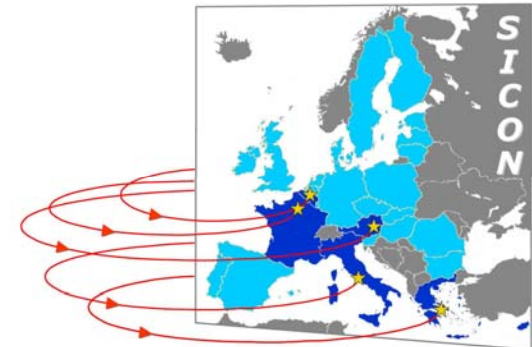


Parameter Estimation and Structural Model Updating Using Modal Methods in the Presence of Nonlinearity

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- 1. Introduction**
- 2. Theoretical Modal Analysis of Nonlinear Systems**
- 3. Nonlinear Experimental Modal Analysis**
- 4. Model Parameter Estimation Techniques**
- 5. Concluding Remarks**



Design of engineering structures relies on

- Numerical predictions → modal analysis (FEM)
- Dynamic testing → experimental modal analysis (EMA)

In the case of **linear structures**, the techniques available for EMA are mature e.g.

- Eigensystem realization algorithm
- Stochastic subspace identification
- Polyreference least-squares complex exponentials frequency domain
- etc

Nonlinearity in Engineering Applications

backlash and friction in
control surfaces and joints

fluid-structure interaction



hardening nonlinearities in
engine-to-ylon connections

composite materials

Many works are reported in the literature on dynamic testing and identification of nonlinear systems but very few address nonlinear phenomena during modal survey tests.

Aim of this presentation

- To extend experimental modal analysis to a practical analogue using the nonlinear normal mode (NNM) theory.
- Validate mathematical models of non-linear structures against experimental data.

Why?

- NNMs offer a solid and rigorous mathematical tool.
- They have a clear conceptual relation to the classical LNMs.
- They are capable of handling strong structural nonlinearity.

1. Introduction

2. Theoretical Modal Analysis of Nonlinear Systems

- **Nonlinear Normal Modes (NNMs)**
- **Numerical Computation of NNMs**
- **Frequency-Energy Plot**

3. Nonlinear Experimental Modal Analysis

4. Model Parameter Estimation Techniques

5. Concluding Remarks

MDOF system

$$\mathbf{M} \ddot{\mathbf{x}} + \mathbf{C} \dot{\mathbf{x}} + \mathbf{K} \mathbf{x} + \mathbf{f}_{NL}(\mathbf{x}, \dot{\mathbf{x}}) = \mathbf{p}(t)$$



Vector of nonlinear forces

- **Dynamic analysis**

Prediction of the responses using a numerical integration procedure (e.g. Newmark's schema)

- **Modal analysis of the MDOF system (with no damping)**

In the linear case

$$\mathbf{M} \ddot{\mathbf{x}} + \mathbf{K} \mathbf{x} = 0$$

Structural eigenproblem

$$\mathbf{K} \Phi_j = \omega_j^2 \mathbf{M} \Phi_j \quad j = 1, \dots, n$$

ω_j j^{th} natural frequency

In the nonlinear case

$$\mathbf{M} \ddot{\mathbf{x}} + \mathbf{K} \mathbf{x} + \mathbf{f}_{NL}(\mathbf{x}, \dot{\mathbf{x}}) = 0$$

Use of the concept of nonlinear normal modes (NNMs) which is a rigorous extension of the concept of eigenmodes to nonlinear systems.

Definitions

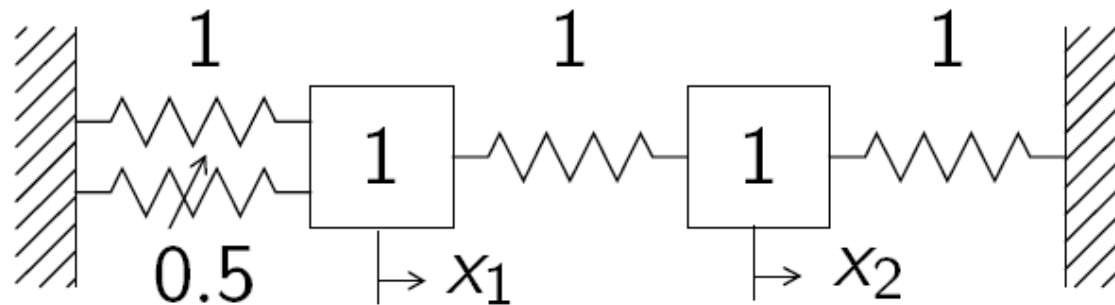
Two definitions of an NNM in the literature:

1. Targeting a straightforward nonlinear extension of the linear normal mode (LNM) concept, **Rosenberg** defined an NNM motion as a vibration in unison of the system (i.e., a synchronous periodic oscillation).
2. To provide an extension of the NNM concept to damped systems, **Shaw and Pierre** defined an NNM as a two-dimensional invariant manifold in phase space. Such a manifold is invariant under the flow (i.e., orbits that start out in the manifold remain in it for all time), which generalizes the invariance property of LNMs to nonlinear systems.

In the **present study**, an NNM motion is defined as a (non-necessarily synchronous) periodic motion of the undamped mechanical system

→ this extended definition is particularly attractive when targeting a numerical computation of the NNMs.

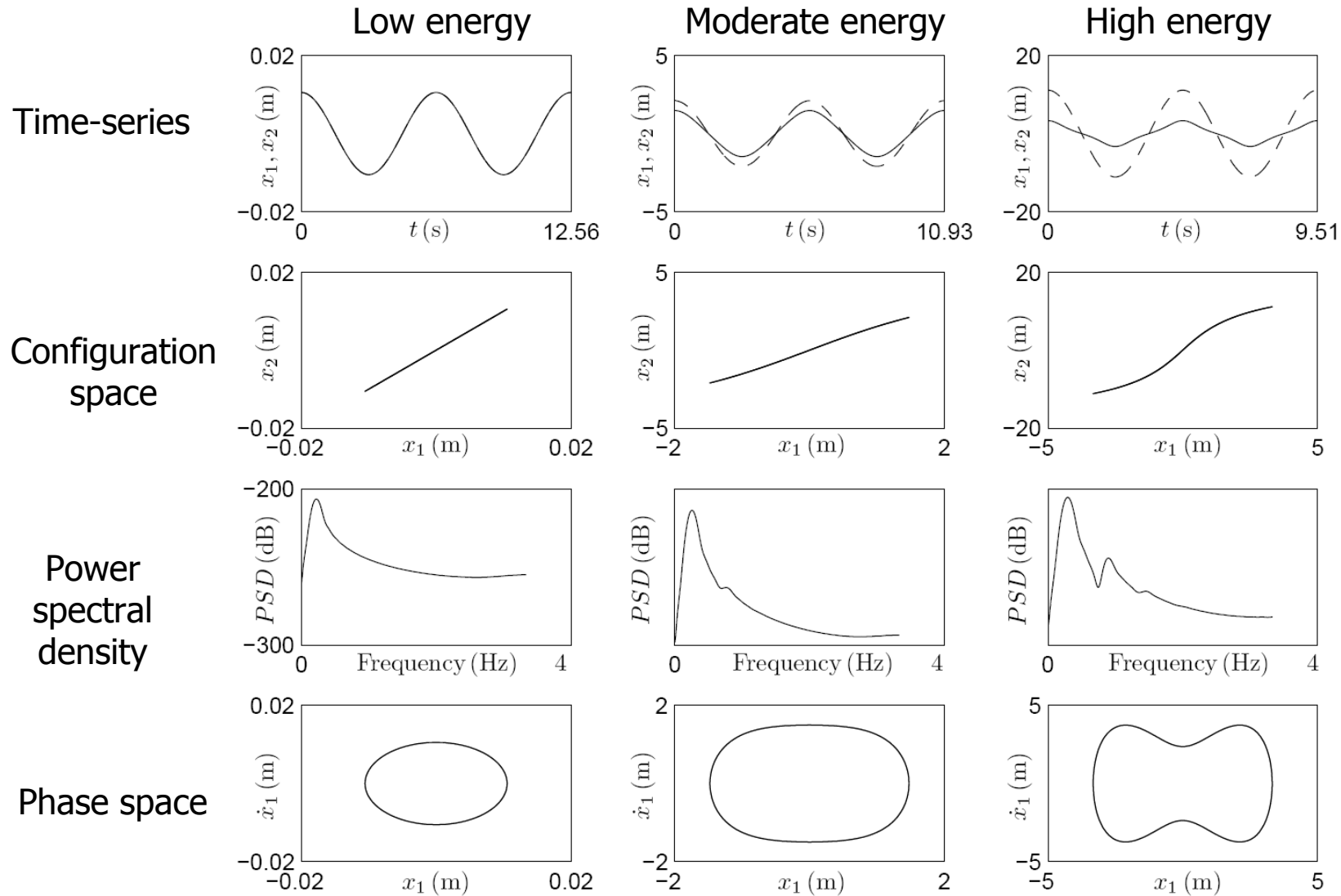
Illustrative example: 2 DOF-system with a cubic stiffness



$$\ddot{x}_1 + (2x_1 - x_2) + 0.5x_1^3 = 0$$

$$\ddot{x}_2 + (2x_2 - x_1) = 0$$

In-Phase NNMs for Increasing Energy



General equation of the nonlinear system (with no damping)

$$\mathbf{M} \ddot{\mathbf{x}} + \mathbf{K} \mathbf{x} + \mathbf{f}_{NL}(\mathbf{x}, \dot{\mathbf{x}}) = 0$$

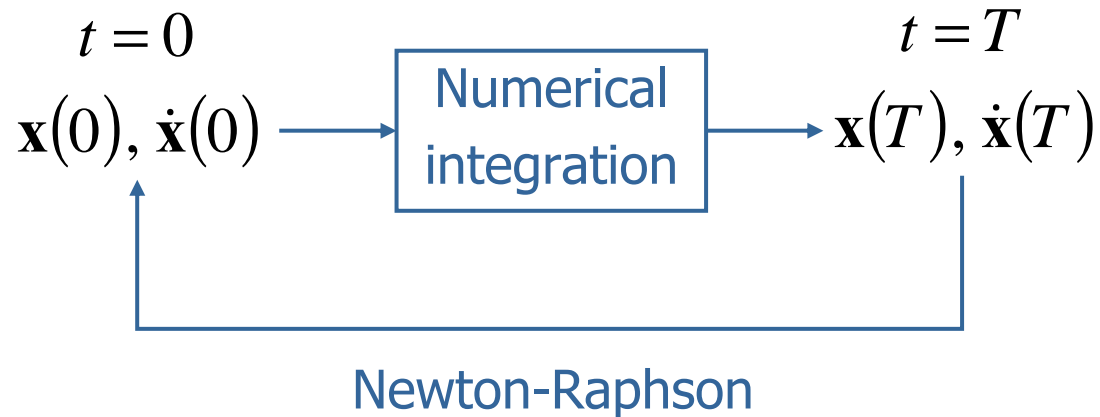


Vector of nonlinear forces

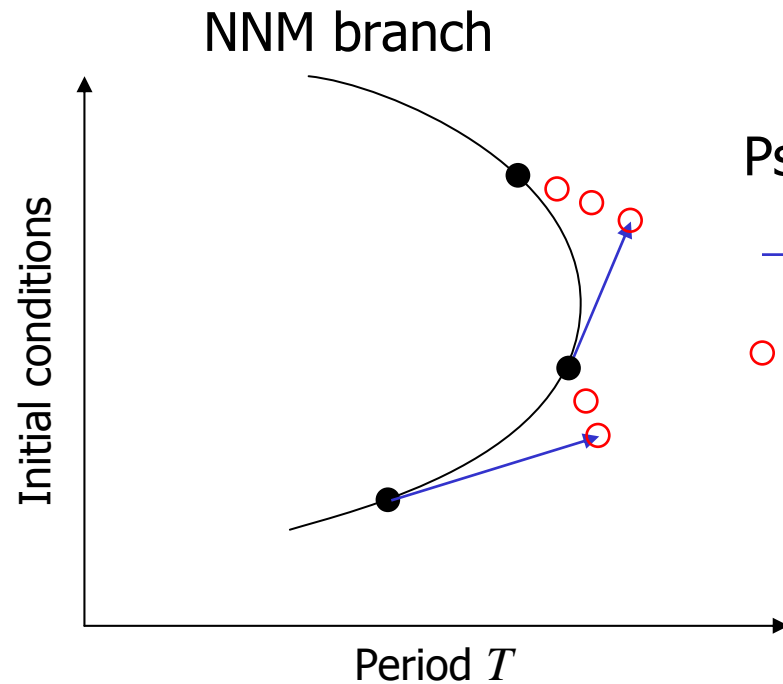
The numerical computation of NNMs relies on two main techniques, namely a ***shooting procedure*** and a method for the ***continuation of periodic solutions***.

