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# Wavelet methods to study the pointwise regularity of the generalized Rosenblatt process ${ }^{* \dagger}$ 

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#### Abstract

We identify three types of pointwise behaviour in the regularity of the (generalized) Rosenblatt process. This extends to a non Gaussian setting previous results known for the (fractional) Brownian motion. On this purpose, fine bounds on the increments of the Rosenblatt process are needed. Our analysis is essentially based on various wavelet methods.


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## 1 Introduction

Precise study of path behaviour, and in particular regularity, of stochastic processes is a classical research field, initiated in the 1920 s by the works of Wiener [47]. It lies in between probability and (harmonic) analysis and a common strategy is to mix probabilistic arguments with analytical tools. Pioneer works concerned Brownian motion. Among them, one can cite Paley and Wiener's expansion [48] using Fourier series, Lévy's representation [30] obtained with some techniques of interpolation theory or, more recently, Kahane's expansion [25] in the Schauder basis.

In the last decades, the emergence of wavelet analysis allowed to obtain series expansions for many stochastic processes. Let $\psi: \mathbb{R} \rightarrow \mathbb{R}$ be a smooth function satisfying the admissibility condition [35]

$$
\begin{equation*}
\int_{\mathbb{R}} \frac{|\hat{\psi}(\xi)|}{|\xi|} d \xi<\infty \tag{1.1}
\end{equation*}
$$

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where $\hat{\psi}$ is the Fourier transform of $\psi$. As such it generates an orthonormal basis of $L^{2}(\mathbb{R})$. More precisely, any function $f \in L^{2}(\mathbb{R})$ can be decomposed as

$$
\begin{equation*}
f=\sum_{j \in \mathbb{Z}} \sum_{k \in \mathbb{Z}} c_{j, k} \psi\left(2^{j} \cdot-k\right), \tag{1.2}
\end{equation*}
$$

where

$$
\begin{equation*}
c_{j, k}=2^{j} \int_{\mathbb{R}} f(x) \psi\left(2^{j} x-k\right) d x . \tag{1.3}
\end{equation*}
$$

It is noteworthy that the expansion (1.2) holds true in many function spaces. We refer to the seminal books [15, 35, 34] for more details and proofs of these facts. Multifractal analysis has demonstrated the efficiency of wavelet methods to study uniform and pointwise Hölder regularity of functions both from a theoretical [9, 10, 13, 20, 21, 23] and a practical point of view $[2,3,11,16,19,22,37,45,46]$.

Now, let us consider a probability space $(\Omega, \mathcal{A}, \mathbb{P})$ and a real-valued stochastic process $X$ definitionined on it. If, for all $\omega \in \Omega$ the sample path $t \mapsto X(t, \omega)$ belongs to $L^{2}(\mathbb{R})$, one can apply expansion (1.2) to it. This way, one definitionines a sequence of random wavelet coefficients $\left(c_{j, k}(\omega)\right)_{j, k \in \mathbb{Z}}$. For instance, if $X=B_{H}$ is the fractional Brownian motion of Hurst index $H \in(0,1)$ and if $\psi$ is a sufficiently regular wavelet, one has [36, 22]

$$
\begin{equation*}
B_{H}=\sum_{j \in \mathbb{N}} \sum_{k \in \mathbb{Z}} 2^{-H j} \xi_{j, k} \psi_{H+1 / 2}\left(2^{j} \cdot-k\right)+R \tag{1.4}
\end{equation*}
$$

where $(R(t, \cdot))_{t \in \mathbb{R}^{+}}$is a process with alsmost surely $C^{\infty}$ sample paths, $\left(\xi_{j, k}\right)_{j \in \mathbb{N}, k \in \mathbb{Z}}$ is a sequence of independent $\mathcal{N}(0,1)$ random variables and $\psi_{H+1 / 2}$ is a fractional antiderivative of $\psi$, see Section 2 for a precise definitioninition.

In [18], Esser and Loosveldt undertook a systematic study of Gaussian wavelet series. Thanks to (1.4), it applies in particular to the fractional Brownian motion and leads to the following theorem.
Theorem 1.1. For all $H \in(0,1)$, there exists an event $\Omega_{H}$ of probability 1 satisfying the following assertions for all $\omega \in \Omega_{H}$ and every non-empty interval I of $\mathbb{R}$.

- For almost every $t \in I$,

$$
0<\limsup _{s \rightarrow t} \frac{\left|B_{H}(t, \omega)-B_{H}(s, \omega)\right|}{|t-s|^{H} \sqrt{\log \log |t-s|^{-1}}}<+\infty .
$$

Such points are called ordinary points.

- There exists a dense set of points $t \in I$ such that

$$
0<\limsup _{s \rightarrow t} \frac{\left|B_{H}(t, \omega)-B_{H}(s, \omega)\right|}{|t-s|^{H} \sqrt{\log |t-s|^{-1}}}<+\infty .
$$

Such points are called rapid points.

- There exists a dense set of points $t \in I$ such that

$$
0<\limsup _{s \rightarrow t} \frac{\left|B_{H}(t, \omega)-B_{H}(s, \omega)\right|}{|t-s|^{H}}<+\infty .
$$

Such points are called slow points.
Note that Theorem 1.1 extends some well-known results of Kahane concerning the Brownian motion [25]. The "ordinary", "rapid" and "slow" terminology is inspired by them. Let us justify it. From a measure-theoretical point of view, the modulus of continuity $x \mapsto|x|^{H} \sqrt{\log \log |x|^{-1}}$ is the most frequent among the points of sample paths.

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Thus, it is natural to refer it to ordinary. Now, $|x|^{H} \sqrt{\log \log |x|^{-1}}=o\left(|x|^{H} \sqrt{\log |x|^{-1}}\right)$ if $x \rightarrow 0^{+}$and thus points for which $x \mapsto|x|^{H} \sqrt{\log |x|^{-1}}$ is the pointwise modulus of continuity are refereed to rapid. On the other side, points for which $x \mapsto|x|^{H}$ is the pointwise modulus of continuity are referred to slow because $|x|^{H}=o\left(|x|^{H} \sqrt{\log \log |x|^{-1}}\right)$ if $x \rightarrow 0^{+}$.

Now, let us turn to the stochastic process we will deal with in this paper. The Rosenblatt process appears naturally as a limit of normalized sums of long-range dependent random variables [17]. Like the fractional Brownian motion, it belongs to the class of Hermite processes, fractional Brownian motion being of order 1 while Rosenblatt process is of order 2. Both are selfsimilar stochastic processes with stationary increments and are characterized by a parameter $H$, called the Hurst exponent. However, unlike the fractional Brownian motion, the Rosenblatt process is not Gaussian. Does it make a big difference regarding ordinary, rapid and slow points? In other words, can Theorem 1.1 be extended to cover the non Gaussian Rosenblatt process?

For the last fifteen years the Rosenblatt process has received a significantly increasing interest in both theoretical and practical lines of research. Due to its self-similarity, its applications are numerous across a multitude of fields, including internet traffic [12] and turbulence [40, 28]. From a statistical point of view, estimating the value of the Hurst index $H$ is important for practical applications and various estimators exist, see [8, 43]. Also, from a mathematical point of view the Rosenblatt process has received a lot of interest since its inception in [39]. Its distribution, still not known in explicit form, was studied first in [1] and more recently in [33] and [44].

In this paper, we even consider a generalization of the Rosenblatt process, as definitionined and studied in [32]. It depends on two parameters $H_{1}, H_{2} \in\left(\frac{1}{2}, 1\right)$ which are such that $H_{1}+H_{2}>\frac{3}{2}$. The generalized Rosenblatt process $\left\{R_{H_{1}, H_{2}}(t, \cdot)\right\}_{t \in \mathbb{R}_{+}}$is definitionined as a double Wiener-Itô integral of a kernel function $K_{H_{1}, H_{2}}$ with respect to a given Brownian motion. More precisely, consider a standard two-sided Brownian motion $B$, and set

$$
\begin{equation*}
R_{H_{1}, H_{2}}(t, \cdot)=\int_{\mathbb{R}^{2}}^{\prime} K_{H_{1}, H_{2}}\left(t, x_{1}, x_{2}\right) d B\left(x_{1}\right) d B\left(x_{2}\right), \tag{1.5}
\end{equation*}
$$

where $\int_{\mathbb{R}^{2}}^{\prime}$ denotes integration over $\mathbb{R}^{2}$ excluding the diagonal. The kernel function in (1.5) is expressed, for all $\left(t, x_{1}, x_{2}\right)$ on $\mathbb{R}_{+} \times \mathbb{R}^{2}$, by

$$
K_{H_{1}, H_{2}}\left(t, x_{1}, x_{2}\right)=\frac{1}{\Gamma\left(H_{1}-\frac{1}{2}\right) \Gamma\left(H_{2}-\frac{1}{2}\right)} \int_{0}^{t}\left(s-x_{1}\right)_{+}^{H_{1}-\frac{3}{2}}\left(s-x_{2}\right)_{+}^{H_{2}-\frac{3}{2}} d s,
$$

where $\Gamma$ stands for the usual Gamma Euler function, and where for $(x, \alpha) \in \mathbb{R}^{2}, x_{+}^{\alpha}:=$ $x^{\alpha} \mathbb{1}_{[0,+\infty)}$ Note that the (standard) Rosenblatt process is the process $\left\{R_{H, H}(t, \cdot)\right\}_{t \in \mathbb{R}_{+}}$ for $H \in(3 / 4,1)$. The generalized Rosenblatt process $\left\{R_{H_{1}, H_{2}}(t, \cdot)\right\}_{t \in \mathbb{R}_{+}}$is non-Gaussian, belongs to the second Wiener chaos, and has the following basic properties:
(1) Continuity: the trajectories of the Rosenblatt process $R_{H_{1}, H_{2}}$ are continuous.
(2) Stationary increments: $R_{H_{1}, H_{2}}$ has stationary increments; that is, the distribution of the process $\left\{R_{H_{1}, H_{2}}(t+s, \cdot)-R_{H_{1}, H_{2}}(s, \cdot)\right\}_{t \in \mathbb{R}_{+}}$does not depend on $s \geqslant 0$.
(3) Self-similarity: $R_{H_{1}, H_{2}}$ is self-similar with exponent $H_{1}+H_{2}-1$; that is, the processes $\left\{R_{H_{1}, H_{2}}(c t, \cdot)\right\}_{t \in \mathbb{R}_{+}}$and $\left\{c^{H_{1}+H_{2}-1} R_{H_{1}, H_{2}}(t, \cdot)\right\}_{t \in \mathbb{R}_{+}}$have the same distribution for all $c>0$.

In [6], Ayache and Esmili presented a wavelet-type representation of the generalized Rosenblatt process, very similar to the one given in [36] for fractional Brownian motion, except for the use of integrals of two-dimensional wavelet bases. This representation

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is the starting point of this paper. It is one of our key tools to prove the following Theorem 1.2 which is the main result of this paper.
Theorem 1.2. For all $H_{1}, H_{2} \in\left(\frac{1}{2}, 1\right)$ such that $H_{1}+H_{2}>\frac{3}{2}$, there exists an event $\Omega_{H_{1}, H_{2}}$ of probability 1 satisfying the following assertions for all $\omega \in \Omega_{H_{1}, H_{2}}$ and every non-empty interval $I$ of $\mathbb{R}$.

- For almost every $t \in I$,

$$
\begin{equation*}
0<\limsup _{s \rightarrow t} \frac{\left|R_{H_{1}, H_{2}}(t, \omega)-R_{H_{1}, H_{2}}(s, \omega)\right|}{|t-s|^{H_{1}+H_{2}-1} \log \log |t-s|^{-1}}<+\infty . \tag{1.6}
\end{equation*}
$$

Such points are called ordinary points.

- There exists a dense set of points $t \in I$ such that

$$
\begin{equation*}
0<\limsup _{s \rightarrow t} \frac{\left|R_{H_{1}, H_{2}}(t, \omega)-R_{H_{1}, H_{2}}(s, \omega)\right|}{|t-s|^{H_{1}+H_{2}-1} \log |t-s|^{-1}}<+\infty . \tag{1.7}
\end{equation*}
$$

Such points are called rapid points.

- There exists a dense set of points $t \in I$ such that

$$
\begin{equation*}
\limsup _{s \rightarrow t} \frac{\left|R_{H_{1}, H_{2}}(t, \omega)-R_{H_{1}, H_{2}}(s, \omega)\right|}{|t-s|^{H_{1}+H_{2}-1}}<+\infty . \tag{1.8}
\end{equation*}
$$

Such points are called slow points.
Theorem 1.2 shows in particular that slow, ordinary and rapid points are not specific to Gaussian processes.
Remark 1.3. Let us compare Theorems 1.1 and 1.2. If $X$ denotes both the fractional Brownian motion or the generalized Rosenblatt process, we see that the asymptotic behaviour of $|X(t, \omega)-X(s, \omega)|$ is always comparable to a modulus of continuity of the form $|t-s|^{\alpha} \theta(|t-s|)$, with $\alpha$ corresponding to the self-similarity exponent of $X$ and $\theta$ a potential logarithmic correction. For the ordinary points, $\theta$ is an iterated logarithm. More precisely, for the fractional Brownian motion, we have $\theta(|t-s|)=\sqrt{\log \log |t-s|^{-1}}$ while, for the generalized Rosenblatt process, $\theta(|t-s|)=\log \log |t-s|^{-1}$. The same feature appears for the rapid points: in the case of the fractional Brownian motion we have $\theta(|t-s|)=\sqrt{\log |t-s|^{-1}}$ and for the generalized Rosenblatt process we have $\theta(|t-s|)=\log |t-s|^{-1}$. Therefore, the only difference between the corresponding logarithmic corrections is the square root that is used for the fractional Brownian motion and not for the generalized Rosenblatt process. It comes from the estimates that can be done on the tails of the distribution of random variables in various Wiener chaoses, see Theorems 3.14 and 3.15 below. Concerning the slow points, there is no logarithmic correction, $\theta=1$ in both case. Unfortunately, contrary to the fractional Brownian motion, we did not manage to show the positiveness of the limit in (1.8). In fact, we would need to find an almost-sure uniform lower modulus of continuity for the generalized Rosenblatt process and to be able to judge its optimality, which seems to be a difficult task. This is discussed in details in Remark 5.1 below.

Our strategy to prove Theorem 1.2 is as follows. First, in Section 3 we derive upper-bounds for the oscillations $\left|R_{H_{1}, H_{2}}(t, \omega)-R_{H_{1}, H_{2}}(s, \omega)\right|$ that are sharp enough to imply the finiteness of the limits (1.6), (1.7) and (1.8). This is done by means of the wavelet-type expansion given in [6], see Theorem 3.3 below. Then, in Section 4, we give lower bounds for the so-called wavelet-leaders, see Section 2, of the generalized Rosenblatt process on a given compactly supported wavelet basis. This will prove the positiveness of the limits (1.6), (1.7). In particular, we use different bases depending

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on whether we deal with the finiteness of the limits in Theorem 1.2 or with their strict positiveness. This is very different from [18] where the authors always work with the same wavelet. The reason is that the expression (3.3) in Theorem 3.3 below is not a wavelet series: it involves additional quantities. Therefore, standard arguments linking wavelet coefficients and regularity of the associated functions can no longer be used.

There is a priori no obstacles to extend our results in Section 4 to any Hermite process. On the contrary, extending the results of Section 3 does not seem obvious at all. This is because a wavelet-type expansion of arbitrary Hermite process is still missing but also because our strategy relies on arguments which are specific to the two-dimensional feature of the Rosenblatt process, see Lemma 2.1 for instance.

Notations used through this paper are rather standard except, maybe, that if $s, t$ are two real numbers, $\int_{[s, t]}$ stands for $\int_{s}^{t}$ if $s \leqslant t$ and $-\int_{s}^{t}=\int_{t}^{s}$ otherwise.

## 2 Some important facts involving wavelets

In this section, we gather all the facts concerning wavelets that we will strongly use all along this article. First, an immediate but important consequence of the admissibility condition (1.1) is that, if the wavelet $\psi \in L^{1}(\mathbb{R})$, its first moment always vanishes, i.e.

$$
\begin{equation*}
\int_{\mathbb{R}} \psi(x) d x=0 . \tag{2.1}
\end{equation*}
$$

This condition is met for all the wavelets we consider in this paper.
First, while dealing with the upper bounds for the limits in Theorem 1.2, we will use a wavelet-type expansion of the generalized Rosenblatt process. It is given in [6] by means of Meyer wavelet: $\psi$ belongs to the Schwartz class $\mathcal{S}(\mathbb{R})$, and its Fourier transform is compactly supported, see [29]. In particular, for all $H \in(1 / 2,1), \psi_{H}$, the fractional antiderivative $\widehat{\psi_{H}}$ of order $H-1 / 2$ of $\psi$ is well-definitionined by means of its Fourier transform as

$$
\begin{equation*}
\widehat{\psi_{H}}(0)=0 \quad \text { and } \quad \widehat{\psi_{H}}(\xi)=(i \xi)^{-\left(H-\frac{1}{2}\right)} \widehat{\psi}(\xi), \forall \xi \neq 0 . \tag{2.2}
\end{equation*}
$$

It also belongs to the Schwartz class $\mathcal{S}(\mathbb{R})$, see [4, 6, 41] for instance. Moreover, some standard facts from distribution theory [41, 4] give us the explicit formula

$$
\psi_{H}(t)=\frac{1}{\Gamma\left(H-\frac{1}{2}\right)} \int_{\mathbb{R}}(t-x)_{+}^{H-\frac{3}{2}} \psi(x) d x .
$$

From (2.2), we see that $\operatorname{supp}\left(\widehat{\psi_{H}}\right)=\operatorname{supp}(\widehat{\psi})$ which is the key fact to establish the following lemma, gathering facts already proved in [6].
Lemma 2.1. Let $H_{1}, H_{2} \in\left(\frac{1}{2}, 1\right)$. If $\left(j_{1}, j_{2}, k_{1}, k_{2}\right) \in \mathbb{Z}^{4}$ are such that $\left|j_{1}-j_{2}\right|>1$, then the integral $I_{j_{1}, j_{2}}^{k_{1}, k_{2}}:=\int_{\mathbb{R}} \psi_{H_{1}}\left(2^{j_{1}} x-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} x-k_{2}\right) d x$ vanishes. Moreover, for all $L>0$, there exists a constant $C_{L}>0$ such that for all $\left(j, k_{1}, k_{2}\right) \in \mathbb{Z}^{3}$, we have

$$
\begin{align*}
&\left|I_{j+1, j}^{k_{1}, k_{2}}\right|=\left|2^{-j} \int_{\mathbb{R}} e^{-i\left(k_{1}-2 k_{2}\right) \xi} \widehat{\psi_{H_{1}}}(\xi) \widehat{\widehat{\psi_{H_{2}}}(2 \xi)} d \xi\right| \leqslant C_{L} \frac{2^{-j}}{\left(3+\left|k_{1}-2 k_{2}\right|\right)^{L}},  \tag{2.3}\\
&\left|I_{j, j}^{k_{1}, k_{2}}\right|=\left\lvert\, 2^{-j} \int_{\mathbb{R}} e^{-i\left(k_{1}-k_{2}\right) \xi \widehat{\psi_{H_{1}}}(\xi) \widehat{\widehat{\psi_{H_{2}}}(\xi)} d \xi \left\lvert\, \leqslant C_{L} \frac{2^{-j}}{\left(3+\left|k_{1}-k_{2}\right|\right)^{L}}\right.}\right.  \tag{2.4}\\
&\left|I_{j, j+1}^{k_{1}, k_{2}}\right|=\left\lvert\, 2^{-j} \int_{\mathbb{R}} e^{-i\left(2 k_{1}-k_{2}\right) \xi \widehat{\psi_{H_{1}}}(2 \xi) \widehat{\widehat{\psi_{H_{2}}}(\xi)} d \xi \left\lvert\, \leqslant C_{L} \frac{2^{-j}}{\left(3+\left|2 k_{1}-k_{2}\right|\right)^{L}} .\right.} .\right. \tag{2.5}
\end{align*}
$$

When dealing with the lower bounds for the limits in Theorem 1.2, we use Daubechies compactly supported wavelets [14]. Note that, if $\operatorname{supp}(\Psi) \subseteq[-N, N]$, for a positive integer $N$, then, for all $(j, k) \in \mathbb{N} \times \mathbb{Z}$, by an obvious change of variable, the wavelet
coefficient (1.3) can be written $c_{j, k}=\int_{-N}^{N} f\left(\frac{x+k}{2 j}\right) \Psi(x) d x$. But now, using the first vanishing moment property (2.1) of $\Psi$, for all $t \in \mathbb{R}$, one can write

$$
\begin{equation*}
c_{j, k}=\int_{-N}^{N}\left(f\left(\frac{x+k}{2^{j}}\right)-f(t)\right) \Psi(x) d x \tag{2.6}
\end{equation*}
$$

Since the value of this integral does not depend on $t$, equality (2.6) somehow means that one can freely choose $t$, according to the context in which we use the wavelet coefficients.

Since $\Psi$ is compactly supported, $\Psi\left(2^{j} \cdot-k\right)$ is localized around the dyadic interval $\lambda_{j, k}:=\left[\frac{k}{2^{j}}, \frac{k+1}{2^{j}}\right)$ and it is therefore common to index wavelets with these intervals. For simplicity, we sometimes omit any references to the indices $j$ and $k$ for such intervals by writing $\lambda=\lambda_{j, k}$, and $k=s(\lambda)$. Then $\Psi_{\lambda}$ stands for the translated and dilated wavelet $\Psi\left(2^{j} .-k\right)$, where $k=s(\lambda)$ and $2^{-j}$ is the side length of the dyadic interval $\lambda$. Similarly, $c_{\lambda}$ is the quantity $c_{j, k}$. The notation $\Lambda_{j}$ stands for the set of dyadic intervals $\lambda$ of $\mathbb{R}$ with side length $2^{-j}$. The unique dyadic interval from $\Lambda_{j}$ containing the point $t_{0} \in \mathbb{R}$ is denoted $\lambda_{j}\left(t_{0}\right)$. The set of dyadic intervals is $\Lambda:=\cup_{j \in \mathbb{N}} \Lambda_{j}$. Two dyadic intervals $\lambda$ and $\lambda^{\prime}$ are adjacent if there exist $j \in \mathbb{N}$ such that $\lambda, \lambda^{\prime} \in \Lambda_{j}$ and $\operatorname{dist}\left(\lambda, \lambda^{\prime}\right)=0$. The set of dyadic intervals adjacent to $\lambda$ is denoted by $3 \lambda$. In this setting, one definitionines the wavelet leader [21] of a function at $t_{0}$ and of scale $j$ by

$$
\begin{equation*}
d_{j}\left(t_{0}\right)=\max _{\lambda \in 3 \lambda_{j}\left(t_{0}\right)} \sup _{\lambda^{\prime} \subseteq \lambda}\left|c_{\lambda}^{\prime}\right| . \tag{2.7}
\end{equation*}
$$

Note that the dependence on $t_{0}$ only comes from the dyadic intervals involved in (2.7). Then, if $\operatorname{supp}(\Psi) \subseteq[-N, N]$, from (2.6), with $t=t_{0}$, one can write

$$
\begin{equation*}
d_{j}\left(t_{0}\right) \leqslant 2 N \sup _{s \in\left(t_{0}-2^{-j}(N+2), t_{0}+2^{-j}(N+2)\right)} \mid f(s)-f\left(t_{0}\right)\| \| \psi \|_{L^{\infty}} . \tag{2.8}
\end{equation*}
$$

When we study stochastic processes, the wavelet leaders are random variables $d_{j}(t, \omega)$. Inequality (2.8) with some easy computations implies that in order to obtain the positiveness of the limit (1.6), it suffices to show that for all $\omega \in \Omega_{H_{1}, H_{2}}$ and all open intervals $I \subseteq \mathbb{R}^{+}$, for almost every $t \in I$,

$$
\begin{equation*}
0<\limsup _{j \rightarrow+\infty} \frac{d_{j}(t, \omega)}{2^{-j\left(H_{1}+H_{2}-1\right)} \log (j)} . \tag{2.9}
\end{equation*}
$$

Similarly, to prove the positiveness of the limit (1.7), we just have to show that for all $\omega \in \Omega_{H_{1}, H_{2}}$ and all open intervals $I \subseteq \mathbb{R}^{+}$, there exists a dense set of points $t \in I$ such that

$$
\begin{equation*}
0<\limsup _{j \rightarrow+\infty} \frac{d_{j}(t, \omega)}{2^{-j\left(H_{1}+H_{2}-1\right)} j} . \tag{2.10}
\end{equation*}
$$

Remark 2.2. Let us mention that wavelet leaders can not be used to prove the finiteness of the limits in Theorem 1.2 because they do not precisely characterize the pointwise regularity, see for instance [27, 31] for more details.

## 3 Upper bounds for oscillations

Starting from now and until the end of the paper, we fix $H_{1}, H_{2} \in\left(\frac{1}{2}, 1\right)$ such that $H_{1}+H_{2}>\frac{3}{2}$. In this section, we show the finiteness of the limits (1.6), (1.7) and (1.8). Concerning the rapid points, we will in fact show a stronger result, obtaining an almost sure uniform modulus of continuity for the generalized Rosenblatt process.

We use a wavelet-type expansion of the generalized Rosenblatt process. It relies on the following random variables.

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Definition 3.1. For all $\left(j_{1}, j_{2}, k_{1}, k_{2}\right) \in \mathbb{Z}^{4}$, let $\varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}}$ be the random variable defined by $2^{\frac{j_{1}+j_{2}}{2}} \int_{\mathbb{R}^{2}}^{\prime} \psi\left(2^{j_{1}} x_{1}-k_{1}\right) \psi\left(2^{j_{2}} x_{2}-k_{2}\right) d B\left(x_{1}\right) d B\left(x_{2}\right)$.
Remark 3.2. For all $\left(j_{1}, j_{2}, k_{1}, k_{2}\right) \in \mathbb{Z}^{4}$, we have ([6, Proposition 2.3])

$$
\begin{equation*}
\varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}}=\left(2^{\frac{j_{1}}{2}} \int_{\mathbb{R}} \psi\left(2^{j_{1}} x-k_{1}\right) d B(x)\right)\left(2^{\frac{j_{2}}{2}} \int_{\mathbb{R}} \psi\left(2^{j_{2}} x-k_{2}\right) d B(x)\right) \tag{3.1}
\end{equation*}
$$

for $j_{1} \neq j_{2}$ or $k_{1} \neq k_{2}$, and

$$
\begin{equation*}
\varepsilon_{j_{1}, j_{1}}^{k_{1}, k_{1}}=\left(2^{\frac{j_{1}}{2}} \int_{\mathbb{R}} \psi\left(2^{j_{1}} x-k_{1}\right) d B(x)\right)^{2}-1 \tag{3.2}
\end{equation*}
$$

for $j_{1}=j_{2}$ and $k_{1}=k_{2}$. Using the fact that $\left(2^{j / 2} \psi\left(2^{j} .-k\right)\right)_{(j, k) \in \mathbb{Z}^{2}}$ forms an orthonormal basis of $L^{2}(\mathbb{R})$, and elementary properties of Wiener integral, we know that $\left(2^{j / 2} \int_{\mathbb{R}} \psi\left(2^{j} x-\right.\right.$ $k) d B(x))_{(j, k) \in \mathbb{Z}^{2}}$ is a family of iid $\mathcal{N}(0,1)$ random variables. So the random variables $\varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}}$ and $\varepsilon_{j_{1}^{\prime}, j_{2}^{\prime}}^{k_{1}^{\prime}, k_{2}^{\prime}}$ are independent as long as $\left\{\left(j_{1}, k_{1}\right),\left(j_{2}, k_{2}\right)\right\} \cap\left\{\left(j_{1}^{\prime}, k_{1}^{\prime}\right),\left(j_{2}^{\prime}, k_{2}^{\prime}\right)\right\}=\emptyset$.

