



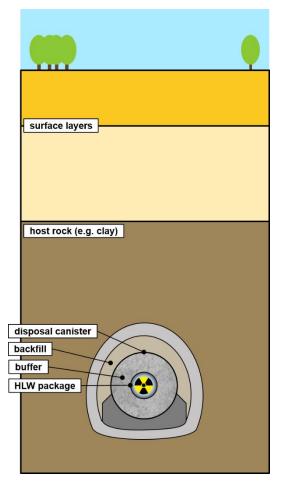
## Multiscale modelling of gas flow in Boom Clay

5<sup>th</sup> International Symposium on Unsaturated soil Mechanics and Waste Disposal 15 – 17 September 2023 • Tongji University (China)

Gilles Corman, Frédéric Collin

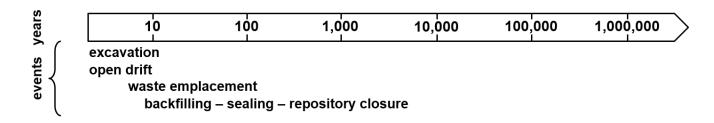


The project leading to this application has received funding from the European Union's Horizon 2020 research and innovation programme under grant agreement n° 847593.

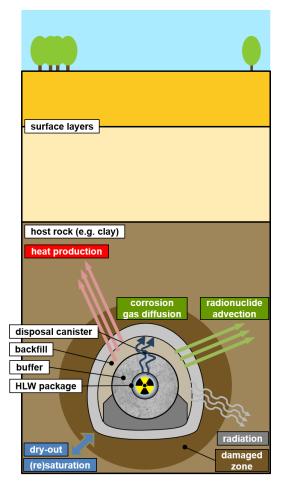


Conceptual scheme of a deep geological repository.

#### Geological disposal of radioactive wastes



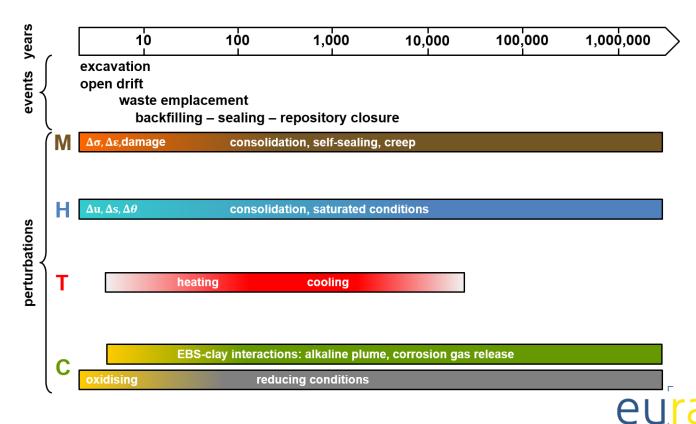


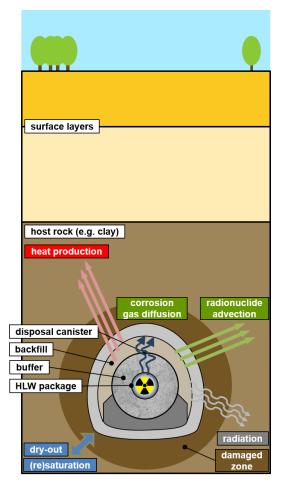


Conceptual scheme of a deep geological repository.

#### Geological disposal of radioactive wastes

Complex multi-physical (THMC) processes

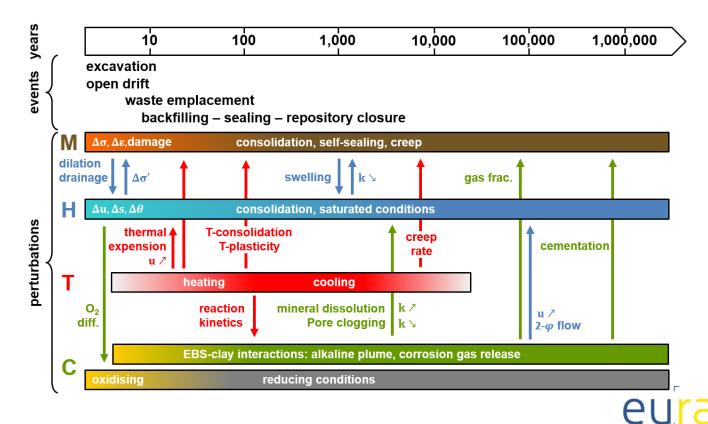


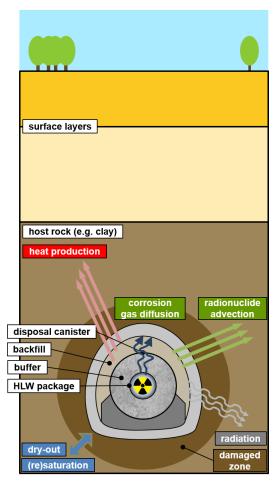


Conceptual scheme of a deep geological repository.

#### Geological disposal of radioactive wastes

- Complex multi-physical (THMC) processes
- Interactions between processes

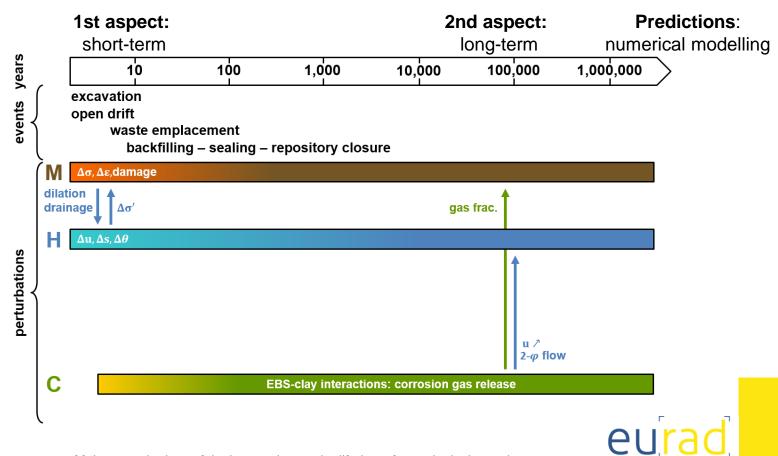




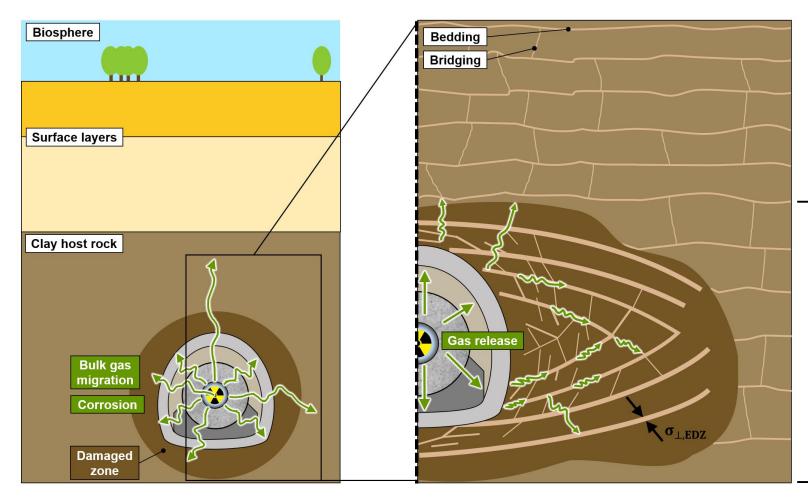
Conceptual scheme of a deep geological repository.

#### Geological disposal of radioactive wastes

- ► Complex multi-physical (**THMC**) processes
- ► Interactions between processes



#### Gas migration issue



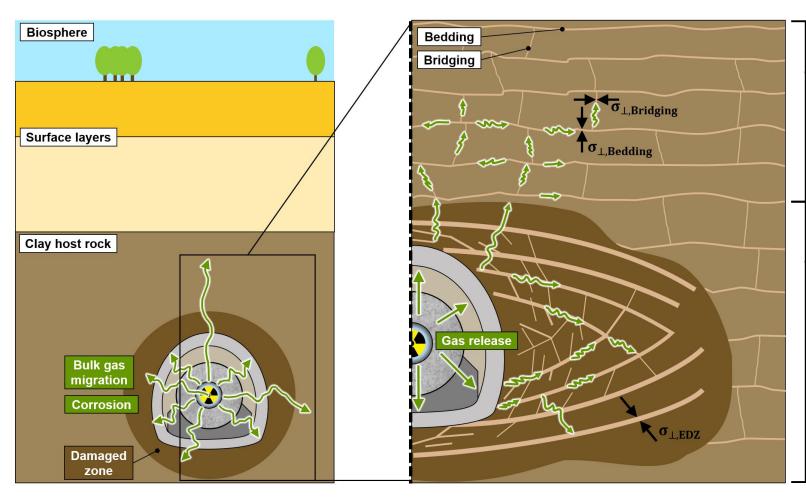
### Expected gas transport modes in the EDZ and the sound rock, from ONDRAF/NIRAS (2016).