- **Shooting method**

The shooting method consists in finding, in an iterative way, the initial conditions $\mathbf{x}(0)$, $\dot{\mathbf{x}}(0)$ and the period T inducing an isolated periodic motion (i.e., an NNM motion) of the conservative system.



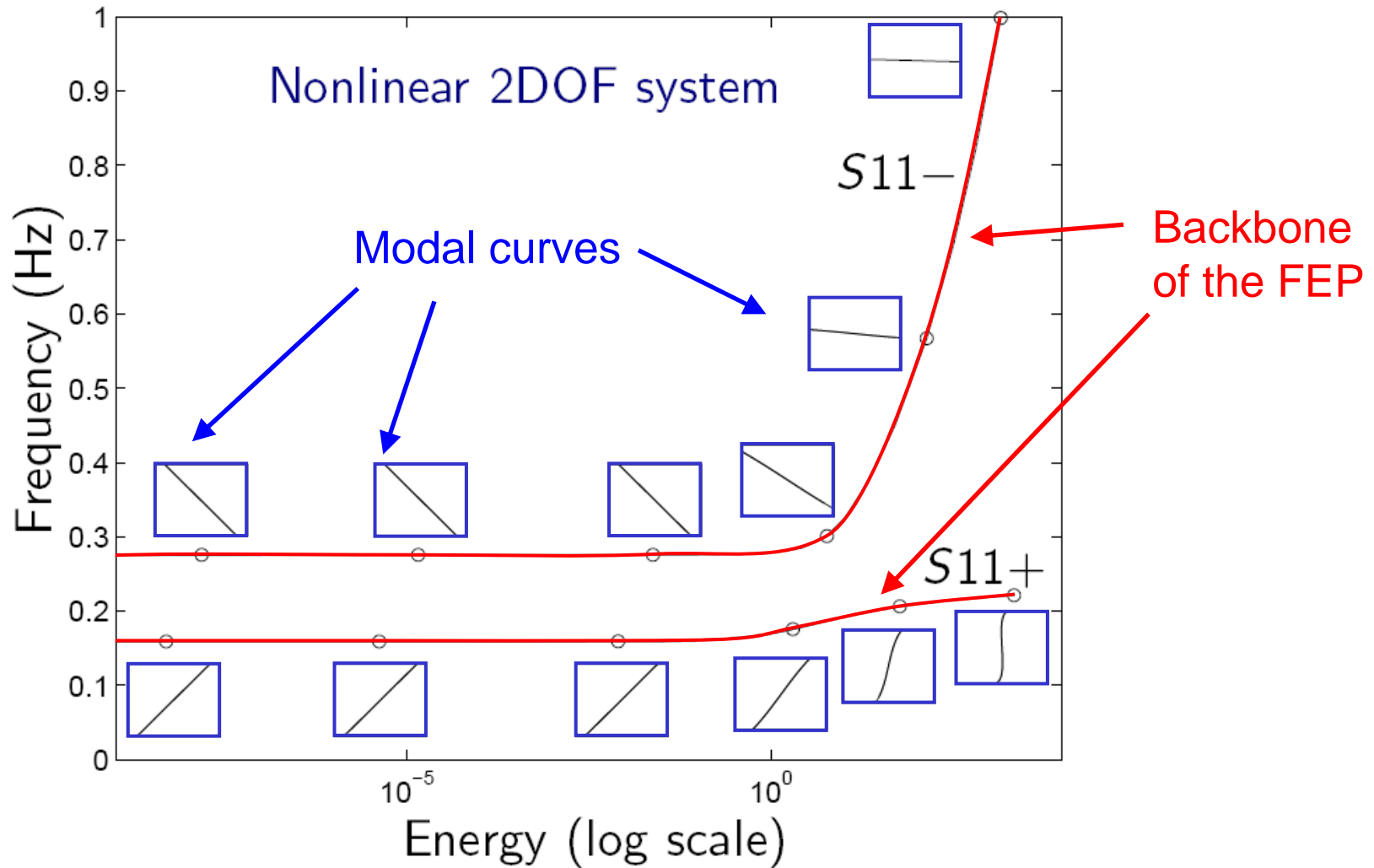
- Pseudo-arclength continuation method



Pseudo-arclength continuation method:

- predictor step tangent to the branch
- ○ ○ corrector step perpendicular to the predictor step (shooting)

Frequency-Energy Plot (FEP)





1. Introduction

2. Theoretical Modal Analysis of Nonlinear Systems

3. Nonlinear Experimental Modal Analysis

- **Phase Separation Methods**

- **Proper Orthogonal Decomposition**

- **Phase Resonance Methods**

- **Nonlinear Normal Mode Testing**

4. Model Parameter Estimation Techniques

5. Concluding Remarks

Experimental Modal Analysis (EMA)

Linear systems

Theoretical Approach

Finite Element Model

$$\mathbf{M} \ddot{\mathbf{x}} + \mathbf{K} \mathbf{x} = \mathbf{0}$$

Eigenvalue problem

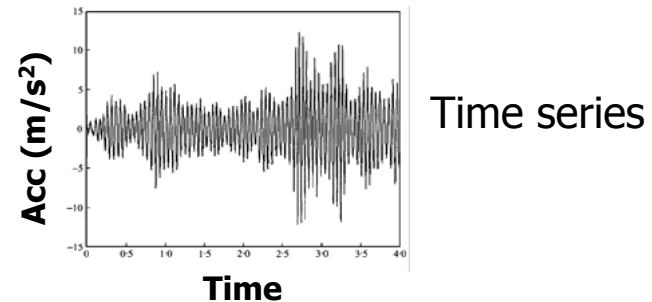
$$\mathbf{K} \Phi_j = \omega_j^2 \mathbf{M} \Phi_j$$

Natural frequencies (ω_j^2)

Mode shapes (Φ_j)

Experimental Approach

Response Measurements



Identification methods

EMA for linear systems is now mature and widely used in structural engineering → well established techniques [1], [2].

Nonlinear systems

Theoretical Approach

Finite Element Model

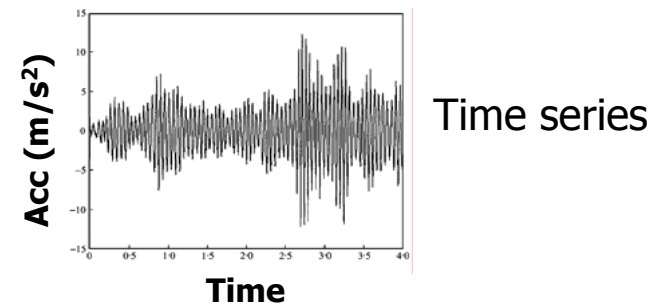
$$\mathbf{M} \ddot{\mathbf{x}} + \mathbf{K} \mathbf{x} + \mathbf{f}_{NL}(\mathbf{x}, \dot{\mathbf{x}}) = 0$$

Numerical NNM computation

NNM frequencies
NNM modal curves

Experimental Approach

Response Measurements



Experimental NNM extraction

EMA for nonlinear systems is **still a challenge**.

There are two main techniques for EMA.

1. Phase separation methods

Several modes are excited at once using either broadband excitation (e.g., hammer impact and random excitation) or swept-sine excitation in the frequency range of interest.

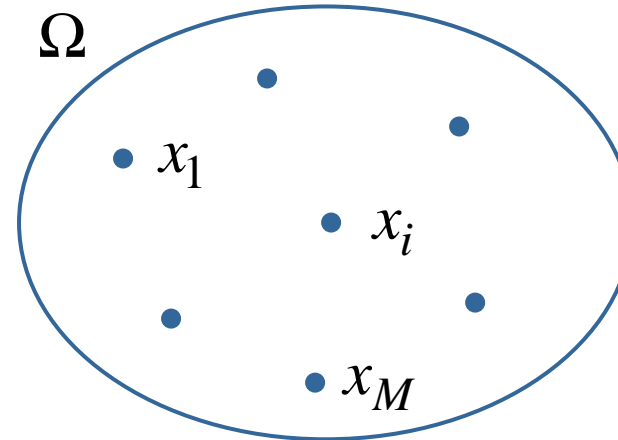
- in the nonlinear case, extraction of individual NNMs is not possible generally, because modal superposition is no longer valid.
- use of the proper orthogonal decomposition (POD) method to extract features from the time series .

Remark

- All structures encountered in practice are nonlinear to some degree.
- If a nonlinear structure is excited with a broadband excitation signal (e.g. random force), then the results will appear linear → experimental modal analysis will lead to an updated linearized model !

Proper Orthogonal Decomposition (POD)

Instrumented structure



N snapshots

$$\mathbf{X} = \begin{matrix} \underbrace{\hspace{10em}} \\ \left[\begin{array}{ccc} x_1(t_1) & \cdots & x_1(t_N) \\ \vdots & \ddots & \vdots \\ x_M(t_1) & \cdots & x_M(t_N) \end{array} \right] \end{matrix} \left. \vphantom{\begin{matrix} \underbrace{\hspace{10em}} \\ \left[\begin{array}{ccc} x_1(t_1) & \cdots & x_1(t_N) \\ \vdots & \ddots & \vdots \\ x_M(t_1) & \cdots & x_M(t_N) \end{array} \right] \end{matrix}} \right\} \begin{matrix} M \\ \text{measurement} \\ \text{co-ordinates} \end{matrix}$$

$\mathbf{X} = [\mathbf{x}(t_1) \quad \dots \quad \mathbf{x}(t_N)]$ is the observation matrix

Proper Orthogonal Decomposition (POD)

The $M \times M$ correlation matrix \mathbf{R} is built

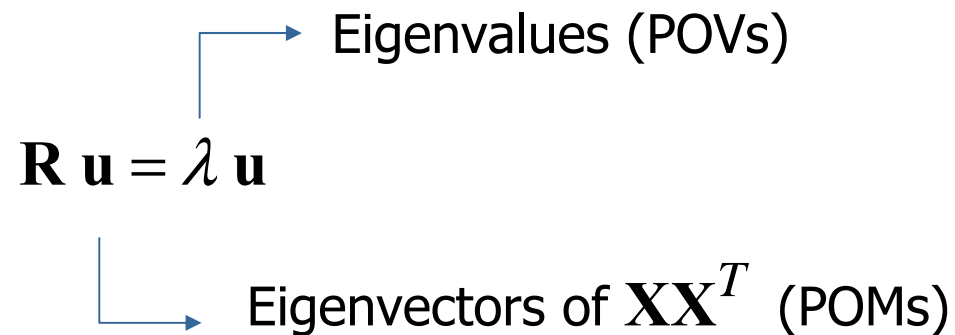
$$\mathbf{R} = \frac{1}{M} \mathbf{X} \mathbf{X}^T$$