The following theorem, proved in [6], gives the wavelet-type expansion we use in this section.
Theorem 3.3. Let $\psi$ be the Meyer wavelet and $I$ be any compact interval of $\mathbb{R}_{+}$. Almost surely, the random series

$$
\begin{equation*}
\sum_{\left(j_{1}, j_{2}, k_{1}, k_{2}\right) \in \mathbb{Z}^{4}} 2^{j_{1}\left(1-H_{1}\right)+j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}} \int_{0}^{t} \psi_{H_{1}}\left(2^{j_{1}} x-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} x-k_{2}\right) d x \tag{3.3}
\end{equation*}
$$

converges uniformly to $R_{H_{1}, H_{2}}$ on the interval $I$.
Remark 3.4. Any open interval in $\mathbb{R}$ can be written as a countable union of dyadic intervals $\left(\lambda_{j, k}\right)_{j \in \mathbb{N}, k \in \mathbb{Z}}$. Then, to prove Theorem 1.2, it is sufficient to show that, for all $j \in \mathbb{N}, k \in \mathbb{Z}$, there exist an event $\Omega_{j, k}$ of probability 1 such that, for all $\omega \in \Omega_{j, k}$, almost every $t \in \lambda_{j, k}$ is ordinary and there exist $t_{r} \in \lambda_{j, k}$ which is rapid and $t_{s} \in \lambda_{j, k}$ which is slow. For notational simplicity, we will only do the proofs in full details for $\lambda_{0,0}=[0,1)$. In fact, after dilatation and translation, our proofs hold true for any arbitrary dyadic interval.

### 3.1 Rapid points

Let us first focus on rapid points. we prove that $x \mapsto|x|^{H_{1}+H_{2}-1} \log |x|^{-1}$ is almost surely a uniform modulus of continuity for $R_{H_{1}, H_{2}}$.
Proposition 3.5. There exists an event $\Omega_{\mathrm{rap}}$ of probability 1 such that for all $\omega \in \Omega_{\mathrm{rap}}$ there exists $C_{R}(\omega)>0$ such that, for all $t, s \in(0,1)$, we have

$$
\begin{equation*}
\left|R_{H_{1}, H_{2}}(t, \omega)-R_{H_{1}, H_{2}}(s, \omega)\right| \leqslant C_{R}(\omega)|t-s|^{H_{1}+H_{2}-1} \log |t-s|^{-1} . \tag{3.4}
\end{equation*}
$$

Let us set, for all $s, t \in(0,1)$ and $\left(j_{1}, j_{2}, k_{1}, k_{2}\right) \in \mathbb{Z}^{4}$,

$$
I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]=\int_{[t, s]} \psi_{H_{1}}\left(2^{j_{1}} x-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} x-k_{2}\right) d x .
$$

All along this section, if $s, t \in(0,1)$ are given, $n$ always refers to the unique positive integer such that $2^{-n-1}<|t-s| \leqslant 2^{-n}$. Our proof consists in writing

$$
\begin{equation*}
\left.\mid R_{H_{1}, H_{2}}(t, \cdot)\right)-R_{H_{1}, H_{2}}(s, \cdot)\left|=\left|\sum_{\left(j_{1}, j_{2}, k_{1}, k_{2}\right) \in \mathbb{Z}^{4}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}} I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right|\right. \tag{3.5}
\end{equation*}
$$

and splitting the sum in the right-hand side into subsums determined according to the position of $j_{1}$ and $j_{2}$ with respect to $n$. To bound from above some of these subsums the following lemma is key.

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Lemma 3.6. [6, Lemma 2.4.] There exist an event $\Omega^{*}$ of probability 1 and a positive random variable $C_{1}$ with finite moment of any order, such that, for all $\omega \in \Omega^{*}$ and for each $\left(j_{1}, j_{2}, k_{1}, k_{2}\right) \in \mathbb{Z}^{4}$,

$$
\begin{equation*}
\left|\varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}}(\omega)\right| \leqslant C_{1}(\omega) \sqrt{\log \left(3+\left|j_{1}\right|+\left|k_{1}\right|\right)} \sqrt{\log \left(3+\left|j_{2}\right|+\left|k_{2}\right|\right)} . \tag{3.6}
\end{equation*}
$$

In view of Lemma 3.6, we set $L_{j_{1}, j_{2}}^{k_{1}, k_{2}}=\sqrt{\log \left(3+\left|j_{1}\right|+\left|k_{1}\right|\right)} \sqrt{\log \left(3+\left|j_{2}\right|+\left|k_{2}\right|\right)}$. As a first step, Lemmata 3.7 to 3.12 are devoted to bound some deterministic series whose general term is $2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} L_{j_{1}, j_{2}}^{k_{1}, k_{2}}\left|I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right|$.

This first lemma will be useful to bound the subsums in the right-hand side of (3.5) for $j_{1}<n$ and $j_{2}<n$.
Lemma 3.7. There exists a deterministic constant $C>0$ such that, for all $t, s \in(0,1)$, we have

$$
\sum_{j_{1}<n} \sum_{j_{2}<n} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} L_{j_{1}, j_{2}}^{k_{1}, k_{2}}\left|I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right| \leqslant C|t-s|^{H_{1}+H_{2}-1} \log |t-s|^{-1} .
$$

Proof. Let us start by considering, for all $\left(j_{1}, j_{2}\right) \in \mathbb{Z}^{2}$, the series

$$
\begin{aligned}
R_{j_{1}, j_{2}} & : t \mapsto \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} L_{j_{1}, j_{2}}^{k_{1}, k_{2}} \int_{0}^{t}\left|\psi_{H_{1}}\left(2^{j_{1}} x-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} x-k_{2}\right)\right| d x \text { and } \\
R_{j_{1}, j_{2}}^{\prime} & : t \mapsto \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} L_{j_{1}, j_{2}}^{k_{1}, k_{2}}\left|\psi_{H_{1}}\left(2^{j_{1}} t-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} t-k_{2}\right)\right| .
\end{aligned}
$$

The fast decay of the fractional antiderivatives of $\psi$ allows us to write, for all $H \in\left\{H_{1}, H_{2}\right\}$ and for all $x \in \mathbb{R}$

$$
\begin{equation*}
\left|\psi_{H}(x)\right| \leqslant C(1+|x|)^{-4} . \tag{3.7}
\end{equation*}
$$

Moreover, according to [6, Lemma 4.2] for all $L>1$ there exists $C>0$ such that, for all $j \in \mathbb{Z}$ and $x \in \mathbb{R}$

$$
\begin{equation*}
\sum_{k \in \mathbb{Z}} \frac{\sqrt{\log (3+|j|+k)}}{\left(3+\left|2^{j} x-k\right|\right)^{L}} \leqslant C \sqrt{\log \left(3+|j|+2^{j}|x|\right)} \tag{3.8}
\end{equation*}
$$

Therefore, if $K$ is any compact set of $\mathbb{R}_{+}$, if $M=\sup _{K}$, for all $t \in K$, we have

$$
\begin{aligned}
\left|R_{j_{1}, j_{2}}(t)\right| & \leqslant C \int_{0}^{t} \sqrt{\log \left(3+\left|j_{1}\right|+2^{j_{1}}|x|\right)} \sqrt{\log \left(3+\left|j_{2}\right|+2^{j_{2}}|x|\right)} d x \\
& \leqslant C M \sqrt{\log \left(3+\left|j_{1}\right|+2^{j_{1}} M\right)} \sqrt{\log \left(3+\left|j_{2}\right|+2^{j_{2}} M\right)}
\end{aligned}
$$

The same arguments can be applied to $R_{j_{1}, j_{2}}^{\prime}$, which means that both series converge uniformly on any compact set of $\mathbb{R}_{+}$. From this, we can use mean value theorem: for all $\left(j_{1}, j_{2}\right) \in \mathbb{Z}^{2}$ there is $\xi\left(j_{1}, j_{2}\right) \in[s, t]$ such that

$$
\begin{align*}
& \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} L_{j_{1}, j_{2}}^{k_{1}, k_{2}}\left|I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right| \\
& \leqslant|t-s| \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} L_{j_{1}, j_{2}}^{k_{1}, k_{2}}\left|\psi_{H_{1}}\left(2^{j_{1}} \xi-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} \xi-k_{2}\right)\right| . \tag{3.9}
\end{align*}
$$

Now, we use the fast decay of the fractional antiderivatives of $\psi$ (3.7) and inequality (3.8) to bound (3.9) from above: for all $j_{1}, j_{2}<n$,

$$
\sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} L_{j_{1}, j_{2}}^{k_{1}, k_{2}}\left|\psi_{H_{1}}\left(2^{j_{1}} \xi-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} \xi-k_{2}\right)\right|
$$

$$
\begin{aligned}
& \leqslant C\left(\sum_{k_{1} \in \mathbb{Z}} \frac{\sqrt{\log \left(3+\left|j_{1}\right|+\left|k_{1}\right|\right)}}{\left(3+\left|2^{j_{1}} \xi-k_{1}\right|\right)^{4}}\right)\left(\sum_{k_{2} \in \mathbb{Z}} \frac{\sqrt{\log \left(3+\left|j_{2}\right|+\left|k_{2}\right|\right)}}{\left(3+\left|2^{j_{2}} \xi-k_{2}\right|\right)^{4}}\right) \\
& \leqslant C \sqrt{\log \left(3+\left|j_{1}\right|+2^{j_{1}}|\xi|\right)} \sqrt{\log \left(3+\left|j_{2}\right|+2^{j_{2}}|\xi|\right)} \\
& \leqslant C \sqrt{\log \left(3+\left|j_{1}\right|+2^{j_{1}}\right)} \sqrt{\log \left(3+\left|j_{2}\right|+2^{j_{2}}\right)},
\end{aligned}
$$

as $\xi \in(0,1)$. Let us then remark that

$$
\begin{align*}
& \sum_{j_{1}<n} 2^{j_{1}\left(1-H_{1}\right)} \sqrt{\log \left(3+\left|j_{1}\right|+2^{j_{1}}\right)} \\
& \quad=\sum_{j_{1} \leqslant 0} 2^{j_{1}\left(1-H_{1}\right)} \sqrt{\log \left(3+\left|j_{1}\right|+2^{j_{1}}\right)}+\sum_{j_{1}=0}^{n-1} 2^{j_{1}\left(1-H_{1}\right)} \sqrt{\log \left(3+\left|j_{1}\right|+2^{j_{1}}\right)} \\
& \quad \leqslant C+\sum_{j_{1}=0}^{n-1} 2^{j_{1}\left(1-H_{1}\right)} \sqrt{\log \left(3+\left|j_{1}\right|+2^{j_{1}}\right)} \leqslant C 2^{n\left(1-H_{1}\right)} \sqrt{n}, \tag{3.10}
\end{align*}
$$

as $1-H_{1}>0$. The same can be applied to the sum over $j_{2}$ and we finally get

$$
\begin{aligned}
& \sum_{j_{1}<n} \sum_{j_{2}<n} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} L_{j_{1}, j_{2}}^{k_{1}, k_{2}}\left|I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right| \\
& \leqslant C|t-s| \sum_{j_{1}<n} \sum_{j_{2}<n} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \sqrt{\log \left(3+\left|j_{1}\right|+2^{j_{1}}\right)} \sqrt{\log \left(3+\left|j_{2}\right|+2^{j_{2}}\right)} \\
& \leqslant C|t-s| 2^{n\left(2-H_{1}-H_{2}\right)} n \leqslant C|t-s|^{H_{1}+H_{2}-1} \log |t-s|^{-1} .
\end{aligned}
$$

Lemmata 3.10 and 3.11 will help finding an upper bound for the subsums in the right-hand side of (3.5) with $j_{1}<n \leqslant j_{2}$ or $j_{2}<n \leqslant j_{1}$ as well as the ones where $n \leqslant j_{1} \leqslant j_{2}$ and $n \leqslant j_{2} \leqslant j_{1}$. Let us define the following partition of $\mathbb{Z}$, which determines the relative positions of $\left[k_{2} 2^{-j_{2}},\left(k_{2}+1\right) 2^{-j_{2}}\right)$ and $[s, t]$.
Definition 3.8. For all $j_{2} \in \mathbb{N}$, we set

$$
\begin{aligned}
\mathbb{Z}_{j_{2}}^{<}(t, s) & =\left\{k_{2} \in \mathbb{Z}: k_{2} 2^{-j_{2}}<\min \{t, s\}\right\}, \\
\mathbb{Z}_{j_{2}}^{>}(t, s) & =\left\{k_{2} \in \mathbb{Z}: k_{2} 2^{-j_{2}}>\max \{t, s\}\right\}, \\
\text { and } \mathbb{Z}_{j_{2}}[t, s] & =\mathbb{Z} \backslash\left(\mathbb{Z}_{j_{2}}^{<}(t, s) \cup \mathbb{Z}_{j_{2}}^{>}(t, s)\right) .
\end{aligned}
$$

Remark 3.9. Note that we have $\# \mathbb{Z}_{j_{2}}[t, s] \leqslant 2^{j_{2}-n}+1$.
Let us also observe that for all $a, b>0$,

$$
\begin{equation*}
\log (3+a+b) \leqslant \log (3+a) \log (3+b) \tag{3.11}
\end{equation*}
$$

Lemma 3.10. There exists a deterministic constant $C>0$ such that, for all $t, s \in(0,1)$ and $j_{1} \leqslant j_{2}$, the quantities

$$
\begin{align*}
& \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}(t, s)} L_{j_{1}, j_{2}}^{k_{1}, k_{2}}\left|I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right|  \tag{3.12}\\
& \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}(t, s)} L_{j_{1}, j_{2}}^{k_{1}, k_{2}}\left|I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right| \tag{3.13}
\end{align*}
$$

are bounded from above by $C \sqrt{\log \left(3+\left|j_{1}\right|+2^{j_{1}}\right)} \sqrt{\log \left(3+\left|j_{2}\right|+2^{j_{2}}\right)} 2^{-j_{2}}$.

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Proof. Let us bound (3.12), the proof for (3.13) being similar. From the fast decay of the fractional antiderivatives of $\psi$ (3.7), inequalities (3.8) and (3.11) for $j_{1} \leqslant j_{2}$, we have

$$
\begin{aligned}
(3.12) \leqslant & C \int_{[s, t]}\left(\sum_{k_{1} \in \mathbb{Z}} \frac{\sqrt{\log \left(3+\left|j_{1}\right|+\left|k_{1}\right|\right)}}{\left(3+\left|2^{j_{1}} x-k_{1}\right|\right)^{4}}\right)\left(\sum_{k_{2} \in \mathbb{Z}_{j_{2}}(t, s)} \frac{\sqrt{\log \left(3+\left|j_{2}\right|+\left|k_{2}\right|\right)}}{\left(3+\left|2^{j_{2}} x-k_{2}\right|\right)^{4}}\right) d x \\
\leqslant & C \sqrt{\log \left(3+\left|j_{1}\right|+2^{j_{1}}\right)} \sqrt{\log \left(3+\left|j_{2}\right|+2^{j_{2}}\right)} \\
& \int_{[s, t]} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}(t, s)} \frac{\sqrt{\log \left(3+\left|2^{j_{2}} x-k_{2}\right|\right)}}{\left(3+\left|2^{j_{2}} x-k_{2}\right|\right)^{4}} d x .
\end{aligned}
$$

For all $x \in[s, t]$ the mapping $\left.y \mapsto 2+2^{j_{2}} x-2^{j_{2}} \min \{s, t\}+y\right)^{-3}$ is decreasing and thus

$$
\begin{align*}
\int_{[s, t]} & \sum_{k_{2} \in \mathbb{Z}_{j_{2}}^{<}(t, s)} \frac{\sqrt{\log \left(3+\left|2^{j_{2}} x-k_{2}\right|\right)}}{\left(3+\left|2^{j_{2}} x-k_{2}\right|\right)^{4}} d x \\
& \leqslant \int_{[s, t]} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}(t, s)} \frac{d x}{\left(3+2^{j_{2}} x-k_{2}\right)^{3}} \\
& \leqslant \int_{[s, t]} \sum_{m=0}^{+\infty} \frac{d x}{\left(3+2^{j_{2}} x-2^{j_{2}} \min \{s, t\}+m\right)^{3}} \\
& \leqslant \int_{[s, t]} \int_{0}^{+\infty} \frac{d x d y}{\left(2+2^{j_{2}} x-2^{j_{2}} \min \{s, t\}+y\right)^{3}} \leqslant C 2^{-j_{2}} . \tag{3.14}
\end{align*}
$$

This bound leads to $(3.12) \leqslant C \sqrt{\log \left(3+\left|j_{1}\right|+2^{j_{1}}\right)} \sqrt{\log \left(3+\left|j_{2}\right|+2^{j_{2}}\right)} 2^{-j_{2}}$.
Lemma 3.11. There exists a deterministic constant $C>0$ such that, for all $t, s \in(0,1)$ and $j_{1} \leqslant j_{2}$, the quantities

$$
\begin{aligned}
& \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}[t, s]} L_{j_{1}, j_{2}}^{k_{1}, k_{2}}\left|\int_{-\infty}^{\min \{s, t\}} \psi_{H_{1}}\left(2^{j_{1}} x-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} x-k_{2}\right) d x\right| \\
& \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}[t, s]} L_{j_{1}, j_{2}}^{k_{1}, k_{2}}\left|\int_{\max \{s, t\}}^{+\infty} \psi_{H_{1}}\left(2^{j_{1}} x-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} x-k_{2}\right) d x\right|
\end{aligned}
$$

are bounded from above by $C \sqrt{\log \left(3+\left|j_{1}\right|+2^{j_{2}}\right)} \sqrt{\log \left(3+\left|j_{2}\right|+2^{j_{2}}\right)} 2^{-j_{2}}$.
Proof. Let us assume that $s \leqslant t$, the argument for $t<s$ being similar. As $j_{2} \geqslant j_{1}$, we have, by inequality (3.8),

$$
\begin{aligned}
& \int_{-\infty}^{s}\left(\sum_{k_{1} \in \mathbb{Z}} \frac{\sqrt{\log \left(3+\left|j_{1}\right|+\left|k_{1}\right|\right)}}{\left(3+\left|2^{j_{1}} x-k_{1}\right|\right)^{4}}\right)\left(\sum_{k_{2} \in \mathbb{Z}_{j_{2}}[t, s]} \frac{\sqrt{\log \left(3+\left|j_{2}\right|+\left|k_{2}\right|\right)}}{\left(3+\left|2^{j_{2}} x-k_{2}\right|\right)^{4}}\right) d x \\
& \quad \leqslant C_{L} \int_{-\infty}^{s} \sqrt{\log \left(3+\left|j_{1}\right|+2^{j_{1}}|x|\right)}\left(\sum_{k_{2} \in \mathbb{Z}_{j_{2}}[t, s]} \frac{\sqrt{\log \left(3+\left|j_{2}\right|+\left|k_{2}\right|\right)}}{\left(3+\left|2^{j_{2}} x-k_{2}\right|\right)^{4}}\right) d x \\
& \quad \leqslant C_{L} \int_{-\infty}^{s} \sqrt{\log \left(3+\left|j_{1}\right|+2^{j_{2}}|x|\right)}\left(\sum_{k_{2} \in \mathbb{Z}_{j_{2}}[t, s]} \frac{\sqrt{\log \left(3+\left|j_{2}\right|+\left|k_{2}\right|\right)}}{\left(3+\left|2^{j_{2}} x-k_{2}\right|\right)^{4}}\right) d x .
\end{aligned}
$$

For all $k_{2} \in \mathbb{Z}_{j_{2}}[t, s],\left|k_{2}\right| \leqslant 2^{j_{2}}$, we have, using (3.11),

$$
\log \left(3+\left|j_{1}\right|+2^{j_{2}}|x|\right) \leqslant \log \left(3+\left|j_{1}\right|+2^{j_{2}}\right) \log \left(3+\left|2^{j_{2}} x-k_{2}\right|\right) \text { and }
$$

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$$
\log \left(3+\left|j_{2}\right|+\left|k_{2}\right|\right) \leqslant \log \left(3+j_{2}+2^{j_{2}}\right) \log \left(3+\left|2^{j_{2}} x-k_{2}\right|\right) .
$$

Thus, it only remains us to deal with

$$
\int_{-\infty}^{s} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}[t, s]} \frac{d x}{\left(3+\left|2^{j_{2}} x-k_{2}\right|\right)^{3}} .
$$

But, for all $x \leqslant s$ and $k_{2} \in \mathbb{Z}_{j_{2}}[t, s],\left|2^{j_{2}} x-k_{2}\right|=k_{2}-2^{j_{2}} x$ and then, using the same method as in (3.14), we get

$$
\begin{equation*}
\int_{-\infty}^{s} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}[t, s]} \frac{d x}{\left(3+k_{2}-2^{j_{2}} x\right)^{3}} \leqslant C 2^{-j_{2}} \tag{3.15}
\end{equation*}
$$

which finally leads to

$$
\begin{align*}
& \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}[t, s]} L_{j_{1}, j_{2}}^{k_{1}, k_{2}}\left|\int_{-\infty}^{s} \psi_{H_{1}}\left(2^{j_{1}} x-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} x-k_{2}\right) d x\right| \\
& \quad \leqslant C \sqrt{\log \left(3+\left|j_{1}\right|+2^{j_{2}}\right)} \sqrt{\log \left(3+j_{2}+2^{j_{2}}\right)} 2^{-j_{2}} \tag{3.16}
\end{align*}
$$

We get in the same way,

$$
\begin{aligned}
& \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}[t, s]} L_{j_{1}, j_{2}}^{k_{1}, k_{2}}\left|\int_{t}^{+\infty} \psi_{H_{1}}\left(2^{j_{1}} x-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} x-k_{2}\right) d x\right| \\
& \quad \leqslant C \sqrt{\log \left(3+\left|j_{1}\right|+2^{j_{2}}\right)} \sqrt{\log \left(3+j_{2}+2^{j_{2}}\right)} 2^{-j_{2}}
\end{aligned}
$$

Next Lemma will be used to bound the subsums of (3.5) with $j_{1}<n \leqslant j_{2}$ or $j_{2}<n \leqslant$ $j_{1}$.
Lemma 3.12. There exists a deterministic constant $C>0$ such that, for all $t, s \in(0,1)$, the quantities

$$
\begin{aligned}
& R^{<\geqslant n}[t, s]:=\sum_{j_{1}<n} \sum_{j_{2} \geqslant n} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} L_{j_{1}, j_{2}}^{k_{1}, k_{2}}\left|I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right| \\
& R^{\geqslant<n}[t, s]:=\sum_{j_{1} \geqslant n} \sum_{j_{2}<n} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} L_{j_{1}, j_{2}}^{k_{1}, k_{2}}\left|I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right|
\end{aligned}
$$

are bounded from above by $C|t-s|^{H_{1}+H_{2}-1} \log |t-s|^{-1}$.
Proof. As $R^{<\geqslant n}[t, s]$ and $R^{\geqslant<n}[t, s]$ can clearly be treated symmetrically, we restrict our attention to $R^{<\geqslant n}[t, s]$. One sees that

$$
\begin{align*}
& \sum_{j_{1}<n} \sum_{j_{2} \geqslant n} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} L_{j_{1}, j_{2}}^{k_{1}, k_{2}} I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s] \\
& =\sum_{j_{1}<n} \sum_{j_{2} \geqslant n} \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}(t, s)} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} L_{j_{1}, j_{2}}^{k_{1}, k_{2}} I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]  \tag{3.17}\\
& +\sum_{j_{1}<n} \sum_{j_{2} \geqslant n} \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}^{>}(t, s)} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} L_{j_{1}, j_{2}}^{k_{1}, k_{2}} I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]  \tag{3.18}\\
& +\sum_{j_{1}<n} \sum_{j_{2} \geqslant n} \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}[t, s]} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} L_{j_{1}, j_{2}}^{k_{1}, k_{2}} I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s] . \tag{3.19}
\end{align*}
$$

For (3.17), we use Lemma 3.10 to get

$$
|(3.17)| \leqslant C \sum_{j_{1}<n} 2^{j_{1}\left(1-H_{1}\right)} \sqrt{\log \left(3+\left|j_{1}\right|+2^{j_{1}}\right)} \sum_{j_{2} \geqslant n} 2^{-j_{2} H_{2}} \sqrt{\log \left(3+\left|j_{2}\right|+2^{j_{2}}\right)} .
$$

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The sum over over $j_{1}$ is bounded just as in (3.10) while, for the sum over $j_{2}$, we have

$$
\begin{align*}
\sum_{j_{2} \geqslant n} 2^{-j_{2} H_{2}} \sqrt{\log \left(3+\left|j_{2}\right|+2^{j_{2}}\right)} & \leqslant \sum_{j_{2} \geqslant n} 2^{-j_{2} H_{2}} \sqrt{\log \left(3+2^{j_{2}+1}\right)} \\
& \leqslant C 2^{-n H_{2}} \sqrt{n} \tag{3.20}
\end{align*}
$$

We bound (3.18) in exactly the same way.
For (3.19), let us again assume $s \leqslant t$, then we write

$$
\begin{align*}
I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s] & =\int_{\mathbb{R}} \psi_{H_{1}}\left(2^{j_{1}} x-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} x-k_{2}\right) d x \\
& -\int_{-\infty}^{s} \psi_{H_{1}}\left(2^{j_{1}} x-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} x-k_{2}\right) d x \\
& -\int_{t}^{+\infty} \psi_{H_{1}}\left(2^{j_{1}} x-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} x-k_{2}\right) d x . \tag{3.21}
\end{align*}
$$

Since $j_{1}<n$ and $j_{2} \geqslant n$, recalling Lemma 2.1, the sum

$$
\sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}[t, s]} L_{j_{1}, j_{2}}^{k_{1}, k_{2}} \int_{\mathbb{R}} \psi_{H_{1}}\left(2^{j_{1}} x-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} x-k_{2}\right) d x
$$

vanishes except maybe when $\left(j_{1}, j_{2}\right)=(n-1, n)$. In this case, note that $\# \mathbb{Z}_{n}[t, s] \leqslant 2$ and, for all $k_{2} \in \mathbb{Z}_{n}[t, s],\left|k_{2}\right| \leqslant 2^{n}$. Then, by Lemma 2.1 and inequality (3.8), we get

$$
\begin{aligned}
& \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{n}[t, s]} L_{n-1, n}^{k_{1}, k_{2}}\left|I_{n-1, n}^{k_{1}, k_{2}}\right| \\
& \leqslant C 2^{-n} \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{n}[t, s]} \frac{\sqrt{\log \left(3+n-1+\left|k_{1}\right|\right)} \sqrt{\log \left(3+n+\left|k_{2}\right|\right.}}{\left(3+\left|2 k_{1}-k_{2}\right|\right)^{4}} \\
& \leqslant C 2^{-n} \sum_{k_{2} \in \mathbb{Z}_{n}[t, s]} \sqrt{\log \left(3+n-1+\left|\frac{k_{2}}{2}\right|\right)} \sqrt{\log \left(3+n+\left|k_{2}\right|\right)} \\
& \leqslant C 2^{-n} n
\end{aligned}
$$

Now, using Lemma 3.11, we also get

$$
\begin{align*}
& \left|\sum_{j_{1}<n} \sum_{j_{2} \geqslant n} \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}[t, s]} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} L_{j_{1}, j_{2}}^{k_{1}, k_{2}} \int_{-\infty}^{s} \psi_{H_{1}}\left(2^{j_{1}} x-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} x-k_{2}\right) d x\right| \\
& \leqslant \sum_{j_{1}<n} \sum_{j_{2} \geqslant n} 2^{j_{1}\left(1-H_{1}\right)} 2^{-j_{2} H_{2}} \sqrt{\log \left(3+\left|j_{1}\right|+2^{j_{2}}\right)} \sqrt{\log \left(3+\left|j_{2}\right|+2^{j_{2}}\right)} \\
& \leqslant \sum_{j_{1}<0} \sum_{j_{2} \geqslant n} 2^{j_{1}\left(1-H_{1}\right)} 2^{-j_{2} H_{2}} \sqrt{\log \left(3+\left|j_{1}\right|\right)} \log \left(3+2^{j_{2}+1}\right) \\
& \quad+\sum_{j_{1}=0}^{n-1} \sum_{j_{2} \geqslant n} 2^{j_{1}\left(1-H_{1}\right)} 2^{-j_{2} H_{2}} \sqrt{\log \left(3+2^{j_{2}+1}\right)} \sqrt{\log \left(3+2^{j_{2}+1}\right)} \\
& \leqslant C 2^{n\left(1-H_{1}-H_{2}\right)} n \tag{3.22}
\end{align*}
$$

The series

$$
\left|\sum_{j_{1}<n} \sum_{j_{2} \geqslant n} \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}[t, s]} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} L_{j_{1}, j_{2}}^{k_{1}, k_{2}} \int_{-\infty}^{s} \psi_{H_{1}}\left(2^{j_{1}} x-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} x-k_{2}\right) d x\right|
$$

is bounded in exactly the same way and the conclusion follows.

