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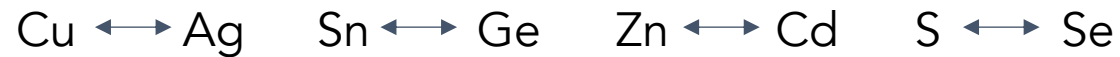
Relevance of Ge incorporation on the physics of deep defects in kesterite materials

T. Ratz^{1,2}, N-D. Nguyen¹, G. Brammertz³, B. Vermang^{2,3,4}, J-Y. Raty¹

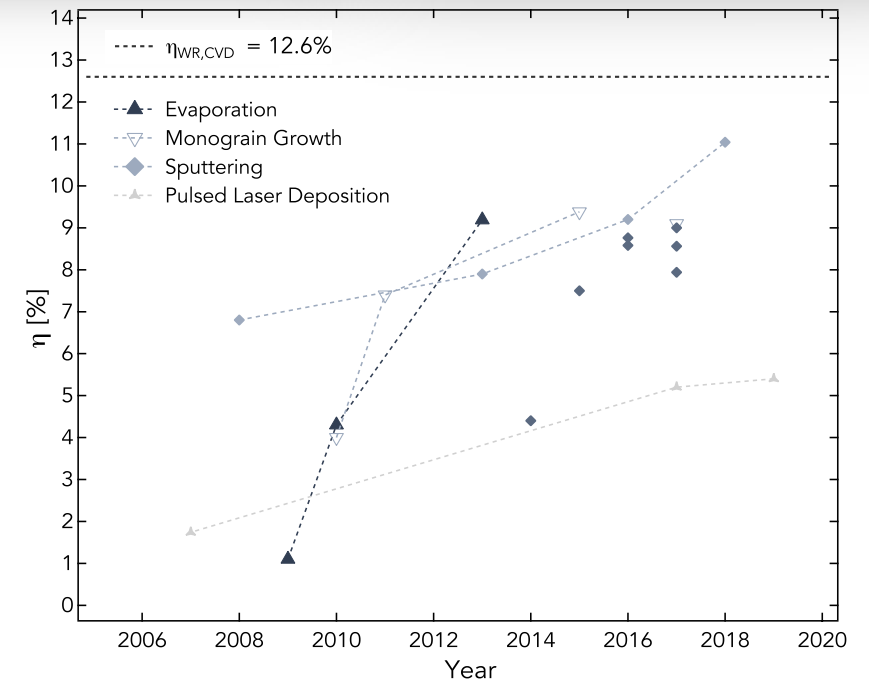
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1. Introduction
2. Kesterite phase diagrams
3. Point defects in Ge-doped and Ge-alloyed kesterites
4. Lattice distortion upon defect incorporation
5. Conclusion

- Limitations of the $\text{Cu}_2\text{ZnSnS}_4$ kesterite-based solar cell efficiency
- Bulk point defects have been pointed out as efficient recombination centres (Sn_{Zn} and/or $[\text{Cu}_{\text{Zn}} + \text{Sn}_{\text{Zn}}]$)
- Cationic substitution as a possible route for efficiency improvement



What is the impact of Ge incorporation on the physics of deep defect in kesterite ?



2

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Relevance of Ge incorporation to control the physical behaviour of point defects in kesterite†

Thomas Ratz, ^{a,b} Ngoc Duy Nguyen, ^a Guy Brammertz, ^c Bart Vermang ^{bcd} and Jean-Yves Raty^a