The eigenvalue problem is solved

$$\mathbf{R} \mathbf{u} = \lambda \mathbf{u}$$

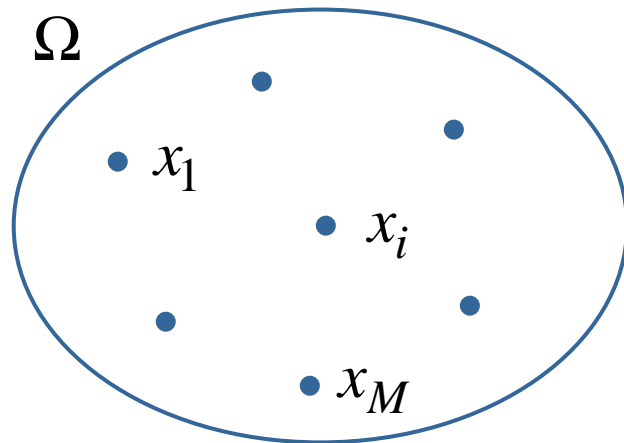
Eigenvalues (POVs)

Eigenvectors of $\mathbf{X} \mathbf{X}^T$ (POMs)



Proper Orthogonal Decomposition (POD)

Computation of the POMs using SVD



M measurement co-ordinates

N time samples

$$\mathbf{X}_{M \times N} = \begin{bmatrix} x_1(t_1) & \cdots & x_1(t_N) \\ \vdots & \ddots & \vdots \\ x_M(t_1) & \cdots & x_M(t_N) \end{bmatrix}$$

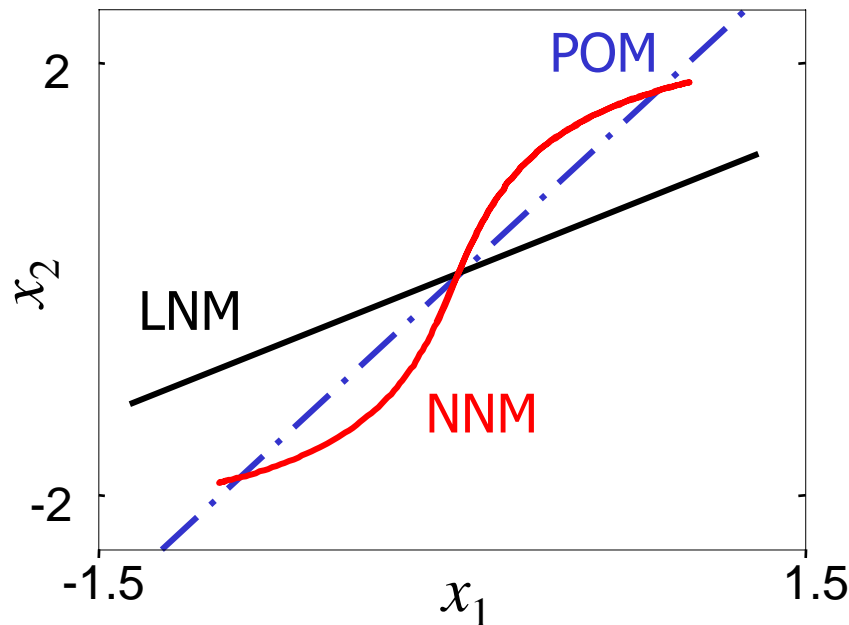
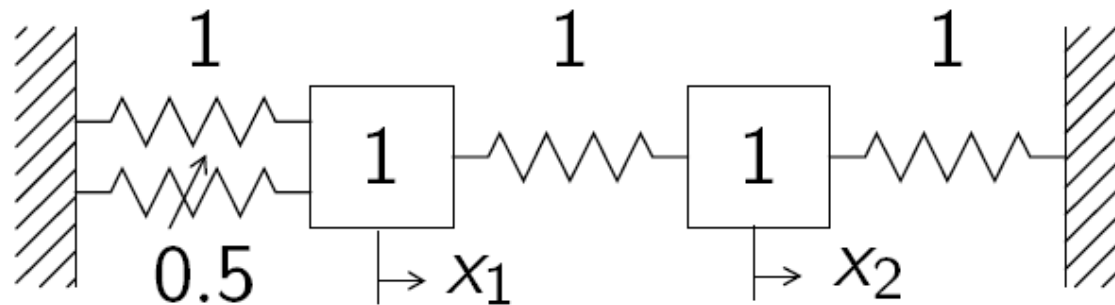
Using SVD

$$\mathbf{X}_{M \times N} = \mathbf{U}_{M \times M} \Sigma_{M \times N} \mathbf{V}_{N \times N}^T$$

$\nearrow \text{diag}(\sqrt{\lambda_i}) \quad (\lambda_i \equiv \text{POV})$
 \searrow Eigenvectors of $\mathbf{X}\mathbf{X}^T$ (POM)

Geometric Interpretation of the POMs

Comparison of LNM, NNM and POM on the 2 DOF example



First mode

The POM is the best linear representation of the nonlinear normal mode.

Key idea: Application of the POD to Features Extraction

Linear Systems

$$\mathbf{M} \ddot{\mathbf{x}} + \mathbf{C} \dot{\mathbf{x}} + \mathbf{K} \mathbf{x} = \mathbf{p}(t)$$

Deterministic approach

Eigenvalue problem :

$$(\mathbf{K} - \omega^2 \mathbf{M}) \Phi = \mathbf{0}$$

Response :

$$\mathbf{x}(t) = \sum_{i=1}^n \eta_i(t) \Phi_{(i)}$$

Spatial information

Natural frequencies

$$\eta_i = A_i \cos(\omega_i t) + B_i \sin(\omega_i t)$$

Nonlinear Systems

$$\mathbf{M} \ddot{\mathbf{x}} + \mathbf{C} \dot{\mathbf{x}} + \mathbf{K} \mathbf{x} + \mathbf{f}_{NL}(\mathbf{x}, \dot{\mathbf{x}}) = \mathbf{p}(t)$$

Statistical approach

Proper Orthogonal Decomposition :

$$\mathbf{X} = \mathbf{U} \Sigma \mathbf{V}^T$$

Response :

$$\mathbf{x}(t) = \sum_{j=1}^n a_j(t) \mathbf{u}_{(j)}$$

Spatial information

Time information
→ Instantaneous frequencies

2. Phase resonance methods (Normal mode testing)

One of the normal mode at a time is excited using multi-point sine excitation at the corresponding natural frequency. The modes are identified one by one.

→ can be extended to nonlinear structures according to the invariance property of NNMs:

« If the motion is initiated on one specific NNM, the remaining NNMs remain quiescent for all time. »

Remark

- Expensive and difficult.
- Extremely accurate mode shapes → a way to identify NNMs (but still a research topic).

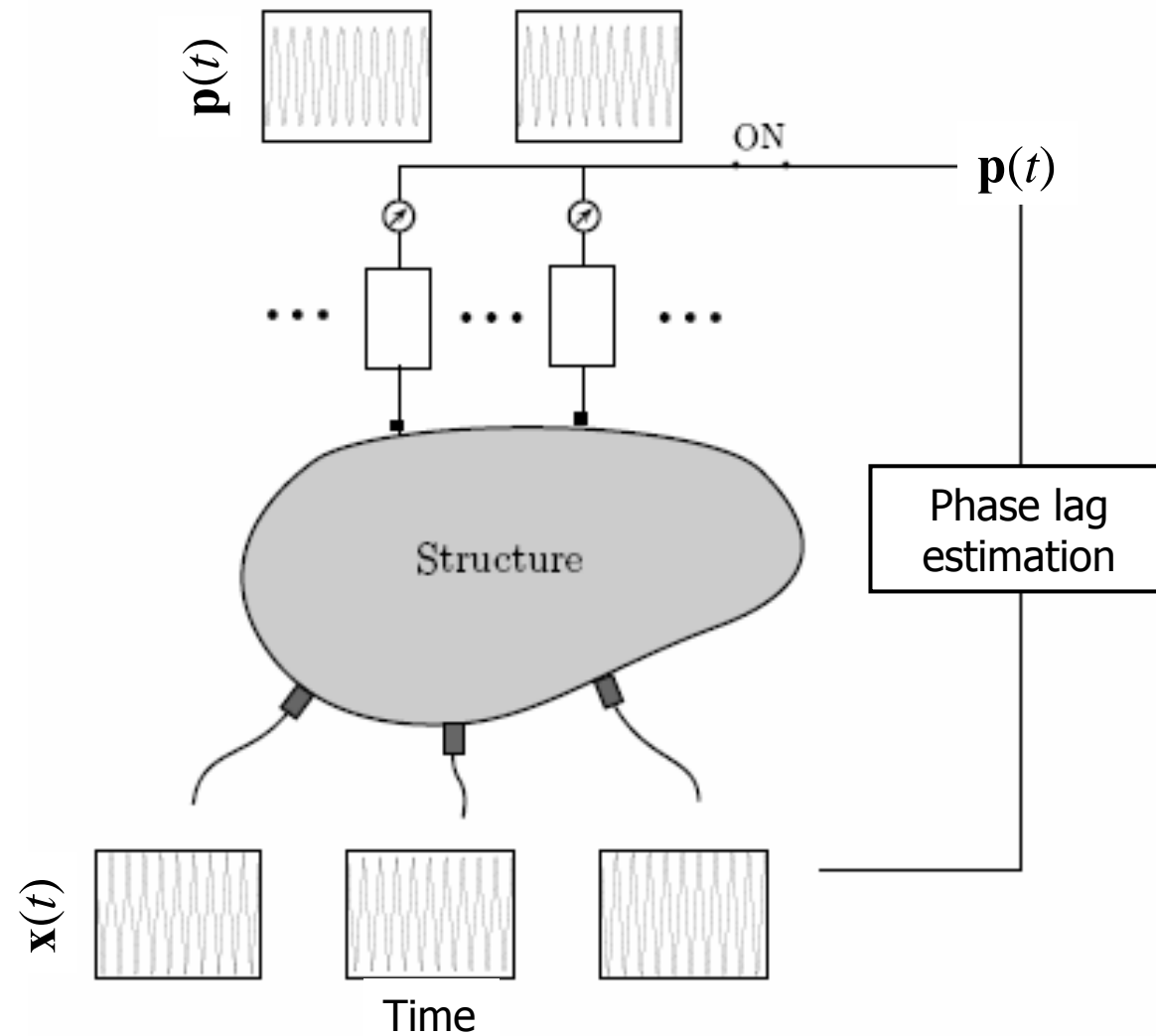
Fundamental properties

1. Forced responses of nonlinear systems at resonance occur in the neighborhood of NNMs [3].
2. According to the invariance property, motions that start out in the NNM manifold remain in it for all time [4].
3. The effect of weak to moderate damping on the transient dynamics is purely parasitic. The free damped dynamics closely follows the NNM of the underlying undamped system [5, 6, 7]

The proposed method for nonlinear EMA relies on a two-step approach that extracts the NNM modal curves and their frequencies of oscillation.

Step 1: NNM Force Appropriation

Objective: isolate a single NNM



Step 1: NNM Force Appropriation

Consider the forced response of a nonlinear structure with linear viscous damping

$$\mathbf{M} \ddot{\mathbf{x}} + \mathbf{C} \dot{\mathbf{x}} + \mathbf{K} \mathbf{x} + \mathbf{f}_{NL}(\mathbf{x}) = \mathbf{p}(t)$$

It is assumed here that the nonlinear restoring force contains only stiffness nonlinearities.

Appropriate excitation

For a given NNM motion $\mathbf{x}_{nnm}(t)$ the equations of motion of the forced and damped system lead to the appropriate excitation


$$\mathbf{p}_{nnm}(t) = \mathbf{C} \dot{\mathbf{x}}_{nnm}(t)$$


This relationship shows that the appropriate excitation is periodic and has the same frequency components as the corresponding NNM motion (i.e., generally including multiharmonic components).