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It remains us to bound the subsums of (3.5) with $j_{1} \geqslant n$ and $j_{2} \geqslant n$. For this, let us define some random variables associated with dyadic intervals.
Definition 3.13. If $\lambda$ is a dyadic interval of scale $n$, we define, for all $j \geqslant n$, the indexation sets

$$
\begin{aligned}
& S_{j}^{0}(\lambda):=\left\{\left(k^{(1)}, K^{(1)}, k^{(2)}, K^{(2)}\right) \in \mathbb{Z}^{4}: \frac{k^{(1)}}{2^{j}}, \frac{K^{(1)}}{2^{j}}, \frac{k^{(2)}}{2^{j}}, \frac{K^{(2)}}{2^{j}} \in \lambda\right\}, \\
& S_{j}^{1}(\lambda):=\left\{\left(k^{(1)}, K^{(1)}, k^{(2)}, K^{(2)}\right) \in \mathbb{Z}^{4}: \frac{k^{(1)}}{2^{j+1}}, \frac{K^{(1)}}{2^{j+1}}, \frac{k^{(2)}}{2^{j}}, \frac{K^{(2)}}{2^{j}} \in \lambda\right\}, \\
& S_{j}^{2}(\lambda):=\left\{\left(k^{(1)}, K^{(1)}, k^{(2)}, K^{(2)}\right) \in \mathbb{Z}^{4}: \frac{k^{(1)}}{2^{j}}, \frac{K^{(1)}}{2^{j}}, \frac{k^{(2)}}{2^{j+1}}, \frac{K^{(2)}}{2^{j+1}} \in \lambda\right\}
\end{aligned}
$$

and consider the random variables, for $\left(k^{(1)}, K^{(1)}, k^{(2)}, K^{(2)}\right) \in S_{j}^{0}(\lambda)$,

$$
\begin{equation*}
{ }_{j}^{0} \sum_{k^{(1)}, K^{(1)}}^{k^{(2)}, K^{(2)}}:=\sum_{k^{(1)} \leqslant k_{1} \leqslant K^{(1)}} \sum_{k^{(2)} \leqslant k_{2} \leqslant K^{(2)}} \varepsilon_{j, j}^{k_{1}, k_{2}} I_{j, j}^{k_{1}, k_{2}} \tag{3.23}
\end{equation*}
$$

, for $\left(k^{(1)}, K^{(1)}, k^{(2)}, K^{(2)}\right) \in S_{j}^{1}(\lambda)$,

$$
\begin{equation*}
{ }_{j}^{1} \sum_{k^{(1)}, K^{(1)}}^{k^{(2)}, K^{(2)}}:=\sum_{k^{(1)} \leqslant k_{1} \leqslant K^{(1)}} \sum_{k^{(2)} \leqslant k_{2} \leqslant K^{(2)}} \varepsilon_{j+1, j}^{k_{1}, k_{2}} I_{j+1, j}^{k_{1}, k_{2}} \tag{3.24}
\end{equation*}
$$

and, for $\left(k^{(1)}, K^{(1)}, k^{(2)}, K^{(2)}\right) \in S_{j}^{2}(\lambda)$,

$$
\begin{equation*}
{ }_{j}^{2} \sum_{k^{(1)}, K^{(1)}}^{k^{(2)}, K^{(2)}}:=\sum_{k^{(1)} \leqslant k_{1} \leqslant K^{(1)}} \sum_{k^{(2)} \leqslant k_{2} \leqslant K^{(2)}} \varepsilon_{j, j+1}^{k_{1}, k_{2}} I_{j, j+1}^{k_{1}, k_{2}} . \tag{3.25}
\end{equation*}
$$

The idea behind the definition of these random variables is, as $|t-s| \leqslant 2^{-n}, s \in 3 \lambda_{n}(t)$ and thus any sum of the form

$$
\begin{equation*}
\sum_{k_{1} \in \mathbb{Z}_{j}[t, s]} \sum_{k_{2} \in \mathbb{Z}_{\ell}[t, s]} \varepsilon_{j, \ell}^{k_{1}, k_{2}} I_{j, \ell}^{k_{1}, k_{2}} \tag{3.26}
\end{equation*}
$$

for $\ell \in\{j, j+1\}$ can be written as the sum of random variables (3.23), (3.24) or (3.25) for some $\left(k^{(1)}, K^{(1)}, k^{(2)}, K^{(2)}\right)$ belonging to at most two $S_{j}^{\ell}(\lambda)(\ell \in\{0,1,2\})$ with $\lambda \in \lambda_{n}(t)$. Indeed,

- if $t$ and $s$ both belong to $\lambda_{n}(t)$ then we only need to rewrite (3.26) in the form (3.23), (3.24) or (3.25) for $\left(k^{(1)}, K^{(1)}, k^{(2)}, K^{(2)}\right) \in S_{j}^{\ell}\left(\lambda_{n}(t)\right)$;
- if $s \in \lambda$ with $\lambda \in 3 \lambda_{n}(t) \backslash \lambda_{n}(t)$ then we need to consider a first sum indexed by a quadruple of $S_{j}^{\ell}\left(\lambda_{n}(t)\right)$ and a second indexed by a quadruple of $S_{j}^{\ell}(\lambda)$.
The reason why we decide to put $\lambda$ instead of $3 \lambda$ in the definition of the sets $S_{j}^{\ell}(\lambda)$ is that if, for all $n \in \mathbb{N}$ and for all $\lambda \in \Lambda_{n}$ and $j \geqslant n$, we define the random variable
we want $\Xi_{j}(\lambda)$ to be independent of $\Xi_{j}\left(\lambda^{\prime}\right)$ as long as $\lambda \cap \lambda^{\prime}=\emptyset$. Moreover, from the definitions of the random variables (3.24), (3.23) and (3.25), the remarks below Theorem 3.3 and the explicit expressions (2.3), (2.4) and (2.5), the law of $\Xi_{j}(\lambda)$ does not depend on $\lambda \in \Lambda_{n}$ but only on $j-n$.

Wavelet methods to study the pointwise regularity of the generalized Rosenblatt process

Let us remind the reader that the first order Wiener chaos (with respect to Brownian motion $B$ ) is the closed linear subspace of $L^{2}(\Omega)$ consisting of random variables of type $\int_{\mathbb{R}} f(x) d B(x)$, for $f \in L^{2}(\mathbb{R})$. The second order Wiener chaos is the closed linear subspace of $L^{2}(\Omega)$ consisting of random variables of type $\int_{\mathbb{R}^{2}}^{\prime} f(x, y) d B(x) d B(y)$, for $f \in L^{2}\left(\mathbb{R}^{2}\right)$. Fractional Brownian motion belongs to first order Wiener chaos while the (generalized) Rosenblatt process belongs to the second order Wiener chaos. The key results to estimate the random variables $\Xi_{j}$ are [24, Theorem 6.7 and Theorem 6.12] that we recall here.
Theorem 3.14. There exists a strictly positive universal deterministic constant $\stackrel{\star}{C}$ such that, for every random variable $X$ belonging to the second order Wiener chaos and for each real number $y \geqslant 2$, one has $\mathbb{P}\left(|X| \geqslant y\|X\|_{L^{2}(\Omega)}\right) \leqslant \exp (-\stackrel{\star}{C} y)$.
Theorem 3.15. If $X$ is a random variable belonging to the second order Wiener chaos, there exist $a, b, y_{0}>0$ such that, for all $y \geqslant y_{0}$, one has $\exp (-a y) \leqslant \mathbb{P}(|X| \geqslant y) \leqslant$ $\exp (-b y)$.
Remark 3.16. As stated in [24], the constants $a, b$ in Theorem 3.15 are not universal and depend on the law of $X$. Note that $b$ can be recovered from Theorem 3.14 and thus is universal on the unit sphere in $L^{2}(\Omega)$.

Lemma 3.17. There exists a deterministic constant $C>0$ such that, for all $n \in \mathbb{N}$, $\lambda \in \Lambda_{n}, j \geqslant n, \ell \in\{0,1,2\}$ and $\left(k^{(1)}, K^{(1)}, k^{(2)}, K^{(2)}\right) \in S_{j}^{\ell}(\lambda)$, we have

$$
\left\|{ }_{j}^{\ell} \sum_{k^{(1)}, K^{(1)}}^{k^{(2)}, K^{(2)}}\right\|_{L^{2}(\Omega)} \leqslant C 2^{\frac{-j-n}{2}} .
$$

Proof. Following an idea from [6, Lemma 2.21], we write

$$
\left\|{ }_{j}^{\ell} \sum_{k^{(1)}, K^{(1)}}^{k^{(2)}, K^{(2)}}\right\|_{L^{2}(\Omega)} \leqslant \sum_{\mathcal{R} \in\{<,>,=\}}\| \|_{j}^{\ell} \sum_{\mathcal{R}}^{k^{(2)}, K^{(2)}, K^{(1)}} \|_{L^{2}(\Omega)}
$$

where ${ }_{j}^{\ell} \sum_{k^{(1)}, K^{(1)} \mathcal{R}}^{k^{(2)}, K^{(2)}}$ is the subsum of (3.24), (3.23) or (3.25) in which $k_{1} \mathcal{R} k_{2}$. By doing so, we make sure that two random variables $\varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}}$ and $\varepsilon_{j_{1}^{\prime}, j_{2}^{\prime}}^{k_{1}^{\prime}, k_{2}^{\prime}}$ appearing in this subsum are uncorrelated except when $\left(k_{1}, k_{2}\right)=\left(k_{1}^{\prime}, k_{2}^{\prime}\right)$. Then from Lemma 2.1, we have for $\ell=0$ (the argument being the same for $\ell=1$ or $\ell=2$ ), for all $\mathcal{R} \in\{<,>,=\}$,

$$
\begin{aligned}
\left\|{ }_{j}^{\ell} \sum_{\mathcal{R}}^{k^{(2)}, K^{(2)}, K^{(1)}}\right\|_{L^{2}(\Omega)}^{2} & =\sum_{k^{(1)} \leqslant k_{1} \leqslant K^{(1)}} \sum_{k^{(2)} \leqslant k_{2} \leqslant K^{(2)}, k_{1} \mathcal{R} k_{2}} \mathbb{E}\left[\left(\varepsilon_{j, j}^{k_{1}, k_{2}}\right)^{2}\right]\left(I_{j, j}^{k_{1}, k_{2}}\right)^{2} \\
& \leqslant \sum_{k^{(1)} \leqslant k_{1} \leqslant K^{(1)}} \sum_{k_{2} \in \mathbb{Z}} \frac{2^{-2 j}}{\left(3+\left|k_{1}-k_{2}\right|\right)^{8}} .
\end{aligned}
$$

Since $\#\left\{k_{1} \in \mathbb{Z}: k^{(1)} \leqslant k_{1} \leqslant K^{(1)}\right\} \leqslant 2^{j-n}$, we conclude that

$$
\begin{equation*}
\left\|{ }_{j}^{\ell} \sum_{k^{(1)}, K^{(1)}}^{k^{(2)}, K^{(2)}}\right\|_{L^{2}(\Omega)} \leqslant C 2^{\frac{-j-n}{2}} . \tag{3.28}
\end{equation*}
$$

Lemma 3.18. There exist an event $\widetilde{\Omega}$ of probability 1 and a positive random variable $C_{2}$ with finite moment of any order such that, on $\widetilde{\Omega}$

$$
\begin{equation*}
\forall n \in \mathbb{N}, \forall \lambda \subseteq[0,1], \lambda \in \Lambda_{n}, \forall j \geqslant n, \Xi_{j}(\lambda) \leqslant C_{2}(j-n+1) n . \tag{3.29}
\end{equation*}
$$

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Proof. Let us take $\theta>0$ and consider, for all $n \in \mathbb{N}$ the event

$$
A_{n}:=\left\{\forall \lambda \subseteq[0,1], \lambda \in \Lambda_{n}, \forall j \geqslant n, \Xi_{j}(\lambda) \leqslant \theta(j-n+1) n\right\} .
$$

If $A_{n}^{c}$ stands for the complementary set of $A_{n}$ in $\Omega$, we have, of course,

$$
\mathbb{P}\left(A_{n}^{c}\right)=\mathbb{P}\left(\exists \lambda \subseteq[0,1], \lambda \in \Lambda_{n}: \exists j \geqslant n \text { s. t. } \Xi_{j}(\lambda) \geqslant \theta(j-n+1) n\right) .
$$

But, for all $\lambda \subseteq[0,1], \lambda \in \Lambda_{n}, j \geqslant n, \ell \in\{0,1,2\}$ and $\left(k^{(1)}, K^{(1)}, k^{(2)}, K^{(2)}\right) \in S_{j}^{\ell}(\lambda)$ we have, by Theorem 3.14, if $\theta \geqslant 2$,

As, for all $j \geqslant n, \# S_{j}^{\ell}(\lambda) \leqslant 2^{4(j-n)}$ and $\#\left\{\lambda \subseteq[0,1]: \lambda \in \Lambda_{n}\right\}=2^{n}$, we get

$$
\begin{aligned}
\mathbb{P}\left(A_{n}^{c}\right) & \leqslant C 2^{n} \sum_{j \geqslant n} 2^{4(j-n)} \exp (-\stackrel{\star}{C} \theta(j-n+1) n) \\
& \leqslant C 2^{n} \exp (-\stackrel{\star}{C} \theta n) \sum_{j \geqslant n} 2^{4(j-n)} \exp (-\stackrel{\star}{C} \theta(j-n))
\end{aligned}
$$

for a deterministic constant $C>0$. Therefore, if we take $\theta>4 \log (2) / \stackrel{\star}{C}$, the conclusion follows from the Borel-Cantelli Lemma.

Lemma 3.19. Let $\Omega^{*}$ and $\widetilde{\Omega}$ be the events of probability 1 given by Lemmata 3.6 and 3.18 respectively. There exists a positive random variable $C_{3}$ with finite moment of any order such that, on $\Omega^{*} \cap \widetilde{\Omega}$, for all $t, s \in(0,1)$ the random variable

$$
\begin{equation*}
\left|\sum_{j_{1} \geqslant n} \sum_{j_{2}>n} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}} I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right| \tag{3.30}
\end{equation*}
$$

is bounded from above by $C_{3}|t-s|^{H_{1}+H_{2}-1} \log |t-s|^{-1}$.
Proof. We start by splitting the sums in (3.30) in two parts:

$$
\begin{align*}
& \sum_{j_{1} \geqslant n} \sum_{j_{2} \geqslant j_{1}} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}} I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s] \text { and } \\
& \sum_{j_{2} \geqslant n} \sum_{j_{1}>j_{2}} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}} I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s] . \tag{3.31}
\end{align*}
$$

We only focus on the first sums, as the argument is symmetric in $j_{1}$ and $j_{2}$. As in Lemma 3.12 we write

$$
\begin{align*}
& \sum_{j_{1} \geqslant n} \sum_{j_{2} \geqslant j_{1}} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}} I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s] \\
& =\sum_{j_{1} \geqslant n} \sum_{j_{2} \geqslant j_{1}} \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}^{<}(t, s)} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}} I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]  \tag{3.32}\\
& +\sum_{j_{1} \geqslant n} \sum_{j_{2} \geqslant j_{1}} \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}^{>}(t, s)} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}} I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s] \tag{3.33}
\end{align*}
$$

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$$
\begin{equation*}
+\sum_{j_{1} \geqslant n} \sum_{j_{2} \geqslant j_{1}} \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}[t, s]} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}} I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s] . \tag{3.34}
\end{equation*}
$$

To bound (3.32), we use inequality (3.6) and Lemma 3.10 to get

$$
\begin{aligned}
|(3.32)| & \leqslant C C_{1} \sum_{j_{1} \geqslant n} \sum_{j_{2} \geqslant j_{1}} 2^{j_{1}\left(1-H_{1}\right)} 2^{-j_{2} H_{2}} \sqrt{\log \left(3+\left|j_{1}\right|+2^{j_{1}}\right)} \sqrt{\log \left(3+\left|j_{2}\right|+2^{j_{2}}\right)} \\
& \leqslant C C_{1} \sum_{j_{1} \geqslant n} 2^{j_{1}\left(1-H_{1}-H_{2}\right)} \sqrt{j_{1}} \leqslant C C_{1} 2^{n\left(1-H_{1}-H_{2}\right)} n
\end{aligned}
$$

by applying twice inequality (3.20). The sum (3.33) is bounded in exactly the same way.
To bound (3.34), we use once again the equality (3.21). First we have, by inequality (3.6) and Lemma 3.11,

$$
\begin{align*}
& \sum_{j_{1} \geqslant n} \sum_{j_{2} \geqslant j_{1}} \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}[t, s]} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}} \int_{-\infty}^{s} \psi_{H_{1}}\left(2^{j_{1}} x-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} x-k_{2}\right) d x \mid \\
& \leqslant C C_{1} \sum_{j_{1} \geqslant n} \sum_{j_{2} \geqslant j_{1}} 2^{j_{1}\left(1-H_{1}\right)} 2^{-j_{2} H_{2}} \sqrt{\log \left(3+\left|j_{1}\right|+2^{j_{2}}\right)} \sqrt{\log \left(3+\left|j_{2}\right|+2^{j_{2}}\right)} \\
& \leqslant C C_{1} 2^{n\left(1-H_{1}-H_{2}\right)} n . \tag{3.35}
\end{align*}
$$

We bound

$$
\left|\sum_{j_{1} \geqslant n} \sum_{j_{2} \geqslant j_{1}} \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}[t, s]} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}} \int_{t}^{+\infty} \psi_{H_{1}}\left(2^{j_{1}} x-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} x-k_{2}\right) d x\right|
$$

in the same way.
It only remains to find an estimate for

$$
\sum_{j_{1} \geqslant n} \sum_{j_{2} \geqslant j_{1}} \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}[t, s]} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}} I_{j_{1}, j_{2}}^{k_{1}, k_{2}}
$$

and thus, recalling Lemma 2.1, we reduce the problem to first bound, for $j \geqslant n$ and $\ell \in\{j, j+1\}$, the sums

$$
\begin{array}{r}
\sum_{k_{1} \in \mathbb{Z}_{j}^{\curvearrowright}(t, s)} \sum_{k_{2} \in \mathbb{Z}_{\ell}[t, s]} \varepsilon_{j, \ell}^{k_{1}, k_{2}} I_{j, \ell}^{k_{1}, k_{2}}, \\
\sum_{k_{1} \in \mathbb{Z}_{j}^{>}(t, s)} \sum_{k_{2} \in \mathbb{Z}_{\ell}[t, s]} \varepsilon_{j, \ell}^{k_{1}, k_{2}} I_{j, \ell}^{k_{1}, k_{2}}, \\
\sum_{k_{1} \in \mathbb{Z}_{j}[t, s]} \sum_{k_{2} \in \mathbb{Z}_{\ell}[t, s]} \varepsilon_{j, \ell}^{k_{1}, k_{2}} I_{j, \ell}^{k_{1}, k_{2}} \tag{3.38}
\end{array}
$$

on $\Omega^{*} \cap \widetilde{\Omega}$. Let us consider (3.36) with $\ell=j$, the argument for $\ell=j+1$ and (3.37) being similar. Using again Lemmata 2.1 and 3.6, we have on $\Omega^{*} \cap \widetilde{\Omega}$, since for all $k_{2} \in \mathbb{Z}_{j}[t, s]$, $\left|k_{2}\right| \leqslant 2^{j}$, for $j \geqslant n$,

$$
\begin{aligned}
|(3.36)| & \leqslant C C_{1} \sum_{k_{1} \in \mathbb{Z}_{j}^{<}(t, s)} \sum_{k_{2} \in \mathbb{Z}_{j}[t, s]} \frac{2^{-j}}{\left(3+\left|k_{1}-k_{2}\right|\right)^{4}} \sqrt{\log \left(3+j+\left|k_{1}\right|\right)} \sqrt{\log \left(3+j+\left|k_{2}\right|\right)} \\
& \leqslant C C_{1} \sum_{k_{1} \in \mathbb{Z}_{j}^{<}(t, s)} \sum_{k_{2} \in \mathbb{Z}_{j}[t, s]} \frac{2^{-j} \sqrt{j}}{\left(3+k_{2}-k_{1}\right)^{4}} \sqrt{\log \left(3+j+\left|k_{1}\right|\right)} \\
& \leqslant C C_{1} \sum_{k_{1} \in \mathbb{Z}_{j}^{<}(t, s)} \sum_{m=0}^{+\infty} \frac{2^{-j} \sqrt{j}}{\left(3+2^{j} \min \{s, t\}+m-k_{1}\right)^{4}} \sqrt{\log \left(3+j+\left|k_{1}\right|\right)}
\end{aligned}
$$

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$$
\begin{align*}
& \leqslant C C_{1} \sum_{k_{1} \in \mathbb{Z}_{j}^{<}(t, s)} 2^{-j} \sqrt{j} \int_{0}^{+\infty} \frac{d y}{\left(2+2^{j} \min \{s, t\}+y-k_{1}\right)^{4}} \sqrt{\log \left(3+j+\left|k_{1}\right|\right)} \\
& \leqslant C C_{1} 2^{-j} \sqrt{j} \sum_{k_{1} \in \mathbb{Z}_{j}^{<}(t, s)} \frac{\sqrt{\log \left(3+j+\left|k_{1}\right|\right)}}{\left(2+2^{j} \min \{s, t\}-k_{1}\right)^{3}} \\
& \leqslant C C_{1} 2^{-j} \sqrt{j} \sqrt{\log \left(3+j+2^{j} \min \{s, t\}\right)} \leqslant C C_{1} 2^{-j} j . \tag{3.39}
\end{align*}
$$

It follows that

$$
\left|\sum_{j \geqslant n} 2^{j\left(2-H_{1}-H_{2}\right)} \sum_{\ell=j}^{j+1} \sum_{k_{1} \in \mathbb{Z}_{j}^{<}(t, s)} \sum_{k_{2} \in \mathbb{Z}_{\ell}[t, s]} \varepsilon_{j, \ell}^{k_{1}, k_{2}} I_{j, \ell}^{k_{1}, k_{2}}\right| \leqslant C C_{1} 2^{n\left(1-H_{1}-H_{2}\right)} n
$$

and, similarly,

$$
\left|\sum_{j \geqslant n} 2^{j\left(2-H_{1}-H_{2}\right)} \sum_{\ell=j}^{j+1} \sum_{k_{1} \in \mathbb{Z}_{j}^{\supset}(t, s)} \sum_{k_{2} \in \mathbb{Z}_{\ell}[t, s]} \varepsilon_{j, \ell}^{k_{1}, k_{2}} I_{j, \ell}^{k_{1}, k_{2}}\right| \leqslant C C_{1} 2^{n\left(1-H_{1}-H_{2}\right)} n .
$$

The bound for (3.38) is obtained using (3.29) and (3.28) which lead to

$$
\begin{aligned}
& \left|\sum_{j \geqslant n} 2^{j\left(2-H_{1}-H_{2}\right)} \sum_{\ell=j}^{j+1} \sum_{k_{1} \in \mathbb{Z}_{j}[t, s]} \sum_{k_{2} \in \mathbb{Z}_{\ell}[t, s]} \varepsilon_{j, \ell}^{k_{1}, k_{2}} I_{j, \ell}^{k_{1}, k_{2}}\right| \\
& \leqslant C C_{2} \sum_{j \geqslant n} 2^{j\left(\frac{3}{2}-H_{1}-H_{2}\right)} 2^{-\frac{n}{2}}(j-n+1) n \\
& \leqslant C C_{2} 2^{n\left(\frac{3}{2}-H_{1}-H_{2}\right)} 2^{-\frac{n}{2}} n \\
& =C C_{2} 2^{n\left(1-H_{1}-H_{2}\right)} n
\end{aligned}
$$

as $\frac{3}{2}<H_{1}+H_{2}$.
Putting all of these together we get that (3.30) is bounded from above by

$$
C \max \left\{C_{1}, C_{2}\right\}|t-s|^{H_{1}+H_{2}-1} \log |t-s|^{-1}
$$

on $\Omega^{*} \cap \widetilde{\Omega}$.
We now prove the main result of this subsection.
Proof of Proposition 3.5. Let us consider $\omega$ in the event $\Omega^{*} \cap \widetilde{\Omega}$ of probability 1, where $\Omega^{*}$ and $\widetilde{\Omega}$ are given by Lemmata 3.6 and 3.18 respectively.

If $t, s \in(0,1)$, we write

$$
\begin{aligned}
\mid R_{H_{1}, H_{2}}( & t, \omega)-R_{H_{1}, H_{2}}(s, \omega) \mid \\
& \leqslant\left|\sum_{j_{1}<n} \sum_{j_{2}<n} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}}(\omega) I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right| \\
& +\left|\sum_{j_{1}<n} \sum_{j_{2} \geqslant n} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}}(\omega) I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right| \\
& +\left|\sum_{j_{1} \geqslant n} \sum_{j_{2}<n} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}}(\omega) I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right| \\
& +\left|\sum_{j_{1} \geqslant n} \sum_{j_{2} \geqslant n} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}}(\omega) I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right| .
\end{aligned}
$$

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The first sum is bounded from above by Lemmata 3.6 and 3.7, the second and the third one are bounded from above by Lemmata 3.6 and 3.12 and the last one is bounded from above by Lemma 3.19.

Remark 3.20. Starting from now and until the end of this section, one can focus on the process

$$
\left\{R_{H_{1}, H_{2}}^{\prime}(t)=\sum_{j_{1}=0}^{+\infty} \sum_{j_{2}=0}^{+\infty} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}} I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[0, \cdot]\right\}
$$

because almost surely, it is the most irregular part of $R_{H_{1}, H_{2}}$. Indeed, using different estimates obtained in this subsection, one can see that, almost surely, there exists a constant $C>0$ such that, for all $s, t \in(0,1)$,

$$
\begin{aligned}
& \left|\sum_{j_{1}<0} \sum_{j_{2}<0} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}} I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right| \leqslant C|t-s|, \\
& \left|\sum_{j_{1}<0} \sum_{j_{2}>0} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}} I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right| \leqslant C|t-s|^{-H_{2}} \log |t-s|^{-1} \\
& \left|\sum_{j_{1}>0} \sum_{j_{2}<0} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}} I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right| \leqslant C|t-s|^{-H_{1}} \log |t-s|^{-1}
\end{aligned}
$$

and we conclude because $H_{1}+H_{2}-1<\min \left\{H_{1}, H_{2}\right\}<1$.