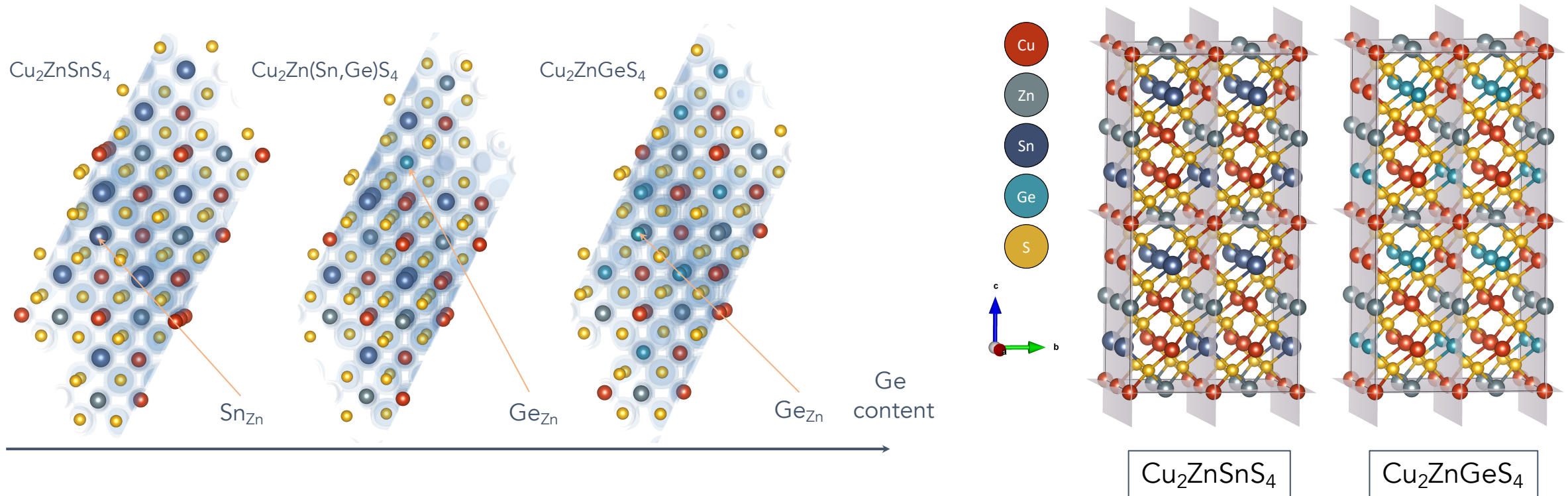
Point defect calculations using DFT

From a perfect kesterite crystal to a defected system – supercell approach (64-atom calculations)

3

Motivation: reduction of the V_{OC} deficit following Ge incorporation into kesterite materials

Goal: Study of the point defect physical behaviour in Ge doped and alloyed kesterites



Point defect calculations using DFT

Point defect formation energy:

$$\Delta H_F(\alpha, q, \mu_i, E_F) = \underbrace{E(\alpha, q) - E(host)}_{\text{Total energy difference between the defected and pristine supercell}} + \underbrace{\sum_i n_i (E_i - \mu_i)}_{\text{Energy cost for particles (elements) exchange}} + \underbrace{q[\epsilon_{VBM}(host) + E_F]}_{\text{Energy cost for charges exchange}}$$

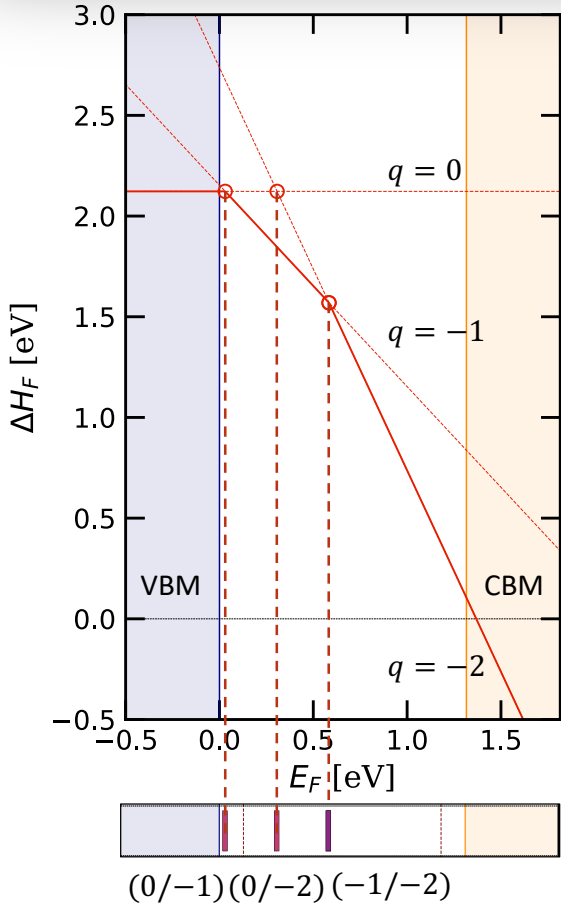
Total energy difference between the defected and pristine supercell

Energy cost for particles (elements) exchange

Energy cost for charges exchange

- Identification of point defects types
 - Acceptor/donor
 - Recombination center

$$\epsilon_{q,q'}(\alpha) = \frac{\Delta H_F(\alpha, q) - \Delta H_F(\alpha, q')}{q' - q}$$