Step 1: NNM Force Appropriation

An NNM motion is now expressed as a Fourier cosine series

$$\mathbf{x}_{nnm}(t) = \sum_{k=1}^{\infty} \mathbf{x}_k^{nnm} \cos(k \omega t)$$

 fundamental pulsation
of the NNM motion

 amplitude vector of the kth harmonic

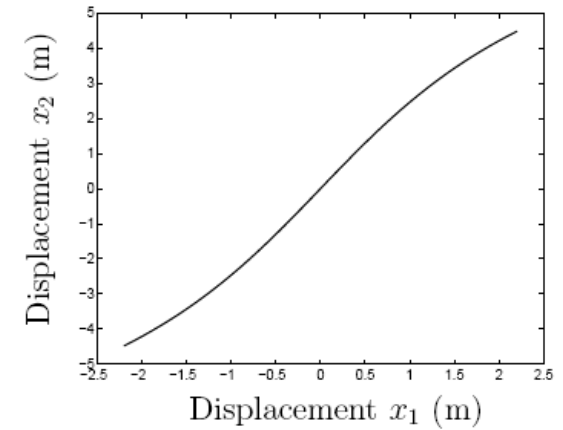
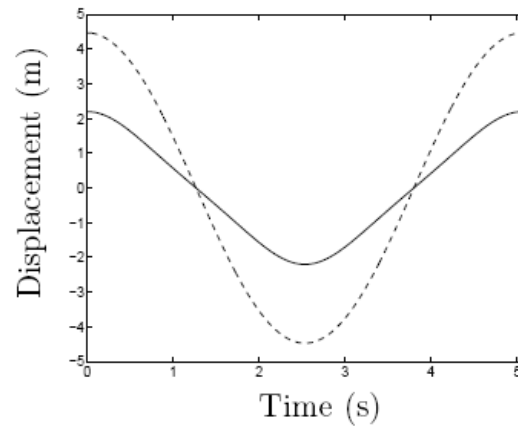
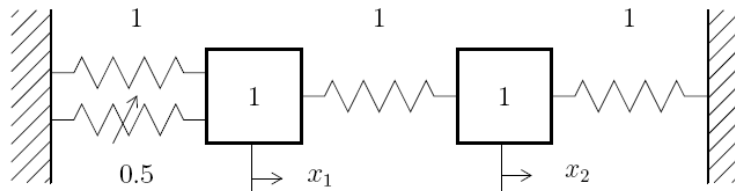
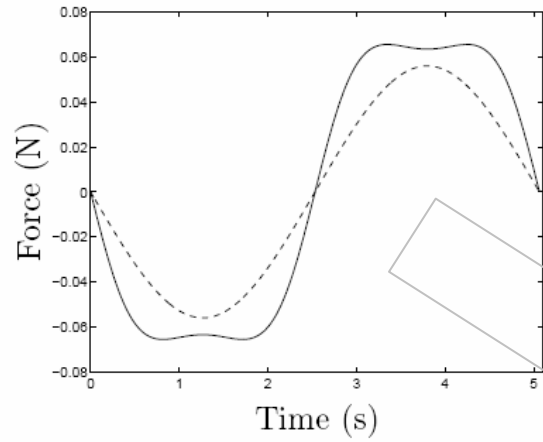
This type of motion is referred to as monophasic NNM motion due to the fact that the displacements of all DOFs reach their extreme values simultaneously.

The appropriate excitation is given by

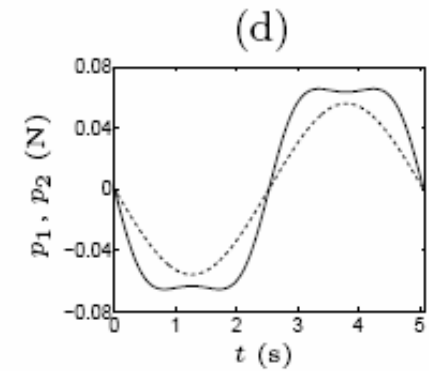
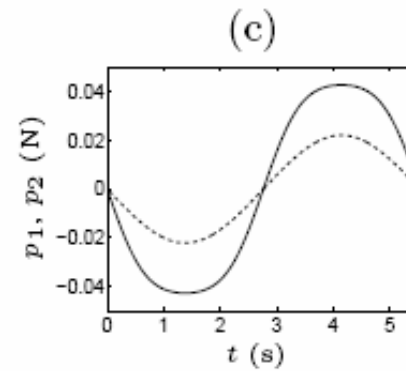
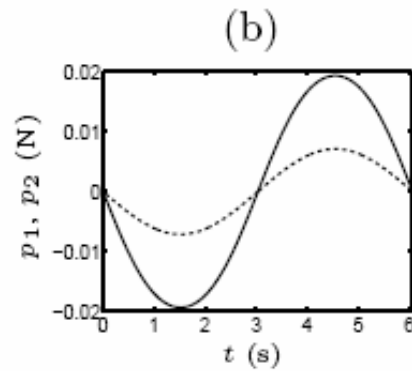
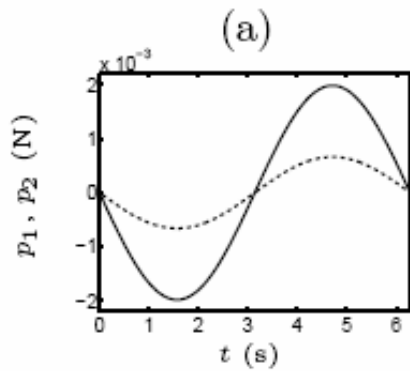
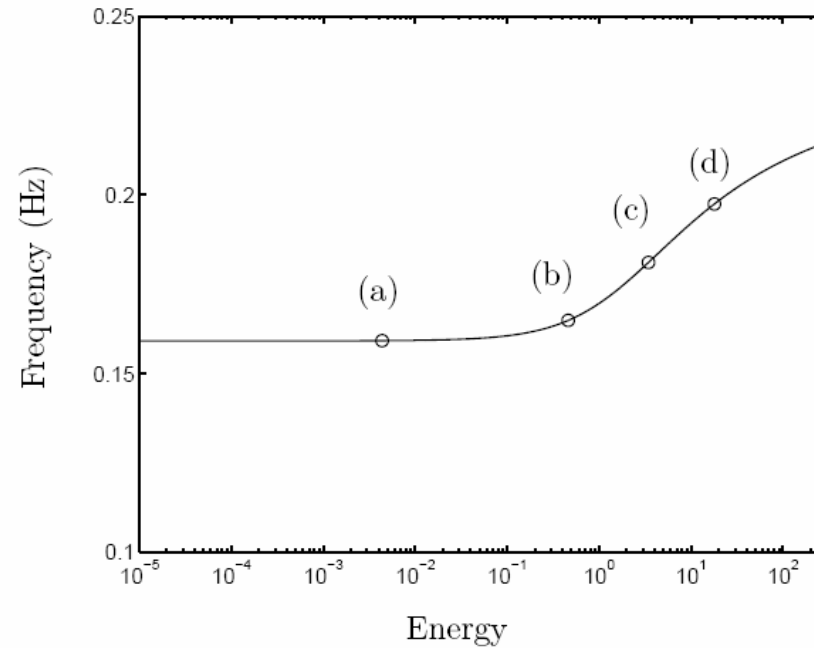
$$\mathbf{p}_{nnm}(t) = - \sum_{k=1}^{\infty} \mathbf{C} \mathbf{x}_k^{nnm} k \omega \sin(k \omega t)$$

→ the excitation of a monophasic NNM is thus characterized by a phase lag of 90° of each harmonics with respect to the displacement response.

Step 1: NNM Force Appropriation



Step 1: NNM Force Appropriation



Step 1: NNM Force Appropriation

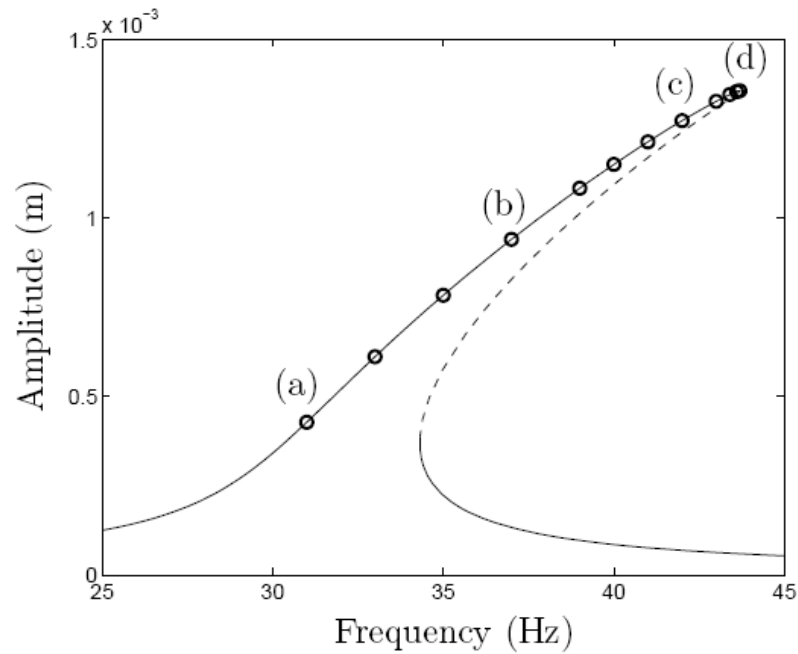
NNM Indicator

Phase lag quadrature criterion:

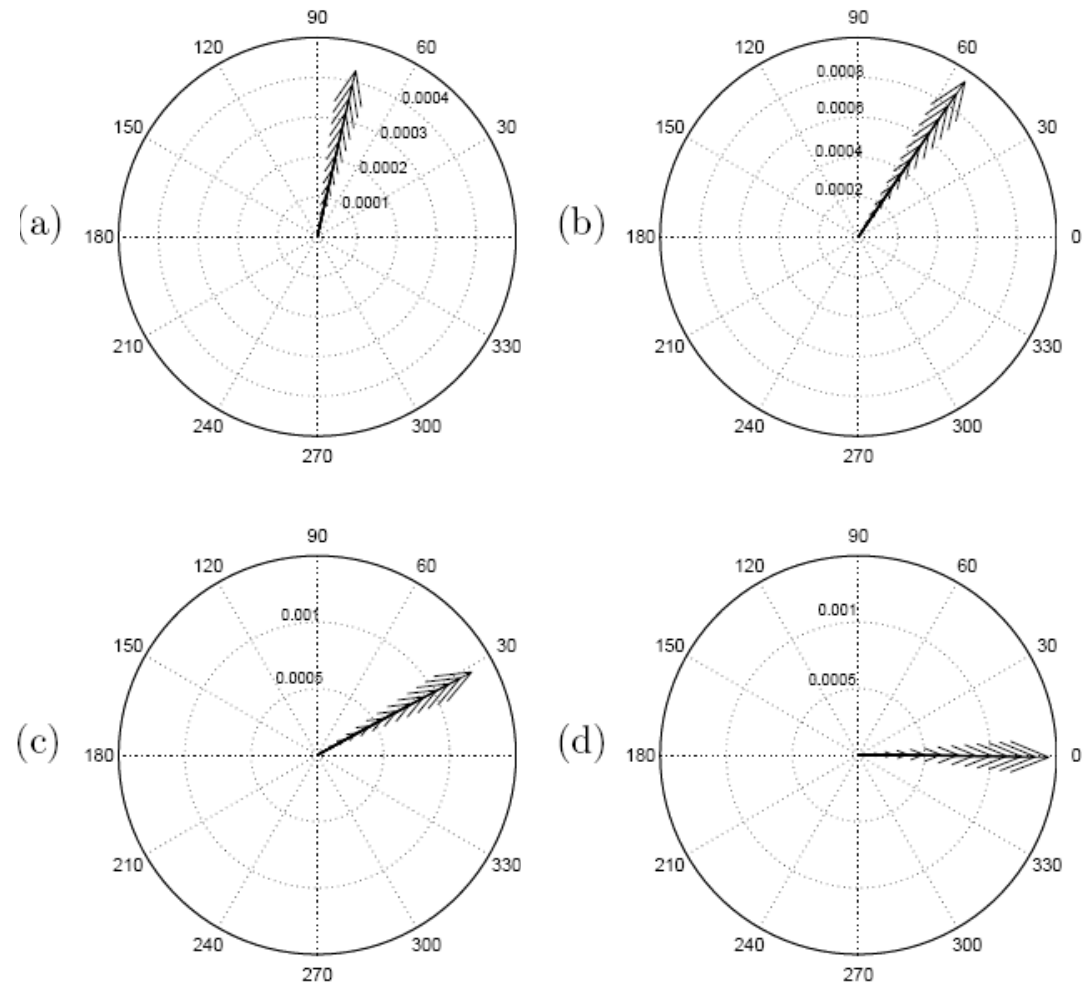
A **linear** structure vibrates according to one of its **LNMs** if the degrees of freedom have a **phase lag of 90°** with respect to the excitation.