### **Excavation damaged zone (EDZ)**

 Governed by the hydraulic properties modifications induced by fracturation



#### Gas migration issue



### Expected gas transport modes in the EDZ and the sound rock, from ONDRAF/NIRAS (2016).

#### Sound rock layers

- Governed by the rock structure <u>at a micro-level</u>
- ▶ Multi-Scale Model

#### **Excavation damaged zone (EDZ)**

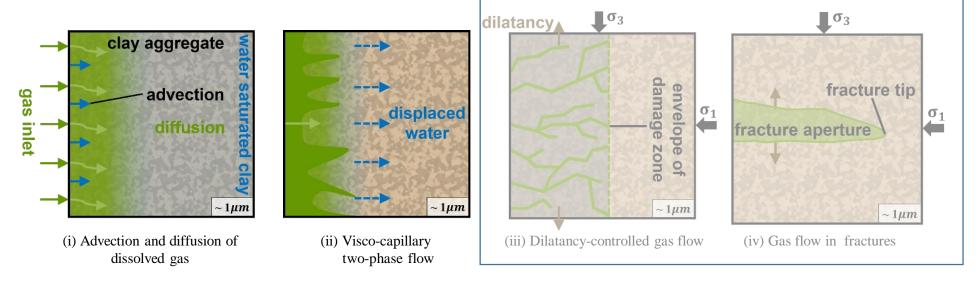
 Governed by the hydraulic properties modifications induced by fracturation



### Content

- Context
- 2 From experimental evidence to modelling
- Multi-scale modelling approach
- 4 Modelling gas injection experiment
- **6** Conclusions

#### **Background**



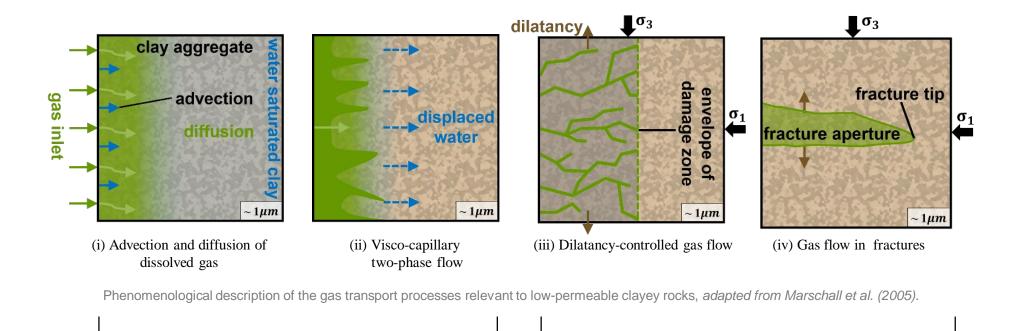
Phenomenological description of the gas transport processes relevant to low-permeable clayey rocks, adapted from Marschall et al. (2005).

**Classical HM two-phase flow models** 



**Classical HM two-phase flow models** 

#### **Background**

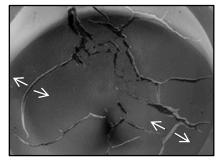


Supported by experimental data

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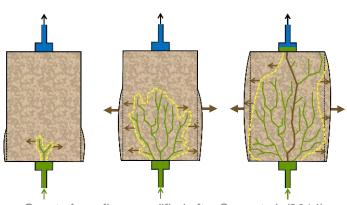
#### **Laboratory experiments**

#### **Clay-rich material**



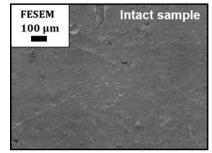
Gas-induced fracturing, Wiseall et al. (2015)

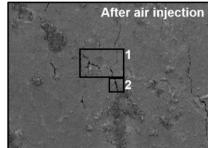
#### **Callovo-Oxfordian claystone**

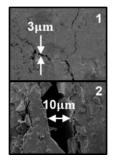


Onset of gas flow, modified after Cuss et al. (2014)

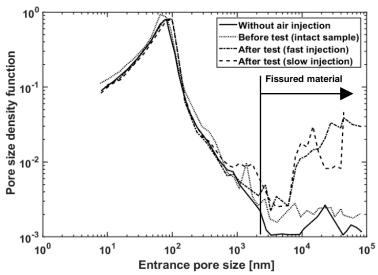
#### **Boom Clay**





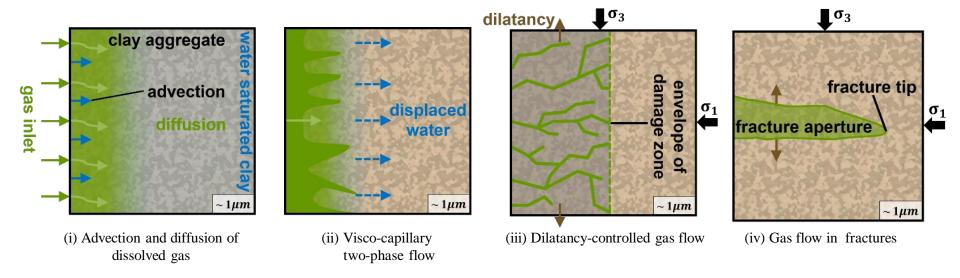


Changes in Boom Clay pore size distribution after air injection, and corresponding FESEM images with zooms on the detected fissures, modified after Gonzalez-Blanco et al. (2022)





#### **Background**



Phenomenological description of the gas transport processes relevant to low-permeable clayey rocks, adapted from Marschall et al. (2005).

Classical HM two-phase flow models

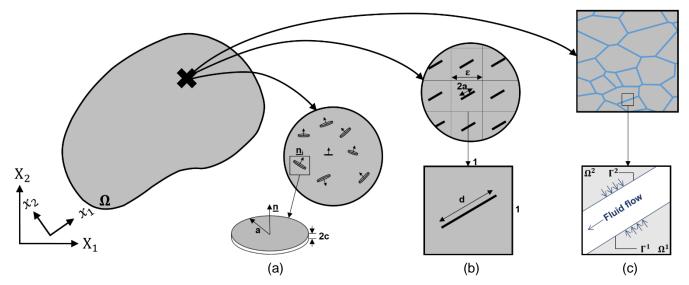
#### Supported by experimental data

- Natural heterogeneities represent preferred weaknesses for the process of opening discrete gas-filled pathway
- Introduce stronger coupling between gas flow and mechanical behaviour into the models.
  - gas flow and mechanical

#### **Advanced HM models**

#### Micro-macro based models

- ► Combines the benefits from large- and small-scale modelling strategies
- Explicit description of all the constituents on their specific length scale through a REV definition



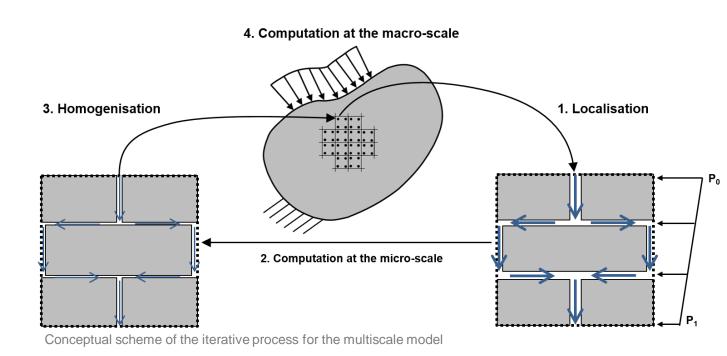


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#### **Overview**

- Macro-to-micro scale transition: Localisation of the macro-scale deformations to the micro-scale
- Resolution of the boundary value problem at the micro-scale
- Micro-to-macro scale transition: Homogenisation of the micro-scale stresses to compute the macroscopic quantities
- Resolution of the boundary value problem at the macro-scale



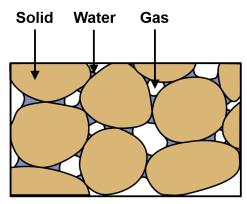
#### **Hybrid developed tool**

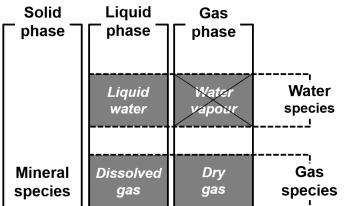
- Complete hydraulic system implemented and solved at the micro-scale
- Mechanical effects addressed at the macro-scale and implicitly integrated at the lower scale through HM couplings



#### Model formulation at the macroscopic scale

Clay material treated as a porous medium





**Balance** equations

Momentum

$$\frac{\partial \sigma_{ij}}{\partial x_j} + \rho g_i = 0$$

Water

$$\dot{M}_w + \frac{\partial f_{w,i}}{\partial x_i} - Q_w = 0$$
Liquid water

Gas

$$\underbrace{\dot{M}_{g} + \frac{\partial f_{g,i}}{\partial x_{i}}}_{\text{Dry gas}} + \underbrace{\dot{M}_{dg} + \frac{\partial f_{dg,i}}{\partial x_{i}}}_{\text{Dissolved gas}} - Q_{g} = 0$$

Constitutive equations

Total stress definition

$$\sigma_{ij} = \sigma'_{ij} + b_{ij} \left[ S_{r_w} p_w^M + (1 - S_{r_w}) p_g^M \right] \delta_{ij}$$

Variation of solid density

$$\frac{\dot{\rho}_{s}}{\rho_{s}} = \frac{(b_{ij} - \phi)(S_{r}^{w}\dot{p}_{w} + S_{r}^{g}\dot{p}_{g}) + \dot{\sigma}'}{(1 - \phi)K_{s}}$$