- Correction terms:
- Electrostatic interaction between image defects
 - Band alignment between charged defects and the host supercell

Thermodynamic conditions

- $\mu_{Cu} < 0, \mu_{Zn} < 0, \mu_X < 0, \mu_S < 0$

Avoid pure phases

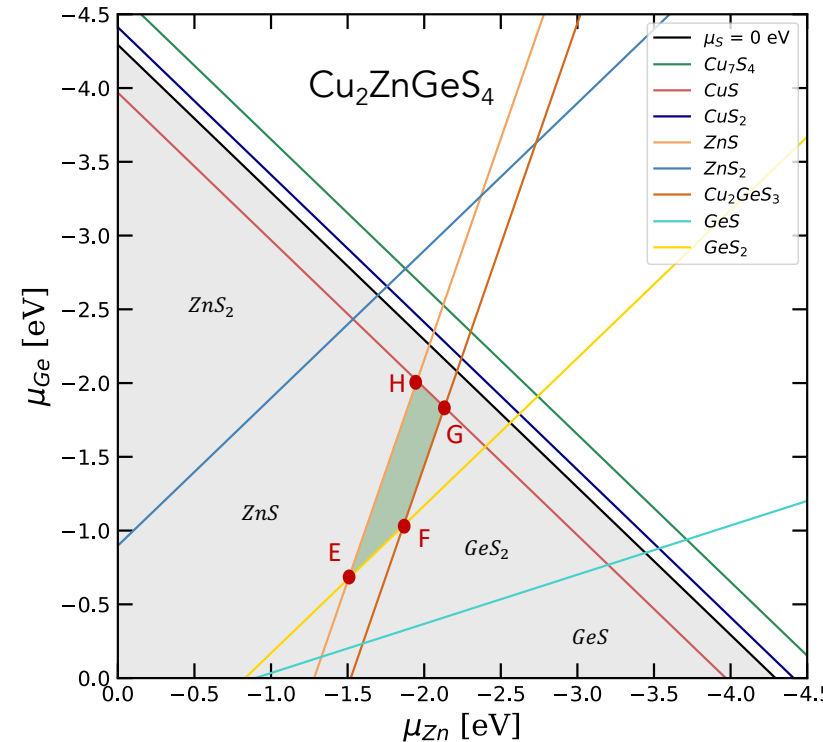
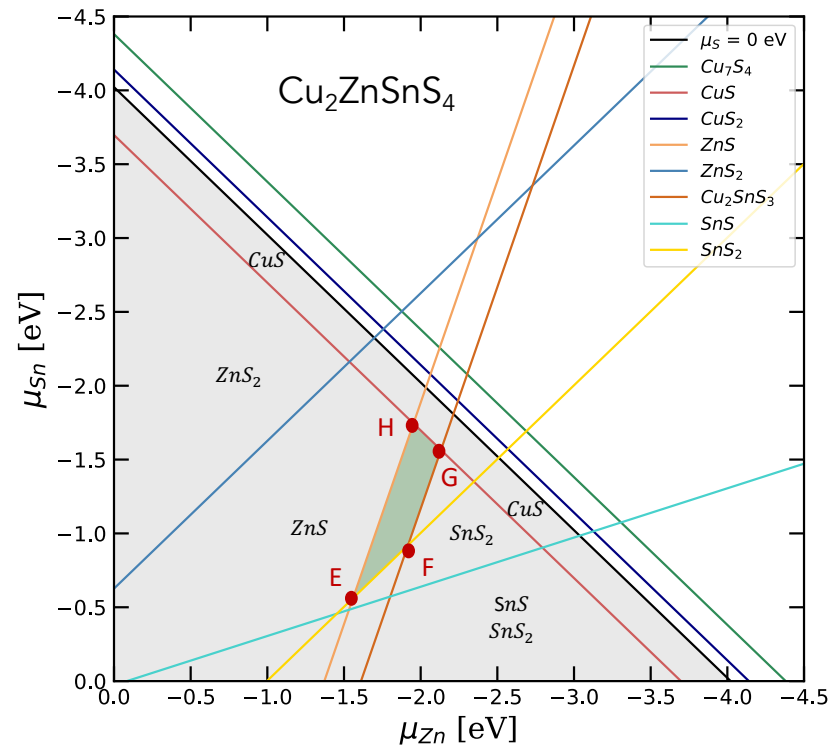
5