A **nonlinear** structure vibrates according to one of its **NNMs** if the degrees of freedom have a **phase lag of 90°** with respect to the excitation.

Step 1: NNM Force Appropriation



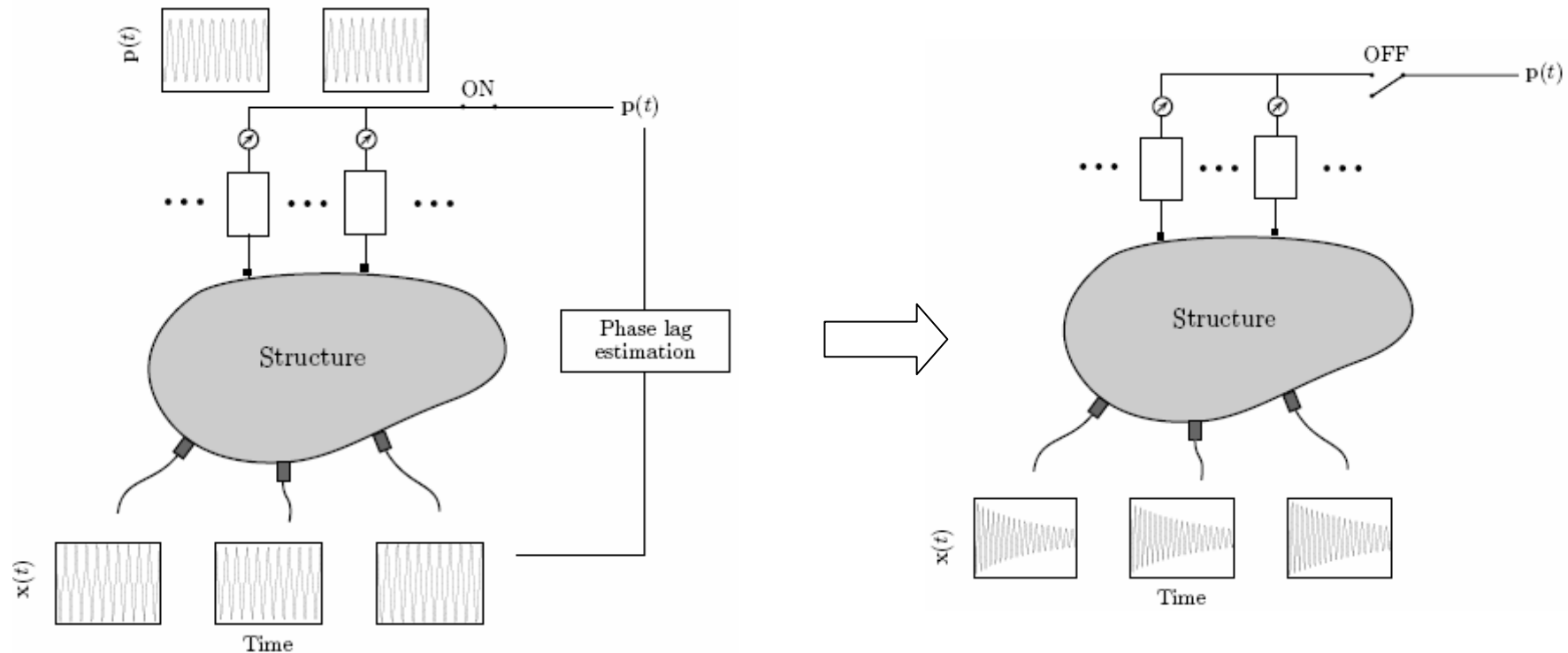
The excitation phase is 90°



Step 2: NNM Free Decay Identification

Turn off the excitation and track the NNM according to the invariance principle:

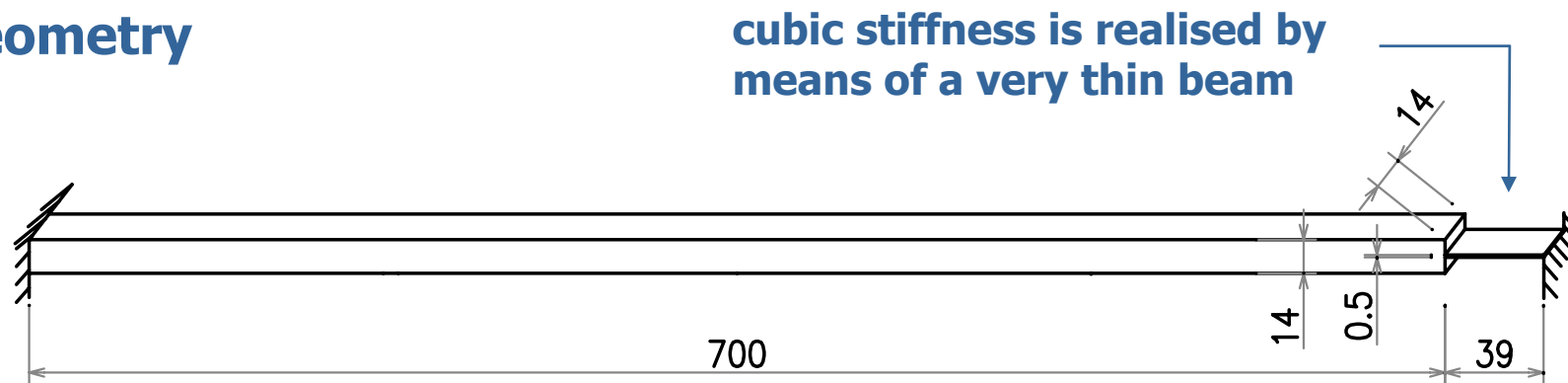
« If the motion is initiated on one specific NNM, the remaining NNMs remain quiescent for all time. »



Nonlinear EMA (Illustrative Example)

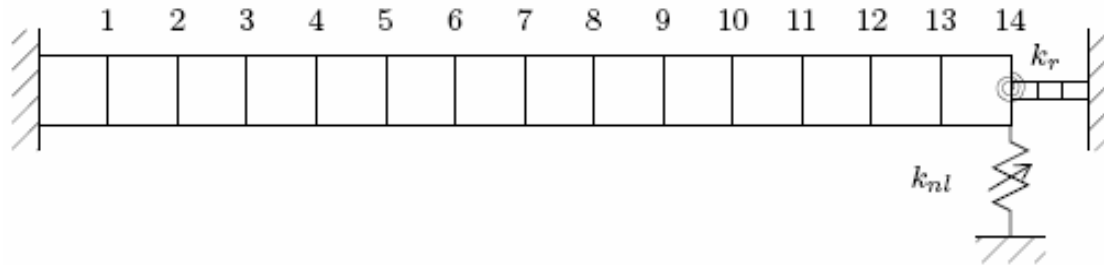
Numerical experiments of a nonlinear beam (defined as benchmark in the framework of the European COST Action F3 « Structural Dynamics » [8]).

Geometry



For weak excitation, the system behaviour may be considered as linear. When the excitation level increases, the thin beam exhibits large displacements and a nonlinear geometric effect is activated resulting in a stiffening effect at the end of the main beam.

Finite Element model

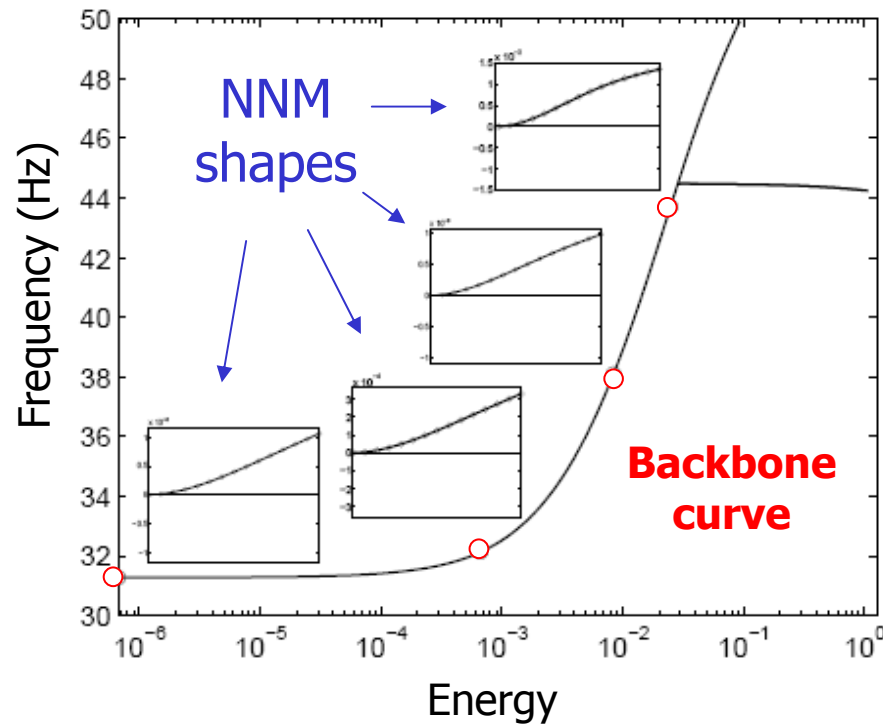


The thin beam is represented by two equivalent grounded springs: one in translation ($k_{nl} + k_l$) and one in rotation (k_r).

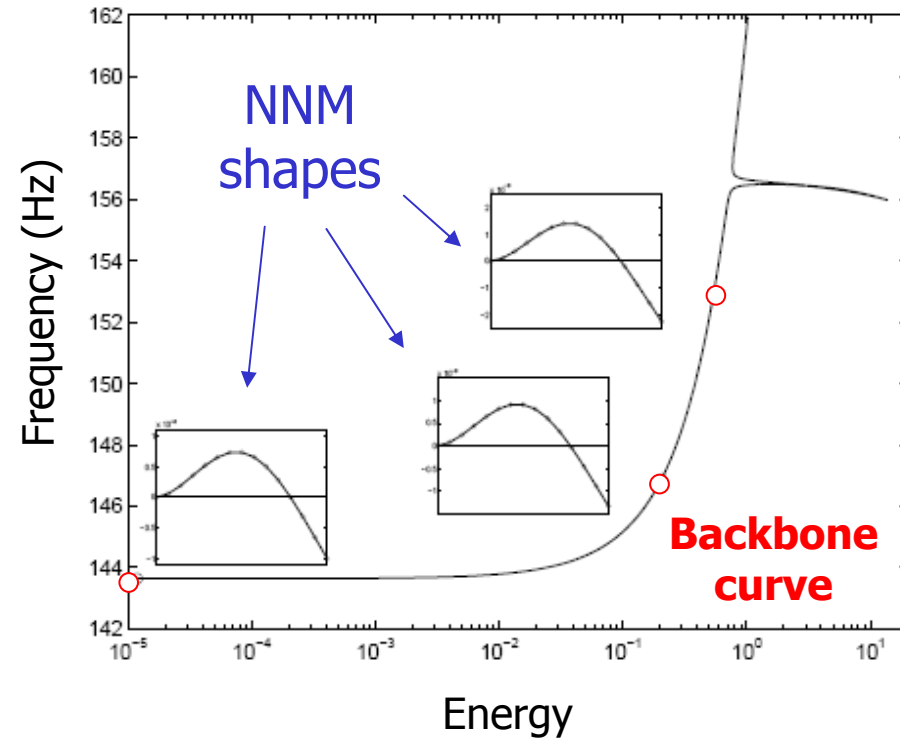
Young's modulus (N/m ²)	Density (kg/m ³)	Nonlinear parameter k_{nl} (N/m ³)
$2.05 \cdot 10^{11}$	7800	$8 \cdot 10^9$

Theoretical frequency-energy plots

First NNM



Second NNM



Simulated experiments

Linear proportional damping is considered.