### 3.2 Ordinary points

Let us now focus on the almost sure finiteness of the limit (1.7) for almost every point. The main idea behind our method is that wavelets which contribute the most in $\left|R_{H_{1}, H_{2}}(t, \cdot)-R_{H_{1}, H_{2}}(s, \cdot)\right|$ are the ones with associated dyadic intervals "close" to the interval $[t, s]$. Thus, we aim at proving the following Proposition.
Proposition 3.21. There exists an event $\Omega_{\text {ord }}$ of probability 1 such that for all $\omega \in \Omega_{\text {ord }}$, for almost every $t \in(0,1)$,

$$
\limsup _{s \rightarrow t} \frac{\left|R_{H_{1}, H_{2}}(t, \omega)-R_{H_{1}, H_{2}}(s, \omega)\right|}{|t-s|^{H_{1}+H_{2}-1} \log \log |t-s|^{-1}}<+\infty .
$$

As in [18], for all $j \in \mathbb{N}$, we denote by $k_{j}(t)$ the unique integer such that $t \in$ $\left[k_{j}(t) 2^{-j},\left(k_{j}(t)+1\right) 2^{-j}\right)$. In other words, $k_{j}(t)=s\left(\lambda_{j}(t)\right)$. If $t \in(0,1)$ is fixed, applying Lemma 3.6 to the sequence of random variables $\left(\xi_{j_{1}, j_{2}}^{k_{1}^{\prime}, k_{2}^{\prime}}\right)_{\left(j_{1}, j_{2}, k_{1}^{\prime}, k_{2}^{\prime}\right) \in \mathbb{Z}^{4}}$ defined by

$$
\xi_{j_{1}, j_{2}}^{k_{1}^{\prime}, k_{2}^{\prime}}=\varepsilon_{j_{1}, j_{2}}^{k_{1}^{\prime}+k_{j_{1}(t)}, k_{2}^{\prime}+k_{j_{2}(t)}}
$$

we deduce the existence of $\Omega_{t}^{*}$, an event of probability 1 , and $C_{t, 1}$, a positive random variable with finite moment of any order, such that, for all $\omega \in \Omega_{t}^{*}$ and for each $\left(j_{1}, j_{2}, k_{1}, k_{2}\right) \in \mathbb{Z}^{4}$, one has

$$
\begin{equation*}
\left|\varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}}(\omega)\right| \leqslant C_{t, 1}(\omega) \sqrt{\log \left(3+\left|j_{1}\right|+\left|k_{1}-k_{j_{1}}(t)\right|\right)} \sqrt{\log \left(3+\left|j_{2}\right|+\left|k_{2}-k_{j_{2}}(t)\right|\right)} . \tag{3.40}
\end{equation*}
$$

In view of this fact, let us set, for $t \in(0,1)$ and $\left(j_{1}, j_{2}, k_{1}, k_{2}\right) \in \mathbb{N}^{2} \times \mathbb{Z}^{2}$

$$
L_{j_{1}, j_{2}}^{k_{1}, k_{2}}(t)=\sqrt{\log \left(3+j_{1}+\left|k_{1}-k_{j_{1}}(t)\right|\right)} \sqrt{\log \left(3+j_{2}+\left|k_{2}-k_{j_{2}}(t)\right|\right)}
$$

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In what follows, we show how to modify Lemmata 3.7 to 3.19 from the previous subsection, using $L_{j_{1}, j_{2}}^{k_{1}, k_{2}}(t)$ instead of $L_{j_{1}, j_{2}}^{k_{1}, k_{2}}$. Before all, we need the following Lemma which is inspired by results from [18] that can be extended in our case.
Lemma 3.22. For all $L>2$ there exists a constant $C_{L}>0$ such that, for all $n \in \mathbb{N}$ and $t, s \in(0,1)$ such that $2^{-n-1}<|t-s| \leqslant 2^{-n}$, for all $x \in[s, t]$

1. For all $0 \leqslant j<n$

$$
\sum_{k \in \mathbb{Z}} \frac{\sqrt{\log \left(3+j+\left|k-k_{j}(t)\right|\right.}}{\left(3+\left|2^{j} x-k\right|\right)^{L}} \leqslant C_{L} \sqrt{\log (3+j)}
$$

2. For all $j \geqslant n$

$$
\sum_{k \in \mathbb{Z}} \frac{\sqrt{\log \left(3+j+\left|k-k_{j}(t)\right|\right.}}{\left(3+\left|2^{j} x-k\right|\right)^{L}} \leqslant C_{L} \sqrt{j-n+1} \sqrt{\log (3+j)}
$$

Proof. For all $j \in \mathbb{N}, k \in \mathbb{Z}$ and $x \in[s, t]$, observe that

$$
\begin{equation*}
\left|k-k_{j}(t)\right| \leqslant\left|k-2^{j} x\right|+\left|2^{j} x-2^{j} t\right|+\left|2^{j} t-k_{j}(t)\right| \leqslant\left|k-2^{j} x\right|+2^{j-n}+1 \tag{3.41}
\end{equation*}
$$

If $0 \leqslant j<n$, then it follows from (3.41) that $\left|k-k_{j}(t)\right| \leqslant\left|2^{j} x-k\right|+2$ which allow us to write, thanks to inequality (3.11),

$$
\begin{aligned}
\frac{\sqrt{\log \left(3+j+\left|k-k_{j}(t)\right|\right.}}{\left(3+\left|2^{j} x-k\right|\right)} & \leqslant \sqrt{\log (3+j)} \frac{\sqrt{\log \left(5+\left|2^{j} x-k\right|\right)}}{\left|2^{j} x-k\right|+3} \\
& \leqslant C \sqrt{\log (3+j)} .
\end{aligned}
$$

where $C:=\sup _{x \geqslant 0}\left(\frac{\sqrt{\log (5+x)}}{x+3}\right)$ and we conclude using the boundedness of the function

$$
\begin{equation*}
\xi \mapsto \sum_{k \in \mathbb{Z}} \frac{1}{(1+|\xi-k|)^{M}} \tag{3.42}
\end{equation*}
$$

for all $M>1$.
Now, if $j \geqslant n$, from (3.41) we get $\left|k-k_{j}(t)\right| \leqslant\left|2^{j} x-k\right|+2^{j-n+1}$ and thus, again by inequality (3.11),

$$
\begin{aligned}
\frac{\sqrt{\log \left(3+j+\left|k-k_{j}(t)\right|\right.}}{\left(3+\left|2^{j} x-k\right|\right)} & \leqslant \sqrt{\log \left(3+2^{j-n+1}\right)} \sqrt{\log (3+j)} \frac{\sqrt{\log \left(3+\left|2^{j} x-k\right|\right)}}{\left|2^{j} x-k\right|+3} \\
& \leqslant C^{\prime} \sqrt{j-n+1} \sqrt{\log (3+j)} .
\end{aligned}
$$

where $C^{\prime}:=\sqrt{3} \sup _{x \geqslant 0}\left(\frac{\sqrt{\log (3+x)}}{x+3}\right)$ and the conclusion comes again from the boundedness of the function in (3.42) for all $M>1$

Lemma 3.23. There exists a deterministic constant $C>0$ such that, for all $t, s \in(0,1)$ we have

$$
\begin{aligned}
& \sum_{0 \leqslant j_{1}<n} \sum_{0 \leqslant j_{2}<n} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} L_{j_{1}, j_{2}}^{k_{1}, k_{2}}(t)\left|I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right| \\
& \leqslant C|t-s|^{H_{1}+H_{2}-1} \log \log |t-s|^{-1}
\end{aligned}
$$

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Proof. If $\xi \in[s, t]$, we get from the fast decay of the fractional antiderivatives of $\psi$ (3.7) and inequality (3.40), for $0 \leqslant j_{1}, j_{2}<n$,

$$
\begin{aligned}
& \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} L_{j_{1}, j_{2}}^{k_{1}, k_{2}}(t)\left|\psi_{H_{1}}\left(2^{j_{1}} \xi-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} \xi-k_{2}\right)\right| \\
\leqslant & C C_{1}\left(\sum_{k_{1} \in \mathbb{Z}} \frac{\sqrt{\log \left(3+j_{1}+\left|k_{1}-k_{j_{1}}(t)\right|\right)}}{\left(3+\left|2^{j_{1}} \xi-k_{1}\right|\right)^{4}}\right)\left(\sum_{k_{2} \in \mathbb{Z}} \frac{\sqrt{\log \left(3+j_{2}+\left|k_{2}-k_{j_{2}}(t)\right|\right)}}{\left(3+\left|2^{j_{2}} \xi-k_{2}\right|\right)^{4}}\right) .
\end{aligned}
$$

These last two sums are bounded by the first point of Lemma 3.22. Using

$$
\begin{equation*}
\sum_{j_{1}=0}^{n-1} 2^{j_{1}\left(1-H_{1}\right)} \sqrt{\log \left(3+j_{1}\right)} \leqslant C 2^{n\left(1-H_{1}\right)} \sqrt{\log (n)} \tag{3.43}
\end{equation*}
$$

instead of (3.10), we conclude, just as in Lemma 3.7, that the desired inequality holds.
Lemma 3.24. There exists a deterministic constant $C>0$ such that, for all $t, s \in(0,1)$ and $0 \leqslant j_{1}<n \leqslant j_{2}$, the quantities

$$
\begin{align*}
& \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}^{<}(t, s)} L_{j_{1}, j_{2}}^{k_{1}, k_{2}}(t)\left|I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right|  \tag{3.44}\\
& \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}(t, s)} L_{j_{1}, j_{2}}^{k_{1}, k_{2}}(t)\left|I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right| \tag{3.45}
\end{align*}
$$

are bounded from above by $C \sqrt{j_{2}-n+1} \sqrt{\log \left(3+j_{1}\right)} \sqrt{\log \left(3+j_{2}\right)} 2^{-j_{2}}$.
Proof. Let us prove the bound for (3.44), the argument for (3.45) being similar. We have, by the first part of Lemma 3.22, for $0 \leqslant j_{1}<n \leqslant j_{2}$,

$$
(3.44) \leqslant C \sqrt{\log \left(3+j_{1}\right)} \int_{[s, t]} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}^{<}(t, s)} \frac{\sqrt{\log \left(3+j_{2}+\left|k_{2}-k_{j_{2}}(t)\right|\right)}}{\left(3+\left|2^{j_{2}} x-k_{2}\right|\right)^{4}} d x
$$

and, as for all $k_{2} \in \mathbb{Z}_{j_{2}}^{<}(t, s)$ and $x \in[s, t]$ we have

$$
\left|k_{2}-k_{j_{2}}(t)\right| \leqslant\left|2^{j_{2}} x-k_{2}\right|+\left|k_{j_{2}}(t)-2^{j_{2}} x\right| \leqslant\left|2^{j_{2}} x-k_{2}\right|+2^{j_{2}-n}+1
$$

and, by inequality (3.11),

$$
\sqrt{\log \left(3+j_{2}+\left|k_{2}-k_{j_{2}}(t)\right|\right)} \leqslant C \sqrt{j_{2}-n+1} \sqrt{\log \left(3+j_{2}\right)} \sqrt{\log \left(3+\left|2^{j_{2}} x-k_{2}\right|\right)}
$$

it just remains us to use the bound (3.14) to write

$$
\begin{equation*}
(3.44) \leqslant C \sqrt{j_{2}-n+1} \sqrt{\log \left(3+j_{1}\right)} \sqrt{\log \left(3+j_{2}\right)} 2^{-j_{2}} \tag{3.46}
\end{equation*}
$$

Lemma 3.25. There exists a deterministic constant $C>0$ such that, for all $t, s \in(0,1)$ and $0 \leqslant j_{1}<n \leqslant j_{2}$, the quantities

$$
\begin{align*}
& \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}[t, s]} L_{j_{1}, j_{2}}^{k_{1}, k_{2}}(t)\left|\int_{-\infty}^{\min \{s, t\}} \psi_{H_{1}}\left(2^{j_{1}} x-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} x-k_{2}\right) d x\right|  \tag{3.47}\\
& \quad \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}[t, s]} L_{j_{1}, j_{2}}^{k_{1}, k_{2}}(t)\left|\int_{\max \{s, t\}}^{+\infty} \psi_{H_{1}}\left(2^{j_{1}} x-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} x-k_{2}\right) d x\right| \tag{3.48}
\end{align*}
$$

are bounded from above by $C \sqrt{\log \left(3+j_{1}\right)} \sqrt{\log \left(3+j_{2}\right)} \sqrt{j_{2}-n+1} 2^{-j_{2}}$.

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Proof. Again we assume $s \leqslant t$. First, using the fast decay of the fractional antiderivatives of $\psi$ (3.7), (3.47) is bounded from above by

$$
\begin{equation*}
\int_{-\infty}^{s} \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}[t, s]} \frac{\sqrt{\log \left(3+j_{1}+\left|k_{1}-k_{j_{1}}(t)\right|\right)}}{\left(3+\left|2^{j_{1}} x-k_{1}\right|\right)^{4}} \frac{\sqrt{\log \left(3+j_{2}+\left|k_{2}-k_{j_{2}}(t)\right|\right)}}{\left(3+\left|2^{j_{2}} x-k_{2}\right|\right)^{4}} d x . \tag{3.49}
\end{equation*}
$$

Observe that, for all $k_{1} \in \mathbb{Z}, k_{2} \in \mathbb{Z}_{j_{2}}[t, s]$ and $x \in(-\infty, s]$, we have, as $j_{1}<n \leqslant j_{2}$,

$$
\begin{aligned}
\left|2^{j_{1}} x-k_{j_{1}}(t)\right| & \leqslant\left|2^{j_{1}} x-2^{j_{1}-j_{2}} k_{2}\right|+\left|2^{j_{1}-j_{2}} k_{2}-2^{j_{1}} t\right|+\left|2^{j_{1}} t-k_{j_{1}}(t)\right| \\
& \leqslant\left|2^{j_{2}} x-k_{2}\right|+2
\end{aligned}
$$

and therefore $\left|k_{1}-k_{j_{1}}(t)\right| \leqslant\left|2^{j_{1}} x-k_{1}\right|+\left|2^{j_{2}} x-k_{2}\right|+2$ while $\left|k_{2}-k_{j_{2}}(t)\right| \leqslant\left|k_{2}-2^{j_{2}} t\right|+$ $\left|2^{j_{2}} t-k_{j_{2}}(t)\right| \leqslant 2^{j_{2}-n}+1$. It allows to write, thanks to inequality (3.11), the boundedness of the function (3.42) and inequality (3.15)

$$
\begin{aligned}
|(3.49)| & \leqslant C \sqrt{\log \left(3+j_{1}\right)} \sqrt{\log \left(3+j_{2}\right)} \sqrt{j_{2}-n+1} \int_{-\infty}^{s} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}[t, s]} \frac{d x}{\left(3+\left|2^{j_{2}} x-k_{2}\right|\right)^{3}} \\
& \leqslant C \sqrt{\log \left(3+j_{1}\right)} \sqrt{\log \left(3+j_{2}\right)} \sqrt{j_{2}-n+1} 2^{-j_{2}} .
\end{aligned}
$$

We bound the second sums in the same way.
Lemma 3.26. There exists a deterministic constant $C>0$ such that, for all $t, s \in(0,1)$, the quantities

$$
\begin{aligned}
& \sum_{0 \leqslant j_{1}<n} \sum_{j_{2} \geqslant n} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} L_{j_{1}, j_{2}}^{k_{1}, k_{2}}(t)\left|I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right| \\
& \sum_{j_{1} \geqslant n} \sum_{0 \leqslant j_{2}<n} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} L_{j_{1}, j_{2}}^{k_{1}, k_{2}}(t)\left|I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right|
\end{aligned}
$$

are bounded from above by $C|t-s|^{H_{1}+H_{2}-1} \log \log |t-s|^{-1}$.
Proof. The proof is exactly the same as the one of Lemma 3.12 except that we use Lemmata 3.24 and 3.25 instead of Lemmata 3.10 and 3.11 respectively and that we conclude using again (3.43) instead of (3.10) and

$$
\begin{equation*}
\sum_{j_{2}=n}^{+\infty} 2^{-j_{2} H_{2}} \sqrt{j_{2}-n+1} \sqrt{\log \left(3+j_{2}\right)} \leqslant C^{\prime} 2^{-n H_{2}} \sqrt{\log (n)} \tag{3.50}
\end{equation*}
$$

instead of (3.20).
Lemma 3.27. There exists a deterministic constant $C>0$ such that, for all $t, s \in$ $(0,1)$ and $n \leqslant j_{1} \leqslant j_{2}$, the quantities (3.44) and (3.45) are bounded from above by $C \sqrt{j_{2}-n+1} \sqrt{j_{1}-n+1} \sqrt{\log \left(3+j_{1}\right)} \sqrt{\log \left(3+j_{2}\right)} 2^{-j_{2}}$.

Proof. The proof is exactly the same as for Lemma 3.24 except that, here, we use the second part of Lemma 3.22 instead of the first one.

Lemma 3.28. There exists a deterministic constant $C>0$ such that, for all $t, s \in$ $(0,1)$ and $n \leqslant j_{1} \leqslant j_{2}$ the quantities (3.47) and (3.48) are bounded from above by $C \sqrt{j_{1}-n+1} \sqrt{j_{2}-n+1} \sqrt{\log \left(3+j_{1}\right)} \sqrt{\log \left(3+j_{2}\right)} 2^{-j_{2}}$.

Proof. The proof is exactly the same as for Lemma 3.25 except that, here, we use the second part of Lemma 3.22 instead of the first one.

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Just as we did for the rapid points, it remains us to bound the random variables $\Xi_{j}(\lambda)$. Here, we don't want anymore to show the existence of an uniform modulus but only a pointwise modulus of continuity at a fixed point of interest $t$. Therefore, we just have to bound, for all $n \in \mathbb{N}$ the random variables $\Xi_{j}(\lambda)$ for $j \geqslant n$ and $\lambda \in 3 \lambda_{n}(t)$. We thus have the following result.
Lemma 3.29. For all $t \in(0,1)$, there exist an event $\widetilde{\Omega_{t}}$ of probability 1 and a positive random variable $C_{t, 2}$ with finite moment of any order such that, on $\widetilde{\Omega_{t}}$,

$$
\begin{equation*}
\forall n \in \mathbb{N}, \forall \lambda \in 3 \lambda_{n}(t), \forall j \geqslant n, \Xi_{j}(\lambda) \leqslant C_{t, 2}(j-n+1) \log (n) \tag{3.51}
\end{equation*}
$$

Proof. If $t \in(0,1)$ is fixed and $\theta>0$, let us define the event

$$
A_{n}(t)=\left\{\forall \lambda \in 3 \lambda_{n}(t) \forall j \geqslant n, \Xi_{j}(\lambda) \leqslant \theta(j-n+1) \log (n)\right\}
$$

Similarly to Lemma 3.18, we get

$$
\begin{aligned}
\mathbb{P}\left(A_{n}(t)^{c}\right) & \leqslant C \sum_{j \geqslant n} 2^{4(j-n)} \exp (-\stackrel{\star}{C} \theta(j-n+1) \log (n)) \\
& \leqslant C \exp (-\stackrel{\star}{C} \theta \log (n)) \sum_{j \geqslant n} 2^{4(j-n)} \exp (-\stackrel{\star}{C} \theta(j-n))
\end{aligned}
$$

for a determistic constant $C>0$. Therefore, if we take again $\theta>4 \log (2) / \stackrel{\star}{C}$ then Borel-Cantelli Lemma implies the existence of an event $\widetilde{\Omega_{t}}$ of probability 1 and $C_{t, 2}$ a positive random variable of finite moment of any order such that, on $\widetilde{\Omega_{t}}$, assertion (3.51) holds.

Lemma 3.30. If $t \in(0,1)$, let $\Omega_{t}^{*}$ be the event of probability 1 where inequality (3.40) holds and $\widetilde{\Omega_{t}}$ be the event of probability 1 given by Lemma 3.29. There exists a positive random variable $C_{t, 3}$ with finite moment of any order such that, on $\Omega_{t}^{*} \cap \widetilde{\Omega_{t}}$, for all $s \in(0,1)$ the random variable

$$
\begin{equation*}
\left|\sum_{j_{1} \geqslant n} \sum_{j_{2}>n} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}} I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right| \tag{3.52}
\end{equation*}
$$

is bounded from above by $C_{t, 3}|t-s|^{H_{1}+H_{2}-1} \log |t-s|^{-1}$.
Proof. Again, we use the split (3.31) and we only do the details for the first sum. We deal with the series (3.32) and (3.33) in the same way that in Lemma 3.19 but using inequality (3.40) and Lemmata 3.27 and 3.28 and finally inequality (3.50).

For (3.34), first, by Lemma 3.28 and inequality (3.40), we have, on $\Omega_{t}^{*} \cap \widetilde{\Omega_{t}}$

$$
\begin{align*}
& \left|\sum_{j_{1} \geqslant n} \sum_{j_{2} \geqslant j_{1}} \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}[t, s]} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}} \int_{-\infty}^{s} \psi_{H_{1}}\left(2^{j_{1}} x-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} x-k_{2}\right) d x\right| \\
& \leqslant C C_{t, 1} \sum_{j_{1} \geqslant n} \sum_{j_{2} \geqslant j_{1}} 2^{j_{1}\left(1-H_{1}\right)} 2^{-j_{2} H_{2}} \sqrt{j_{1}-n+1} \sqrt{j_{2}-n+1} \sqrt{\log \left(3+j_{1}\right)} \sqrt{\log \left(3+j_{2}\right)} \\
& \leqslant C C_{1} 2^{n\left(1-H_{1}-H_{2}\right)} \log (n) . \tag{3.53}
\end{align*}
$$

We bound

$$
\left|\sum_{j_{1} \geqslant n} \sum_{j_{2} \geqslant j_{1}} \sum_{k_{1} \in \mathbb{Z}} \sum_{k_{2} \in \mathbb{Z}_{j_{2}}[t, s]} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}} \int_{-\infty}^{s} \psi_{H_{1}}\left(2^{j_{1}} x-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} x-k_{2}\right) d x\right|
$$

on $\Omega_{t}^{*} \cap \widetilde{\Omega_{t}}$ exactly in the same way.
To finish the proof, again, we have to bound (3.36), (3.37) and (3.38) for $\ell \in\{j, j+1\}$ (with $j \geqslant n$ ) on $\Omega_{t}^{*} \cap \widetilde{\Omega_{t}}$. For (3.36), in the case $\ell=j$, one can note that, for all $k_{2} \in \mathbb{Z}_{j}[t, s]$, $\left|k_{2}-k_{j}(t)\right| \leqslant 2^{j-n}+1$ and, for all $k_{1} \in \mathbb{Z}_{j}^{<}(t, s),\left|k_{1}-k_{j}(t)\right| \leqslant\left|2^{j} \min \{t, s\}-k_{1}\right|+2^{j-n}+1$. Using the same tricks as in (3.39), we get, on $\Omega_{t}^{*} \cap \widetilde{\Omega_{t}}$

$$
\begin{aligned}
|(3.36)| & \leqslant C C_{t, 1}(j-n+1) \log (3+j) 2^{-j} \sum_{k_{1} \in \mathbb{Z}_{j}^{<}(t, s)} \frac{1}{\left(2+\left|2^{j} \min \{s, t\}-k_{1}\right|\right)^{3}} \\
& \leqslant C(j-n+1) \log (3+j) 2^{-j} .
\end{aligned}
$$

The bounds for (3.37) and in the case $\ell=j+1$ are obtained in the same way. Finally to bound (3.38), we use (3.51) and (3.28) and get on $\Omega_{t}^{*} \cap \widetilde{\Omega_{t}}$

$$
\begin{aligned}
& \left|\sum_{j \geqslant n} 2^{j\left(2-H_{1}-H_{2}\right)} \sum_{\ell=j}^{j+1} \sum_{k_{1} \in \mathbb{Z}_{j}[t, s]} \sum_{k_{2} \in \mathbb{Z}_{\ell}[t, s]} \varepsilon_{j, \ell}^{k_{1}, k_{2}} I_{j, \ell}^{k_{1}, k_{2}}\right| \\
& \leqslant C C_{t, 2} \sum_{j \geqslant n} 2^{j\left(\frac{3}{2}-H_{1}-H_{2}\right)} 2^{-\frac{n}{2}}(j-n+1) \log (n) \\
& \leqslant C C_{t, 2} 2^{n\left(1-H_{1}-H_{2}\right)} \log (n) .
\end{aligned}
$$

We conclude that (3.52) is bounded from above by $C \max \left\{C_{t, 1}, C_{t, 2}\right\}|t-s|^{H_{1}+H_{2}-1} \log \mid t-$ $\left.s\right|^{-1}$ on $\Omega_{t}^{*} \cap \widetilde{\Omega_{t}}$.

We can now prove Proposition 3.21.
Proof of Proposition 3.21. Let us fix $t \in(0,1)$ and consider $\omega \in \Omega_{t}^{*} \cap \widetilde{\Omega_{t}}$. For all $s \in(0,1)$, we write ${ }^{1}$

$$
\begin{aligned}
\mid R_{H_{1}, H_{2}}^{\prime}( & (t, \omega)-R_{H_{1}, H_{2}}^{\prime}(s, \omega) \mid \\
& \leqslant\left|\sum_{0 \leqslant j_{1}<n} \sum_{0 \leqslant j_{2}<n} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}}(\omega) I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right| \\
& +\left|\sum_{0 \leqslant j_{1}<n} \sum_{j_{2} \geqslant n} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}}(\omega) I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right| \\
& +\left|\sum_{j_{1} \geqslant n} \sum_{0 \leqslant j_{2}<n} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}}(\omega) I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right| \\
& +\left|\sum_{j_{1} \geqslant n} \sum_{j_{2} \geqslant n} \sum_{\left(k_{1}, k_{2}\right) \in \mathbb{Z}^{2}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}}(\omega) I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]\right| .
\end{aligned}
$$

We bound from above the first sum by inequality (3.40) and Lemma 3.23, the second and the third sums by inequality (3.40) and Lemma 3.26 and the last sum by Lemma 3.30.

Using inequalities (3.50) and (3.53) and Remark 3.20, one can finally write that for all $t \in(0,1)$, for all $\omega$ in the event of probability $1 \Omega_{t}^{*} \cap \widetilde{\Omega_{t}}$

$$
\limsup _{s \rightarrow t} \frac{\left|R_{H_{1}, H_{2}}(t, \omega)-R_{H_{1}, H_{2}}(s, \omega)\right|}{|t-s|^{H_{1}+H_{2}-1} \log \log |t-s|^{-1}}<+\infty
$$

and we conclude by Fubini Theorem.