- $2\mu_{Cu} + \mu_{Zn} + \mu_X + 4\mu_S = \Delta H_F(CZXS)$

Kesterite materials is desired

- $\sum_i n_i \mu_i < \Delta H_F(\text{Secondaries}) = E(A_x B_y) - xE(A) - yE(B)$

No secondary phases

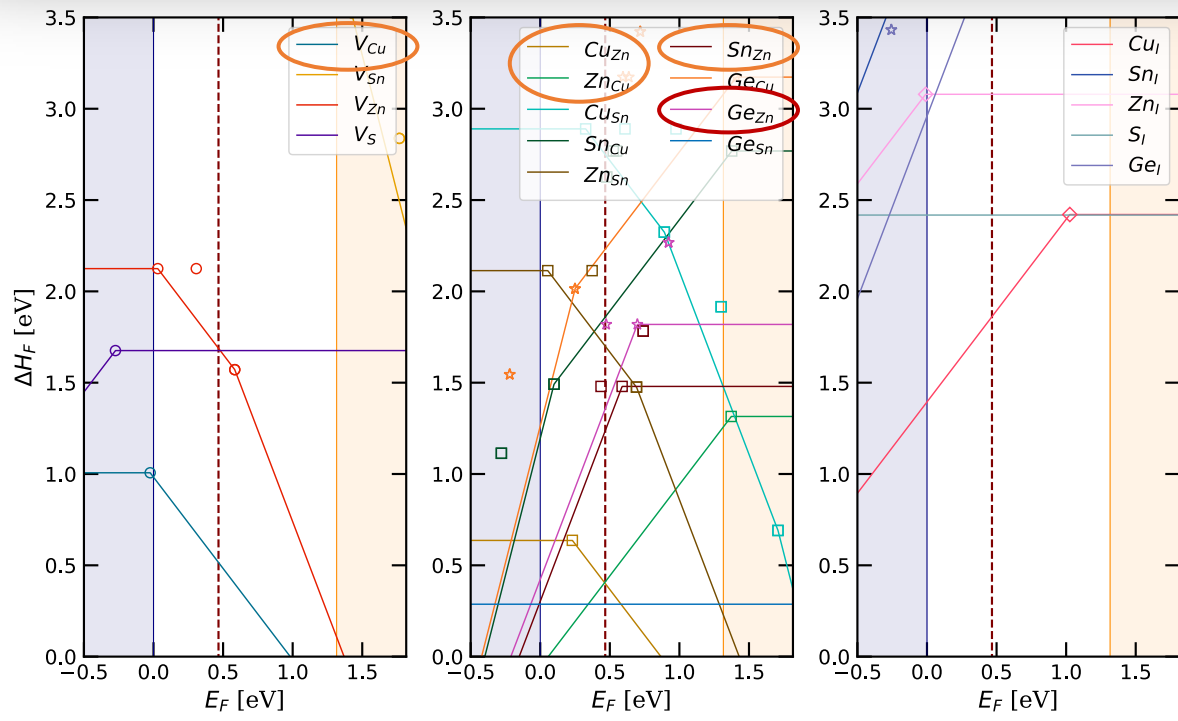


Kesterite phase
 Pure phase

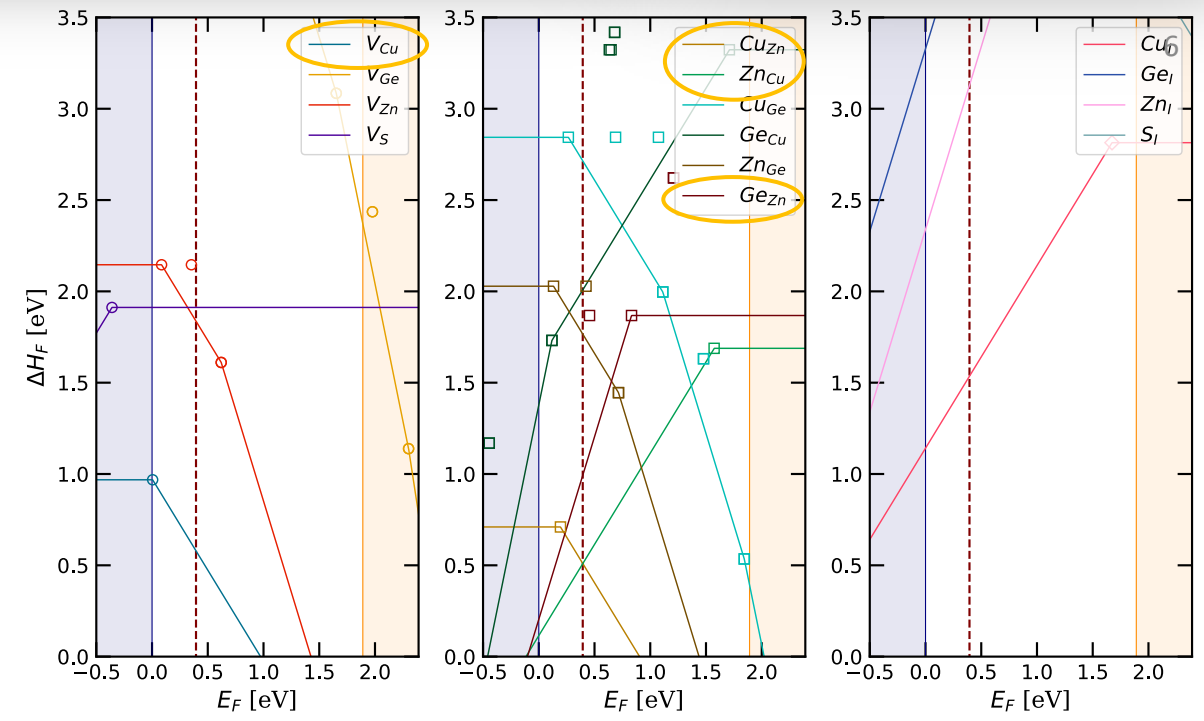
$\mu_i \rightarrow 0$: rich in element i
 $|\mu_i| \rightarrow \infty$: poor in element i

$\mu_{Cu} = -0.55 \text{ eV}$

Point defect formation energies (point E)