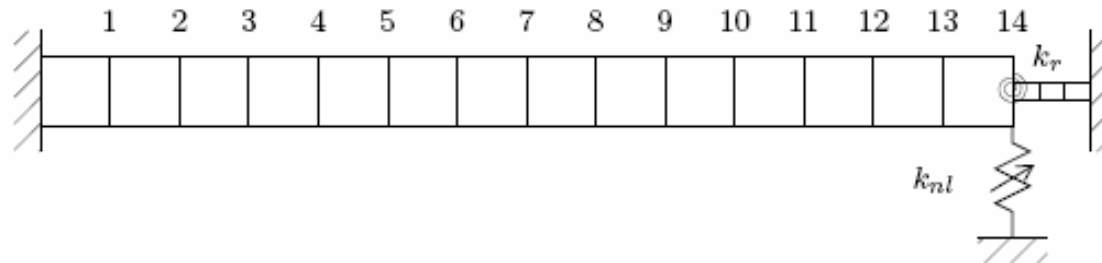
$$\mathbf{C} = 3 \cdot 10^{-7} \mathbf{K} + 5 \mathbf{M}$$

It corresponds to moderate damping; for instance, the modal damping ratio is equal to 1.28% for the first linear normal mode.

Imperfect force appropriation

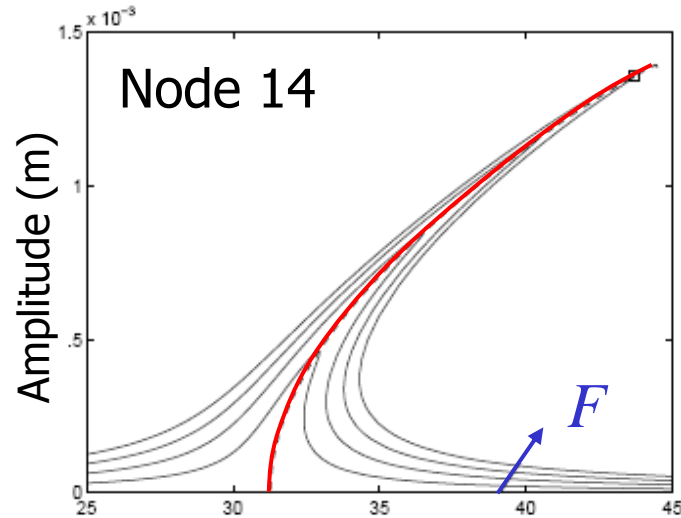
From a practical viewpoint, it is useful to study the quality of imperfect force appropriation consisting of a single-point mono-harmonic excitation, i.e., using a single shaker with no harmonics of the fundamental frequency.

The harmonic force $p(t) = F \sin(\omega t)$ is applied to node 4 of the main beam.



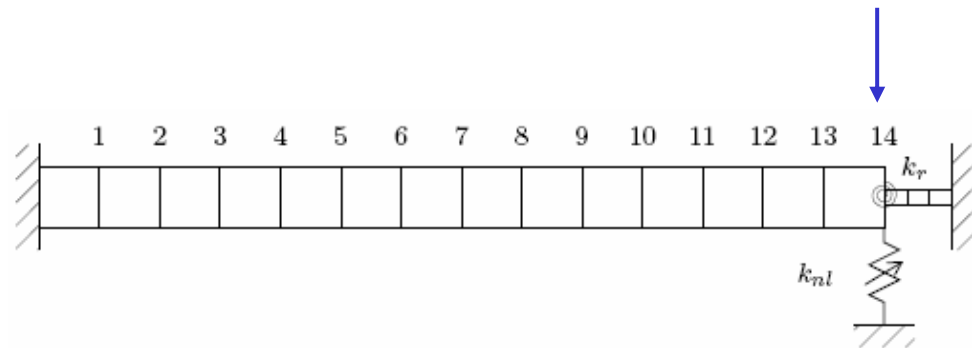
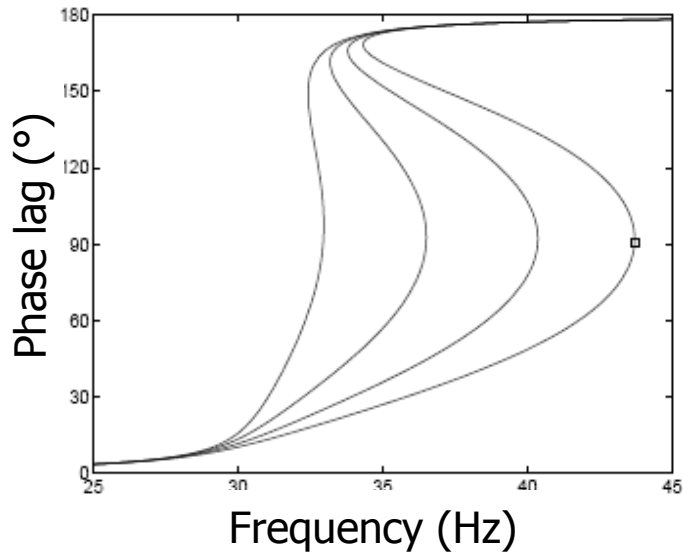
Nonlinear EMA (Illustrative Example)

Nonlinear forced frequency responses to the first resonant frequency

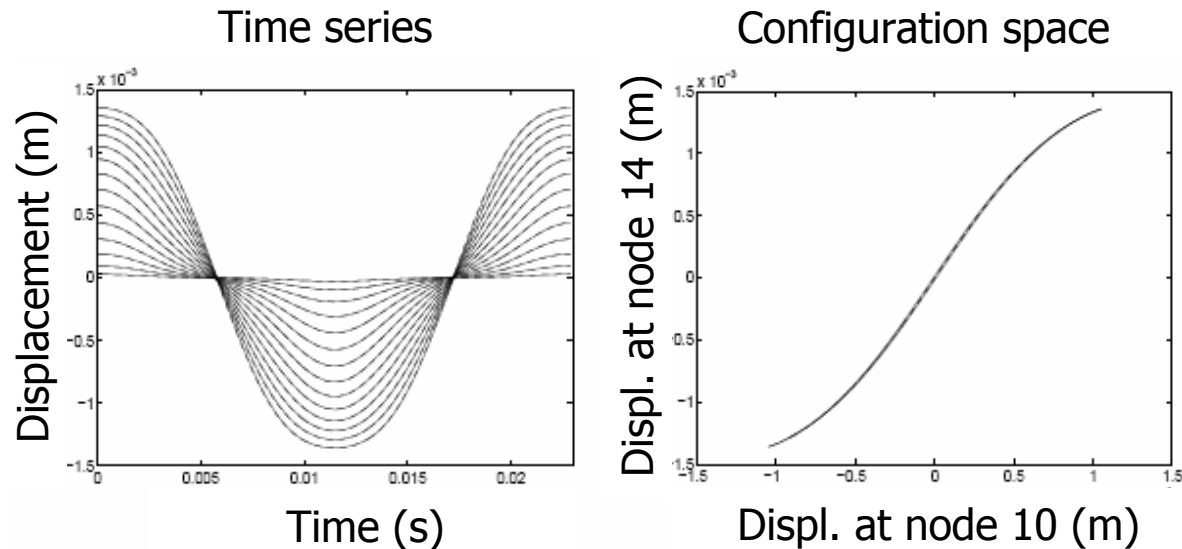


backbone of the first undamped NNM

4 different forcing amplitudes: 1N, 2N, 3N, 4N.



$$F = 4N$$



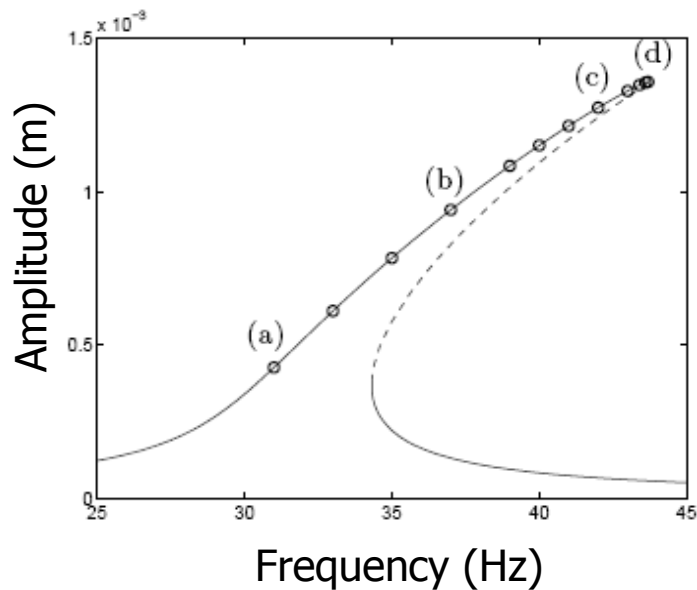
Observations

- The phase lag quadrature criterion is fulfilled close to resonant frequencies.
- Forced responses at resonance occur in the neighbourhood of NNMs.
- Imperfect appropriation can isolate the NNM of interest (the beam has well-separated modes).

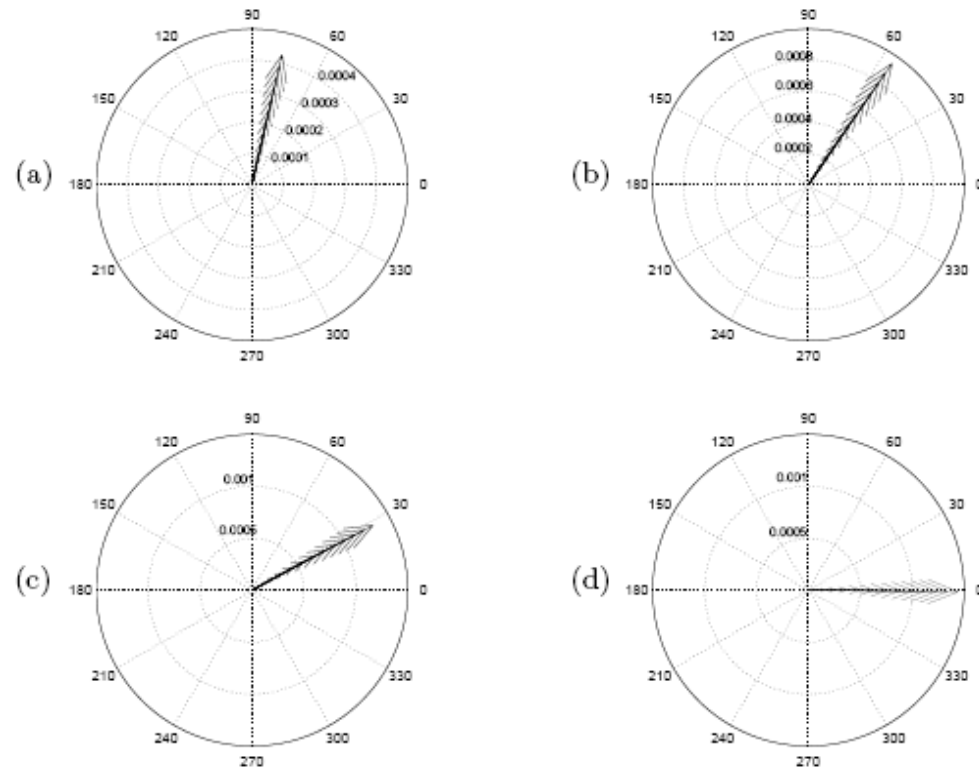
These findings also hold for the second beam NNM.

Nonlinear EMA (Illustrative Example)

Stepped sine excitation procedure for carrying out the NNM force appropriation ($F = 4\text{N}$) of the damped nonlinear beam.



