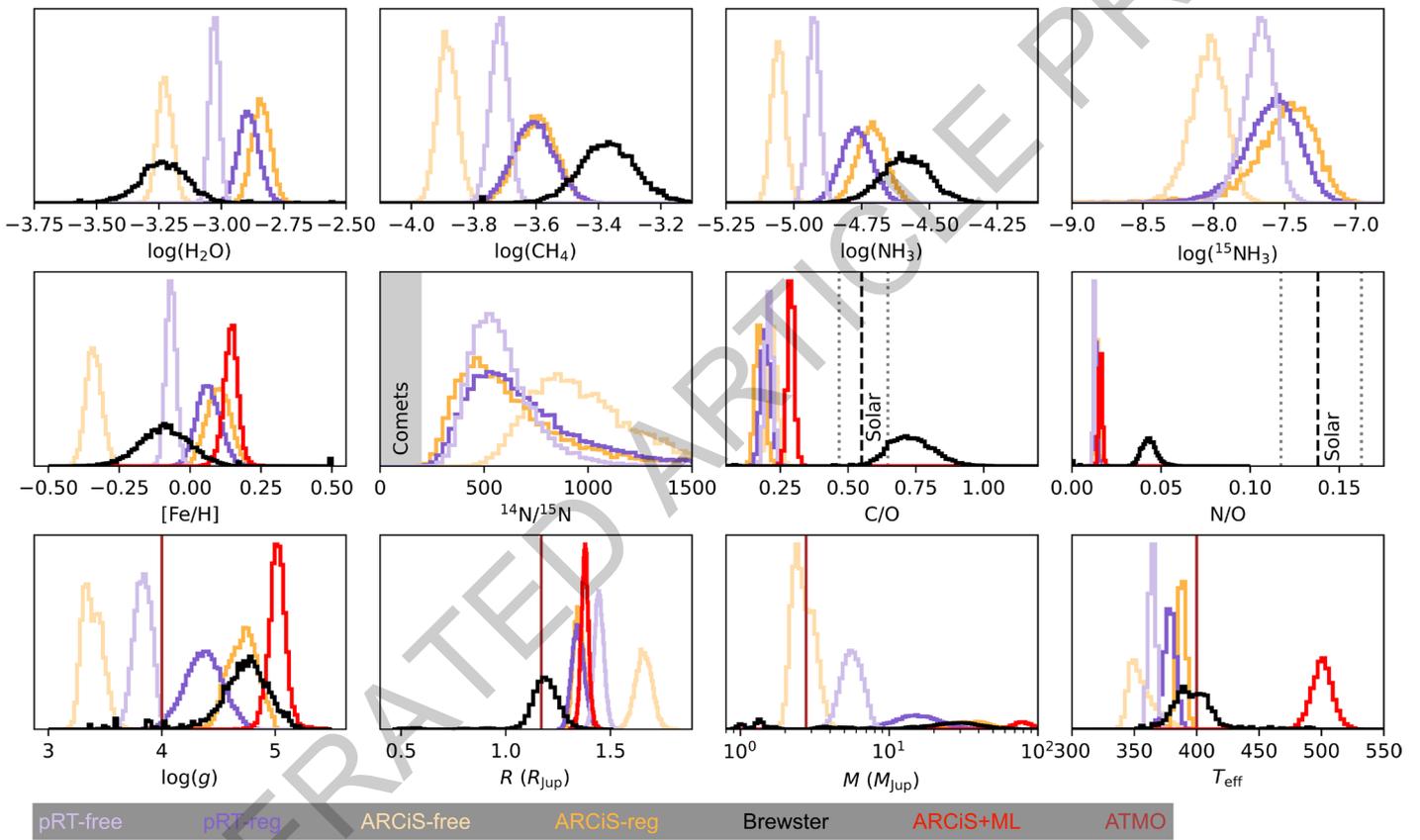




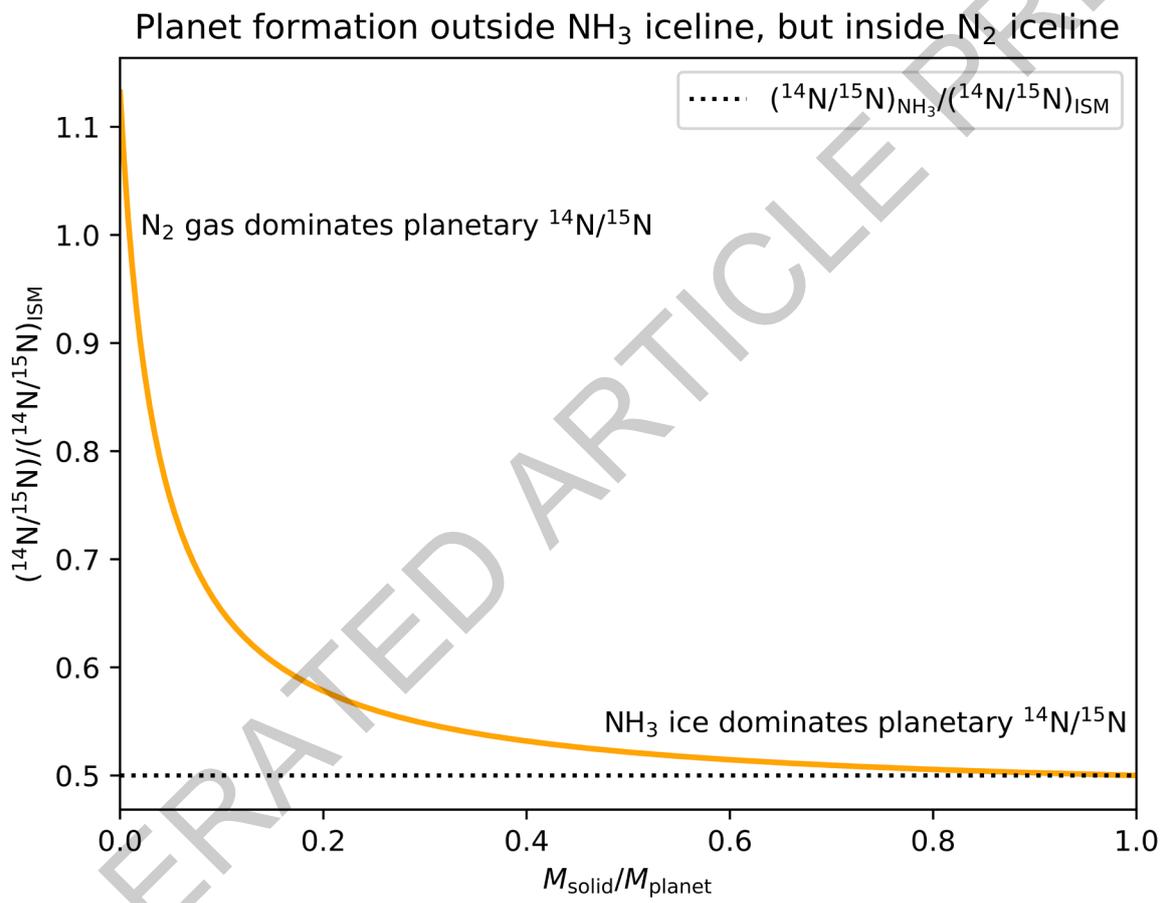
Extended Data Fig. 1



**Extended Data Fig. 2**



Extended Data Fig. 3



Extended Data Fig. 4

Quantity	Retrieval codes					Self-consistent codes	
	pRT-free	pRT-reg	ARCiS-free	ARCiS-reg	Brewster	ARCiS+ML	ATMO
$T_{\text{eff}}$ (K)	$364^{+3}_{-3}$	$379^{+4}_{-4}$	$354^{+13}_{-7}$	$388^{+3}_{-4}$	$397^{+16}_{-14}$	$501^{+8}_{-8}$	400
$\log(g)$	$3.83^{+0.08}_{-0.08}$	$4.37^{+0.16}_{-0.17}$	$3.38^{+0.1}_{-0.08}$	$4.72^{+0.12}_{-0.14}$	$4.74^{+0.19}_{-0.21}$	$5.02^{+0.07}_{-0.07}$	4.0
$R$ ( $R_{\text{Jup}}$ )	$1.45^{+0.02}_{-0.02}$	$1.35^{+0.03}_{-0.03}$	$1.65^{+0.04}_{-0.04}$	$1.35^{+0.03}_{-0.03}$	$1.19^{+0.06}_{-0.06}$	$1.38^{+0.02}_{-0.02}$	1.17
$M$ ( $M_{\text{Jup}}$ )	$5.76^{+1.04}_{-0.89}$	$17.27^{+6.86}_{-5.01}$	$2.69^{+0.58}_{-0.40}$	$38.57^{+10.92}_{-9.66}$	$31.50^{+14.30}_{-11.60}$	$79.99^{+12.79}_{-10.87}$	2.76
$R_{\text{bin}}$ ( $R_{\text{Jup}}$ )	$1.02^{+0.02}_{-0.02}$	$0.95^{+0.02}_{-0.02}$	$1.17^{+0.03}_{-0.03}$	$0.95^{+0.02}_{-0.02}$	$0.84^{+0.04}_{-0.04}$	$0.98^{+0.01}_{-0.01}$	0.83