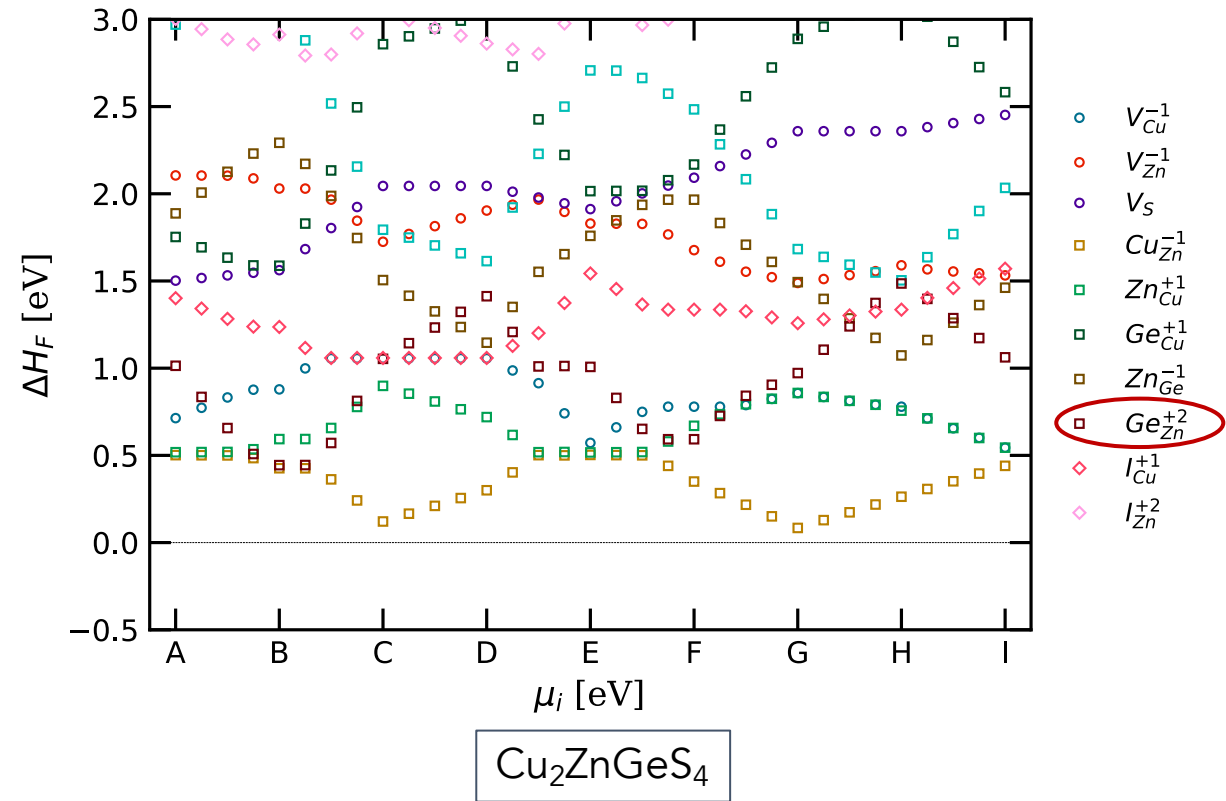
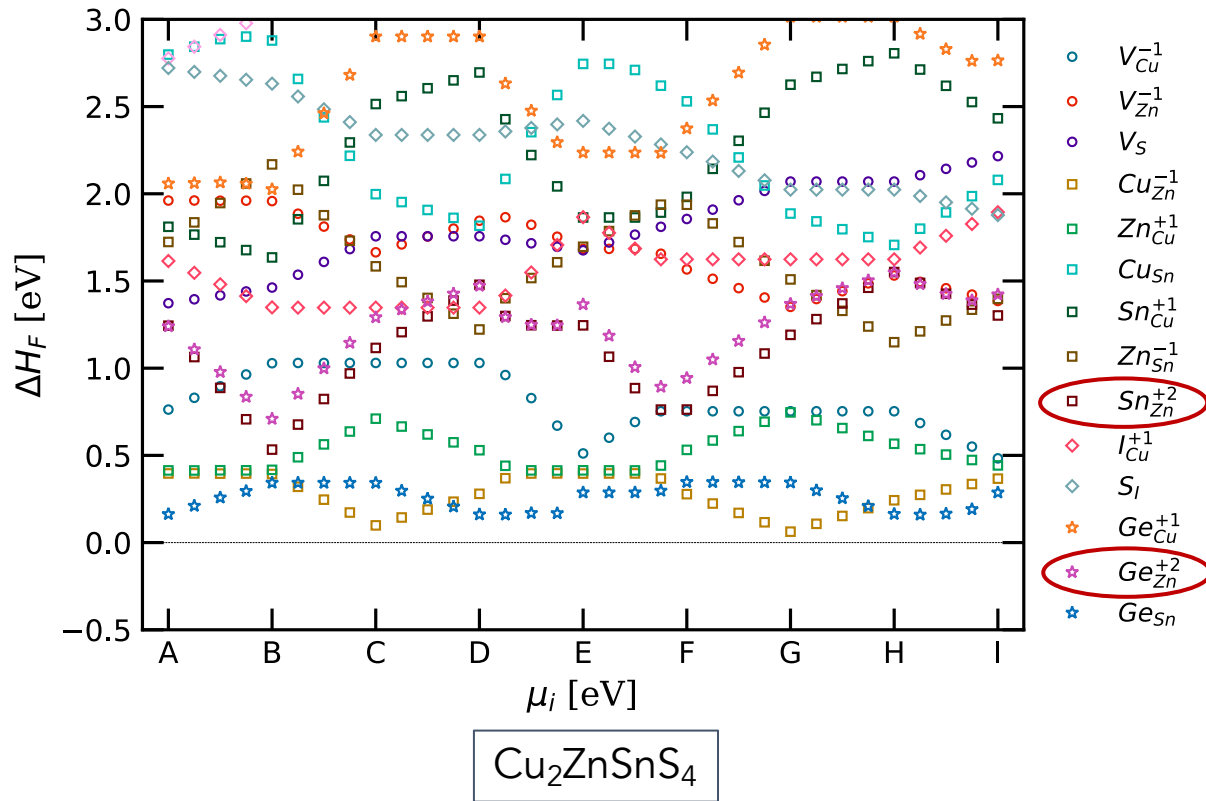
$\text{Cu}_2\text{ZnSnS}_4$



$\text{Cu}_2\text{ZnGeS}_4$

- Fermi energy level under thermodynamic equilibrium condition
- p-type intrinsic conductivity of kesterites
- V_{Cu} , Cu_{Zn} , Zn_{Cu} , X_{Zn} ($X=\text{Sn}, \text{Ge}$) low formation energies