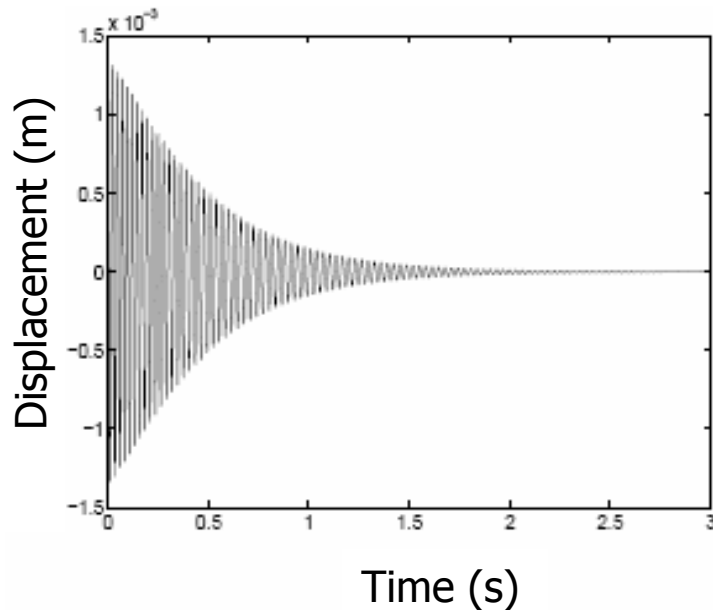
Responses along the branch close to the first resonance



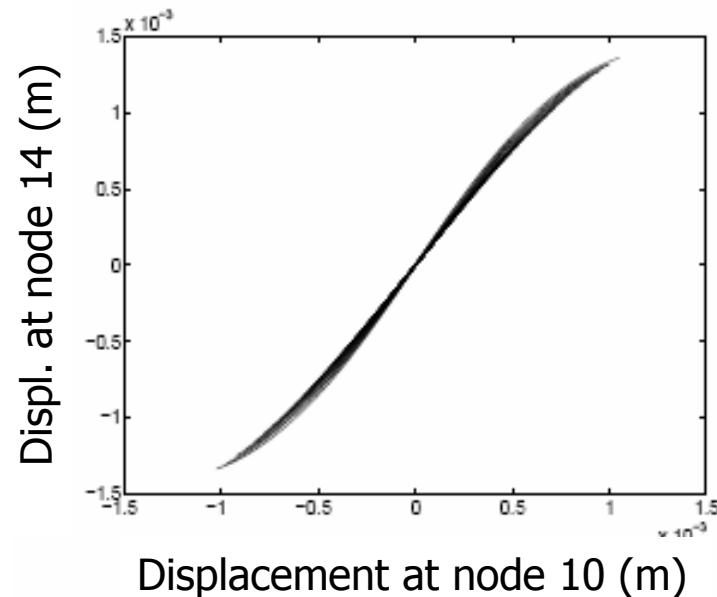
Phase scatter diagrams of the complex Fourier coefficients of the displacements corresponding to the fundamental frequency for the responses (a), (b), (c) and (d).

NNM free decay identification

Free response of the damped nonlinear beam initiated from the imperfect appropriated forced response



Time series of the displacement at the tip of the main beam (node 14).

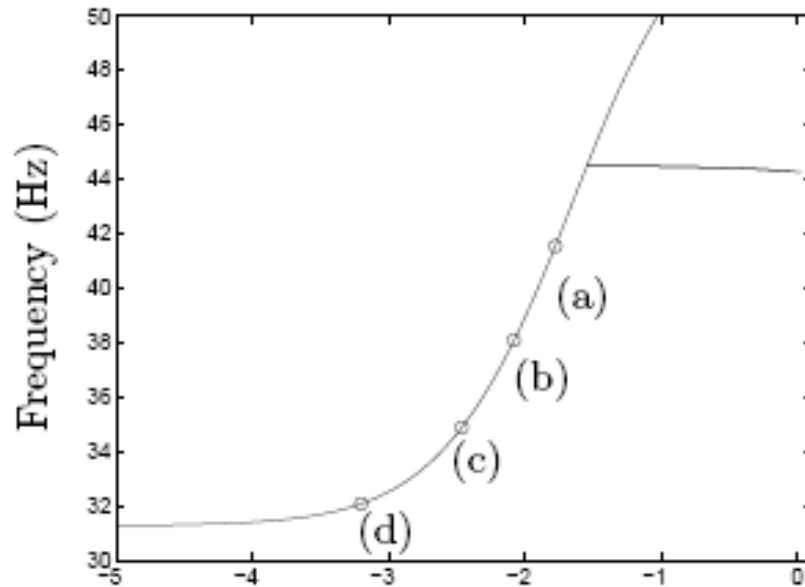


Motion in the configuration space composed of the displacements at nodes 10 and 14.

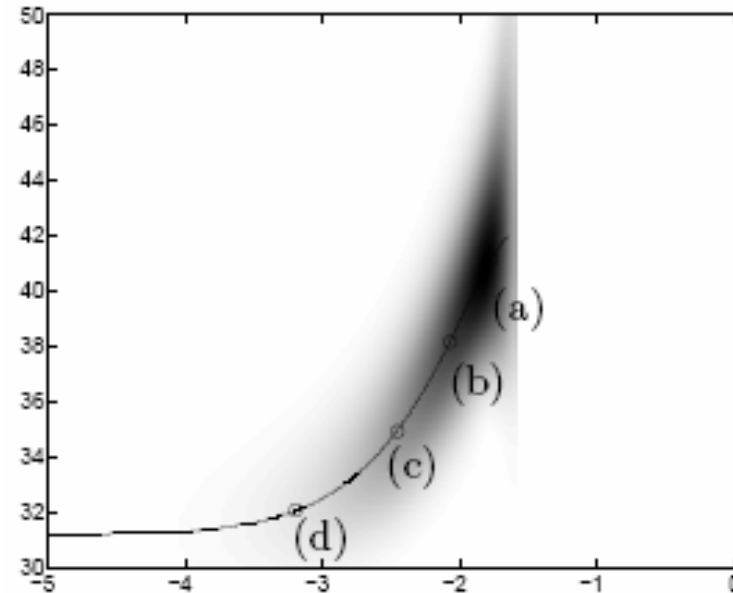
Nonlinear EMA (Illustrative Example)

Frequency-energy plot of the first NNM of the nonlinear beam.

Theoretical FEP

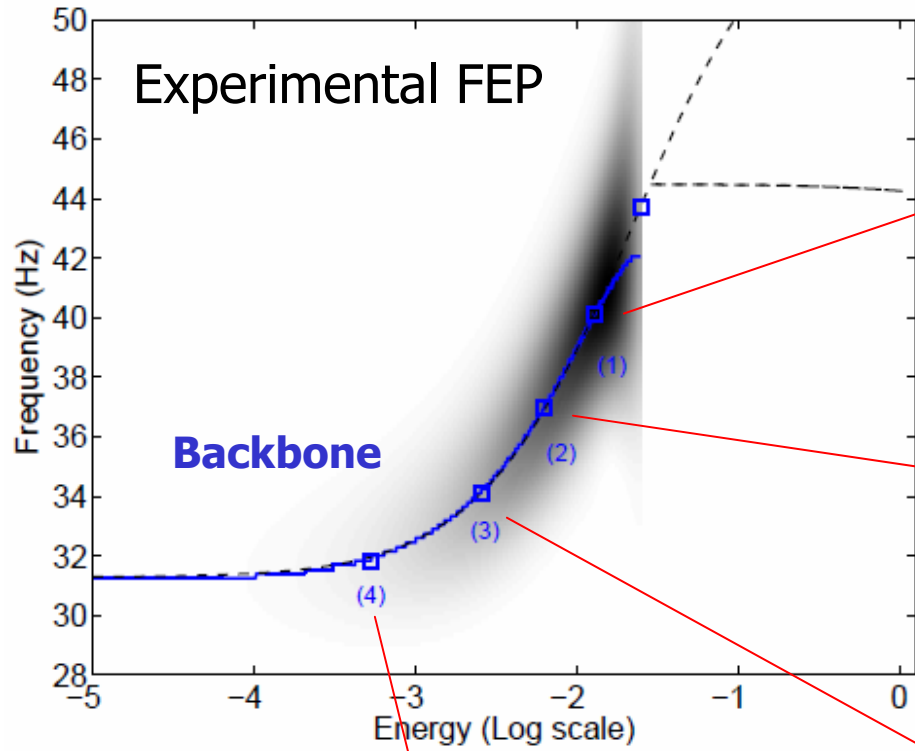


Experimental FEP



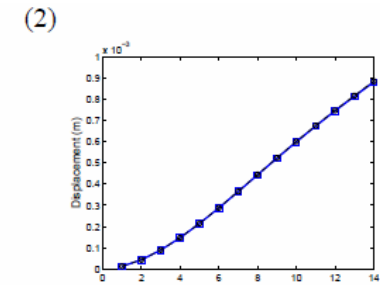
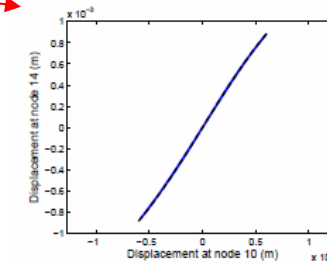
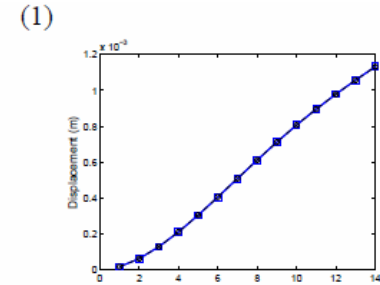
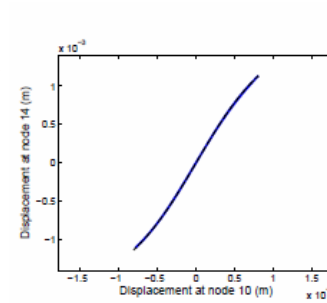
This FEP was calculated from the time series of the free damped response using the CWT. The solid line is the ridge of the transform.

Nonlinear EMA (Illustrative Example)

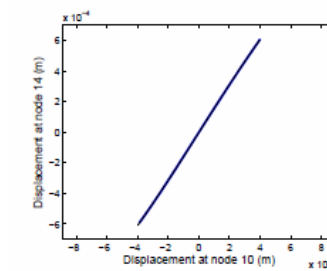
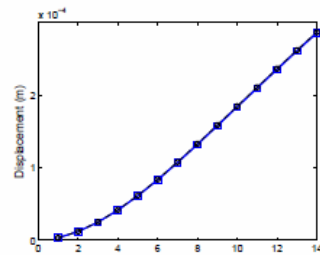
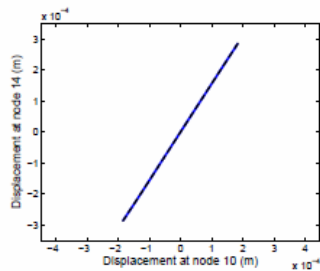


Modal curves

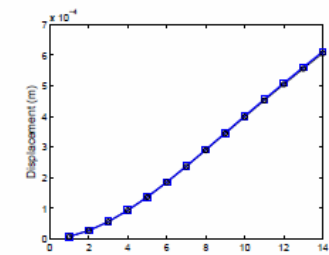
Modal shapes



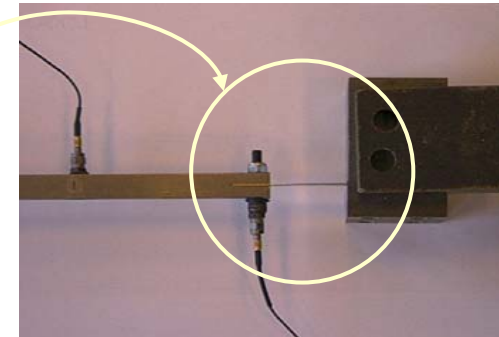
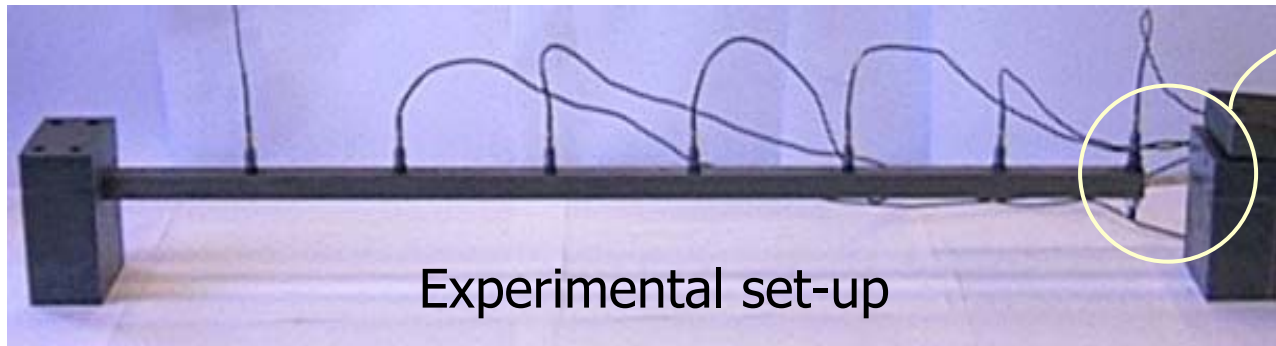
(4)