$M_{\text{bin}}$ ( $M_{\text{Jup}}$ )	$2.88^{+0.52}_{-0.45}$	$8.63^{+3.43}_{-2.51}$	$1.34^{+0.29}_{-0.2}$	$19.29^{+5.46}_{-4.83}$	$15.75^{+7.15}_{-5.8}$	$39.99^{+6.4}_{-5.43}$	1.38
[M/H]	$-0.07^{+0.02}_{-0.02}$	$0.06^{+0.04}_{-0.04}$	$-0.34^{+0.03}_{-0.03}$	$0.10^{+0.04}_{-0.04}$	$-0.08^{+0.10}_{-0.09}$	$0.14^{+0.02}_{-0.03}$	solar
C/O	$0.21^{+0.01}_{-0.01}$	$0.20^{+0.02}_{-0.02}$	$0.22^{+0.02}_{-0.02}$	$0.17^{+0.02}_{-0.02}$	$0.73^{+0.08}_{-0.10}$	$0.29^{+0.01}_{-0.01}$	solar
(N/O) / (N/O) <sub>⊙</sub>	$0.091^{+0.004}_{-0.004}$	$0.096^{+0.006}_{-0.006}$	$0.107^{+0.006}_{-0.006}$	$0.097^{+0.006}_{-0.006}$	$0.314^{+0.036}_{-0.032}$	$0.117^{+0.006}_{-0.007}$	solar
$^{14}\text{N}/^{15}\text{N}$	$560^{+165}_{-115}$	$642^{+365}_{-192}$	$949^{+322}_{-208}$	$591^{+432}_{-190}$	–	–	–
$\log(\text{H}_2\text{O})$	$-3.03^{+0.02}_{-0.02}$	$-2.89^{+0.04}_{-0.04}$	$-3.23^{+0.03}_{-0.03}$	$-2.84^{+0.04}_{-0.04}$	$-3.23^{+0.10}_{-0.10}$	chem	chem
$\log(\text{CH}_4)$	$-3.72^{+0.03}_{-0.03}$	$-3.61^{+0.07}_{-0.06}$	$-3.88^{+0.03}_{-0.03}$	$-3.60^{+0.06}_{-0.06}$	$-3.36^{+0.10}_{-0.09}$	chem	chem
