

Anharmonic Contributions To Lattice Thermal Conductivity



José Batista^{1*}, Olle Hellman² & Matthieu Verstraete¹ *jpabatista@uliege.be

¹ nanomat/Q-MAT/CESAM and European Theoretical Spectroscopy Facility, Université de Liège, B-4000 Liège, Belgium

European Theoretical Spectroscopy Facility ² Department of Molecular Chemistry and Material Science, Weizmann Institute of Science, Rehovot 76100, Israel

1. Introduction

Anharmonic thermal transport in insulators can be approached in several ways of varying complexity [1, 2], from independent phonon models to a more rigorous treatment with inter-branch coupling and memory included [3]. All current approaches rely, however, on the renormalization of the harmonic expression to include anharmonic effects. In this work we use an explicitly anharmonic heat current to derive the thermal transport coefficients, and find their higher order anharmonic contributions.

2. Quasi-harmonic Approaches

3. Spectral functions

- Green-Kubo: $\kappa = \frac{V}{k_B T^2} \int_0^\infty \langle \mathbf{J}(t) \mathbf{J}(0) \rangle dt$.
- Harmonic heat current: $\mathbf{J}_{ha} = \sum_{ab} \mathbf{v}_{ab} \omega_a B_a A_b$, $\mathbf{q} = (\mathbf{q}, s_q), B_q = a_q a_{\bar{q}}^{\dagger}, A_q = a_q + a_{\bar{q}}^{\dagger}$.
- Thermal transport coefficients: $\kappa = \frac{1}{2V} \sum_{ab} \mathbf{v}_{ab} \otimes \mathbf{v}_{ab} c_{ab} \tau_{ab}$.
- Anharmonicity is introduced in τ as:

 $- \operatorname{RTA} [4]: \tau_{a} = \frac{1}{2\Gamma_{a}} \to \operatorname{diagonal} \text{ in momentum and mode.}$ $- \operatorname{QHGK} [1]: \tau_{\mathbf{a}s_{a}s_{b}} = \frac{\Gamma_{\mathbf{a}s_{a}} + \Gamma_{\mathbf{a}s_{b}}}{(\omega_{\mathbf{a}s_{a}} - \omega_{\mathbf{a}s_{b}})^{2} + (\Gamma_{\mathbf{a}s_{a}} + \Gamma_{\mathbf{a}s_{b}})^{2}} \to \operatorname{no} \operatorname{longer} \operatorname{diagonal} \operatorname{in mode.}$

- Γ calculated via Fermi golden rule \rightarrow no memory effects.

• Anharmonicity renormalizes harmonic results, but no purely anharmonic contributions.



4. Anharmonic Heat Current

- AHC: $\mathbf{J}_{anh} = \sum_{abc} \mathbf{v}_{abc} \omega_a B_a A_b A_c$.
- \mathbf{v}_{abc} generalizes \mathbf{v}_{ab} to the 3-phonon

6. Anharmonic Contributions

• In order to evaluate the Kubo integral, we transform the correlation functions into phonon spectral functions $J(\Omega)$:

case.

- $\mathbf{J} = \mathbf{J}_{ha} + \mathbf{J}_{anh}$.
- $\kappa \propto \langle (\mathbf{J}_{ha}(t) + \mathbf{J}_{anh}(t))(\mathbf{J}_{ha} + \mathbf{J}_{anh}) \rangle.$
- Purely anharmonic contributions will appear.

5. Decoupling Scheme

- 4-point correlation function $\langle B_a(t)A_b(t)B_cA_d\rangle$ leads to the quasi-harmonic results.
- 5 and 6-point correlation functions lead to new contributions.
- Decoupling scheme:

 $\langle B_a(t)A_b(t)A_c(t)B_dA_eA_f\rangle \approx$ $\langle B_a(t)B_d\rangle\langle A_b(t)A_e\rangle\langle A_c(t)A_f\rangle$ +other 2-pt

$\langle X_a(t)Y_b^{\dagger}\rangle = \int J_{ab}^{XY}(\Omega)(n(\Omega)+1)e^{\Omega\tau}d\Omega$.

• By convolution of these spectral functions we obtain the purely anharmonic contributions to κ :



$$\kappa_{anh-anh}^{2-pt} \propto J_{-(\mathbf{a}+\mathbf{b})s_{c}s_{f}}^{AA}(\Omega) * J_{\mathbf{a}s_{a}s_{d}}^{BB}(\Omega') * J_{\mathbf{b}s_{b}s_{e}}^{AA}(\Omega'')$$

- Simple 2-momentum dependency.
- Main contribution to anharmonic κ .
- $\kappa_{anh-anh}^{3-pt} \propto (S_{ab}^{\text{III}}(\Omega) J_{(\mathbf{a}+\mathbf{b})s_{\mu}s_{d}}^{AB}(\Omega)) * \\ (S_{-(\mathbf{a}+\mathbf{b})s_{c}}^{\text{I}}(\Omega') J_{-\mathbf{e}s_{\nu}s_{e}}^{AA}(\Omega')) * \\ (J_{-(\mathbf{a}+\mathbf{b}-\mathbf{e})s_{\eta}s_{f}}^{AA}(\Omega''))$
 - 3-momentum dependency, $J(\Omega)$ filtered by $S(\Omega)$ and $S(\Omega)$.
- $\kappa_{anh-ha} \propto (S_{ab}^{\mathrm{III}}(\Omega) J_{(\mathbf{a}+\mathbf{b})s_{\mu}s_{d}}^{AB}(\Omega)) * \qquad \mid \kappa_{ha-anh} \propto (\mathcal{S}_{a}^{\mathrm{IV}}(\Omega) J_{-\mathbf{c}s_{\mu}s_{c}}^{AB}(\Omega)) *$

 $+\langle B_a(t)A_b(t)B_d\rangle\langle A_c(t)A_eA_f\rangle$ +other 3-pt.

• Since we can write all 3-point functions in terms of 2-point ones, all contributions will still be written in terms of dressed-phonons.

References

- [1] L. Isaeva et al. Nature Communications, 10(1):3853, Aug 2019.
- [2] M. Simoncelli et al. Nature Physics, 15(8):809–813, Aug 2019.
- [3] D. Dangić et al. npj Computational Materials, 7, 12 2021.
- [4] G.D. Mahan. Many-Particle Physics. Springer US, 1990.

 $anh-ha \propto \left(S_{ab}^{AA}(\Omega)J_{(\mathbf{a}+\mathbf{b})s_{\mu}s_{d}}^{AA}(\Omega)\right) * \left(J_{-(\mathbf{a}+\mathbf{b})s_{c}s_{e}}^{AA}(\Omega')\right)$

 $-2-\text{momentum dependency, corrects} -2-\text{momentum} \\ \kappa_{ha}. -2-\text{momentum} \\ J(\Omega) \text{ filtered } \mathbf{I}$

 $\left(J_{(\mathbf{a}+\mathbf{c})s_{\nu}s_{d}}^{AA}(\Omega')\right)*\left(J_{-\mathbf{a}s_{b}s_{e}}^{AA}(\Omega'')\right)$

- 2-momentum dependency, but with $J(\Omega)$ filtered by $\mathcal{S}(\Omega)$.

7. Conclusions and Outlook

• Pure anharmonic contributions appear as new combinations of single phonon spectral functions. Memory is always present.

• Correction to the 2 phonon expression arises from cross-correlation of \mathbf{J}_{ha} and \mathbf{J}_{anh} .