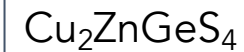
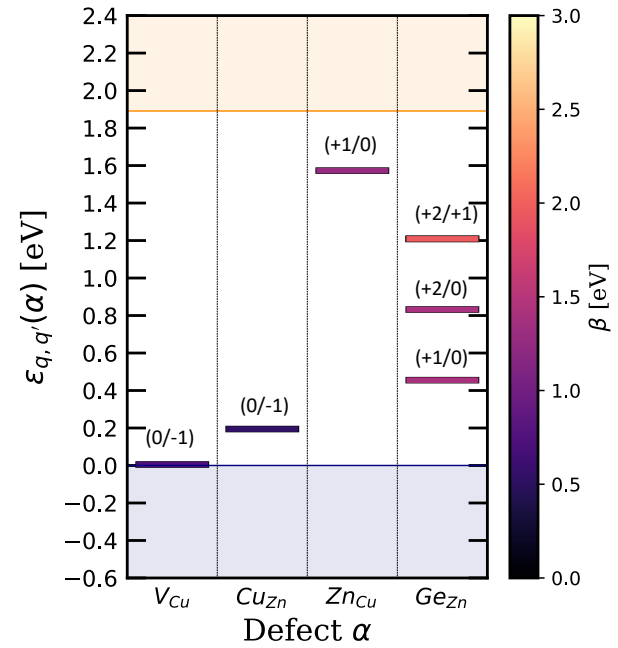
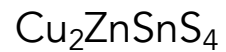
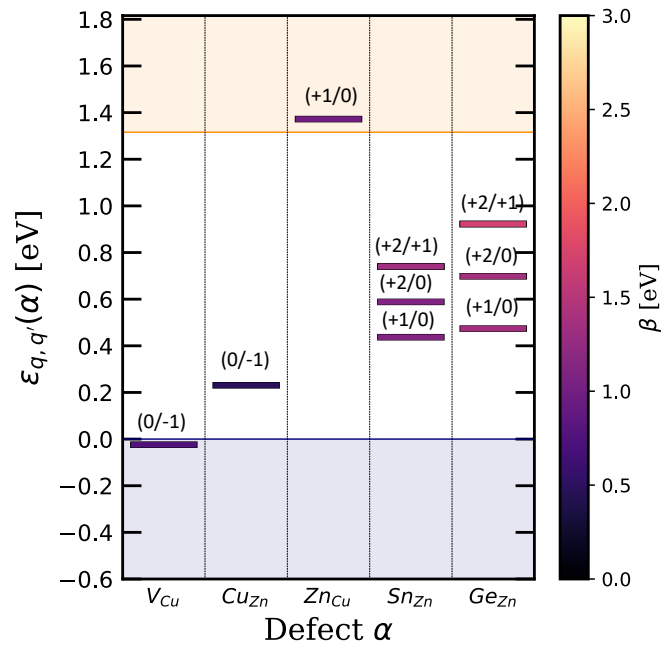
Evolution of the point defect formation energies according to the synthesis conditions:



- $V_{\text{Cu}}, \text{Cu}_{\text{Zn}}, \text{Zn}_{\text{Cu}}, \text{X}_{\text{Zn}}$ ($\text{X}=\text{Sn}, \text{Ge}$) low formation energies
- Ge_{Sn} most abundant doping defect
- From Sn_{Zn} to the dopant Ge_{Zn} to the Ge_{Zn} substitutional defect in the Ge-based kesterite

Physical behaviour of the defects previously identified as abundant in the materials

8



- V_{Cu} , Cu_{Zn} , Zn_{Cu} acting respectively as acceptor and donor for the later
- X_{Zn} ($X=\text{Sn},\text{Ge}$) acting as recombination centres
- The formation energy is not the only key ingredient impacting the physical behaviour of a defect

What is a deep defect? Combining Shockley-Read-Hall statistics with multiphonon recombination theory

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Transition between two states: Ψ_1 the initial band state and Ψ_2 , the final defect state

$$W_{1,2} = \frac{2\pi}{\hbar} |\langle \Psi_1 | H_{int} | \Psi_2 \rangle|^2 V N(E_2)$$

electron-lattice interaction

Article
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Effective and Noneffective Recombination Center Defects in $\text{Cu}_2\text{ZnSnS}_4$: Significant Difference in Carrier Capture Cross Sections