$\log(\text{NH}_3)$	$-4.93^{+0.02}_{-0.02}$	$-4.77^{+0.06}_{-0.06}$	$-5.06^{+0.03}_{-0.02}$	$-4.71^{+0.05}_{-0.05}$	$-4.58^{+0.10}_{-0.10}$	chem	chem
$\log(^{15}\text{NH}_3)$	$-7.67^{+0.10}_{-0.11}$	$-7.58^{+0.17}_{-0.21}$	$-8.03^{+0.11}_{-0.13}$	$-7.49^{+0.18}_{-0.25}$	–	–	–
$\log(K_{zz})$	–	–	–	–	–	$9.01^{+0.19}_{-0.19}$	7

Extended Data Table 1

Quantity	Retrievals	Self-consistent codes
$T_{\text{eff}}$ (K)	$378^{+13}_{-18}$	$450 \pm 101$
$\log(g)$	$4.34^{+0.42}_{-0.88}$	$4.5 \pm 1.0$
$R$ ( $R_{\text{Jup}}$ )	$1.37^{+0.26}_{-0.13}$	$1.27 \pm 0.21$
$M$ ( $M_{\text{Jup}}$ )	$15.8^{+22.82}_{-12.64}$	41*
$R_{\text{bin}}$ ( $R_{\text{Jup}}$ )	$0.97^{+0.18}_{-0.09}$	$0.90 \pm 0.15$
$M_{\text{bin}}$ ( $M_{\text{Jup}}$ )	$7.91^{+11.33}_{-6.34}$	21*
[M/H]	$-0.05^{+0.15}_{-0.27}$	no comb
C/O	$0.21^{+0.45}_{-0.03}$	no comb
(N/O) / (N/O) $_{\odot}$	$0.10^{+0.19}_{-0.01}$	no comb
$^{14}\text{N}/^{15}\text{N}$	$673^{+393}_{-212}$	–
$\log(\text{H}_2\text{O})$	$-3.03^{+0.18}_{-0.21}$	chem
$\log(\text{CH}_4)$	$-3.65^{+0.21}_{-0.21}$	chem
$\log(\text{NH}_3)$	$-4.79^{+0.15}_{-0.25}$	chem
$\log(^{15}\text{NH}_3)$	$-7.68^{+0.24}_{-0.34}$	–
$\log(K_{zz})$	–	$8 \pm 2$

Extended Data Table 2