#### Macro-to-micro scale transition: Localisation

Decomposition of the micro-kinematics:

• Macro-pressure fields  $(\Box^M)$  of water and gas must be identical to the micro-quantities  $(\Box^m)$  for any point of the material

$$p_w^M(\hat{P}) = p_w^m(\hat{P}) \qquad \qquad p_g^M(\hat{P}) = p_g^m(\hat{P})$$

• For any point P close to  $\hat{P}$ , at the macroscopic scale:

$$p_w^M(P) \approx p_w^M(\hat{P}) + \frac{\partial p_w^M(\hat{P})}{\partial x_j} \left( x_j - \hat{x}_j \right) \qquad p_g^M(P) \approx p_g^M(\hat{P}) + \frac{\partial p_g^M(\hat{P})}{\partial x_j} \left( x_j - \hat{x}_j \right)$$

Higher-order terms neglected

at the microscopic scale:

$$p_w^m(P) \approx p_w^M(\hat{P}) + \frac{\partial p_w^M(\hat{P})}{\partial x_j} \left( x_j - \hat{x}_j \right) + p_w^f(\hat{P}) \qquad p_g^m(P) \approx p_g^M(\hat{P}) + \frac{\partial p_g^M(\hat{P})}{\partial x_j} \left( x_j - \hat{x}_j \right) + p_g^f(\hat{P})$$

#### Separation of scales

 Approach restricted to situations where the variations of the macroscopic fields is large compared to the variations of micro-scale fields

$$\frac{\partial p_w^M(\hat{P})}{\partial x_j} \left( x_j - \hat{x}_j \right) + p_w^f(\hat{P}) \ll p_w^M(\hat{P})$$

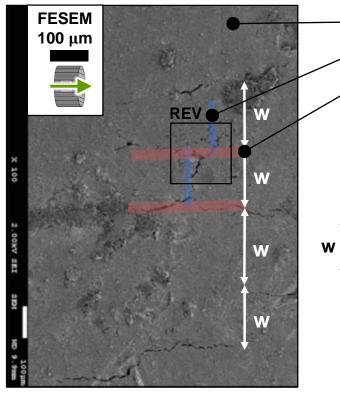
$$\frac{\partial p_g^M(\hat{P})}{\partial x_j} \left( x_j - \hat{x}_j \right) + p_g^f(\hat{P}) \ll p_g^M(\hat{P})$$



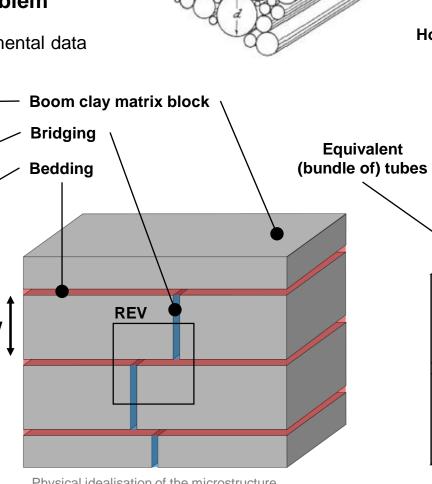
### Multi-scale modelling

#### Micro-scale boundary value problem

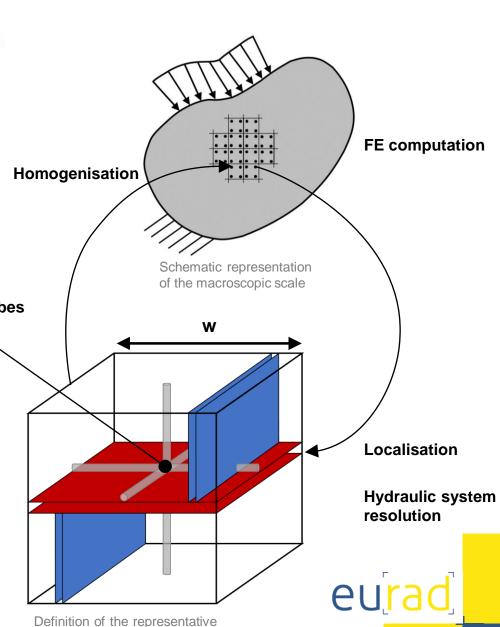
Multi-scale model supported by experimental data



Internal visualisation of a Boom Clay sample using FESEM, from Gonzalez-Blanco (2017).



Physical idealisation of the microstructure.



element volume (REV)

#### Micro-scale boundary value problem

Balance equations at the micro-scale

Gas

$$+ \frac{\partial f_{g_i}^m}{\partial x_i} + \sum_{i=1}^m + \frac{\partial f_{dg_i}^m}{\partial x_i} = 0$$

Water

$$+ \frac{\partial f_{w_i}^m}{\partial x_i} = 0$$

$$\dot{M}_g^m \dot{M}_{dg}^m \dot{M}_w^m$$
 Variations of fluid contents

$$\begin{split} f_{w_i}^m &= \rho_w q_{w_i} \\ f_{g_i}^m &= \rho_g q_{g_i} & \text{Mass flows} \\ f_{dg_i}^m &= \rho_{dg} q_{w_i} + i_{dg_i} \end{split}$$

 Mechanical effects: computed at the macro-scale and transferred to the micro-scale through <u>HM couplings</u>

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#### **Characterisation of the microstructure parameters**

1. Size of the REV

2. Macroporosity

Bedding plane separation  $w = 300 \, \mu m$ 

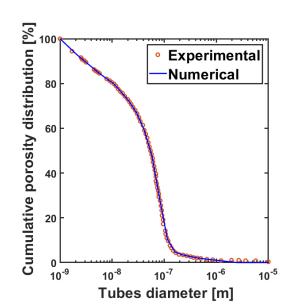
Fitting of the pore size distribution

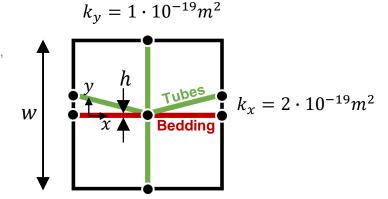
(Tortuosity = Calibration factor)

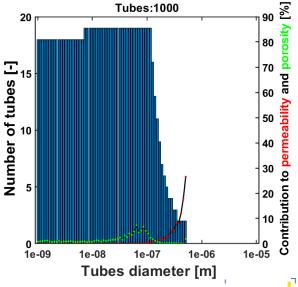
Effect of small-size pores

Experimental estimations of bedding plane separation, from Gonzalez-Blanco (2017)

FESEM	μ-CT
150 – 270 μm	410 – 560 μm







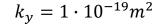
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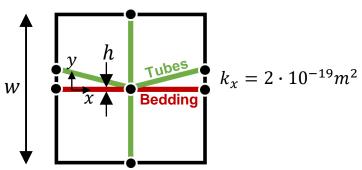
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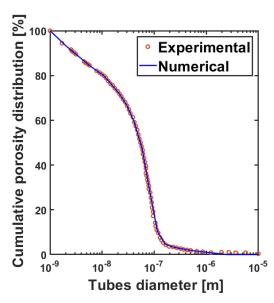


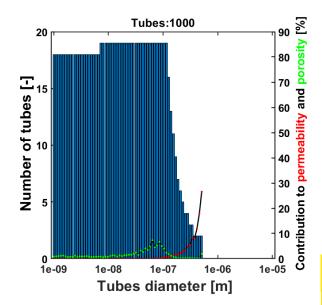
▶ 2. Macroporosity

Fitting of the pore size distribution

Effect of small-size pores

(Tortuosity = Calibration factor)





Fracture aperture

 $k_{x,frac,0} = 10^{-19} m^2$  $\rightarrow h_0 = \sqrt[3]{12 w k}$  Macropores

$$k_{x} = \frac{\pi}{8} \left(\frac{D}{2}\right)^{4} \left(\frac{1}{w^{2}}\right) + \underbrace{\frac{h_{x,f}rac}{h^{2}(\cdot w)} \left(\frac{h}{w}\right)}_{12} \qquad k_{x,f}$$

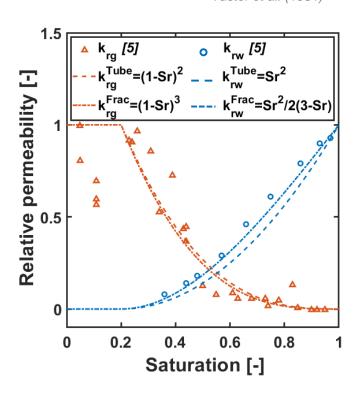
 $k_y = \frac{\pi}{8} \left(\frac{D}{2}\right)^4 \left(\frac{1}{w^2}\right)^4$ 

➤ 3. Intrinsic permeability Effect of large-size pores

#### **Characterisation of the microstructure parameters**

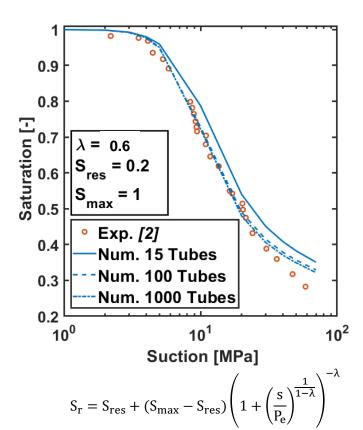
#### ► 4. Relative permeability curves

Yuster et al. (1951)