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### 3.3 Slow points

In this section, we aim at showing that the generalized Rosenblatt process admits slow points: we prove the following Proposition.
Proposition 3.31. There exists an event $\Omega_{\text {slo }}$ of probability 1 such that for all $\omega \in \Omega_{\text {slo }}$ there exist $t \in(0,1)$ such that

$$
\begin{equation*}
\limsup _{s \rightarrow t} \frac{\left|R_{H_{1}, H_{2}}(t, \omega)-R_{H_{1}, H_{2}}(s, \omega)\right|}{|t-s|^{H_{1}+H_{2}-1}}<+\infty . \tag{3.54}
\end{equation*}
$$

In [25], Kahane described a procedure to insure the existence of slow points for the Brownian motion. This procedure was then generalized in [18] to fit for any arbitrary fractional Brownian motion. It consists in showing that for any $m>0$, almost surely, there exist $\mu>0$ and $t \in(0,1)$ such that, if one sets

$$
\begin{equation*}
\Lambda_{j}^{0}(t)=\left\{\lambda \in \Lambda_{j}:|s(\lambda(t))-s(\lambda)| \leqslant 1\right\} \tag{3.55}
\end{equation*}
$$

and, for all $1 \leqslant l$

$$
\begin{equation*}
\Lambda_{j}^{l}(t)=\left\{\lambda \in \Lambda_{j},: 2^{m(l-1)}<|s(\lambda(t))-s(\lambda)| \leqslant 2^{m l}\right\} \tag{3.56}
\end{equation*}
$$

then, for all $\lambda \in \Lambda_{j}^{l}(t)$ we have

$$
\begin{equation*}
\left|\varepsilon_{\lambda}\right| \leqslant 2^{l} \mu, \tag{3.57}
\end{equation*}
$$

where $\varepsilon_{\lambda}$ is the random variable $2^{\frac{j}{2}} \int_{\mathbb{R}} \psi_{\lambda}(x) d B(x)$. In this procedure, if $\mu \in \mathbb{N}$, for all $j, l \in \mathbb{N}_{0}$ and $\lambda \in \Lambda_{j}, \lambda \subseteq[0,1]$, we define

$$
\Lambda_{j, l}(\lambda)=\left\{\lambda^{\prime} \in \Lambda_{j},:\left|s(\lambda)-s\left(\lambda^{\prime}\right)\right| \leqslant 2^{m l}\right\}
$$

and the random set

$$
S_{j, l}^{\mu}=\left\{\lambda^{\prime} \in \Lambda_{j},: 2^{l} \mu<\left|\varepsilon_{\lambda^{\prime}}\right| \leqslant 2^{l+1} \mu\right\} .
$$

Finally we consider the random set

$$
I_{j}^{\mu}=\left\{\lambda \in \Lambda_{j}, \lambda \subseteq[0,1]: \forall l \in \mathbb{N}_{0}, \Lambda_{j, l}(\lambda) \cap S_{j, l}^{\mu}=\emptyset\right\}
$$

and show that almost surely, there exists $\mu \in \mathbb{N}$ such that

$$
S_{\text {low }}^{\mu}=\bigcap_{j \in \mathbb{N}_{0}} \bigcup_{\lambda \in I_{j}^{\mu}} \bar{\lambda} \neq \emptyset
$$

which is equivalent to the fact that, for any $J$

$$
S_{\mathrm{low}, J}^{\mu}=\bigcap_{j \leqslant J} \bigcup_{\lambda \in I_{j}^{\mu}} \bar{\lambda} \neq \emptyset
$$

as $\left(S_{\text {low }, J}^{\mu}\right)_{J}$ is a decreasing sequence of compact sets. To do so, let us denote by $2 S_{\text {low, } J}^{\mu}$ the sets of dyadic intervals of scale $J+1$ obtained by cutting in two the remaining intervals ${ }^{2}$ in $S_{\text {low }, J}^{\mu}$ and remark that $S_{\text {low }, J+1}^{\mu}$ is obtained from $2 S_{\text {low }, J}^{\mu}$ by removing the dyadic intervals $\lambda$ such that $\Lambda_{J+1, l}(\lambda) \cap S_{J+1, l}^{\mu} \neq \emptyset$ for a $l \in \mathbb{N}_{0}$. But now, if $\xi \sim \mathcal{N}(0,1)$, we set, for all such a $l$

$$
p_{l}(\mu)=\mathbb{P}\left(2^{l} \mu<|\xi| \leqslant 2^{l+1} \mu\right) .
$$

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and note that, if $N$ is the number of intervals of $S_{\text {low, } J}^{\mu}$, counting the number of intervals in $2 S_{\text {low }, J}^{\mu} \cap S_{J+1, l}^{\mu}$ is a binomial random variable of parameter $2 N$ and $p_{l}(\mu)$ and this number is thus bounded by

$$
2 N\left(p_{l}(\mu)+(l+1) \sqrt{p_{l}(\mu)\left(1-p_{l}(\mu)\right)}\right)
$$

on an event of probability $1-(l+1)^{-2} N^{-1}$. Therefore, to pass from $S_{\text {low }, J}^{\mu}$ to $S_{\text {low }, J+1}^{\mu}$ we remove at most

$$
2 N \sum_{l=0}^{+\infty}\left(2^{m l+1}+1\right)\left(p_{l}(\mu)+(l+1) \sqrt{p_{l}(\mu)\left(1-p_{l}(\mu)\right)}\right)
$$

intervals with probability greater than $1-N^{-1}$. But if $\mu$ is large enough, as $p_{l}(\mu)$ is of order $\frac{e^{-\left(2^{l} \mu\right)^{2}}}{2^{l} \mu}$, one can make sure that this last term is bounded by $\frac{N}{2}$. So, if $N_{J}^{\mu}$ is the random variable counting the number of subintervals of $S_{\text {low, } J}^{\mu}$, we have

$$
\mathbb{P}\left(\left.N_{J+1}^{\mu} \geqslant \frac{3}{2} N_{J}^{\mu} \right\rvert\, N_{J}^{\mu}=N\right) \geqslant 1-N^{-1}
$$

which leads to the recursive formula

$$
\mathbb{P}\left(N_{J+1}^{\mu} \geqslant\left(\frac{3}{2}\right)^{J+1}\right) \geqslant\left(1-\left(\frac{2}{3}\right)^{J}\right) \mathbb{P}\left(N_{J}^{\mu} \geqslant\left(\frac{3}{2}\right)^{J}\right), \quad \forall J \in \mathbb{N}_{0},
$$

see [18, Lemma 3.6 and Theorem 3.7.]. Finally, we deduce

$$
\begin{equation*}
\mathbb{P}\left(\bigcup_{\mu} \bigcap_{J \in \mathbb{N}_{0}}\left(N_{J}^{\mu} \geqslant 1\right)\right)=1 \tag{3.58}
\end{equation*}
$$

Moreover, we can show that, in this case, $S_{\text {low }}^{\mu} \cap(0,1) \neq \emptyset$. If $\alpha>0$, applying this procedure with $\frac{1}{m}<\alpha$ gives us that any point $t \in S_{\text {low }}^{\mu} \cap(0,1)$ is a slow point of the fractional Brownian motion of exponent $\alpha$.

From formulas (3.1) and (3.2), we see that this procedure is also useful to bound the random variables appearing in the expansion (3.3) of the generalized Rosenblatt process. But, from the proofs of Propositions 3.5 and 3.21 we know that this is not sufficient and we also need to give a bound for the random variables $\Xi_{j}(\lambda)$, for $\lambda \in 3 \lambda_{n}(t), n \in \mathbb{N}$ and $j \geqslant n$. Such dyadic intervals are precisely the ones in the set $\Lambda_{n, 0}\left(\lambda_{n}(t)\right)$ and this fact forces us to consider the following modification of the procedure. For all $j \in \mathbb{N}$, if $l \neq 0$, the sets $S_{j, l}^{\mu}$ remain untouched as well as its associated probability $p_{l}(\mu)$ while for $l=0$ we set

$$
S_{j, 0}^{\mu}=\left\{\lambda^{\prime} \in \Lambda_{j}, \lambda^{\prime} \subseteq[0,1]: \exists j^{\prime} \geqslant j \Xi_{j^{\prime}}\left(\lambda^{\prime}\right)>\left(j^{\prime}-j+1\right) \mu\right\},
$$

with associated probability (which only depends on $\mu$ )

$$
p_{0}(\mu)=\mathbb{P}\left(\exists j^{\prime} \geqslant j \Xi_{j^{\prime}}(\lambda)>\left(j^{\prime}-j+1\right) \mu\right) .
$$

As $\Xi_{j^{\prime}}\left(\lambda_{1}\right)$ is independent of $\Xi_{j^{\prime}}\left(\lambda_{2}\right)$ as long as $\lambda_{1} \cap \lambda_{2}=\emptyset$, for all $J \in \mathbb{N}$, if $N$ is again the number of dyadic intervals of $S_{\text {low }, J}^{\mu}$, the number of such intervals in $2 S_{\text {low }, J}^{\mu} \cap S_{J+1,0}^{\mu}$ is still a binomial random variable of parameter $2 N$ and $p_{0}(\mu)$. Therefore If $\mu$ is large enough, using Theorems 3.14 and 3.15, one can still affirm

$$
2 N \sum_{l=0}^{+\infty}\left(2^{m l+1}+1\right)\left(p_{l}(\mu)+(l+1) \sqrt{p_{l}(\mu)\left(1-p_{l}(\mu)\right)}\right) \leqslant \frac{N}{2}
$$

and the end of the procedure is saved: equality (3.58) still holds. Now, if $t \in S_{\text {low }}^{\mu} \cap(0,1)$ we know that

$$
\begin{equation*}
\forall n \in \mathbb{N}, \forall \lambda \in 3 \lambda_{n}(t), \forall j \geqslant n, \Xi_{j}(\lambda) \leqslant(j-n+1) \mu . \tag{3.59}
\end{equation*}
$$

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Let us remark that, as for all $\lambda \in \Lambda_{n},\left|\varepsilon_{\lambda}^{2}\right| \leqslant 2 \Xi_{n}(\lambda)+1$, we still have, in this case, for all $\lambda \in 3 \lambda_{n}(t),\left|\varepsilon_{\lambda}\right| \leqslant C \mu$, for a deterministic constant $C>0$.

Starting from now we take $m$ such that $1 / m<\min \left\{H_{1}, H_{2}\right\}$ and $2 / m<1-H_{1}-H_{2}$.
In order to use notations (3.55) and (3.56), here after $\lambda_{1}$ (resp. $\lambda_{2}$ ) will always stand for the dyadic interval $\left[k_{1} 2^{-j_{1}},\left(k_{1}+1\right) 2^{-j_{1}}\right)$ (resp. $\left[k_{2} 2^{-j_{2}},\left(k_{2}+1\right) 2^{-j_{2}}\right)$ ) and $\psi_{\lambda_{1}}$ (resp. $\psi_{\lambda_{2}}$ ) will be the associated antiderivative of wavelet $\psi_{H_{1}}\left(2^{j_{1}} \cdot-k_{1}\right)$ (resp. $\psi_{H_{2}}\left(2^{j_{2}} \cdot-k_{2}\right)$ ) and $I_{\lambda_{1}, \lambda_{2}}[t, s]$ will stand for $I_{j_{1}, j_{2}}^{k_{1}, k_{2}}[t, s]$. Finally, $\varepsilon_{\lambda_{1}, \lambda_{2}}$ will stand for $\varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}}$. If $t \in(0,1)$, let $\left(y_{\lambda}(t)\right)_{\lambda \Lambda}$ be the sequence defined by

$$
y_{\lambda}(t)=2^{l} \text { if } \lambda \in \Lambda_{j}^{l}(t) .
$$

Note that, if we apply the preceding procedure, we find $\Omega_{\text {slo }}$ an event of probability 1 such that, for all $\omega \in \Omega_{\text {slo }}$, there exists $\mu$ for which $S_{\text {low }}^{\mu} \cap(0,1) \neq \emptyset$. Then, if $t$ belong to this set, we have, thanks to inequality (3.57) and equalities (3.1) and (3.2)

$$
\begin{equation*}
\left|\varepsilon_{\lambda_{1}, \lambda_{2}}(\omega)\right| \leqslant C \mu^{2} y_{\lambda_{1}}(t) y_{\lambda_{2}}(t) \tag{3.60}
\end{equation*}
$$

for a deterministic constant $C>0$. Again, we need to adapt the Lemmata from previous sections with this alternative upper bound.
Lemma 3.32. There exists a deterministic constant $C>0$ such that, for all $t, s \in(0,1)$ we have

$$
\begin{aligned}
& \quad \sum_{0 \leqslant j_{1}<n} \sum_{0 \leqslant j_{2}<n} \sum_{\lambda_{1} \in \Lambda_{j_{1}}, \lambda_{2} \in \Lambda_{j_{2}}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} y_{\lambda_{1}}(t) y_{\lambda_{2}}(t)\left|I_{\lambda_{1}, \lambda_{2}}[t, s]\right| \\
& \\
& \leqslant C|t-s|^{H_{1}+H_{2}-1} .
\end{aligned}
$$

Proof. If $\xi \in[s, t]$ and $\lambda \in \lambda_{j}^{l}(t)$, for $0 \leqslant j<n$ and $l \geqslant 1$,

$$
\left|2^{j} \xi-s(\lambda)\right| \geqslant|s(\lambda(t))-s(\lambda)|-2>2^{m(l-1)}-2
$$

and so, using the fast decay of the fractional antiderivatives of $\psi$ (3.7) and the definition of $\left(y_{\lambda}\right)_{\lambda \in \Lambda}$, we get for $0 \leqslant j_{1}, j_{2}<n$

$$
\begin{align*}
& \sum_{\lambda_{1} \in \Lambda_{j_{1}, \lambda_{2} \in \Lambda_{j_{2}}} y_{\lambda_{1}}(t) y_{\lambda_{2}}(t)\left|\psi_{\lambda_{1}}(\xi) \psi_{\lambda_{2}}(\xi)\right|}=\sum_{\left(l_{1}, l_{2}\right) \in \mathbb{N}_{0}^{2}} \sum_{\lambda_{1} \Lambda_{j_{1}}^{l_{1}}}(t) \\
\leqslant & \sum_{\lambda_{2} \Lambda_{j_{2}}^{l_{2}}(t)} y_{\lambda_{1}}(t) y_{\lambda_{2}}(t)\left|\psi_{\lambda_{1}}(\xi) \psi_{\lambda_{2}}(\xi)\right| \\
\leqslant & \sum_{\left(l_{1}, l_{2}\right) \in \mathbb{N}_{0}} \sum_{\lambda_{1} \Lambda_{j_{1}}^{l_{1}}(t)} \sum_{\lambda_{2} \Lambda_{j_{2}}^{l_{2}}(t)} \frac{2^{l_{1}+l_{2}}}{\left(3+\left|2^{j_{1}} \xi-k_{1}\right|\right)^{4}\left(3+\left|2^{j_{2}} \xi-k_{2}\right|\right)^{4}} \\
\leqslant & C \sum_{\left(l_{1}, l_{2}\right) \in \mathbb{N}_{0}} \sum_{\lambda_{1} \Lambda_{j_{1}}^{l_{1}}(t)} \sum_{\lambda_{2} \Lambda_{j_{2}}^{l_{2}}(t)} \frac{2^{l_{1}+l_{2}} 2^{-m\left(l_{1}+l_{2}\right)}}{\left(3+\left|2^{j_{1}} \xi-k_{1}\right|\right)^{3}\left(3+\left|2^{j_{2}} \xi-k_{2}\right|\right)^{3}} \\
\leqslant & C \sum_{k_{1} \in \mathbb{Z}} \frac{1}{\left(3+\left|2^{j_{1}} \xi-k_{1}\right|\right)^{3}} \sum_{k_{2} \in \mathbb{Z}} \frac{1}{\left(3+\left|2^{j_{2}} \xi-k_{2}\right|\right)^{3}} \leqslant C . \tag{3.61}
\end{align*}
$$

It leads, just as in Lemmata 3.7 and 3.23, to the desired estimate.
In what follows, we use these notations instead of the one given in Definition 3.8:

$$
\begin{aligned}
& \Lambda_{j_{2}}^{<}(t, s)=\left\{\lambda_{2} \in \Lambda_{j_{2}}: s\left(\lambda_{2}\right) \in \mathbb{Z}_{j_{2}}^{<}(t, s)\right\}, \\
& \Lambda_{j_{2}}^{>}(t, s)=\left\{\lambda_{2} \in \Lambda_{j_{2}}: s\left(\lambda_{2}\right) \in \mathbb{Z}_{j_{2}}^{>}(t, s)\right\}, \\
& \Lambda_{j_{2}}[t, s]=\left\{\lambda_{2} \in \Lambda_{j_{2}}: s\left(\lambda_{2}\right) \in \mathbb{Z}_{j_{2}}[t, s]\right\} .
\end{aligned}
$$

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Lemma 3.33. There exists a deterministic constant $C>0$ such that, for all $t, s \in(0,1)$ and $0 \leqslant j_{1}<n \leqslant j_{2}$, the quantities

$$
\begin{align*}
& \sum_{\lambda_{1} \in \Lambda_{j_{1}}} \sum_{\lambda_{2} \in \Lambda_{j_{2}}^{<}(t, s)} y_{\lambda_{1}}(t) y_{\lambda_{2}}(t)\left|I_{\lambda_{1}, \lambda_{2}}[t, s]\right|  \tag{3.62}\\
& \sum_{\lambda_{1} \in \Lambda_{j_{1}}} \sum_{\lambda_{2} \in \Lambda_{j_{2}}^{>}(t, s)} y_{\lambda_{1}}(t) y_{\lambda_{2}}(t)\left|I_{\lambda_{1}, \lambda_{2}}[t, s]\right| \tag{3.63}
\end{align*}
$$

are bounded by $C 2^{\frac{1}{m}\left(j_{2}-n\right)} 2^{-j_{2}}$.
Proof. Again, we prove the bound for (3.62), the reasoning for (3.63) being similar. Let us remark that, if $j_{2} \geqslant n x \in[s, t]$ and $\lambda_{j_{2}}(x) \in \Lambda_{j_{2}}^{l}(t)$ then, the construction and the definition of $\left(y_{\lambda}(t)\right)_{\lambda \in \Lambda}$ gives that

- $l \leqslant \frac{1}{m}\left(j_{2}-n\right)$, as $|s-t| \leqslant 2^{-n}$,
- if $\lambda \in \Lambda_{j_{2}}^{l_{2}}(x)$ then $\left|y_{\lambda}\right| \leqslant 2^{l_{2}} 2^{l+1} \mu$ while, by definition, if $l_{2} \geqslant 1$

$$
3+\left|2^{j_{2}} x-s(\lambda)\right| \geqslant 2+2^{m\left(l_{2}-1\right)}
$$

Therefore, if we set $D_{j_{2}}^{l}(t)=\bigcup_{\lambda \in \Lambda_{j_{2}}^{l}(t)} \lambda$, we have

$$
\begin{align*}
(3.62) & \leqslant \sum_{\lambda_{1} \in \Lambda_{j_{1}}} \sum_{\lambda_{2} \in \Lambda_{j_{2}}^{<}(t, s)} y_{\lambda_{1}}(t) y_{\lambda_{2}}(t) \int_{[s, t]}\left|\psi_{\lambda_{1}}(x) \psi_{\lambda_{2}}(x)\right| d x  \tag{3.64}\\
& \leqslant \sum_{0 \leqslant l \leqslant \frac{1}{m}\left(j_{2}-n\right)} \sum_{\lambda_{1} \in \Lambda_{j_{1}}} \sum_{\lambda_{2} \in \Lambda_{j_{2}}^{<}(t, s)} y_{\lambda_{1}}(t) y_{\lambda_{2}}(t) \int_{D_{j_{2}}^{l}(t)}\left|\psi_{\lambda_{1}}(x) \psi_{\lambda_{2}}(x)\right| d x
\end{align*}
$$

But, for all $x \in D_{j_{2}}^{l}$, using the same method as in (3.61), but splitting the sums according to the set $\Lambda_{j_{1}}^{l_{1}}(x)$ and $\Lambda_{j_{2}}^{l_{2}}(x)$ on which $y_{\lambda_{1}}(t) y_{\lambda_{2}}(t) \leqslant 2^{l+l_{1}+l_{2}+1}$ we get

$$
\begin{align*}
\sum_{\lambda_{1} \in \Lambda_{j_{1}}} & \sum_{\lambda_{2} \in \Lambda_{j_{2}}^{\circ}(t, s)}\left|\varepsilon_{\lambda_{1}, \lambda_{2}} \| \psi_{\lambda_{1}}(x) \psi_{\lambda_{2}}(x)\right| \\
& \leqslant C 2^{l+1} \sum_{\lambda_{1} \in \Lambda_{j_{1}}} \frac{1}{\left(3+\left|2^{j_{1}} x-k_{1}\right|\right)^{3}} \sum_{\lambda_{2} \in \Lambda_{j_{2}}^{<}(t, s)} \frac{1}{\left(3+\left|2^{j_{2}} x-k_{2}\right|\right)^{3}}  \tag{3.65}\\
& \leqslant C 2^{l+1} \sum_{\lambda_{2} \in \Lambda_{j_{2}}^{<}(t, s)} \frac{1}{\left(3+\left|2^{j_{2}} x-k_{2}\right|\right)^{3}} .
\end{align*}
$$

Finally, using the techniques in (3.14), we get (3.62) $\leqslant C 2^{\frac{1}{m}\left(j_{2}-n\right)} 2^{-j_{2}}$.
Lemma 3.34. There exists a deterministic constant $C>0$ such that, for all $t, s \in(0,1)$ and $0 \leqslant j_{1}<n \leqslant j_{2}$, the quantities

$$
\begin{align*}
& \sum_{\lambda_{1} \in \Lambda_{j_{1}}} \sum_{\lambda_{2} \in \Lambda_{j_{2}}[t, s]} y_{\lambda_{1}}(t) y_{\lambda_{2}}(t)\left|\int_{-\infty}^{\min \{s, t\}} \psi_{H_{1}}\left(2^{j_{1}} x-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} x-k_{2}\right) d x\right|  \tag{3.66}\\
& \sum_{\lambda_{1} \in \Lambda_{j_{1}}} \sum_{\lambda_{2} \in \Lambda_{j_{2}}[t, s]} y_{\lambda_{1}}(t) y_{\lambda_{2}}(t)\left|\int_{\max \{s, t\}}^{+\infty} \psi_{H_{1}}\left(2^{j_{1}} x-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} x-k_{2}\right) d x\right| \tag{3.67}
\end{align*}
$$

are bounded by $C 2^{\frac{1}{m}\left(j_{2}-n\right)} 2^{-j_{2}}$.

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Proof. Again, we assume $s \leqslant t$. If $x \in(-\infty, s]$ is such that $\lambda_{j_{1}}(x) \in \Lambda_{j_{1}}^{l}(s)$, we have, for all $\lambda_{1} \in \Lambda_{j_{1}}^{l_{1}}(x)$ and $\lambda_{2} \in \Lambda_{j_{2}}[t, s] \cap \Lambda_{j_{2}}^{l_{2}}(s)$ (with $j_{1}<n \leqslant j_{2}$ ),

$$
\begin{align*}
\frac{y_{\lambda_{1}}(t) y_{\lambda_{2}}(t)}{\left(3+\left|2^{j_{1}} x-k_{1}\right|\right)^{4}\left(3+\left|2^{j_{2}} x-k_{2}\right|\right)^{4}} & \leqslant C \frac{2^{\frac{1}{m}\left(j_{2}-n\right)+l+l_{1}+l_{2}+1} \mu^{2}}{\left(3+\left|2^{j_{1}} x-k_{1}\right|\right)^{4}\left(3+\left|2^{j_{2}} x-k_{2}\right|\right)^{5}} \\
& \leqslant C \frac{2^{\frac{1}{m}\left(j_{2}-n\right)+l+1}}{\left(3+\left|2^{j_{1}} x-k_{1}\right|\right)^{3}\left(3+\left|2^{j_{2}} x-k_{2}\right|\right)^{4}}  \tag{3.68}\\
& \leqslant C \frac{2^{\frac{1}{m}\left(j_{2}-n\right)}}{\left(3+\left|2^{j_{1}} x-k_{1}\right|\right)^{3}\left(3+\left|2^{j_{2}} x-k_{2}\right|\right)^{3}}
\end{align*}
$$

because $3+\left|2^{j_{2}} x-k_{2}\right|=3+k_{2}-2^{j_{2}} x \geqslant 2+2^{j_{1}}(s-x) \geqslant 2^{m(l-1)}$. Thus we get, using the fast decay of the fractional antiderivatives of the wavelet before splitting the integral over $(-\infty, s]$ into the integral over the sets $(-\infty, s] \cap D_{j_{1}}^{l}(s)$, in the same way as in (3.64), using (3.68) and finally the boundedness of the function (3.42) for $M=3$ and inequality (3.15)

$$
\begin{aligned}
\sum_{\lambda_{1} \in \Lambda_{j_{1}}} & \sum_{\lambda_{2} \in \Lambda_{j_{2}}[t, s]} \int_{-\infty}^{s}\left|\psi_{\lambda_{1}}(x) \psi_{\lambda_{2}}(x)\right| d x \\
& \leqslant C 2^{\frac{1}{m}\left(j_{2}-n\right)} \int_{-\infty}^{s} \sum_{\lambda_{1} \in \Lambda_{j_{1}}} \sum_{\lambda_{2} \in \Lambda_{j_{2}}[t, s]} \frac{d x}{\left(3+\left|2^{j_{1}} x-k_{1}\right|\right)^{3}\left(3+\left|2^{j_{2}} x-k_{2}\right|\right)^{3}} \\
& \leqslant C 2^{\frac{1}{m}\left(j_{2}-n\right)} \int_{-\infty}^{s} \sum_{\lambda_{2} \in \Lambda_{j_{2}}[t, s]} \frac{d x}{\left(3+\left|2^{j_{2}} x-k_{2}\right|\right)^{3}} \leqslant C 2^{\frac{1}{m}\left(j_{2}-n\right)} 2^{-j_{2}} .
\end{aligned}
$$

In the same way we get

$$
\sum_{\lambda_{1} \in \Lambda_{j_{1}}} \sum_{\lambda_{2} \in \Lambda_{j_{2}}[t, s]} y_{\lambda_{1}}(t) y_{\lambda_{2}}(t)\left|\int_{\max \{t, s\}}^{+\infty} \psi_{H_{1}}\left(2^{j_{1}} x-k_{1}\right) \psi_{H_{2}}\left(2^{j_{2}} x-k_{2}\right) d x\right| \leqslant C 2^{-j_{2}} .
$$

Lemma 3.35. There exists a deterministic constant $C>0$ such that, for all $t, s \in(0,1)$, the quantities

$$
\begin{aligned}
& \sum_{0 \leqslant j_{1}<n} \sum_{j_{2} \geqslant n} \sum_{\lambda_{1} \in \Lambda_{j_{1}}, \lambda_{2} \in \Lambda_{j_{2}}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} y_{\lambda_{1}}(t) y_{\lambda_{2}}(t)\left|I_{\lambda_{1}, \lambda_{2}}[t, s]\right| \\
& \sum_{j_{1} \geqslant n} \sum_{0 \leqslant j_{2}<n} \sum_{\lambda_{1} \in \Lambda_{j_{1}}, \lambda_{2} \in \Lambda_{j_{2}}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} y_{\lambda_{1}}(t) y_{\lambda_{2}}(t)\left|I_{\lambda_{1}, \lambda_{2}}[t, s]\right|
\end{aligned}
$$

are bounded by $C|t-s|^{H_{1}+H_{2}-1}$.
Proof. The proof is exactly the same as the one of Lemma 3.12 excepted that we use Lemmata 3.33 and 3.34 instead of Lemmata 3.10 and 3.11 respectively. It leads on one side us to consider the sums

$$
\sum_{j_{1}=0}^{n-1} 2^{j_{1}\left(1-H_{1}\right)} \sum_{j_{2}=n}^{+\infty} 2^{\frac{1}{m}\left(j_{2}-n\right)} 2^{-j_{2} H_{2}}+2^{n\left(1-H_{1}-H_{2}\right)}
$$

which are bounded by $C 2^{n\left(1-H_{1}-H_{2}\right)} \leqslant C|t-s|^{H_{1}+H_{2}-1}$, because $\frac{1}{m}<H_{2}$. On the other side, if we write $I_{\lambda_{1}, \lambda_{2}}$ for $I_{j_{1}, j_{2}}^{k_{1}, k_{2}}$ in Lemma 2.1, we have, from it,

$$
\left|\sum_{\lambda_{1} \in \Lambda_{n-1}} \sum_{\lambda_{2} \in \Lambda_{n}[t, s]} y_{\lambda_{1}}(t) y_{\lambda_{2}}(t) I_{\lambda_{1}, \lambda_{2}}\right|
$$

$$
\begin{align*}
& \leqslant C 2^{-n}\left|\sum_{l_{1}=0}^{+\infty} \sum_{\lambda_{1} \in \Lambda_{n-1}^{l_{1}}(t)} \sum_{\lambda_{2} \in \Lambda_{n}[t, s]} \frac{2^{l_{1}}}{\left(3+\left|2 k_{1}-k_{2}\right|\right)^{4}}\right| \\
& \leqslant C 2^{-n}\left|\sum_{\lambda_{1} \in \Lambda_{n-1}} \sum_{\lambda_{2} \in \Lambda_{n}[t, s]} \frac{1}{\left(3+\left|2 k_{1}-k_{2}\right|\right)^{3}}\right| \leqslant C 2^{-n} . \tag{3.69}
\end{align*}
$$

Lemma 3.36. There exists a deterministic constant $C>0$ such that, for all $t, s \in(0,1)$ and $n \leqslant j_{1} \leqslant j_{2}$ the quantities (3.62) and (3.63) are bounded by $C 2^{\frac{1}{m}\left(j_{1}-n\right)} 2^{\frac{1}{m}\left(j_{2}-n\right)} 2^{-j_{2}}$.

Proof. The proof is essentially the same as for Lemma 3.33 excepted that, now, as $n \leqslant j_{1} \leqslant j_{2}$, we remark that if $x \in D_{j_{2}}^{l}(t)$ for a $0 \leqslant l \leqslant \frac{1}{m}\left(j_{2}-n\right)$ then $x \in D_{j_{1}}^{l^{\prime}}(t)$ for a $0 \leqslant l^{\prime} \leqslant \frac{1}{m}\left(j_{1}-n\right)$.