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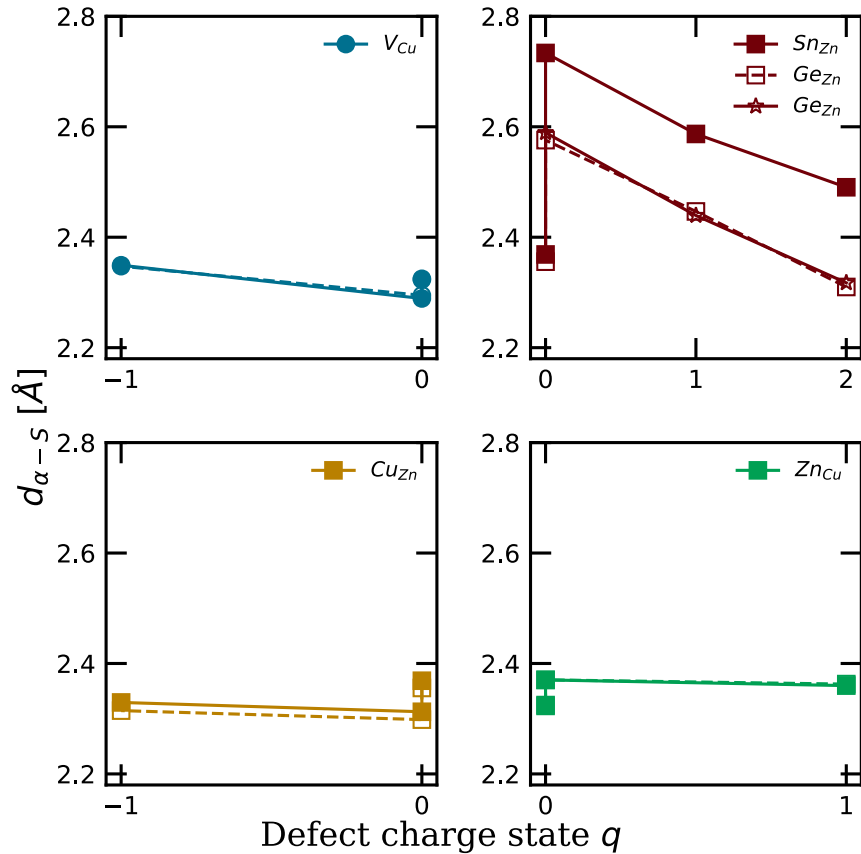
$\text{Sn}_{\text{Zn}}^{+1}$ (+1/0) and $\text{Sn}_{\text{Zn}}^{+2}$ (+2/+1) are not equivalent recombination centres as $\sigma_{\text{Sn}_{\text{Zn}}^{+1}} \ll \sigma_{\text{Sn}_{\text{Zn}}^{+2}}$

The larger the structural relaxation, the larger the carrier capture cross section σ

What is the impact of the Ge incorporation on the lattice distortion ?

Relation between the lattice distortion upon defect incorporation and the defect carrier capture cross section

10



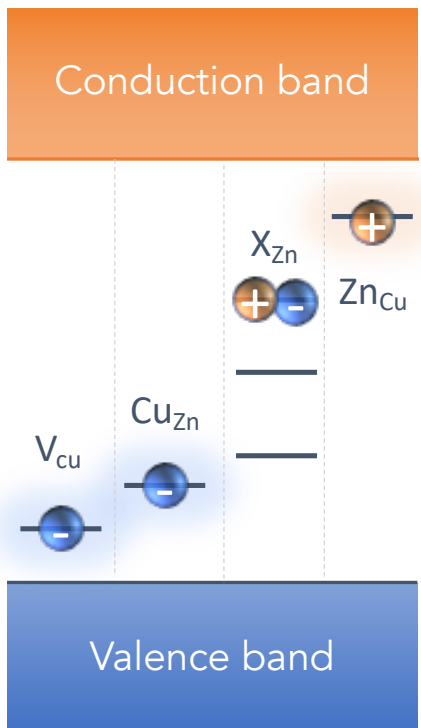
$$e_p = \sigma_p \langle v_p \rangle N_V \exp\left(-\frac{E_T - E_V}{k_B T}\right)$$

Defect emission rate e_p depends on

- The energetic position of the defect within the band gap and with respect to the CBM or VBM
- Carrier capture cross section σ_p related to the electron-phonon interaction

→ Reduction of the lattice distortion induced by the formation of the Ge_{Zn} substitutional defect in comparison to its Sn_{Zn} parent in the Sn-based kesterite

- Similar physical behaviour of intrinsic point defects in Sn-based and Ge-based kesterites
- Via Ge-doping in CZTS, Ge_{Sn} has been identified as the most abundant defect
- Sn_{Zn} and Ge_{Zn} are identified as recombination centres providing energy levels within the materials band gap
- Net reduction of the lattice distortion upon Ge_{Zn} formation both for Ge-doping and alloying in comparison to its Sn counterpart



→ Possible explanation for the V_{OC} improvement reported experimentally

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Thank you for your attention