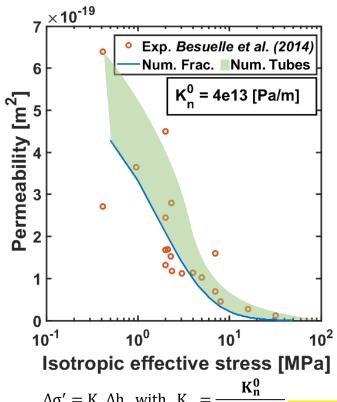
#### ► 5. Retention curve

Van Genuchten (1980)

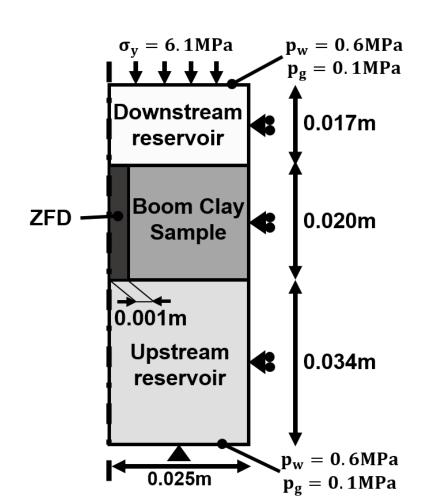


► 6. Normal stiffness of the fracture

Goodman (1976)



#### **Geometry and boundary conditions**



#### **Parameters**

#### Reservoirs

Stiff elements: E = 10000MPa v = 0.3Highly conductive: n = 0.5  $k = 10^{-10}m^2$ 

• Flat retention curve:  $P_{entry} = 0.01MPa$ 

#### **Boom Clay matrix**

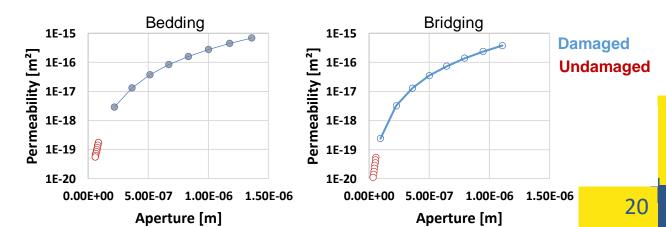
• Mechanical: E = 200 - 400MPa v = 0.33

Hydraulic:

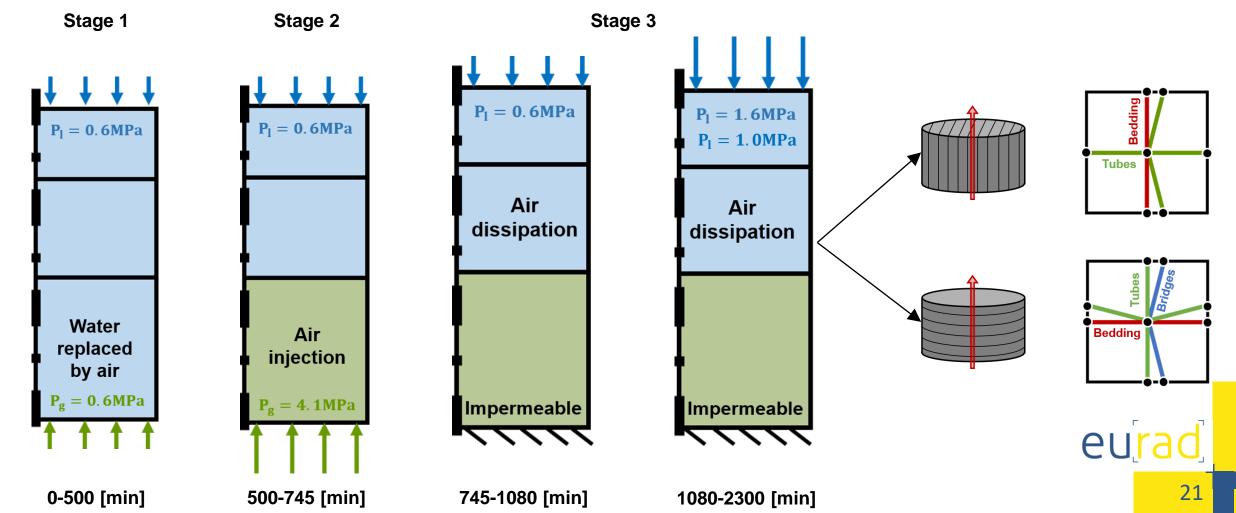
Initial aperture:  $0.80 - 1.27 \cdot 10^{-7} m$ Initial permeability:  $2.0 - 4.0 \cdot 10^{-19} m^2$ 

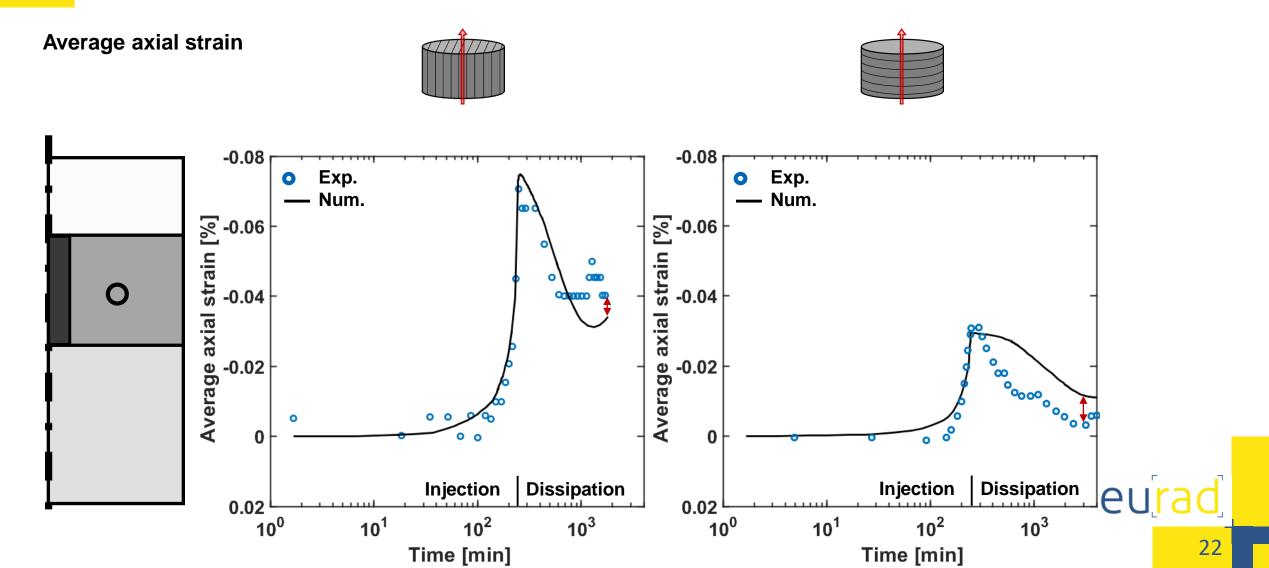
• Initial porosity: 0.363

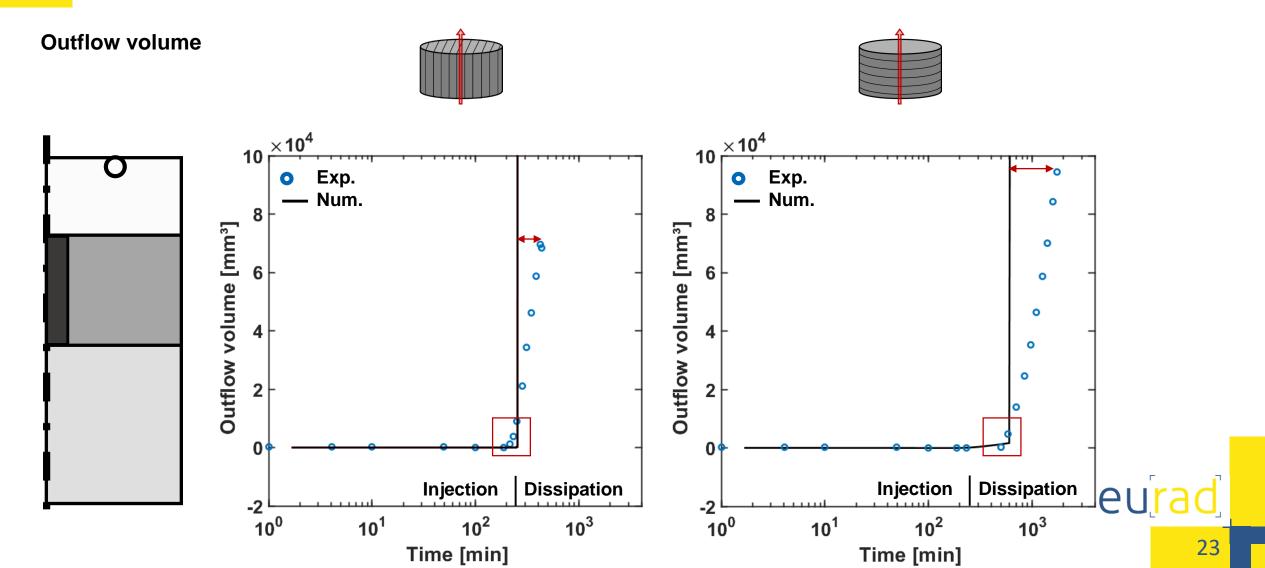
#### **Boom Clay Zone of Fracture Development (ZFD)**