Lemma 3.37. There exists a deterministic constant $C>0$ such that, for all $t, s \in(0,1)$ and $n \leqslant j_{1} \leqslant j_{2}$ the quantities (3.66) and (3.67) are bounded by $C 2^{\frac{1}{m}\left(j_{1}-n\right)} 2^{\frac{1}{m}\left(j_{2}-n\right)} 2^{-j_{2}}$.

Proof. The proof is essentially the same as for Lemma 3.34 and the only modification is the same as in the proof of Lemma 3.36.

This time, the bound for the random variables $\Xi_{j}(\lambda)$ are already considered in the construction and we can directly go to the proof of the main Proposition of this subsection.

Proof of Proposition 3.31. If we apply the procedure with $m$ such that $1 / m<\min \left\{H_{1}, H_{2}\right\}$ and $2 / m<1-H_{1}-H_{2}$, we find an event $\Omega_{\text {slo }}$ of probability 1 such that, for all $\omega \in \Omega_{\text {slo }}$, there is $\mu \in \mathbb{N}$ for which $S_{\text {low }}^{\mu} \cap(0,1) \neq \emptyset$. Then, if $\omega \in \Omega_{\text {slo }}$ and $t \in S_{\text {low }}^{\mu}(\omega) \cap(0,1)$ and $s \in(0,1)$, we write

$$
\begin{align*}
&\left|R_{H_{1}, H_{2}}^{\prime}(t, \omega)-R_{H_{1}, H_{2}}^{\prime}(s, \omega)\right| \\
& \leqslant\left|\sum_{0 \leqslant j_{1}<n} \sum_{0 \leqslant j_{2}<n} \sum_{\lambda_{1} \in \Lambda_{j_{1}, \lambda_{2} \in \Lambda_{j_{2}}}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{\lambda_{1}, \lambda_{2}}(\omega) I_{\lambda_{1}, \lambda_{2}}[t, s]\right| \\
&+\left|\sum_{0 \leqslant j_{1}<n} \sum_{j_{2} \geqslant n} \sum_{\lambda_{\lambda_{1} \in \Lambda_{j_{1}}, \lambda_{2} \in \Lambda_{j_{2}}}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{\lambda_{1}, \lambda_{2}}(\omega) I_{\lambda_{1}, \lambda_{2}}[t, s]\right|  \tag{3.70}\\
&+\left|\sum_{j_{1} \geqslant n} \sum_{0 \leqslant j_{2}<n} \sum_{\lambda_{\lambda_{1} \in \Lambda_{j_{1}}, \lambda_{2} \in \Lambda_{j_{2}}}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{j_{1}, j_{2}}^{k_{1}, k_{2}}(\omega) I_{\lambda_{1}, \lambda_{2}}[t, s]\right| \\
&+\left|\sum_{j_{1} \geqslant n} \sum_{j_{2} \geqslant n} \sum_{\lambda_{1} \in \Lambda_{j_{1}, \lambda_{2} \in \Lambda_{j_{2}}}} 2^{j_{1}\left(1-H_{1}\right)} 2^{j_{2}\left(1-H_{2}\right)} \varepsilon_{\lambda_{1}, \lambda_{2}}(\omega) I_{\lambda_{1}, \lambda_{2}}[t, s]\right| .
\end{align*}
$$

As inequality (3.60) holds, we use Lemma 3.32 to bound the first sum, and Lemma 3.35 to bound the second and the third one. For the last sum, from inequality (3.60) and Lemmata 3.36 and 3.37, it just remains us to find bound for the random variables (3.36), (3.37) and (3.38)with $\ell \in\{j, j+1\}$ on $\Omega_{\text {slo }}$. For (3.36) with $\ell=j$, we have, as in (3.69) and then (3.39)

$$
\left|\sum_{\lambda_{1} \in \Lambda_{j}^{\wedge}(t, s)} \sum_{\lambda_{2} \in \Lambda_{j}[t, s]} \varepsilon_{\lambda_{1}, \lambda_{2}}(\omega) I_{\lambda_{1}, \lambda_{2}}\right|
$$

$$
\begin{aligned}
& \leqslant C 2^{-j} 2^{\frac{2}{m}(j-n)} \mu^{2}\left|\sum_{\lambda_{1} \in \Lambda_{j}^{<}(t, s)} \sum_{\lambda_{2} \in \Lambda_{j}[t, s]} \frac{1}{\left(3+\left|2 k_{1}-k_{2}\right|\right)^{3}}\right| \\
& \leqslant C 2^{-j} 2^{\frac{2}{m}(j-n)} \mu^{2} .
\end{aligned}
$$

The same bound holds when we consider the sums over $\lambda_{1} \in \Lambda_{j}^{>}(t, s)$ or $\lambda_{2} \in \Lambda_{j+1}[t, s]$, i.e. for (3.36) and (3.37). Finally the construction and especially (3.59) insures us that

$$
\left|\sum_{\lambda_{1} \in \Lambda_{j}[t, s]} \sum_{\lambda_{2} \in \Lambda_{j}[t, s]} \varepsilon_{\lambda_{1}, \lambda_{2}}(\omega) I_{\lambda_{1}, \lambda_{2}}\right| \leqslant C(j-n+1) 2^{\frac{-j-n}{2}} \mu .
$$

Therefore, the last term in (3.70) is bounded from above by

$$
\begin{aligned}
& C \mu^{2}\left(\sum_{j_{1} \geqslant n} 2^{j_{1}\left(1-H_{1}\right)} 2^{\frac{1}{m}\left(j_{1}-n\right)} \sum_{j_{2} \geqslant j_{1}} 2^{-j_{2} H_{2}} 2^{\frac{1}{m}\left(j_{2}-n\right)}+\sum_{j \geqslant n} 2^{j\left(\frac{3}{2}-H_{1}-H_{2}\right)}(j-n+1) 2^{-\frac{n}{2}}\right) \\
& \leqslant C \mu^{2}\left(\sum_{j_{1} \geqslant n} 2^{j_{1}\left(1-H_{1}-H_{2}\right)} 2^{\frac{2}{m}\left(j_{1}-n\right)}+2^{n\left(\frac{3}{2}-H_{1}-H_{2}\right)} 2^{-\frac{n}{2}}\right) \\
& \leqslant C \mu^{2} 2^{n\left(1-H_{1}-H_{2}\right)} \leqslant C \mu^{2}|t-s|^{H_{1}+H_{2}-1}
\end{aligned}
$$

and thus inequality (3.54) holds.

## 4 Lower bounds for wavelet leaders

In this section, we show that the limits (1.6) and (1.7) are strictly positive. In [7], the authors used the independence of the increments of the Brownian motion to bound from below its wavelet leaders. But, for the (generalized) Rosenblatt process this nice feature is not met anymore. Nevertheless, following an idea by Ayache in a closely related but different context ${ }^{3}$ [5], we decompose the wavelet coefficients of the generalized Rosenblatt process in two parts. We gain some independence properties in the first part while the second is, in some sense, negligible compared to the first, see Proposition 4.6 below. All along this section, $C$ stands for a deterministic constant whose value may change from a line to another but does not depend on any relevant quantities, and in order to ease notations we set

$$
C_{H_{1}, H_{2}}:=\frac{1}{\Gamma\left(H_{1}-\frac{1}{2}\right) \Gamma\left(H_{2}-\frac{1}{2}\right)}
$$

and for $s, x_{1}, x_{2} \in \mathbb{R}$

$$
f_{H_{1}, H_{2}}\left(s, x_{1}, x_{2}\right)=\left(s-x_{1}\right)_{+}^{H_{1}-3 / 2}\left(s-x_{2}\right)_{+}^{H_{2}-3 / 2}
$$

Let $\Psi$ be a wavelet with compact support included in $[-N, N]$. Using formula (2.6) at $t=k / 2^{j}$, the wavelet coefficient $c_{j, k}$ of the generalized Rosenblatt process is given by

$$
\begin{aligned}
c_{j, k} & =\int_{-N}^{N}\left[R_{H_{1}, H_{2}}\left(\frac{x+k}{2^{j}}\right)-R_{H_{1}, H_{2}}\left(\frac{k}{2^{j}}\right)\right] \Psi(x) d x \\
& =c_{H_{1}, H_{2}} \int_{-N}^{N} \Psi(x) \int_{\mathbb{R}^{2}}^{\prime} \int_{\frac{k}{2^{j}}}^{\frac{x+k}{2^{j}}} f_{H_{1}, H_{2}}\left(s, x_{1}, x_{2}\right) d s d B\left(x_{1}\right) d B\left(x_{2}\right) d x \\
& =c_{H_{1}, H_{2}} \int_{\mathbb{R}^{2}}^{\prime} \int_{-N}^{N} \Psi(x) \int_{\frac{k}{2^{j}}}^{\frac{x+k}{2^{j}}} f_{H_{1}, H_{2}}\left(s, x_{1}, x_{2}\right) d s d x d B\left(x_{1}\right) d B\left(x_{2}\right)
\end{aligned}
$$

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$$
=c_{H_{1}, H_{2}} \int_{A}^{\prime} \int_{-N}^{N} \Psi(x) \int_{\frac{k}{2 j}}^{\frac{x+k}{2 j}} f_{H_{1}, H_{2}}\left(s, x_{1}, x_{2}\right) d s d x d B\left(x_{1}\right) d B\left(x_{2}\right)
$$

where $\left.A:=]-\infty, \frac{k+N}{2^{j}}\right]^{2}$, because, as long as $x \in[-N, N]$ and $s \in\left[k 2^{-J},(k+N) 2^{-j}\right]$, $f\left(s, x_{1}, x_{2}\right)$ vanishes for all $x_{1}, x_{2}$ outside of $A$.
Definition 4.1. Given an integer $M \geqslant 0, c_{j, k}$ can be written as $c_{j, k}=\widetilde{c_{j, k}}{ }^{M}+{\widetilde{c_{j, k}}}^{M}$ where

$$
\begin{equation*}
\widetilde{c}_{j, k}^{M}=c_{H_{1}, H_{2}} \int_{\lambda_{j, k}^{M}}^{\prime} \int_{-N}^{N} \Psi(x) \int_{\frac{k}{2 j}}^{\frac{x+k}{2 j}} f_{H_{1}, H_{2}}\left(s, x_{1}, x_{2}\right) d s d x d B\left(x_{1}\right) d B\left(x_{2}\right) \tag{4.1}
\end{equation*}
$$

with

$$
\left.\left.\lambda_{j, k}^{M}:=\right] \frac{k-N M}{2^{j}}, \frac{k+N}{2^{j}}\right]^{2}
$$

and

$$
{\widetilde{c_{j, k}}}^{M}=c_{H_{1}, H_{2}} \int_{A \backslash \lambda_{j, k}^{M,}}^{\prime} \int_{-N}^{N} \Psi(x) \int_{\frac{k}{2^{j}}}^{\frac{x+k}{2 j}} f_{H_{1}, H_{2}}\left(s, x_{1}, x_{2}\right) d s d x d B\left(x_{1}\right) d B\left(x_{2}\right) .
$$

Remark 4.2. Let us highlight the fact that using time change of variable for Wiener-Itô integrals [38, Theorem 8.5.7], for all $j, k$, we have that $\widetilde{c_{j, k}}{ }^{M}$ is equal in law to the random variable

$$
c_{H_{1}, H_{2}} 2^{-j\left(H_{1}+H_{2}-1\right)} \int_{I_{M}}^{\prime} \int_{-N}^{N} \psi(x) \int_{0}^{x} f_{H_{1}, H_{2}}\left(s, x_{1}, x_{2}\right) d s d x d B\left(x_{1}\right) d B\left(x_{2}\right)
$$

with $I_{M}=(-M N, N]^{2}$, while ${\widetilde{c_{j, k}}}^{M}$ is equal in law to the random variable

$$
c_{H_{1}, H_{2}} 2^{-j\left(H_{1}+H_{2}-1\right)} \int_{I_{M}^{\prime}}^{\prime} \int_{-N}^{N} \psi(x) \int_{0}^{x} f_{H_{1}, H_{2}}\left(s, x_{1}, x_{2}\right) d s d x d B\left(x_{1}\right) d B\left(x_{2}\right)
$$

with $I_{M}^{\prime}=(-\infty, N]^{2} \backslash(-M N, N]^{2}$.
Definition 4.3. For all $(j, k) \in \mathbb{N} \times \mathbb{Z}$ and $M \in \mathbb{N}$ we define the random variables

$$
{\widetilde{\varepsilon_{j, k}}}^{M}:=\frac{{\widetilde{c_{j, k}}}^{M}}{2^{-j\left(H_{1}+H_{2}-1\right)}} \text { and }{\widetilde{\varepsilon_{j, k}}}^{M}:=\frac{{\widetilde{c_{j, k}}}^{M}}{2^{-j\left(H_{1}+H_{2}-1\right)}} .
$$

Remark 4.4. Note that $\widetilde{\varepsilon_{j, k}}{ }^{M}$ and $\widetilde{\varepsilon_{j^{\prime}, k^{\prime}}}{ }^{M}$ are independent when

$$
\begin{equation*}
\lambda_{j, k}^{M} \cap \lambda_{j^{\prime}, k^{\prime}}^{M}=\emptyset . \tag{4.2}
\end{equation*}
$$

Indeed, if $\left(f_{j}\right)_{j}$ is a sequence of real-valued step functions on $\mathbb{R}^{2} \backslash\{(x, x): x \in \mathbb{R}\}$ which converge to the integrand with respect to $d B\left(x_{1}\right) d B\left(x_{2}\right)$ in (4.1) then the integral $\int_{\mathbb{R}^{2}}^{\prime} f_{j}\left(x_{1}, x_{2}\right) d B\left(x_{1}\right) d B\left(x_{2}\right)$ is a polynomial function of a finite number of increments $B\left(t_{2}\right)-B\left(t_{1}\right)$ of the Brownian motion for some $t_{1}, t_{2} \in \lambda_{j, k}^{M}$. Thus $\widetilde{\varepsilon_{j, k}}{ }^{M}$ is measurable with respect to the $\sigma$-algebra generated by these increments

$$
\sigma_{j, k}^{M}:=\sigma\left(\left\{B\left(t_{2}\right)-B\left(t_{1}\right): t_{1}, t_{2} \in \lambda_{j, k}^{M}\right\}\right) .
$$

Using the independence of the increments of the Brownian motion, one concludes that $\sigma_{j, k}^{M}$ and $\sigma_{j^{\prime}, k^{\prime}}^{M}$ are independent as long as condition (4.2) is met and so the same holds for $\widetilde{\varepsilon_{j, k}}{ }^{M}$ and $\widetilde{\varepsilon_{j, k^{\prime}}}{ }^{M}$. Moreover, $\widetilde{\varepsilon_{j_{1}, k_{1}}}{ }^{M}, \ldots, \widetilde{\varepsilon_{j_{n}, k_{n}}}{ }^{M}$ are independent when the following condition is satisfied

$$
\begin{equation*}
\lambda_{j_{i}, k_{i}}^{M} \cap \lambda_{j_{l}, k_{l}}^{M}=\emptyset \text { for all } 1 \leqslant i<l \leqslant n . \tag{4.3}
\end{equation*}
$$

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This leads to defining the following condition.
Definition 4.5. Let $n \geqslant 2$. We say $\lambda_{j_{1}, k_{1}}, \ldots, \lambda_{j_{n}, k_{n}}$ satisfy condition ( $C_{M}$ ) if (4.3) is satisfied.

From Remark 4.2, we know that $\left({\widetilde{\varepsilon_{j, k}}}^{M}\right)_{j, k}$ is a family of identically distributed second order Wiener chaos random variables. Moreover, $\widetilde{\varepsilon_{j_{1}, k_{1}}}{ }^{M}, \ldots, \widetilde{\varepsilon_{j_{n}, k_{n}}}{ }^{M}$ are independent as long as $\lambda_{j_{1}, k_{1}}, \ldots, \lambda_{j_{n}, k_{n}}$ satisfies $\left(C_{M}\right)$.

Our aim now is to provide a lower bound (independent of $M$ ) for the tail behaviour of the random variables $\widetilde{\varepsilon_{j, k}}{ }^{M}$. To this end, we start by proving the following lemma.
Lemma 4.6. There exist three strictly positive deterministic constants $C_{\Psi, H_{1}, H_{2}}$, $C_{\Psi, H_{1}, H_{2}}^{\prime}$ and $C_{\Psi, H_{1}, H_{2}}^{*}$ such that for all $(j, k) \in \mathbb{N} \times \mathbb{Z}$ and $M \geqslant 2$ one has

$$
\begin{aligned}
C_{\Psi, H_{1}, H_{2}} 2^{-j\left(H_{1}+H_{2}-1\right)} \leqslant & \left\|{\widetilde{c_{j, k}} M}_{M}\right\|_{L^{2}(\Omega)} \leqslant C_{\Psi, H_{1}, H_{2}}^{\prime} 2^{-j\left(H_{1}+H_{2}-1\right)} \\
& \| \widetilde{c_{j, k}} M
\end{aligned} \|_{L^{2}(\Omega)} \leqslant C_{\Psi, H_{1}, H_{2}}^{*} 2^{-j\left(H_{1}+H_{2}-1\right)} M^{\max \left\{H_{1}, H_{2}\right\}-1} .
$$

Proof. Let us assume, w.l.o.g. that $H_{1} \geqslant H_{2}$. We define the functions

$$
\begin{aligned}
& \Phi_{1}:\left(x_{1}, x_{2}\right) \mapsto \int_{-N}^{N} \Psi(x) \int_{0}^{x} f_{H_{1}, H_{2}}\left(s, x_{1}, x_{2}\right) d s d x \\
& \Phi_{2}:\left(x_{1}, x_{2}\right) \mapsto \int_{-N}^{N} \Psi(x) \int_{0}^{x} f_{H_{1}, H_{2}}\left(s, x_{2}, x_{1}\right) d s d x
\end{aligned}
$$

and the symmetric function ${ }^{4} \Phi=\frac{1}{2}\left(\Phi_{1}+\Phi_{2}\right)$. By Remark 4.2 we have, using the "Wiener isometry" ${ }^{5}$ [42, Section 5],

$$
\left\|\widetilde{c j, k}^{M}\right\|_{L^{2}(\Omega)}=\sqrt{2} c_{H_{1}, H_{2}} 2^{-j\left(H_{1}+H_{2}-1\right)}\|\Phi\|_{L^{2}\left(I_{M}\right)}
$$

and thus it suffices to take $C_{\Psi, H_{1}, H_{2}}:=\sqrt{2} c_{H_{1}, H_{2}}\|\Phi\|_{L^{2}\left([-N, N]^{2}\right)}$ and $C_{\Psi, H_{1}, H_{2}}^{\prime}:=\sqrt{2} c_{H_{1}, H_{2}}$ $\|\Phi\|_{L^{2}\left((-\infty, N]^{2}\right)}$. Now, still using Remark 4.2 and "Wiener isometry" we have

$$
\begin{aligned}
\left\|{\widetilde{c_{j, k}}}^{M}\right\|_{L^{2}(\Omega)} & =\sqrt{2} c_{H_{1}, H_{2}} 2^{-j\left(H_{1}+H_{2}-1\right)}\|\Phi\|_{L^{2}\left(I_{M}^{\prime}\right)} \\
& \leqslant \sqrt{2} c_{H_{1}, H_{2}} 2^{-j\left(H_{1}+H_{2}-1\right)}\left\|\Phi_{1}\right\|_{L^{2}\left(I_{M}^{\prime}\right)} .
\end{aligned}
$$

Also as

$$
I_{M}^{\prime}=(-\infty, N]^{2} \backslash(-M N, N]^{2} \subset \mathbb{R} \times(-\infty,-M N] \bigcup(-\infty,-M N] \times \mathbb{R}
$$

we write

$$
\begin{aligned}
& \left\|\Phi_{1}\right\|_{L^{2}\left(I_{M}^{\prime}\right)}^{2} \\
& =\int_{I_{M}^{\prime}}\left|\int_{-N}^{N} \Psi(x) \int_{0}^{x}\left(s-x_{1}\right)_{+}^{H_{1}-\frac{3}{2}}\left(s-x_{2}\right)_{+}^{H_{2}-\frac{3}{2}} d s d x\right|^{2} d x_{1} d x_{2} \\
& \leqslant \int_{I_{M}^{\prime}}\left(\int_{-N}^{N}|\Psi(x)| \int_{[0, x]}\left(s-x_{1}\right)_{+}^{H_{1}-\frac{3}{2}}\left(s-x_{2}\right)_{+}^{H_{2}-\frac{3}{2}} d s d x\right)^{2} d x_{1} d x_{2} \\
& \leqslant \int_{\mathbb{R}} \int_{-\infty}^{-M N}\left(\int_{-N}^{N}|\Psi(x)| \int_{[0, x]}\left(s-x_{1}\right)_{+}^{H_{1}-\frac{3}{2}}\left(s-x_{2}\right)_{+}^{H_{2}-\frac{3}{2}} d s d x\right)^{2} d x_{1} d x_{2}
\end{aligned}
$$

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$$
+\int_{\mathbb{R}} \int_{-\infty}^{-M N}\left(\int_{-N}^{N}|\Psi(x)| \int_{[0, x]}\left(s-x_{1}\right)_{+}^{H_{1}-\frac{3}{2}}\left(s-x_{2}\right)_{+}^{H_{2}-\frac{3}{2}} d s d x\right)^{2} d x_{2} d x_{1} .
$$

Let us deal with the first term in the last sum, the second one can be treated similarly by permuting the roles of $H_{1}$ and $H_{2}$ as well as $x_{1}$ and $x_{2}$. As the function $y \mapsto y^{H_{1}-3 / 2}$ is decreasing, one gets

$$
\begin{aligned}
& \int_{\mathbb{R}} \int_{-\infty}^{-M N}\left(\int_{-N}^{N}|\Psi(x)| \int_{[0, x]}\left(s-x_{1}\right)_{+}^{H_{1}-\frac{3}{2}}\left(s-x_{2}\right)_{+}^{H_{2}-\frac{3}{2}} d s d x\right)^{2} d x_{1} d x_{2} \\
& \leqslant\left(\int_{-\infty}^{-M N}\left(-N-x_{1}\right)^{2 H_{1}-3} d x_{1}\right) \times \int_{\mathbb{R}}\left(\int_{-N}^{N}|\Psi(x)|\left|\int_{[0, x]}\left(s-x_{2}\right)_{+}^{H_{2}-3 / 2} d s\right| d x\right)^{2} d x_{2} .
\end{aligned}
$$

Concerning the first integral, we have, as $M \geqslant 2$

$$
\begin{aligned}
\int_{-\infty}^{-N M}\left(-N-x_{1}\right)^{2 H_{1}-3} d x_{1} & =\frac{1}{2-2 H_{1}}(N M-N)_{+}^{2 H_{1}-2} \\
& =\frac{1}{2-2 H_{1}} N^{2 H_{1}-2}(M-1)^{2 H_{1}-2} \\
& \leqslant c \cdot M^{2 H_{1}-2}
\end{aligned}
$$

while, using again the "Wiener isometry",

$$
\begin{aligned}
& \int_{\mathbb{R}}\left(\int_{-N}^{N}|\Psi(x)|\left|\int_{[0, x]}\left(s-x_{2}\right)_{+}^{H_{2}-3 / 2}\right| d x\right)^{2} d x_{2} \\
& \leqslant 2 N\|\Psi\|_{\infty} \sup _{x \in[-N, N]} \int_{\mathbb{R}}\left|\int_{[0, x]}\left(s-x_{2}\right)_{+}^{H_{2}-3 / 2} d s\right|^{2} d x_{2} \\
& =2 N\|\Psi\|_{\infty} \sup _{x \in[-N, N]} \mathbb{E}\left[\left|B_{H_{2}}(x)-B_{H_{2}}(0)\right|^{2}\right] \\
& \leqslant 2 N\|\Psi\|_{\infty} \sup _{x \in[-N, N]} C_{H_{2}}(|x|)^{2 H_{2}} \leqslant c,
\end{aligned}
$$

where $B_{H_{2}}$ denotes the fractional Brownian motion with parameter $H_{2}$. As a result, there exists a positive constant $C_{\Psi, H_{1}, H_{2}}^{*}$ such that, as we suppose $H_{1} \geqslant H_{2}$, one has

$$
\left\|{\widetilde{c_{j, k}}}\right\|_{L^{2}(\Omega)} \leqslant C_{\Psi, H_{1}, H_{2}}^{*} 2^{-j\left(H_{1}+H_{2}-1\right)} M^{H_{1}-\boxminus}
$$

Proposition 4.7. Let $M \in \mathbb{N}$ and $y \in \mathbb{R}^{+}$. If $M$ and $y$ are large enough, then the exists a deterministic constant $c_{2}>0$ (independent of $M$ ) such that

$$
\begin{equation*}
\mathbb{P}\left(\left|{\widetilde{\varepsilon_{j, k}}}^{M}\right|>y\right) \geqslant \exp \left(-c_{2} y\right) \tag{4.4}
\end{equation*}
$$

for $\operatorname{all}(j, k) \in \mathbb{N} \times \mathbb{Z}$
Proof. Fix $y \in \mathbb{R}^{+}$, large enough. By Lemma 4.6, one can remark that as $M \rightarrow+\infty$, $\left(\widetilde{\varepsilon_{j, k}}{ }^{M}\right)_{M}$ converges in $L^{2}(\Omega)$ to the random variable

$$
\varepsilon_{j, k}:=\frac{c_{j, k}}{2^{-j\left(H_{1}+H_{2}-1\right)}}
$$

with, for all $M \in \mathbb{N}$,

$$
\varepsilon_{j, k}-{\widetilde{\varepsilon_{j, k}}}^{M}={\widetilde{\varepsilon_{j, k}}}^{M} .
$$

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By Theorem 3.15, there exists a constant $c_{1}>0$ such that, for all $\lambda \in \Lambda$ and $y$ sufficiently large

$$
\mathbb{P}\left(\left|\varepsilon_{j, k}\right| \geqslant y\right) \geqslant \exp \left(-c_{1} y\right) .
$$

Then, for all $M \in \mathbb{N}$, we have, for all such $\lambda$ and $y$

$$
\begin{aligned}
& \mathbb{P}\left(\left|\widetilde{\varepsilon_{j, k}}{ }^{M}\right| \geqslant y\right) \geqslant \mathbb{P}\left(\left\{\left|\widetilde{\varepsilon_{j, k}}{ }^{M}\right| \geqslant y\right\} \cap\left\{\left|{\widetilde{\varepsilon_{j, k}}}^{M}\right| \leqslant y\right\}\right) \\
& \geqslant \mathbb{P}\left(\left\{\left|\varepsilon_{j, k}\right|-\left|{\widetilde{\varepsilon_{j, k}}}^{M}\right| \geqslant y\right\} \cap\left\{\left|{\widetilde{\varepsilon_{j, k}}}^{M}\right| \leqslant y\right\}\right) \\
& \geqslant \mathbb{P}\left(\left\{\left|\varepsilon_{j, k}\right| \geqslant 2 y\right\} \cap\left\{\left|{\widetilde{\varepsilon_{j, k}}}^{M}\right| \leqslant y\right\}\right) \\
& \geqslant \mathbb{P}\left(\left|\varepsilon_{j, k}\right| \geqslant 2 y\right)-\mathbb{P}\left(\left(\left|{\overline{\varepsilon_{j, k}}}^{M}\right|>y\right) .\right.
\end{aligned}
$$

Using Lemma 4.6 and Theorem 3.14 one has

$$
\begin{aligned}
\mathbb{P}\left(\left|{\widetilde{\varepsilon_{j, k}}}^{M}\right|>y\right) & \leqslant \mathbb{P}\left(\left(\left|{\overline{c_{j, k}}}^{M}\right|>y\left\|{\widetilde{c_{j, k}}}^{M}\right\|_{L^{2}(\Omega)}\left(C_{\Psi, H_{1}, H_{2}}^{*}\right)^{-1} M^{1-\max \left\{H_{1}, H_{2}\right\}}\right)\right. \\
& \leqslant \exp \left(-\stackrel{\left.\stackrel{C}{C}\left(C_{\Psi, H_{1}, H_{2}}^{*}\right)^{-1} M^{1-\max \left\{H_{1}, H_{2}\right\}} y\right) .}{ } .\right.
\end{aligned}
$$

Thus, if $M$ is large enough, one has, as $1-\max \left\{H_{1}, H_{2}\right\}>0$,

$$
\exp \left(-\stackrel{\star}{C}\left(C_{\Psi, H_{1}, H_{2}}^{*}\right)^{-1} M^{1-\max \left\{H_{1}, H_{2}\right\}} y\right) \leqslant \frac{1}{2} \exp \left(-2 c_{1} y\right)
$$

which gives that, for all large enough $y$, one gets

$$
\mathbb{P}\left(\left|{\widetilde{\varepsilon_{j, k}}}^{M}\right|>y\right) \geqslant \exp \left(-c_{2} y\right)
$$

with $c_{2}:=2 c_{1}$. In the sequel, we will implicitly always consider such large enough $M$.
In the following two subsections, Lemmata 4.9 and 4.12 follow the lines of Lemmata 3.6 and 3.8 in [7] respectively, with some subtle modifications as the authors in [7] deal with $\mathcal{N}(0,1)$ random variables while, here, we focus on random variables in the second order Wiener chaos that depend on the parameter $M$. For the sake of completeness and clarity, we write the proofs in full details.