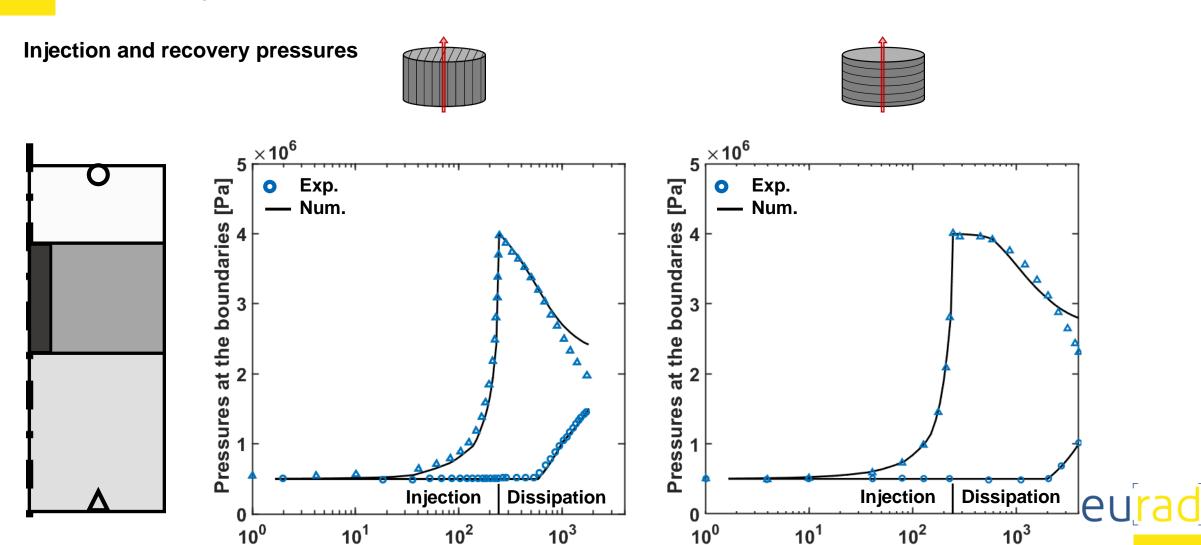


### **Simulation stages**



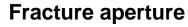






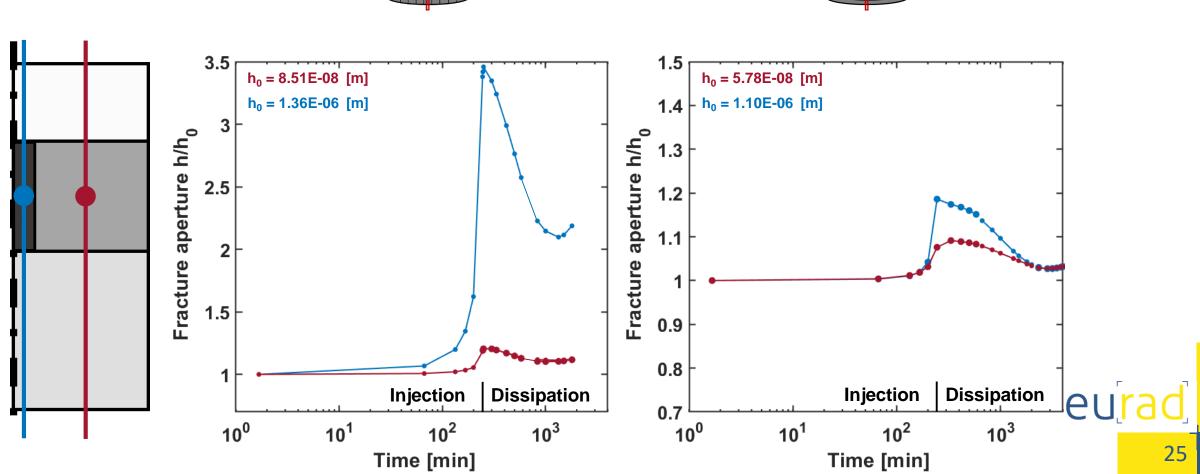
Time [min]

Time [min]

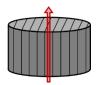


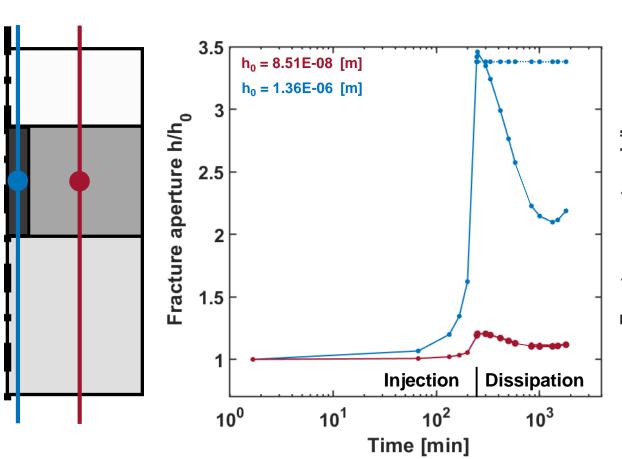


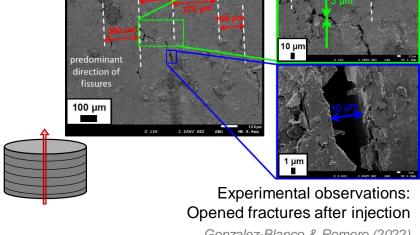




Fracture aperture

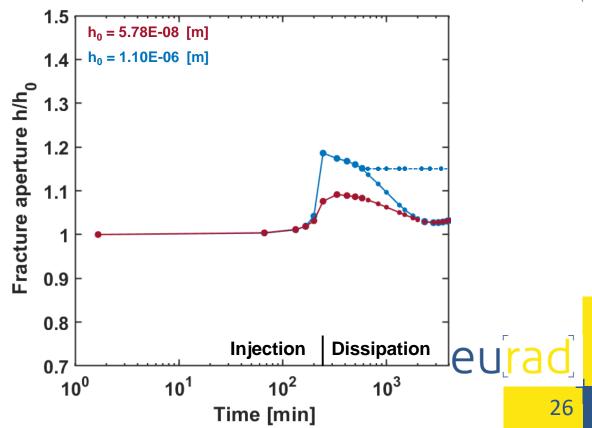


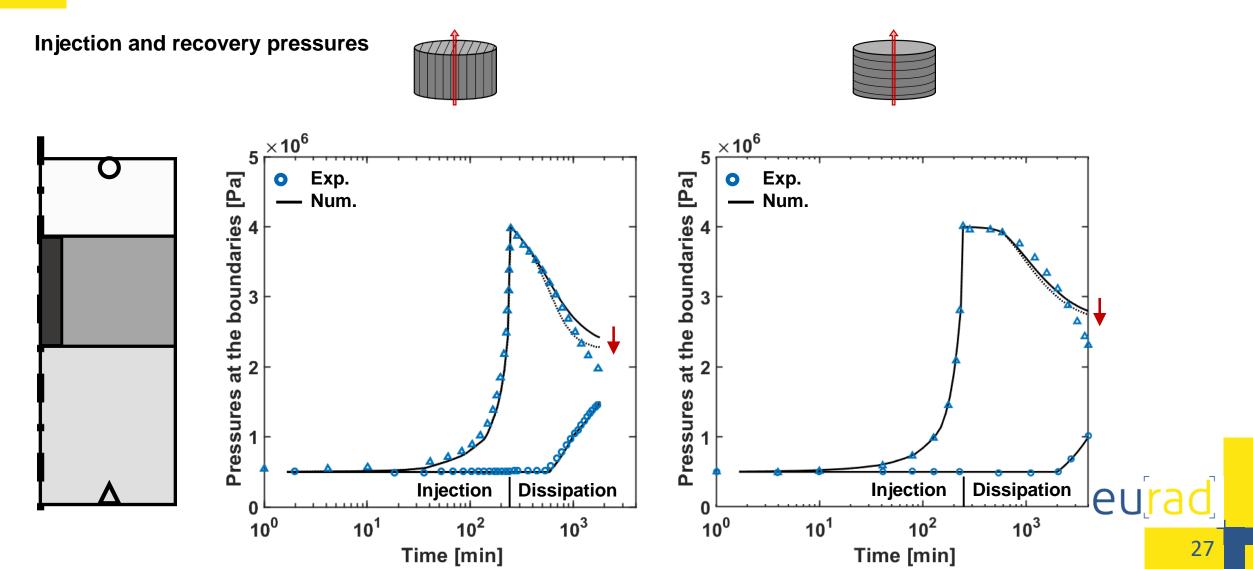




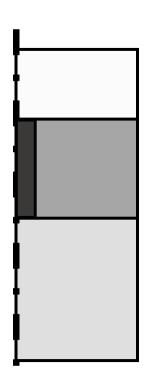
After air injection

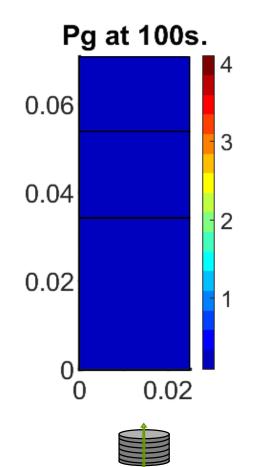
Gonzalez-Blanco & Romero (2022)

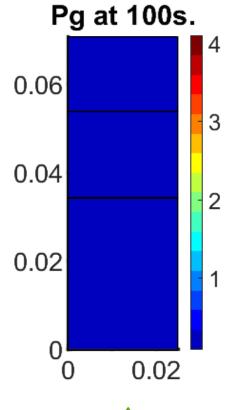




### Injection and recovery pressures



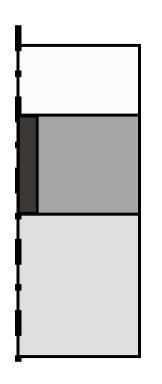


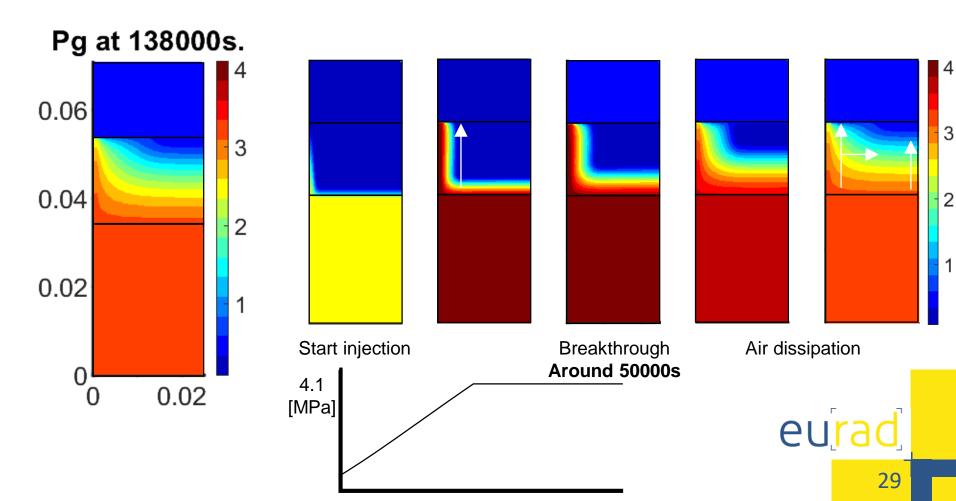




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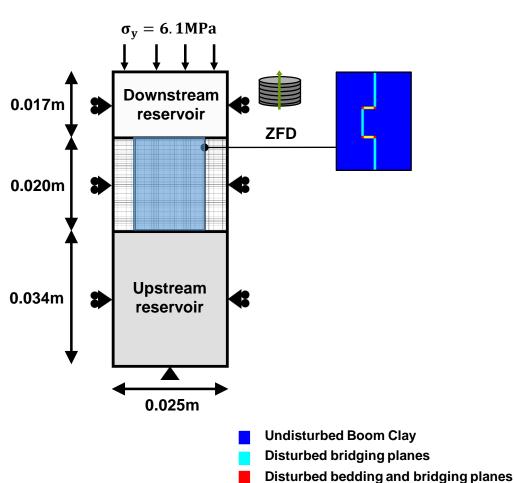
### Injection and recovery pressures

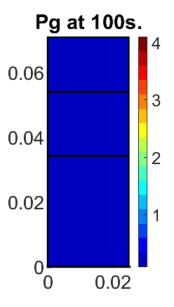




Disturbed bedding planes

#### Effect of the connectivity of the planes

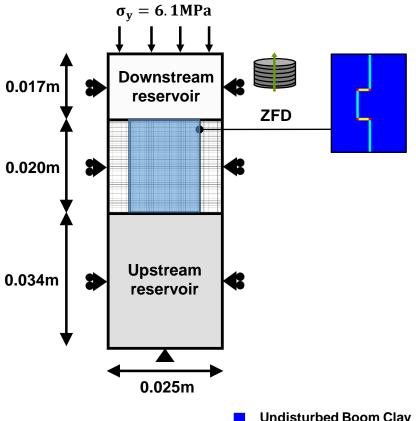


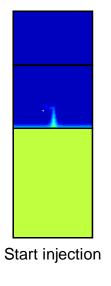


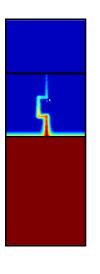
# 4.1

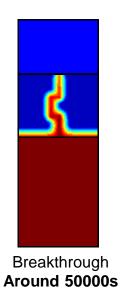
### Gas injection experiment

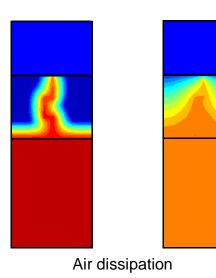
Effect of the connectivity of the planes











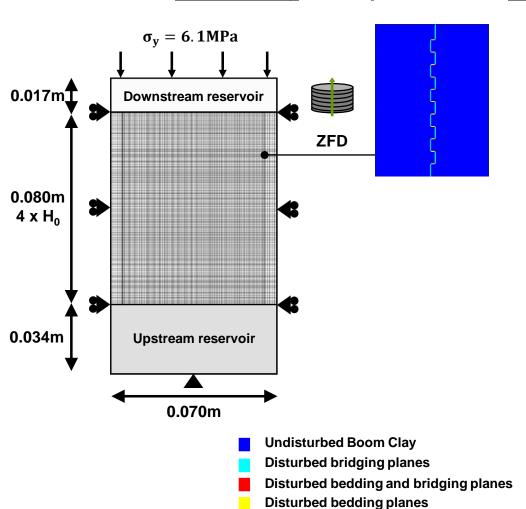
**Undisturbed Boom Clay** 

**Disturbed bridging planes** 

Disturbed bedding and bridging planes

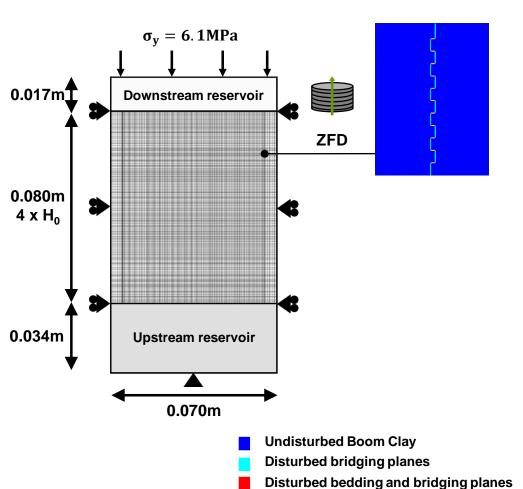
Disturbed bedding planes

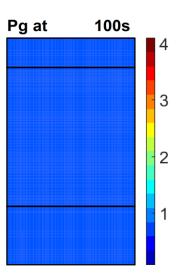
Effect of the connectivity of the planes under up-scaling

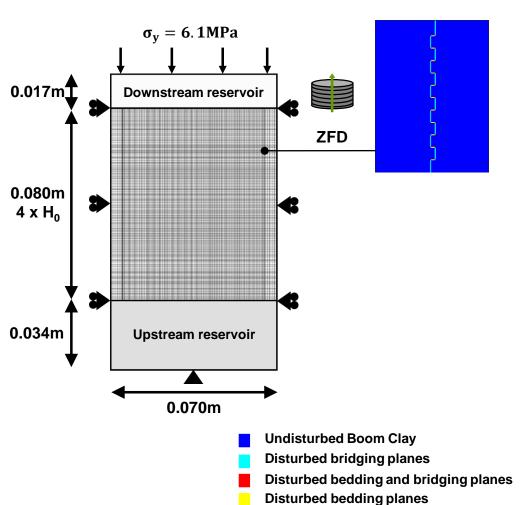


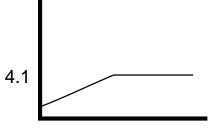
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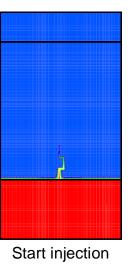
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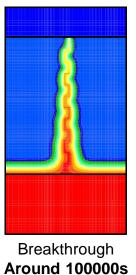