### 4.1 Ordinary points

In this section our aim is to prove the following proposition.
Proposition 4.8. There exists $\Omega_{1}^{*} \subset \Omega$ with probability 1 such that for all $\omega \in \Omega_{1}^{*}$ and Lebesgue almost every $t \in(0,1)$ one has

$$
\begin{equation*}
\limsup _{j \rightarrow+\infty} \frac{d_{j}(t, \omega)}{2^{-j\left(H_{1}+H_{2}-1\right)} \log j}>0 . \tag{4.5}
\end{equation*}
$$

To this end, as a first step, let us state the following lemma concerning the random variable $\widetilde{\varepsilon}_{\lambda}{ }^{M}$. If $\lambda=\lambda_{j, k}$ is a dyadic interval and $m \in \mathbb{N}, S_{\lambda, m}=S_{j, k, m}$ stands for the finite set of cardinality $2^{m}$ whose elements are the dyadic intervals of scale $j+m$ included in $\lambda_{j, k}$, formally speaking $S_{j, k, m}:=\left\{\lambda \in \Lambda_{j+m}: \lambda \subset \lambda_{j, k}\right\}$.
Lemma 4.9. There is a deterministic constant $C>0$ such that the following holds: for all $M \in \mathbb{N}$ and for all $t \in(0,1)$, there exists $\Omega_{t, 1} \subset \Omega$ with probability 1 such that for all $\omega \in \Omega_{t, 1}$ there are infinitely many $j \in \mathbb{N}$ such that

$$
\max _{\lambda^{\prime} \in S_{\lambda,\left[\log _{2}(N M)\right]+2}^{\lambda \in 3 \lambda_{j}(t)}}\left|{\widetilde{\varepsilon_{\lambda^{\prime}}}}^{M}(\omega)\right| \geqslant C \log j .
$$

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Proof. Let us fix $t \in(0,1)$ and $j \in \mathbb{N}$. For any $\lambda \in S_{j, k_{j}(t), m}$, there exists a unique decreasing finite sequence $\left(I_{n}\right)_{0 \leqslant n \leqslant m}$ of decreasing dyadic intervals in the sense of inclusion such that $I_{0}=\lambda_{j, k_{j}(t)}, I_{m}=\lambda$ and $I_{n} \in S_{j, k_{j}(t), n}$. Then, define the sequence $\left(T_{n}\right)_{1 \leqslant n \leqslant m}$ of unique dyadic intervals such that $I_{n-1}=I_{n} \cup T_{n}$. Note that for all $1 \leqslant n \leqslant m$, $T_{n} \in 3 I_{n}$. Moreover, as $\left(I_{n}\right)_{0 \leqslant n \leqslant m}$ is decreasing, $\left(T_{n}\right)_{1 \leqslant n \leqslant m}$ are pairwisely disjoint. Furthermore, for every $n \in\{1, \ldots, m\}$, there exist $T_{n}^{\prime}=\lambda_{j_{n}, k_{n}} \in S_{T_{n},\left\lfloor\log _{2} N M\right\rfloor+2}$ such that $\left(\frac{k_{n}-N M}{2^{j_{n}}}, \frac{k_{n}+N}{2^{j_{n}}}\right) \subset T_{n}$. As a consequence, the associated random variables $\left(\widetilde{\varepsilon}_{T_{n}^{\prime}}^{M}\right)_{1 \leqslant n \leqslant m}$ are independent as the dyadic intervals $\left(T_{n}^{\prime}\right)_{1 \leqslant n \leqslant m}$ satisfies condition $\left(C_{M}\right)$ in Definition 4.5. Next, for a constant $C>0$ to be chosen later, we set

$$
\mathcal{E}_{j, m}(t)=\left\{\omega \in \Omega: \max _{1 \leqslant n \leqslant m}\left|{\widetilde{\varepsilon_{T_{n}^{\prime}}}}^{M}\right| \geqslant C \log (2 m)\right\} .
$$

Note that, as the random variables $\left(\widetilde{\widetilde{T}_{n}^{\prime}}\right)_{1 \leqslant n \leqslant m}$ are independent,

$$
\mathbb{P}\left(\mathcal{E}_{j, m}(t)\right)=1-\prod_{n=1}^{m} \mathbb{P}\left(\left|{\widetilde{\varepsilon_{n}^{\prime}}}^{M}\right|<C \log (2 m)\right)
$$

Recalling (4.4), and the fact that $\log (1-x) \leqslant-x$ if $x \in(0,1)$, one gets, for $m$ is large enough,

$$
\begin{aligned}
\mathbb{P}\left(\mathcal{E}_{j, m}(t)\right) & \geqslant 1-\left(1-\exp \left(-C c_{2} \log (2 m)\right)^{m}\right. \\
& =1-\left(1-\left(\frac{1}{2 m}\right)^{C c_{2}}\right)^{m} \\
& \geqslant 1-\exp \left(\frac{m}{(2 m)^{C c_{2}}}\right) \\
& =1-\exp \left(\frac{m^{1-C c_{2}}}{2^{C c_{2}}}\right)
\end{aligned}
$$

Finally, choosing $C$ such that $0<C c_{2}<1$, one obtain that $\sum_{p \in \mathbb{N}} \mathbb{P}\left(\mathcal{E}_{2^{p}, 2^{p}}(t)\right)=+\infty$. Knowing that the events $\mathcal{E}_{2^{p}, 2^{p}}(t)$ are independent for all $p \in \mathbb{N}$, one concludes using Borel-Cantelli Lemma that

$$
\mathbb{P}\left(\limsup _{m \rightarrow+\infty} \mathcal{E}_{2^{m}, 2^{m}}(t)\right)=1
$$

It follows that for a fixed $t \in \mathbb{R}$, almost surely, there are infinitely many $j \in \mathbb{N}$ such that

$$
\max _{\substack{\lambda^{\prime} \in S_{\lambda,\left\lfloor\log _{2} N M\right\rfloor+2} \\ \lambda \in 3 \lambda_{j}(t)}}\left|{\widetilde{\varepsilon_{\lambda^{\prime}}}}^{M}(\omega)\right| \geqslant C \log j .
$$

Concerning the "non-independent part" of the wavelet coefficients, one can state the following Lemma.

Lemma 4.10. There is a deterministic constant $C^{\prime}>0$ such that, for all $M \in \mathbb{N}$ and for all $t \in(0,1)$, there exists $\Omega_{t, 2} \subset \Omega$ with probability 1 such that for all $\omega \in \Omega_{t, 2}$ there exists $J \in \mathbb{N}$ such that, for all $j \geqslant J$,

$$
\begin{aligned}
& \max _{\lambda^{\prime} \in S_{\lambda,\left\lfloor\log _{2}(N M)\right]+2}}\left|\widetilde{\varepsilon_{\lambda^{\prime}}}{ }^{M}(\omega)\right| \leqslant C^{\prime} M^{\max \left\{H_{1}, H_{2}\right\}-1} \log j . \\
& \lambda \in 3 \lambda_{j}(t)
\end{aligned}
$$

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Proof. Let us fix $t \in(0,1)$. For any $C^{\prime}>0$, for all $j$ sufficiently large and $\lambda \in 3 \lambda_{j}(t)$, we have, by Theorem 3.14,

$$
\begin{aligned}
& \mathbb{P}\left(\exists \lambda^{\prime} \in S_{\lambda,\left\lfloor\log _{2} N M\right\rfloor+2}:\left|\widetilde{\widetilde{\varepsilon}_{\lambda^{\prime}}}{ }^{M}\right| \geqslant C^{\prime} M^{\max \left\{H_{1}, H_{2}\right\}-1} \log j\right) \\
& \left.\leqslant \sum_{\lambda^{\prime} \in S_{\lambda, \mid \log _{2}}{ }_{N M\rfloor+2}} \mathbb{P}\left(\mid \widetilde{\varepsilon_{\lambda^{\prime}}}{ }^{M}\right) \mid \geqslant C^{\prime} M^{\max \left\{H_{1}, H_{2}\right\}-1} \log j\right) \\
& \leqslant \sum_{\lambda^{\prime} \in S_{\lambda,\left\lfloor\log _{2} N M\right\}+2}} \mathbb{P}\left(\left|\widetilde{\varepsilon_{\lambda^{\prime}}}{ }^{M}\right| \geqslant C^{\prime}\left(C_{\Psi, H_{1}, H_{2}}^{*}\right)^{-1}\left\|\widetilde{\varepsilon_{\lambda^{\prime}}}{ }^{M}\right\|_{L^{2}(\Omega)} \log j\right) \\
& \leqslant 4 N M \exp \left(-\stackrel{\star}{C} C^{\prime}\left(C_{\Psi, H_{1}, H_{2}}^{*}\right)^{-1} \log j\right)
\end{aligned}
$$

Thus, for $C^{\prime}>C_{\Psi, H_{1}, H_{2}}^{*} /{ }^{\star}$, the conclusion follows by Borel-Cantelli Lemma.
Proof of Proposition 4.8. The constant $C$ and $C^{\prime}$ of Lemmata 4.9 and 4.10 being deterministic and independent of $M$, on can choose $M$ large enough such that

$$
C-C^{\prime} M^{\max \left\{H_{1}, H_{2}\right\}-1}>0 .
$$

Let us fix $t \in(0,1)$ and consider $\omega \in \Omega_{t, 1} \cap \Omega_{t, 2}$, where the events, of probability $1, \Omega_{t, 1}$ and $\Omega_{t, 2}$ are given by the same Lemmata. For all $J \in \mathbb{N}$, by Lemma 4.9, there exist $j \geqslant J$ and $\lambda^{\prime}(j) \subseteq 3 \lambda_{j}(t)$ of scale $j^{\prime}=j+\lfloor\log N M\rfloor+2$ such that

$$
\left|\widetilde{c_{\lambda^{\prime}(j)}} M(\omega)\right| \geqslant C 2^{-j^{\prime}\left(H_{1}+H_{2}-1\right)} \log j .
$$

If $J$ is large enough, we also have, for all such $j \geqslant J$, by Lemma 4.10,

$$
\left|\widetilde{c_{\lambda^{\prime}(j)}} M(\omega)\right| \leqslant C^{\prime} M^{\max \left\{H_{1}, H_{2}\right\}-1} 2^{-j^{\prime}\left(H_{1}+H_{2}-1\right)} \log j .
$$

From this we deduce that

$$
\begin{aligned}
d_{j}(t, \omega) & \geqslant\left|c_{\lambda^{\prime}(j)}(\omega)\right| \\
& \geqslant\left|\widetilde{c_{\lambda^{\prime}(j)}} M(\omega)\right|-\left|\widetilde{c_{\lambda^{\prime}(j)}} M(\omega)\right| \\
& \geqslant 2^{-j^{\prime}\left(H_{1}+H_{2}-1\right)} \log j\left(C-C^{\prime} M^{\max \left\{H_{1}, H_{2}\right\}-1}\right) \\
& \geqslant 2^{-j\left(H_{1}+H_{2}-1\right)}(4 N M)^{1-H_{1}-H_{2}} \log j\left(C-C^{\prime} M^{\max \left\{H_{1}, H_{2}\right\}-1}\right)
\end{aligned}
$$

Therefore, (4.5) holds true for all $t \in(0,1)$ and $\omega \in \Omega_{t, 1} \cap \Omega_{t, 2}$. The conclusion follows then from Fubini Theorem.

### 4.2 Rapid points

In this section our aim is to prove the following proposition.
Proposition 4.11. There exists $\Omega_{2}^{*} \subset \Omega$ with probability 1 such that, for all $\omega \in \Omega_{2}^{*}$, there exist $t \in(0,1)$ such that

$$
\begin{equation*}
\limsup _{j \rightarrow+\infty} \frac{d_{j}(t, \omega)}{2^{-j\left(H_{1}+H_{2}-1\right)} j}>0 . \tag{4.6}
\end{equation*}
$$

As in the previous subsection, we start by working with the random variables $\widetilde{\varepsilon}_{\lambda}^{M}$.
Lemma 4.12. There exists a deterministic constant $C>0$ such that for all $M$ there is $\Omega_{2} \subset \Omega$ with probability 1 such that for all $\omega \in \Omega_{2}$ there exist $t \in(0,1)$ such that

$$
\begin{equation*}
\limsup _{j \rightarrow+\infty} \frac{\left|{\widetilde{\varepsilon_{\lambda_{j}}(t)}} M(\omega)\right|}{j} \geqslant C . \tag{4.7}
\end{equation*}
$$

Proof. Let us fix $a \in(0,1)$ and $C>0$ to be chosen later on. For every $(j, l) \in \mathbb{N} \times$ $\left\{0, \ldots,\left\lfloor 2^{j(1-a)}\right\rfloor-1\right\}$, we set

$$
S_{j, l}^{M}=\left\{l\left[2^{a j} /(2 N M)\right\rfloor, \ldots,(l+1)\left\lfloor 2^{a j} /(2 N M)\right\rfloor-1\right\}
$$

and consider the event

$$
\mathcal{E}_{j, l}^{M}=\left\{\omega \in \Omega: \max _{k \in S_{j, l}^{M}}\left|\widetilde{j, 2 k N M}^{M}(\omega)\right| \geqslant C j\right\}
$$

Let $j_{0}$ be the smallest integer such that $\left\lfloor 2^{a j} /(2 N M)\right\rfloor \geqslant 1$. If we assume that

$$
\begin{equation*}
\Omega_{2}^{*}=\bigcup_{J \geqslant j_{0}} \bigcap_{j \geqslant J} \bigcap_{l \in\left\{0, \ldots,\left[2^{j(1-a)} \mathrm{l}-1\right\}\right.} \mathcal{E}_{j, l}^{M} \tag{4.8}
\end{equation*}
$$

is an event of probability 1 and we consider $\omega \in \Omega_{2}^{*}$. For every $j \geqslant j_{0}$, denote by

$$
\begin{equation*}
G_{j}^{M}(\omega):=\left(k \in\left\{0, \ldots, 2^{j}-1\right\}:\left|{\widetilde{\varepsilon_{j, k}}}^{M}(\omega)\right| \geqslant C j\right) . \tag{4.9}
\end{equation*}
$$

Moreover, for every $n \geqslant j_{0}$, one considers

$$
\begin{equation*}
O_{n}^{M}(\omega):=\bigcup_{j \geqslant n} U_{j}^{M}(\omega), \quad \text { where } U_{j}^{M}(\omega):=\bigcup_{k \in G_{j}^{M}(\omega)}\left(\frac{k}{2^{j}}, \frac{k+1}{2^{j}}\right) \tag{4.10}
\end{equation*}
$$

If one proves that $O_{n}^{M}(\omega)$ is dense in $(0,1)$, then by Baire's theorem the set $\cap_{n \geqslant j_{0}} O_{n}^{M}(\omega)$ is non-empty and let t be an element of this set. Then for every $n \geqslant j_{0}$, there is $j \geqslant n$ such that $\left|\widetilde{\varepsilon_{\lambda_{j}(t)}}{ }^{M}(\omega)\right| \geqslant C j$, and so desired statement (4.7) is true.

We still have to prove two points:

1. $O_{n}^{M}(\omega)$ is dense in $(0,1)$.
2. $\Omega_{2}^{*}$ is an event of probability 1.

Indeed, starting with statement 1 , consider $t \in(0,1), j \geqslant j_{0}$ and $k$ such that $\lambda_{j}(t)=\lambda_{j, k}$. Then, we have two cases:

Case 1: There is $l \in\left\{0, \ldots,\left\lfloor 2^{j(1-a)}\right\rfloor-1\right\}$ such that

$$
k \in\left\{l\left\lfloor 2^{a j}\right\rfloor, \ldots,(l+1)\left\lfloor 2^{a j}\right\rfloor-1\right\}
$$

Using (4.8) and (4.9), there is $k^{\prime} \in\left\{l\left[2^{a j} /(2 N M)\right\rfloor, \ldots,(l+1)\left[2^{a j} /(2 N M)\right\rfloor-1\right\}$ such that $2 k^{\prime} N M \in G_{j}(\omega)$. Then, by (4.10), $\left(\frac{2 N M k^{\prime}}{2^{j}}, \frac{2 N M k^{\prime}+1}{2^{j}}\right) \subset O_{n}^{M}(\omega)$. which is at distance at most $2^{-j}\left(\left\lfloor 2^{a j}\right\rfloor+2 N M\left\lfloor 2^{a j} /(2 N M)\right\rfloor\right)$ from $t$. Finally, we get that $t$ is at a distance at most $22^{j(a-1)}$ of $U_{j}^{M}(\omega)$.
Case 2: $k \in\left\{\left[2^{j(1-a)}\right\rfloor\left[2^{j a}\right\rfloor, \ldots, 2^{j}-1\right\}$. Again by (4.8) and (4.9), there is $k^{\prime} \in S_{j, l}^{M}$ such that $2 k^{\prime} N M \in G_{j}^{M}(\omega)$, and similarly, we get that $t$ is at a distance at most $c 2^{j(a-1)}$ of $U_{j}^{M}(\omega)$, for some constant $c>0$ depending only on $N, M$ and $a$.

Finally, in both cases $t$ is at a distance at most $c 2^{j(a-1)}$, and so the density follows.
Now for statement 2, in order to prove that $\Omega_{2}^{*}$ has a probability 1 , it is enough to prove that

$$
\begin{equation*}
\mathbb{P}\left(\mathcal{C}\left(\bigcap_{l \in\left\{0, \ldots,\left[2^{j(1-a)}\right\rfloor-1\right\}} \mathcal{E}_{j, l}^{M}\right)\right) \tag{4.11}
\end{equation*}
$$

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is the general term of a convergent series, then the result follows by Borel-Cantelli Lemma. Note that the variables ${\widetilde{\varepsilon_{j, 2 N M}}}^{M}(\omega), k \in S_{j, l}^{M}$ and $l \in\left\{0, \ldots,\left\lfloor 2^{j(1-a)}\right\rfloor-1\right\}$, are independent because for every $k \neq k^{\prime},\left|2 N M k-2 N M k^{\prime}\right| \geqslant 2 N M$ and so $\lambda_{j, 2 M N k}^{M} \cap$ $\lambda_{j, 2 M N k^{\prime}}^{M}=\emptyset$. Consequently, one has

$$
\begin{align*}
& \mathbb{P}\left(\mathcal{C}\left(\bigcap_{\left.l \in\left\{0, \ldots, 2^{j(1-a)}\right\rfloor-1\right\}} \mathcal{E}_{j, l}^{M}\right)\right) \\
= & 1-\mathbb{P}\left(\bigcap_{\left.l \in\left\{0, \ldots, \mid 2^{j(1-a)}\right\rfloor-1\right\}} \mathcal{C}\left(\mathcal{E}_{j, l}^{M}\right)\right) \\
= & 1-\prod_{l \in\left\{0, \ldots,\left\lfloor 2^{j(1-a)}\right\rfloor-1\right\}}\left(1-\prod_{k \in S_{j, l}^{M}} \mathbb{P}\left(\left|{\widetilde{\varepsilon_{j, 2 N M}}}{ }^{M}(\omega)\right|<C j\right)\right) \\
= & 1-\left(1-(1-\mathbb{P}(|\varepsilon| \geqslant C j))^{\left[^{a j} /(2 N M)\right\rfloor}\right)^{\left\lfloor 2^{j(1-a)}\right\rfloor} \\
\leqslant & 1-\exp \left(2^{j(1-a)} \log \left(1-p_{j}\right)\right) \tag{4.12}
\end{align*}
$$

where $\varepsilon$ is a random variable belonging to the Wiener chaos of order 2 distributed according to the $\left(\tilde{\varepsilon_{\lambda}}\right)_{\lambda \in \Lambda}$ and $p_{j}=(1-\mathbb{P}(|\varepsilon| \geqslant C j))^{\left\lfloor 2^{a j} /(2 N M)\right\rfloor}$. Remark that $p_{j}$ is a positive term that tends to 0 as $j \rightarrow+\infty$. Indeed, using the fact that $\log (1-x) \leqslant-x$ if $x \in(0,1)$ together with (4.4), there exist $J \in \mathbb{N}$ such that for all $j \geqslant J$,

$$
\begin{align*}
0 \leqslant p_{j} & \leqslant\left(1-\exp \left(-C c_{2} j\right)\right)^{\left.2^{a j} /(2 N M)\right\rfloor} \\
& \leqslant \exp \left(-\left\lfloor\frac{2^{a j}}{2 N M}\right\rfloor \exp \left(-C c_{2} j\right)\right) \\
& \leqslant \exp \left(-C^{\prime} \exp \left(\log 2^{a j}\right) \exp \left(-C c_{2} j\right)\right) \\
& \leqslant \exp \left(-C^{\prime} \exp j\left(a \log 2-C c_{2}\right)\right) \tag{4.13}
\end{align*}
$$

where $C^{\prime}$ depends only on $N, M$ and $a$ and $c_{2}$ is the constant given in (4.4). It is enough to choose $C$ such that $a \log 2-C c_{2}>0$ to deduce that and so $p_{j} \rightarrow 0$ as $j \rightarrow+\infty$. Similarly, one can get for all $j \geqslant J$

$$
0 \leqslant 2^{j(1-a)} p_{j} \leqslant \exp \left(-C^{\prime} \exp j\left(\log 2-C c_{2}\right)\right)
$$

which indeed shows that $2^{j(1-a)} p_{j}$ tends to 0 as $j \rightarrow+\infty$. Now, using the fact that $\log (1-x)=-x+o(x)$ and $\exp (x)=1+x+o(x)$ as $x \rightarrow 0$, together with (4.12) we obtain that for all $\delta>0$

$$
\mathbb{P}\left(\mathcal{C}\left(\bigcap_{\left.l \in\left\{0, \ldots, \mid 2^{j(1-a)}\right\rfloor-1\right\}} \mathcal{E}_{j, l}^{M}\right)\right) \leqslant 2^{j(1-a)}\left(\delta\left(p_{j}+\delta p_{j}\right)+p_{j}+\delta p_{j}\right)
$$

for $j$ large enough. Using the upper bound in (4.13), one can finally conclude that (4.11) is indeed the general term of a convergent series.

Concerning the random variable $\breve{\varepsilon}_{\lambda}{ }^{M}$, one can give an almost sure upper bound.
Lemma 4.13. There exists a deterministic constant $C^{\prime}>0$ such that for all $M$ there is $\Omega_{2}^{\prime} \subset \Omega$ with probability 1 such that for all $\omega \in \Omega_{2}^{\prime}$ there exist $J \in \mathbb{N}$ such that, for all $j \geqslant J$, for all $\lambda \in \Lambda_{j}, \lambda \subseteq[0,1]$,

$$
\left|{\widetilde{\varepsilon_{\lambda}}}^{M}(\omega)\right| \leqslant C^{\prime} M^{\max \left\{H_{1}, H_{2}\right\}-1} j
$$

Proof. If $C^{\prime}>0$, for all $j$ sufficiently large, we have, by Theorem 3.14

$$
\begin{aligned}
\mathbb{P} & \left(\exists \lambda \in \Lambda_{j}, \lambda \subseteq[0,1]:\left|\check{\varepsilon \lambda}_{\lambda}^{M}(\omega)\right| \geqslant C^{\prime} M^{\max \left\{H_{1}, H_{2}\right\}-1} j\right) \\
& \leqslant \sum_{\lambda \in \Lambda_{j} \lambda \subseteq[0,1]} \mathbb{P}\left(\left|\breve{\varepsilon}_{\lambda}^{M}(\omega)\right| \geqslant C^{\prime} M^{\max \left\{H_{1}, H_{2}\right\}-1} j\right) \\
& \leqslant 2^{j} \exp \left(-\stackrel{\star}{C} C^{\prime}\left(C_{\Psi, H_{1}, H_{2}}^{*}\right)^{-1} j\right)
\end{aligned}
$$

and thus, if $C^{\prime}>\log (2) C_{\Psi, H_{1}, H_{2}}^{*} / \stackrel{\star}{C}$, the conclusion follows by Borel-Cantelli Lemma.
Proof of Proposition 4.11. Again, one can choose $M$ large enough such that

$$
C-C^{\prime} M^{\max \left\{H_{1}, H_{2}\right\}-1}>0,
$$

where $C$ and $C^{\prime}$ are the constant given by Lemmata 4.12 and 4.13 respectively. Let us consider $\omega \in \Omega_{2}^{*}:=\Omega_{2} \cap \Omega_{2}^{\prime}$ where the events, of probability $1, \Omega_{2}$ and $\Omega_{2}^{\prime}$ are giving by the same Lemmata. We use the same notations as in them. First there exist $t \in(0,1)$ such that for all $J \in \mathbb{N}$ there exist $j \geqslant n$ such that

$$
\begin{equation*}
\left|\widetilde{c_{\lambda_{j}(t)}}{ }^{M}(\omega)\right| \geqslant C j 2^{-j\left(H_{1}+H_{2}-1\right)} . \tag{4.14}
\end{equation*}
$$

Moreover, if $J$ is large enough, for all such $j$ we also have

$$
\begin{equation*}
\left|{\widetilde{c_{\lambda_{j}}(t)}}^{M}(\omega)\right| \leqslant C^{\prime} M^{\max \left\{H_{1}, H_{2}\right\}-1} 2^{-j\left(H_{1}+H_{2}-1\right)} j . \tag{4.15}
\end{equation*}
$$

In this case, as in 4.8 we have that for all $J$ great enough, there is $j \geqslant J$ such that

$$
d_{j}(t, \omega) \geqslant 2^{-j\left(H_{1}+H_{2}-1\right)} j\left(C-C^{\prime} M^{\max \left\{H_{1}, H_{2}\right\}-1}\right)
$$

and so one can conclude that (4.6) holds true for all $\omega \in \Omega_{2}^{*}$.

## 5 Proof of the main theorem

Theorem 1.2 is then a straightforward consequence of Propositions 3.5, 3.21, 3.31, 4.8 and 4.11 .

Proof of Theorem 1.2. If $\lambda$ is a dyadic interval, taking into account Remark 3.4, let us denote by $\Omega_{\lambda}$ the event obtained by taking the intersection of all the events of probability 1 induced by Propositions 3.5, 3.21, 3.31, 4.8 and 4.11. Then $\Omega_{R}=\bigcap_{\lambda} \Omega_{\lambda}$ is also an event of probability 1 , as it is a countable intersection of events of probability 1.