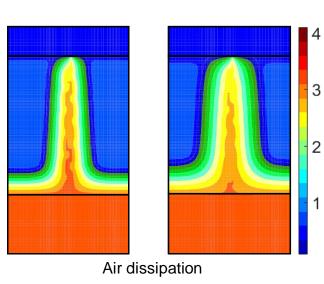




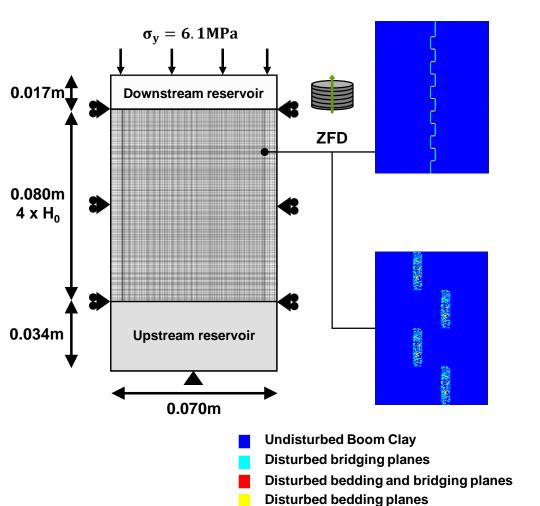


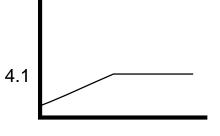


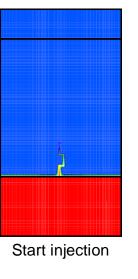


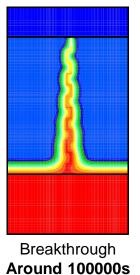


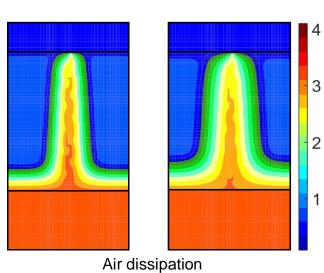
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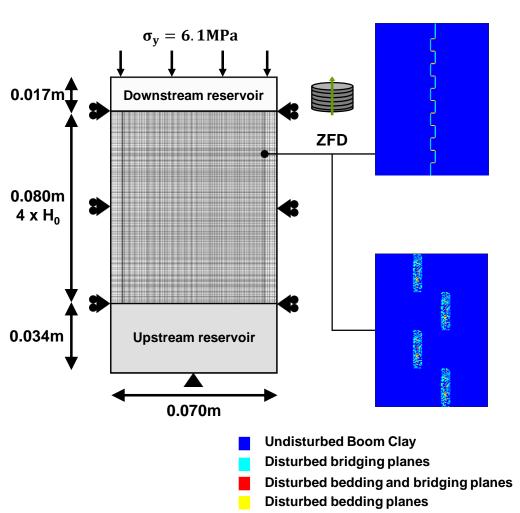


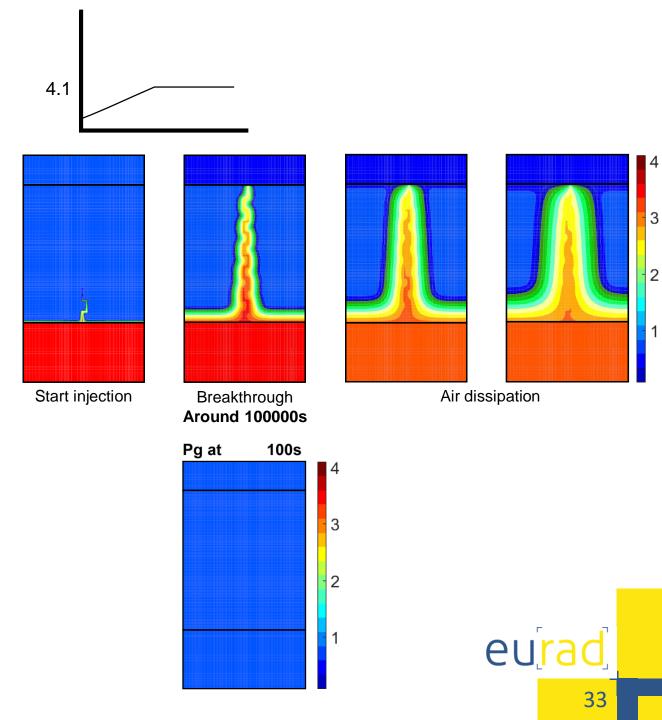


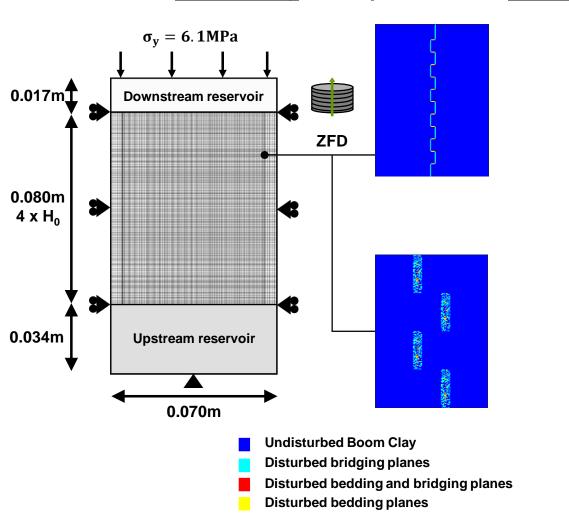


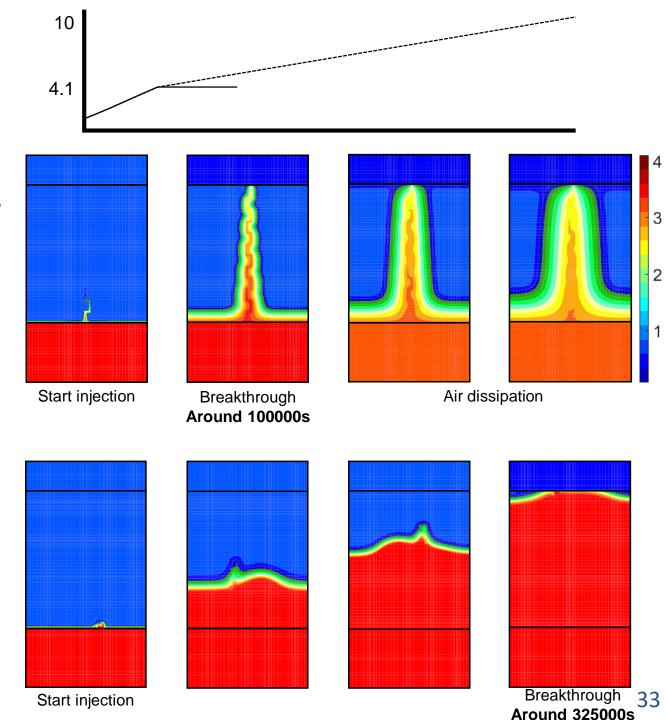


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### Content

- Context
- 2 From experimental evidence to modelling
- Multi-scale modelling approach
- 4 Modelling gas injection experiment
- **6** Conclusions

### **Conclusions**

### We **developed** a multi-scale model able to

- 1. Simply idealise the microstructure of the rock with fractures and tubes
- 2. Reproduce mechanisms inherent to gas migrations in sound rock layers

### We **showed** that

- 1. Macro-pores, bedding planes and bridging planes play different roles in gas flows
- 2. Preferential flow paths can be generated through fractures with weaker properties
- 3. Different gas mechanisms occur in the presence of weaker bridging planes

# Thank you for your attention!

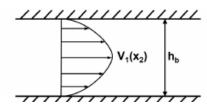
#### Micro-scale boundary value problem

Constitutive equations: Hydraulic problem considering a channel flow model (Navier-Stokes equations)



$$q_{\alpha_i} = -\frac{k_{r_{\alpha}}}{\mu_{\alpha}} \frac{1}{A} \kappa_{frac} \frac{\partial p_{\alpha}}{\partial x_i} = -\frac{k_{r_{\alpha}}}{\mu_{\alpha}} \frac{h_b^3}{12w} \frac{\partial p_{\alpha}}{\partial x_i}$$

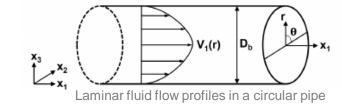
$$q_{\alpha_i} = -\frac{k_{r_{\alpha}}}{\mu_{\alpha}} \frac{1}{A} \kappa_{frac} \frac{\partial p_{\alpha}}{\partial x_i} = -\frac{k_{r_{\alpha}}}{\mu_{\alpha}} \frac{h_b^3}{12w} \frac{\partial p_{\alpha}}{\partial x_i} \qquad q_{\alpha_i} = -\frac{k_{r_{\alpha}}}{\mu_{\alpha}} \frac{1}{A} \kappa_{tube} \frac{\partial p}{\partial x_i} = -\frac{k_{r_{\alpha}}}{\mu_{\alpha}} \frac{D^4}{128w^2} \frac{\partial p}{\partial x_i}$$



Laminar fluid flow profiles between two parallel plates

$$\kappa_{frac} = -\frac{h_b^2}{12} h_b \cdot w$$

$$\kappa_{tube} = -\pi \frac{D^4}{128}$$



W

Gas flow in between of water flows in a fracture space

$$k_{r_w} = \frac{S_r^2}{2} (3 - S_r)$$
$$k_{r_g} = (1 - S_r)^3$$

$$k_{r_w} = S_r^2$$
  
$$k_{r_o} = (1 - S_r)^2$$



Gas flow in between of water flows in a circular pipe



$$i_{dg_i} = -S_{r_w} \, \bar{\mathsf{\tau}} \, D_{dg/w} \, \rho_w \, rac{\partial}{\partial x_i} \left( rac{\mathsf{\rho}_{dg}}{\mathsf{\rho}_w} 
ight)$$