If we consider $\omega \in \Omega_{R}$, for all dyadic interval $\lambda$, first, from Proposition 3.5, for all $t \in \lambda$

$$
\begin{equation*}
\limsup _{s \rightarrow t} \frac{\left|R_{H_{1}, H_{2}}(t, \omega)-R_{H_{1}, H_{2}}(s, \omega)\right|}{|t-s|^{H_{1}+H_{2}-1} \log |t-s|^{-1}}<+\infty \tag{5.1}
\end{equation*}
$$

while, for almost every $t_{o} \in \lambda$, from Propositions 3.21 and 4.8,

$$
0<\limsup _{s \rightarrow t_{o}} \frac{\left|R_{H_{1}, H_{2}}\left(t_{o}, \omega\right)-R_{H_{1}, H_{2}}(s, \omega)\right|}{\left|t_{o}-s\right|^{H_{1}+H_{2}-1} \log \log \left|t_{o}-s\right|^{-1}}<+\infty
$$

Nevertheless, from Proposition 4.11 we also know that there exists $t_{r} \in \lambda$ such that

$$
0<\limsup _{s \rightarrow t_{r}} \frac{\left|R_{H_{1}, H_{2}}\left(t_{r}, \omega\right)-R_{H_{1}, H_{2}}(s, \omega)\right|}{\left|t_{r}-s\right|^{H_{1}+H_{2}-1} \log \left|t_{r}-s\right|^{-1}}
$$

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which, combined with (5.1), gives that, for all such a $t_{r}$,

$$
0<\limsup _{s \rightarrow t_{r}} \frac{\left|R_{H_{1}, H_{2}}\left(t_{r}, \omega\right)-R_{H_{1}, H_{2}}(s, \omega)\right|}{\left|t_{r}-s\right|^{H_{1}+H_{2}-1} \log \left|t_{r}-s\right|^{-1}}<+\infty .
$$

Moreover, from Proposition 3.31, we also know that one can find $t_{\sigma} \in \lambda$ such that

$$
\limsup _{s \rightarrow t_{\sigma}} \frac{\left|R_{H_{1}, H_{2}}\left(t_{\sigma}, \omega\right)-R_{H_{1}, H_{2}}(s, \omega)\right|}{\left|t_{\sigma}-s\right|^{H_{1}+H_{2}-1}}<+\infty .
$$

Thus, we have shown that, on $\Omega_{R}$, for all dyadic interval $\lambda$, almost every points in $\lambda$ are ordinary while there exist rapid and slow points in $\lambda$. Now, if $I$ is a non-empty interval of $\mathbb{R}$, writing the interior of $I$ as a countable union of dyadic intervals, we show that, on $\Omega_{R}$, almost every points in $\lambda$ are ordinary. Finally, if $U$ is an open set with $I \cap U \neq \emptyset$, there is a dyadic interval $\lambda \subseteq I \cap U$. Thus, on $\Omega_{R}, I \cap U$ contains rapid and slow points and the density of these points in $I$ is proved.

Remark 5.1. Unfortunately, our method does not allow us to affirm the positiveness of the limit (1.8), at the opposite of limits (1.6) and (1.7). Indeed, as for almost every $\omega \in \Omega$

$$
\limsup _{s \rightarrow t} \frac{\left|R_{H_{1}, H_{2}}(t, \omega)-R_{H_{1}, H_{2}}(s, \omega)\right|}{|t-s|^{H_{1}+H_{2}-1}}
$$

is finite for some $t$, we would need to show its positiveness for all $t$ and thus the positiveness of the limit

$$
\begin{equation*}
\limsup _{j \rightarrow+\infty} \frac{d_{j}(t, \omega)}{2^{-j\left(H_{1}+H_{2}-1\right)}} \tag{5.2}
\end{equation*}
$$

for all $t$.
Concerning the random variables $\left({\widetilde{\varepsilon_{\lambda}}}^{M}\right)_{\lambda}$, one can obtain a positive result ${ }^{6}$. Indeed, from [24, Theorem 6.9 and Remark 6.10] we know that there exists an universal deterministic constant $\gamma \in[0,1)$ such that, for each random variable $X$ in the Wiener chaos of order 2

$$
\mathbb{P}\left(|X| \leqslant \frac{1}{2}\|X\|_{L^{2}(\Omega)}\right) \leqslant \gamma
$$

As $0 \leqslant \gamma<1$, of course, one can find $\ell_{0} \in \mathbb{N}$ such that

$$
\begin{equation*}
\gamma^{\ell_{0}}<2^{-1} \tag{5.3}
\end{equation*}
$$

Let us go back to the construction starting the proof of Lemma 4.9. If the dyadic interval $\lambda_{j, k}$ and $m \in \mathbb{N}$ are fixed and $S \in \mathcal{S}_{j, k, m}$ we define the sequences of dyadic intervals $\left(I_{n}\right)_{0 \leqslant n \leqslant m}$ and $\left(T_{n}\right)_{1 \leqslant n \leqslant m}$ in the same way: $I_{0}=\lambda_{j, k}, I_{m}=S$ and, for all $1 \leqslant n \leqslant m, I_{n-1}=I_{n} \cup T_{n}$. Now, for any $1 \leqslant n \leqslant m$, there are $\ell_{0}$ dyadic intervals $\left(T_{n}^{\ell}=\lambda_{j_{n}^{(\ell)}, k_{n}^{(\ell)}}\right)_{1 \leqslant \ell \leqslant \ell_{0}}$ in $S_{\left.T_{n}, \log _{2}\left(\ell_{0} N M\right)\right\rfloor+2}$ such that, for all $1 \leqslant \ell \leqslant \ell_{0}$

$$
\left(\frac{k_{n}^{(\ell)}-N M}{2^{j_{n}^{(\ell)}}}, \frac{k_{n}^{(\ell)}+N}{2^{j_{n}^{(\ell)}}}\right) \subseteq T_{n}
$$

and, if $\ell \neq \ell^{\prime}, T_{n}^{\ell} \cap T_{n}^{\ell^{\prime}}=\emptyset$. Therefore, the dyadic intervals $\left(T_{n}^{\ell}\right)_{1 \leqslant n \leqslant m, 1 \leqslant \ell \leqslant \ell_{0}}$ satisfy condition $\left(C_{M}\right)$ in Definition 4.5. From this, for all $S \in \mathcal{S}_{j, k, m}$ we define the Bernouilli random variable

$$
\mathcal{B}_{j, k, m}(S)=\prod_{1 \leqslant n \leqslant m, 1 \leqslant \ell \leqslant \ell_{0}} \mathbf{1}_{\left\{\left|{\widetilde{\varepsilon_{n}^{\ell}}}^{M}\right|<2^{-1} C_{\Psi, H_{1}, H_{2}}\right\}}
$$

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for which, by Proposition 4.6, we have, using the independence of the random variables $\left(\widetilde{\varepsilon_{n}^{\ell}}{ }^{M}\right)_{1 \leqslant n \leqslant m, 1 \leqslant \ell \leqslant \ell_{0}}, \mathbb{E}\left[\mathcal{B}_{j, k, m}(S)\right] \leqslant \gamma^{m \ell_{0}}$. Therefore, if we define the random variable

$$
\mathcal{G}_{j, k, m}=\sum_{S \in \mathcal{S}_{j, k, m}} \mathcal{B}_{j, k, m}(S)
$$

then $\mathbb{E}\left[\mathcal{G}_{j, k, m}\right] \leqslant\left(2 \gamma^{\ell_{0}}\right)^{m}$ and it follows from inequality (5.3) and Fatou Lemma that

$$
\mathbb{E}\left[\liminf _{m \rightarrow+\infty} \mathcal{G}_{j, k, m}\right]=0
$$

As a consequence, $\Omega_{1}=\bigcap_{j \in \mathbb{N}, 0 \leqslant k<2^{j}}\left\{\omega: \liminf _{m \rightarrow+\infty} \mathcal{G}_{j, k, m}(\omega)=0\right\}$ is an event of probability 1.

Now if $\omega \in \Omega_{1}$ and $t \in(0,1)$, we take $j \in \mathbb{N}$ and $k=k_{j}(t)$ and since, for all $m, \mathcal{G}_{j, k_{j}(t), m}$ has values in $\left\{0, \ldots, 2^{m}\right\}$ we conclude that there are infinitely many $m$ for which, for every $S \in \mathcal{S}_{j, k_{j}(t), m}, \mathcal{B}_{j, k, m}(S)=0$. Considering such a $m$ and $S=\lambda_{j+m}(t)$ then we first remark that, for all $1 \leqslant n \leqslant m, I_{n}=\lambda_{j+n}(t)$ and thus $T_{n} \in 3 \lambda_{j+n}(t)$. Now, as $\mathcal{B}_{j, k, m}\left(\lambda_{j+m}(t)\right)=0$, one can find $1 \leqslant n \leqslant m$ and $1 \leqslant \ell \leqslant \ell_{0}$ such that

$$
\left|{\widetilde{\varepsilon_{n}^{\ell}}}^{M}(\omega)\right| \geqslant 2^{-1} C_{\Psi, H_{1}, H_{2}} .
$$

Thus we have showed that, for all $\omega \in \Omega_{1}$ and $t \in(0,1)$ there exist infinitely many $j^{\prime} \in \mathbb{N}$ such that

$$
\max _{\substack{\lambda^{\prime} \in S_{\lambda,\left[\log _{2}\left(\ell_{0} N M\right)\right]+2} \\ \lambda \in 3 \lambda_{j^{\prime}}(t)}}\left|{\widetilde{\lambda^{\prime}}}^{M}(\omega)\right| \geqslant 2^{-1} C_{\Psi, H_{1}, H_{2}} .
$$

To pass to the wavelet leaders, in the spirit of Propositions 4.8 and 4.11 , we would need to get from Borel-Cantelli Lemma an upper bound of

$$
\max _{\lambda^{\prime} \in S_{\lambda,\left[\log _{2}\left(\ell_{0} N M\right)\right\rfloor+2}\left|\widetilde{\varepsilon_{\lambda^{\prime}}}{ }^{M}(\omega)\right|}^{\lambda \in 3 \lambda_{j^{\prime}}(t)} \mid
$$

for all $j$ sufficiently large on an event of probability 1 which does not depend on $t$. Then, as

$$
\left.\begin{array}{l}
\mathbb{P}\left(\exists \lambda \in \Lambda_{j}, \lambda \subseteq[0,1]: \max _{\substack{\prime \prime} S_{\left.\lambda^{\prime}, \log _{2}\left(\ell_{0} N M\right)\right]+2} \mid \widetilde{\varepsilon}_{\lambda^{\prime}} \in 3 \lambda}{ }^{M}(\omega) \mid \geqslant C^{\prime} M^{\max \left\{H_{1}, H_{2}\right\}-1} j\right) \\
\end{array}\right)
$$

if $C^{\prime}>\log (2) C_{\Psi, H_{1}, H_{2}}^{*} / \stackrel{\star}{C}$ this probability is the general term of some convergent series and in this case one can affirm the existence of an event $\Omega_{1}^{\prime}$ of probability 1 such that, for all $\omega \in \Omega_{1}^{\prime}$ there exist $J \in \mathbb{N}$ such that, for all $j \geqslant J$, for all $\lambda \in \Lambda_{j}, \lambda \subseteq[0,1]$,

$$
\max _{\lambda^{\prime \prime} \in S_{\substack{\lambda^{\prime},\left[\log _{2}\left(\ell_{0} N M\right)\right]+2 \\ \lambda^{\prime} \in 3 \lambda}}\left|\widetilde{\varepsilon_{\lambda^{\prime}} M}(\omega)\right| \leqslant C^{\prime} M^{\max \left\{H_{1}, H_{2}\right\}-1} j .}
$$

It seems to be the sharper upper bound that we can hope to find with our constraints and the fact that we don't have any independence property to take advantage of when

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dealing with the random variables $\breve{\varepsilon}_{\lambda}{ }^{M}$. This is insufficient to consider properly limit (5.2). Nevertheless, if, instead of working with an uniform constant $M$ we make it depends on the scale $j$ by setting $M_{j}=\left(4 C^{\prime} C_{\Psi, H_{1}, H_{2}}^{-1} j\right)^{\frac{1}{1-\max \left\{H_{1}, H_{2}\right\}}}$, where $C^{\prime}>\log (2) C_{\Psi, H_{1}, H_{2}}^{*} / \stackrel{\star}{C}$ is the same constant as in Lemma 4.13,

$$
\begin{gathered}
\left.\left.\lambda_{j, k}^{M_{j}}:=\right] \frac{k-N M_{j}}{2^{j}}, \frac{k+N}{2^{j}}\right]^{2} \\
\widetilde{c_{j, k}} M_{j}=c_{H_{1}, H_{2}} \int_{\lambda_{j, k} M_{j}}^{\prime} \int_{-N}^{N} \Psi(x) \int_{\frac{k}{2^{j}}}^{\frac{x+k}{2 j}} f_{H_{1}, H_{2}}\left(s, x_{1}, x_{2}\right) d s d x d B\left(x_{1}\right) d B\left(x_{2}\right)
\end{gathered}
$$

and

$$
{\widetilde{c_{j, k}}}^{M_{j}}=c_{j, k}-{\widetilde{c_{j, k}}}^{M_{j}}
$$

then Proposition 4.6 stills holds if we replace $M$ by $M_{j}$ with $j$ sufficiently large and, by directly adapting what precedes one can find on event $\Omega_{1}^{*}$ of probability 1 such that, for all $\omega \in \Omega_{1}^{*}$ and $t \in(0,1)$ there exist infinitely many $j \in \mathbb{N}$ such that ${ }^{7}$

$$
\max _{\substack{\lambda^{\prime} \in S_{\lambda,\left[\log _{2}\left(l_{0} N M_{j}\right)\right]+2}^{\lambda \in 3 \lambda_{j}(t)}}}\left|{\widetilde{\varepsilon_{\lambda^{\prime}}}}^{M_{j}}(\omega)\right| \geqslant 2^{-1} C_{\Psi, H_{1}, H_{2}} .
$$

while there exist $J \in \mathbb{N}$ such that, for all $j \geqslant J$, for all $\lambda \in \Lambda_{j}, \lambda \subseteq[0,1]$,

$$
\begin{aligned}
\max _{\substack{\prime \prime} S_{\lambda^{\prime},\left[\log _{2}\left(l_{0} N M_{j}\right)\right]+2}^{\lambda^{\prime} \in 3 \lambda}}\left|\widetilde{\varepsilon_{\lambda^{\prime}}} M_{j}(\omega)\right| & \leqslant C^{\prime}\left(M_{j}\right)^{\max \left\{H_{1}, H_{2}\right\}-1} j \\
& \leqslant 4^{-1} C_{\Psi, H_{1}, H_{2}} .
\end{aligned}
$$

As a consequence, as in Proposition 4.8, for all $J \in \mathbb{N}$ there exist $j \geqslant J$ with

$$
d_{j}(t, \omega) \geqslant 2^{-j\left(H_{1}+H_{2}-1\right)}\left(4 C^{\prime} C_{\Psi, H_{1}, H_{2}}^{-1} j\right)^{\frac{1-H_{1}-H_{2}}{1-\max \left\{H_{1}, H_{2}\right\}}}\left(4 \ell_{0} N\right)^{1-H_{1}-H_{2}} 4^{-1} C_{\Psi, H_{1}, H_{2}}
$$

which allows to state that, for all $t \in(0,1)$ and $\omega \in \Omega_{1}$,

$$
\limsup _{j \rightarrow+\infty} \frac{d_{j}(t, \omega)}{2^{-j\left(H_{1}+H_{2}-1\right)} j^{\frac{1-H_{1}-H_{2}}{1-\max \left\{H_{1}, H_{2}\right\}}}}>0
$$

and thus, for all $\omega \in \Omega_{1}$ and for all $t \in(0,1)$,

$$
\limsup _{s \rightarrow t} \frac{\left|R_{H_{1}, H_{2}}(t, \omega)-R_{H_{1}, H_{2}}(s, \omega)\right|}{|t-s|^{H_{1}+H_{2}-1}\left(\log |t-s|^{-1}\right)^{\frac{1-H_{1}-H_{2}}{1-\max \left\{H_{1}, H_{2}\right\}}}}>0 .
$$

In particular, we find an almost sure uniform lower modulus of continuity for the generalized Rosenblatt process, similar to the one established in [26] for the Rosenblatt process. However, we are not able to judge the optimality of this modulus, which seems to be a difficult problem, as already stated in [5, Remark 1.2].

An interesting corollary of Remark 5.1 and Proposition 3.5 is the fact that, almost surely, the pointwise Hölder exponent of the generalized Rosenblatt process is everywhere $H_{1}+H_{2}-1$ and, in particular, it is nowhere differentiable.

Similarly, one can also take $\left(M_{j}=\left(4 C^{\prime} C_{\Psi, H_{1}, H_{2}}^{-1} \log (j)^{\frac{1}{1-\max \left\{H_{1}, H_{2}\right\}}}\right)_{j}\right.$, where $C^{\prime}$ is this time the same constant that in Lemma 4.10 and show, precisely like in this Lemma,

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that there exists a deterministic constant $C^{\prime}>0$ such that, for all $t \in(0,1)$ there exists $\Omega_{t, 2} \subset \Omega$ with probability 1 such that for all $\omega \in \Omega_{t, 2}$ there exist $J \in \mathbb{N}$ such that, for all $j \geqslant J$,

$$
\max _{\substack{\lambda^{\prime} \in S_{\lambda,\left[\log _{2}\left(\ell_{0} N M_{j}\right)\right]+2} \\ \lambda \in 3 \lambda_{j}(t)}}\left|{\widetilde{\varepsilon_{\lambda^{\prime}}}}^{M_{j}}(\omega)\right| \leqslant 4^{-1} C_{\Psi, H_{1}, H_{2}}
$$

and conclude in the same way that there exists an event of probability 1 such that, for all $\omega$ in this event and for almost every $t \in(0,1)$

$$
\limsup _{s \rightarrow t} \frac{\left|R_{H_{1}, H_{2}}(t, \omega)-R_{H_{1}, H_{2}}(s, \omega)\right|}{|t-s|^{H_{1}+H_{2}-1}\left(\log \log |t-s|^{-1}\right)^{\frac{1-H_{1}-H_{2}}{1-\max \left\{H_{1}, H_{2}\right\}}}}>0
$$

## References

[1] J. M. P. Albin. A note on Rosenblatt distributions. Statist. Probab. Lett., 40(1):83-91, 1998. MR1650532
[2] A. Arneodo, Y. d'Aubenton Carafa, E. Bacry, P.V. Graves, J.F. Muzy, and C. Thermes. Wavelet based fractal analysis of DNA sequences. Physica D.: Nonlinear Phenomena, (96):291-320, 1996. MR1368691
[3] A. Arneodo, J.-F. Muzy, and E. Bacry. Wavelets and multifractal formalism for singular signals: application to turbulence data. Physical Review Letters, (67):3515-3518, 1991.
[4] A. Ayache. Multifractional stochastic fields. World Scientific Publishing Co. Pte. Ltd., Hackensack, NJ, 2019. Wavelet strategies in multifractional frameworks. MR3839281
[5] A. Ayache. Lower bound for local oscillations of Hermite processes. Stoch. Process. Their Appl., 130:4593-4607, 2020. MR4108464
[6] A. Ayache and Y. Esmili. Wavelet-type expansion of the generalized Rosenblatt process and its rate of convergence. J. Fourier Anal. Appl., 26(3):Paper No. 51, 35, 2020. MR4110623
[7] A. Ayache, C. Esser, and T. Kleyntssens. Different possible behaviors of wavelet leaders of the Brownian motion. Statist. Probab. Lett., 150:54-60, 2019. MR3922488
[8] J.-M. Bardet and C. A. Tudor. A wavelet analysis of the Rosenblatt process: chaos expansion and estimation of the self-similarity parameter. Stochastic Process. Appl., 120(12):2331-2362, 2010. MR2728168
[9] J. Barral and S. Seuret. From multifractal measures to multifractal wavelet series. J. Fourier Anal Appl., (11):589-614, 2005. MR2182637
[10] F. Bastin, C. Esser, and S. Jaffard. Large deviation spectra based on wavelet leaders. Rev. Mat. Iberoam., 32(3):859-890, 2016. MR3556054
[11] A. Bunde, J. Kropp, and H.J. Schellnhuber. The science of disasters: climate disruptions, heart attacks, and market crashes, volume 2. Springer-Verlag, 2002.
[12] Amit Chaurasia. Performance of synthetic Rosenblatt process under multicore architecture. In 2019 3rd International conference on Electronics, Communication and Aerospace Technology (ICECA), pages 377-381. IEEE, 2019.
[13] M. Clausel and S. Nicolay. Wavelet techniques for pointwise anti-Hölderian irregularity. Constr. Approx., 33:41-75, 2011. MR2747056
[14] I. Daubechies. Orthonormal bases of compactly supported wavelets. Comm. Pure App. Math., 41:909-996, 1988. MR0951745
[15] I. Daubechies. Ten lectures on wavelets. CBMS-NSF Regional Conference Series in Applied Mathematics, 1992. MR1162107
[16] A. Deliège, T. Kleyntssens, and S. Nicolay. Mars topography investigated through the wavelet leaders method: a multidimensional study of its fractal structure. Planet. Space Sci., 136:4658, 2017.

Wavelet methods to study the pointwise regularity of the generalized Rosenblatt process
[17] R.L. Dobrushin and P. Major. Non-central limit theorems for nonlinear functionals of gaussian fields. Z. Wahrsch. Verw. Gebiete, (50):27-52, 1979. MR0550122
[18] C. Esser and L. Loosveldt. Slow, ordinary and rapid points for Gaussian Wavelets Series and application to Fractional Brownian Motions. ALEA, Lat. Am. J. Probab. Math. Stat., 19:1471-1495, 2022.
[19] P. Flandrin. Time-frequency / Time-scale analysis. In Wavelet Analysis and its Applications. Academic Press, 1999. MR1681043
[20] S. Jaffard. Multifractal formalism for functions part I: Results valid for all functions. SIAM J. Math. Anal., (28):944-970, 1997. MR1453315
[21] S. Jaffard. Wavelet techniques in multifractal analysis, fractal geometry and applications: A jubilee of Benoit Mandelbrot. Proceedings of Symposia in Pure Mathematics, 72:91-151, 2004. MR2112122
[22] S. Jaffard, P. Abry, and S. Roux. Function spaces vs. Scaling functions: Some issues inimage classification. Mathematical Image processing, pages 1-40, 2011. MR2867518
[23] S. Jaffard and Y. Meyer. Wavelet methods for pointwise regularity and local oscillations of functions. Mem. Amer. Math. Soc., 123(587):x+110, 1996. MR1342019
[24] S. Janson. Gaussian Hilbert Spaces. Cambridge Tracts in Mathematics. Cambridge University Press, 1997. MR1474726
[25] J.-P. Kahane. Some random series of functions, volume 5 of Cambridge Studies in Advanced Mathematics. Cambridge University Press, Cambridge, second edition, 1985. MR0833073
[26] G. Kerchev, I. Nourdin, E. Saksman, and L. Viitasaari. Local times and sample path properties of the Rosenblatt process. Stoch. Process. Their Appl., 131:498-522, 2021. MR4165649
[27] D. Kreit and S. Nicolay. Generalized pointwise Hölder spaces definitionined via admissible sequences. J. Funct. Spaces, ID 8276258:11, 2018. MR3812831
[28] E. H. Lakhel and A. Tlidi. Existence, uniqueness and stability of impulsive stochastic neutral functional differential equations driven by rosenblatt process with varying-time delays. Random Operators and Stochastic Equations, 27(4):213-223, 2019. MR4036646
[29] P.G. Lemarié and Y. Meyer. Ondelettes et bases hilbertiennes. Rev. Mat. Iberoam., 2:1-18, 1986. MR0864650
[30] P. Lévy. Processus stochastiques et mouvement brownien. Gauthier-Villars, Paris, 1948. MR0190953
[31] L. Loosveldt and S. Nicolay. Generalized spaces of pointwise regularity: Toward a general framework for the WLM. Nonlinearity, 34:6561-6586, 2021. MR4304490
[32] M. Maejima and C.A. Tudor. Selfsimilar processes with stationnary increments in the second Wiener chaos. Probab. Math. Stat., 32(1):167-186, 2012. MR2959876
[33] M. Maejima and C.A. Tudor. On the distribution of the Rosenblatt process. Statist. Probab. Lett., 83(6):1490-1495, 2013. MR3048314
[34] S. Mallat. A Wavelet Tour of Signal Processing. Academic Press, 1999. MR1614527
[35] Y. Meyer and D. Salinger. Wavelets and operators, volume 1. Cambridge university press, 1995. MR1228209
[36] Y. Meyer, F. Sellan, and M. S. Taqqu. Wavelets, generalized white noise and fractional integration: the synthesis of fractional Brownian motion. J. Fourier Anal. Appl., 5(5):465-494, 1999. MR1755100
[37] S. Nicolay, M. Touchon, B. Audit, Y. d'Aubenton Carafa, C. Thermes, A. Arneodo, et al. Bifractality of human DNA strand-asymmetry profiles results from transcription. Phys. Rev. E, 75:032902, 2007.
[38] Bernt Oksendal. Stochastic differential equations: an introduction with applications. Springer Science \& Business Media, 2013. MR2001996
[39] M. Rosenblatt. A central limit theorem and a strong mixing condition. Proc. Nat. Acad. Sci. U.S.A., 42:43-47, 1956. MR0074711
[40] R. Sakthivel, P. Revathi, Y. Ren, and G. Shen. Retarded stochastic differential equations with infinite delay driven by rosenblatt process. Stochastic analysis and applications, 36(2):304323, 2018. MR3750683

Wavelet methods to study the pointwise regularity of the generalized Rosenblatt process
[41] L. Schwartz. Théorie des distributions. Hermann, 1978. MR0209834
[42] M.S. Taqqu. The rosenblatt process. In Richard A. Davis, Keh-Shin Lii, and Dimitris N. Politis, editors, Selected Works of Murray Rosenblatt, pages 29-45. Springer New York, New York, NY, 2011. MR2742596
[43] C.A. Tudor and F.G. Viens. Variations and estimators for self-similarity parameters via Malliavin calculus. Ann. Probab., 37(6):2093-2134, 2009. MR2573552
[44] M.S.. Veillette and M. S. Taqqu. Properties and numerical evaluation of the Rosenblatt distribution. Bernoulli, 19(3):982-1005, 2013. MR3079303
[45] H. Wendt, P. Abry, S. Jaffard, H. Ji, and Z. Shen. Wavelet leader multifractal analysis for texture classification. ICIP, pages 3829-3832, 2009.
[46] H. Wendt, S. Roux, S. Jaffard, and P. Abry. Wavelet leaders and bootstrap for multifractal analysis of images. Signal Processing, (89(6)):1100-1114, 2009.
[47] N. Wiener. Collected works, volume 1. The MIT press, 1976. MR0532698
[48] N. Wiener and R.C. Paley. Fourier transforms in the complex domain. Amer. Math. Soc. Colloq. Pub., 19, 1934. 183pp. MR1451142

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[^1]:    ${ }^{1}$ We recall that $R_{H_{1}, H_{2}}^{\prime}$ is defined in Remark 3.20.

[^2]:    ${ }^{2}$ The interval $\left[k 2^{-j},(k+1) 2^{-j}\right]$ is cut into $\left[(2 k) 2^{-(j+1)},(2 k+1) 2^{-(j+1)}\right]$ and $\left[(2 k+1) 2^{-(j+1)},(2 k+\right.$ 2) $2^{-(j+1)}$ ].

[^3]:    ${ }^{3}$ In [5], Ayache does not consider wavelets at all but directly work on Wiener-Itô integrals.

[^4]:    ${ }^{4}$ The function $\Phi$ is in the fact the symmetrization of $\Phi_{1}$.
    ${ }^{5}$ For $f$ a symmetric function in $L^{2}\left(\mathbb{R}^{2}\right)$, and $I_{2}(f)$ the second order Wiener-Itô integral of $f$, one has $\mathbb{E}\left(I_{2}(f)\right)^{2}=2!\|f\|_{L^{2}\left(\mathbb{R}^{2}\right)}$.

[^5]:    ${ }^{6}$ This result is again a generalization of [7, Lemma 3.3.] where most of the modifications comes from the fact that we are working in the Wiener chaos of order 2.

[^6]:    ${ }^{7}$ The random variables $\widetilde{\varepsilon_{\lambda^{\prime}}} M_{j}$ and $\widetilde{\varepsilon_{\lambda^{\prime}}} M_{j}$ are defined in an obvious way.

[^7]:    ${ }^{1}$ LOCKSS: Lots of Copies Keep Stuff Safe http://www.lockss.org/
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