#### Micro-scale boundary value problem

Constitutive equations: Hydro-mechanical couplings

Stress-dependent evolution of micro-elements aperture

$$\Delta \sigma' = K_n \, \Delta h$$

$$\Delta \sigma' = K \, \Delta D_b$$

$$K_n = \frac{K_n^0}{\left(1 + \frac{\Delta h}{h_0}\right)^2}$$

$$K = \frac{2C}{D_0}$$

Stress-dependent formulation of the transmissivity and the entry pressure of micro-elements

$$\kappa_{frac} = -\frac{h_b^2}{12} h_b \cdot w$$

$$\kappa_{tube} = -\pi \frac{D^4}{128}$$

$$p_e = p_{e_0} \left(\frac{h_{b_0}}{h_b}\right)^m$$

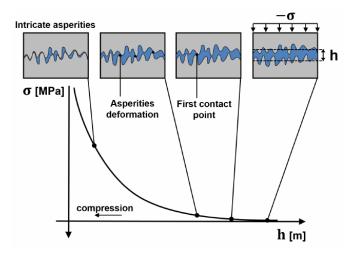
$$p_e = p_{e_0} \left(\frac{D_{b_0}}{D_b}\right)^m$$

$$p_{e_0} = \frac{2\sigma_{GL} cos\theta}{h_b}$$

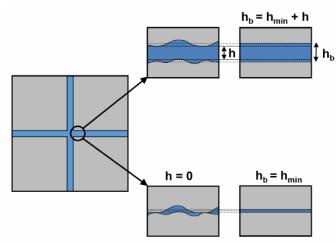
$$p_{e_0} = \frac{2\sigma_{GL} cos\theta}{D_b/2}$$

$$p_{e_0} = \frac{D_{b_0}}{D_b/2}$$

$$p_{e_0} = \frac{D_{b_0}}{D_b/2}$$



Constitutive law describing the normal behaviour of a rough rock joint, *Cerfontaine (2015)* 

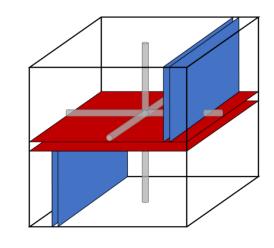


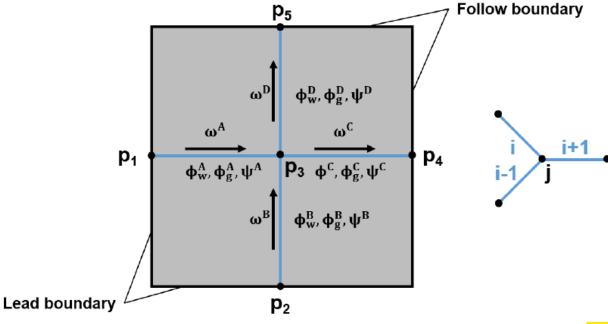
Definitions of the hydraulic and the mechanical aperture in reality (left) and in the modelling (right), *Marinelli (2016)* 

#### Micro-scale boundary value problem

General principles for numerical resolution of the hydraulic system

- Hydraulic network respecting these conditions:
  - Anti-symmetric boundary fluxes
  - Macroscopic pressure gradient between the boundaries
- Hydraulic problem established through mass balance on each node (j)
- Hydraulic problem solved
  - For a given configuration
  - Under steady-state conditions
  - By applying the macro-pressure to one node





Example of a channel network with the mass balance on node j



#### Micro-scale boundary value problem

General principles for numerical resolution of the hydraulic system

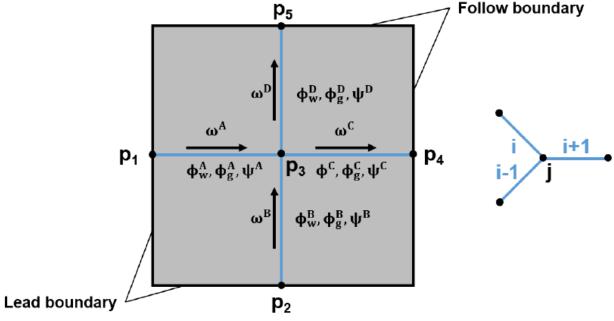
- Hydraulic network respecting these conditions:
  - Anti-symmetric boundary fluxes
  - Macroscopic pressure gradient between the boundaries
- => Channel (fracture or tube) mass fluxes of water and gas

$$\omega_{w} = -\underbrace{\frac{\rho_{w}k_{r_{w}}}{\mu_{w}}\kappa\frac{\partial p_{w}^{m}}{\partial s}}_{\text{Advection of liquid water}}$$

$$\omega_{g} = -\underbrace{\frac{\rho_{g}k_{r_{g}}}{\mu_{g}}\kappa\frac{\partial p_{g}^{m}}{\partial s}}_{\text{Advection of gaseous gas}} - \underbrace{H_{g}\frac{\rho_{g}k_{r_{w}}}{\mu_{w}}\kappa\frac{\partial p_{w}^{m}}{\partial s}}_{\text{Advection of dissolved gas}}$$

$$-S_{r_{w}}\bar{\tau}D_{dg/w}\frac{H_{g}}{\rho_{w}}\left(\frac{\rho_{w}\rho_{g,0}}{p_{g,0}}\frac{\partial p_{g}^{m}}{\partial s} - \frac{\rho_{g}\rho_{w,0}}{\chi_{w}}\frac{\partial p_{w}^{m}}{\partial s}\right)$$

Diffusion of dissolved gas



Example of a channel network with the mass balance on node j



#### Micro-scale boundary value problem

General principles for numerical resolution of the hydraulic system

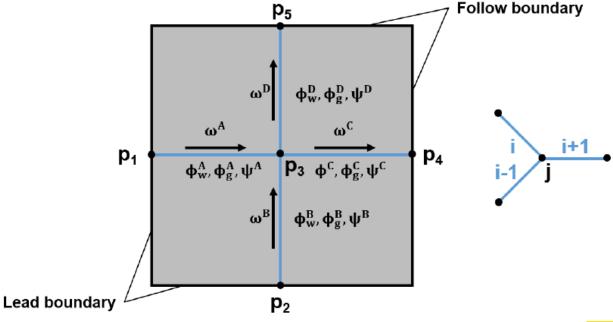
- Hydraulic problem established through mass balance on each node (j)
  - Mass conservation principle, i.e. for each node of the network, the sum of the input flows is equal to the sum of the output flows

$$\frac{d\omega_{\alpha}^{i}}{ds^{i}} = 0 \qquad \Leftrightarrow \qquad \omega_{\alpha}^{i-1} + \omega_{\alpha}^{i} + \omega_{\alpha}^{i+1} = 0$$
 
$$\alpha = w, g \qquad \text{Liquid or gaseous phase}$$

Well-posed hydraulic system to solve

$$\left[G_{ww}\right]\left\{p_w^m\right\} = 0 \qquad \left[G_{gg}\right]\left\{p_g^m\right\} + \left[G_{gw}\right]\left\{p_w^m\right\} = 0$$

- For a given configuration
- Under steady-state conditions
- By applying the macro-pressure to one node



Example of a channel network with the mass balance on node j



#### Micro-to-macro scale transition: Homogenisation

Fluid fluxes

$$f_{w_{i}}^{M} \frac{\partial p_{w}^{\star,M}}{\partial x_{i}} = \frac{1}{\Omega} \int_{\Omega} f_{w_{i}}^{m} \frac{\partial p_{w}^{\star,M}}{\partial x_{i}} d\Omega = \frac{1}{\Omega} \int_{\Gamma} \bar{q}_{w}^{m} p_{w}^{\star,M} d\Gamma$$
$$= \frac{1}{\Omega} \frac{\partial p_{w}^{\star,M}}{\partial x_{i}} \int_{\Gamma} \bar{q}_{w}^{m} x_{i} d\Gamma$$
$$= \frac{1}{\Omega} \int_{\Gamma} \bar{q}_{w}^{m} x_{i} d\Gamma$$

$$f_{g_i}^M + f_{dg_i}^M = \frac{1}{\Omega} \int_{\Gamma} \bar{q}_g^m x_i d\Gamma$$

Fluid masses: total amount of fluids inside the fractures and tubes

$$M_w^M = \frac{1}{\Omega} \int_{\Omega_w^{int}} \rho_w d\Omega$$
$$= \rho_w S_{r_w} \phi_n$$

$$M_g^M = M_g^m + M_{dg}^m$$

$$= \frac{1}{\Omega} \left( \int_{\Omega_g^{int}} \rho_g d\Omega + \int_{\Omega_w^{int}} \rho_{dg} d\Omega \right)$$

$$= \rho_g (1 - S_{r_w}) \phi_n + \rho_{dg} S_{r_w} \phi_n$$

#### Macro-scale boundary value problem

Under matrix form:

$$\begin{bmatrix} \begin{bmatrix} K_{ww}^{M} \end{bmatrix}_{(3 \times 3)} & \begin{bmatrix} K_{wg}^{M} \end{bmatrix}_{(3 \times 3)} \\ \begin{bmatrix} K_{gw}^{M} \end{bmatrix}_{(3 \times 3)} & \begin{bmatrix} K_{wg}^{M} \end{bmatrix}_{(3 \times 3)} \end{bmatrix} \begin{cases} \begin{cases} \delta \nabla p_{w}^{M} \\ \delta p_{w}^{M} \end{cases} \\ \begin{cases} \delta \nabla p_{g}^{M} \\ \delta p_{g}^{M} \end{cases} \end{cases}_{(3)} \end{cases} = \begin{cases} \begin{cases} \delta f_{w}^{M} \\ \delta \dot{M}_{w}^{M} \end{cases}_{(3)} \\ \begin{cases} \delta f_{g}^{M} \\ \delta \dot{M}_{g}^{M} \end{cases}_{(3)} \end{cases}$$

Summarized as:

$$\left[A^{M}\right]_{(10\times10)}\left\{\delta U^{M}\right\}_{(10)}=\left\{\delta \varSigma^{M}\right\}_{(10)}$$