(3)



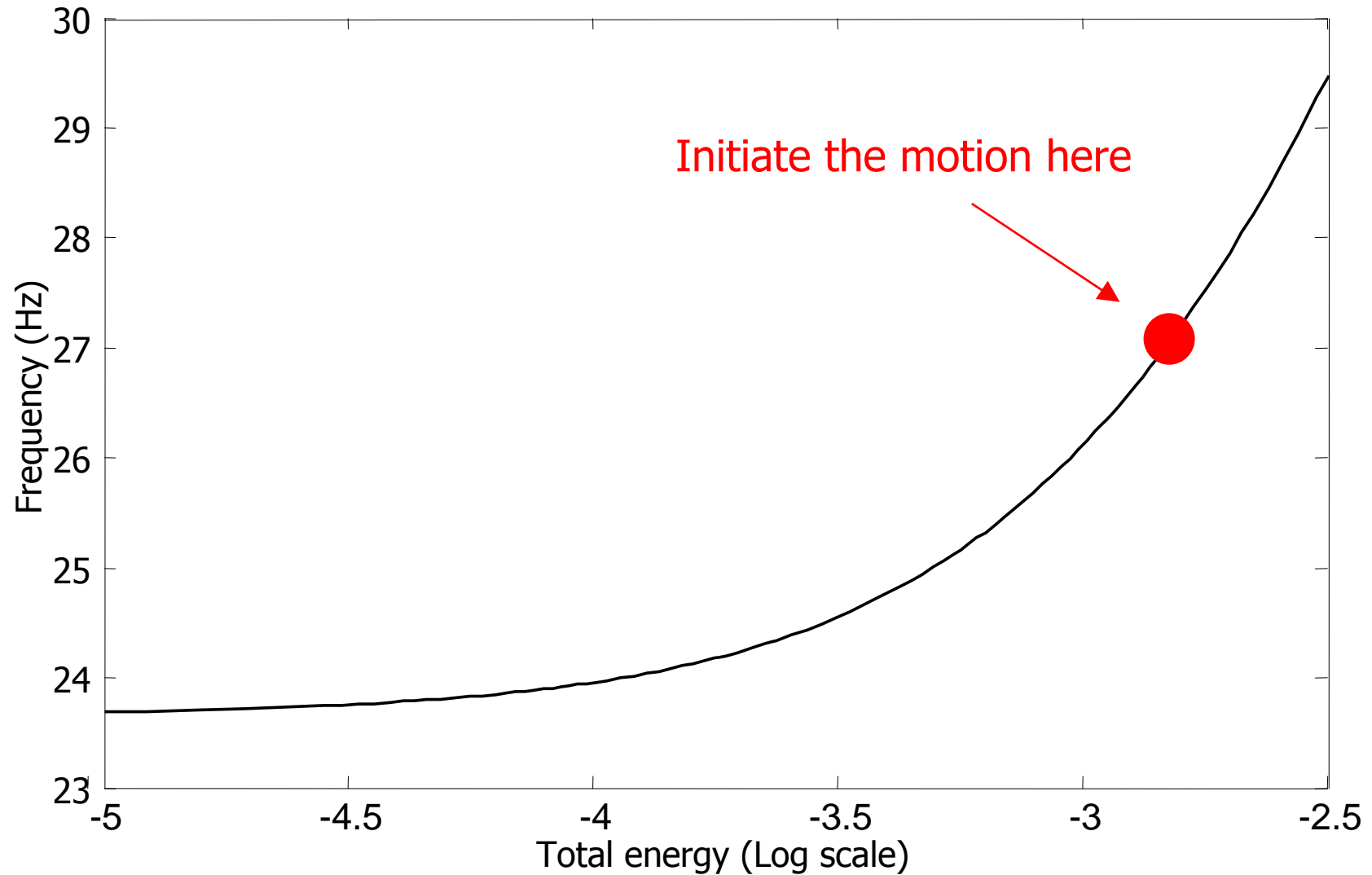
Benchmark of the European COST Action F3 « Structural Dynamics ».



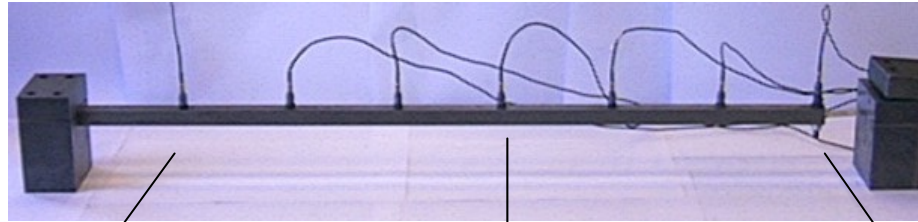
Test conditions

- Harmonic excitation of the nonlinear beam.
- Response measured using seven accelerometers.
- Very preliminary results.

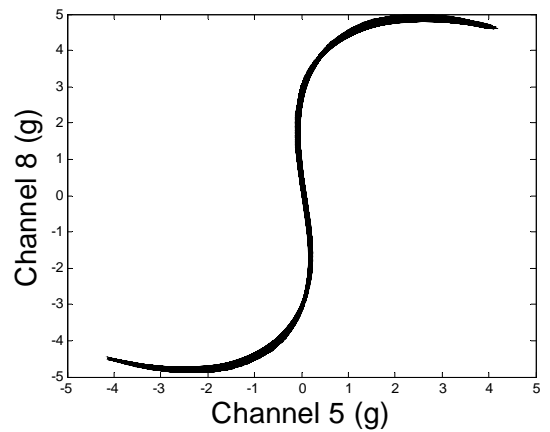
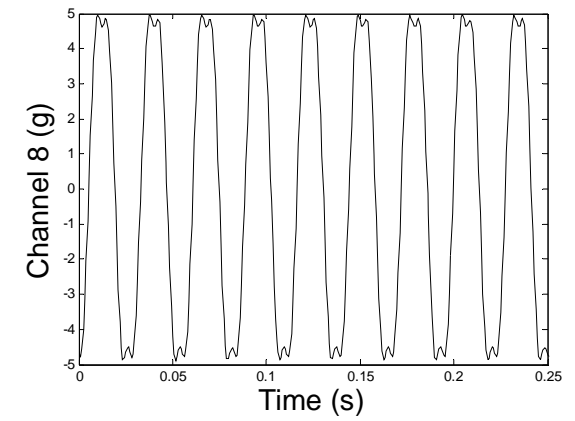
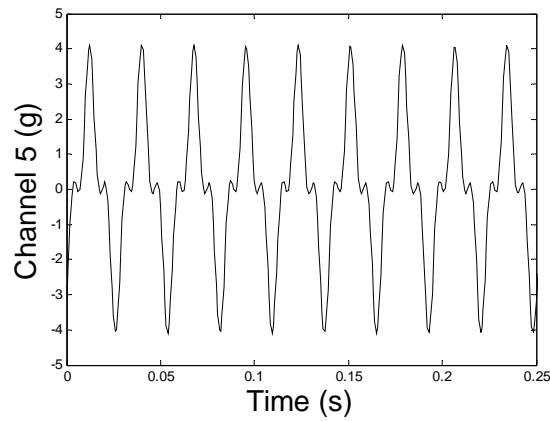
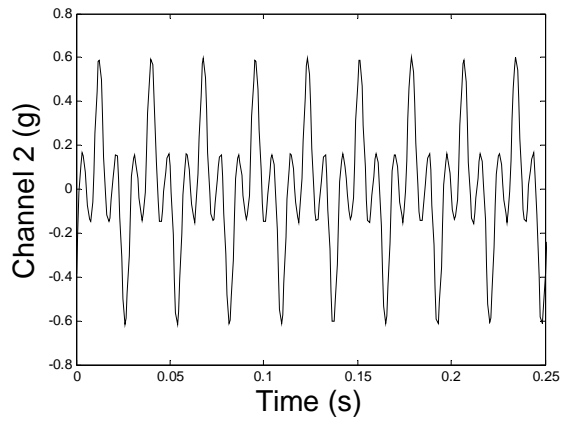
Excitation of the 1st NNM of the beam



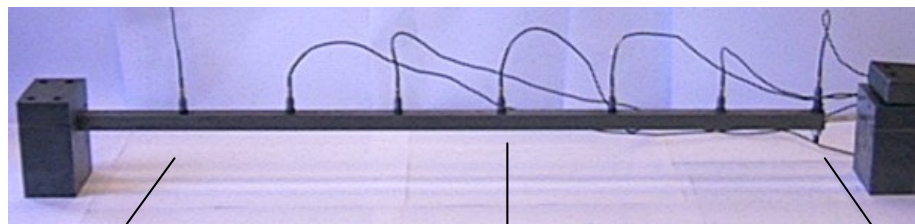
Experimental Demonstration



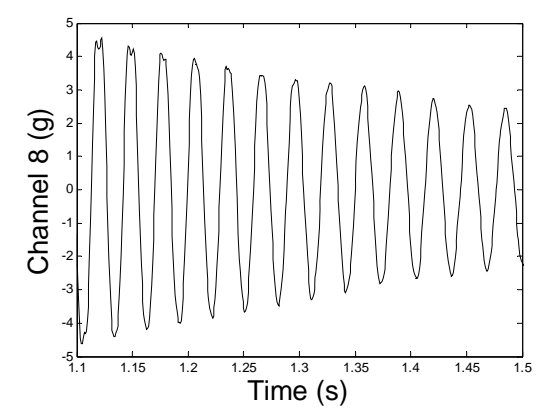
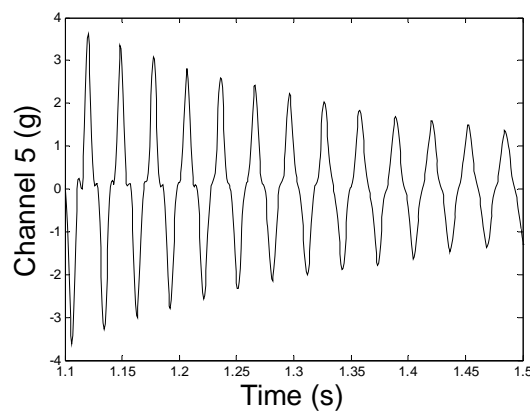
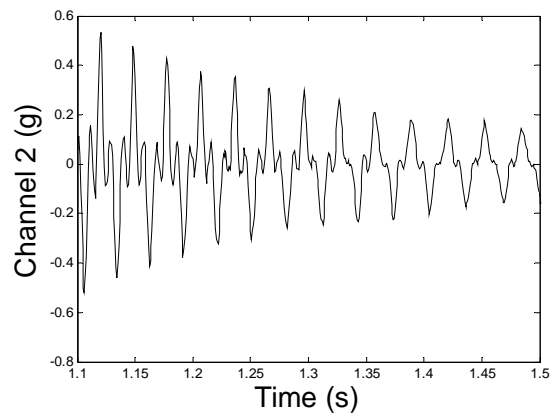
Sustained harmonic
excitation

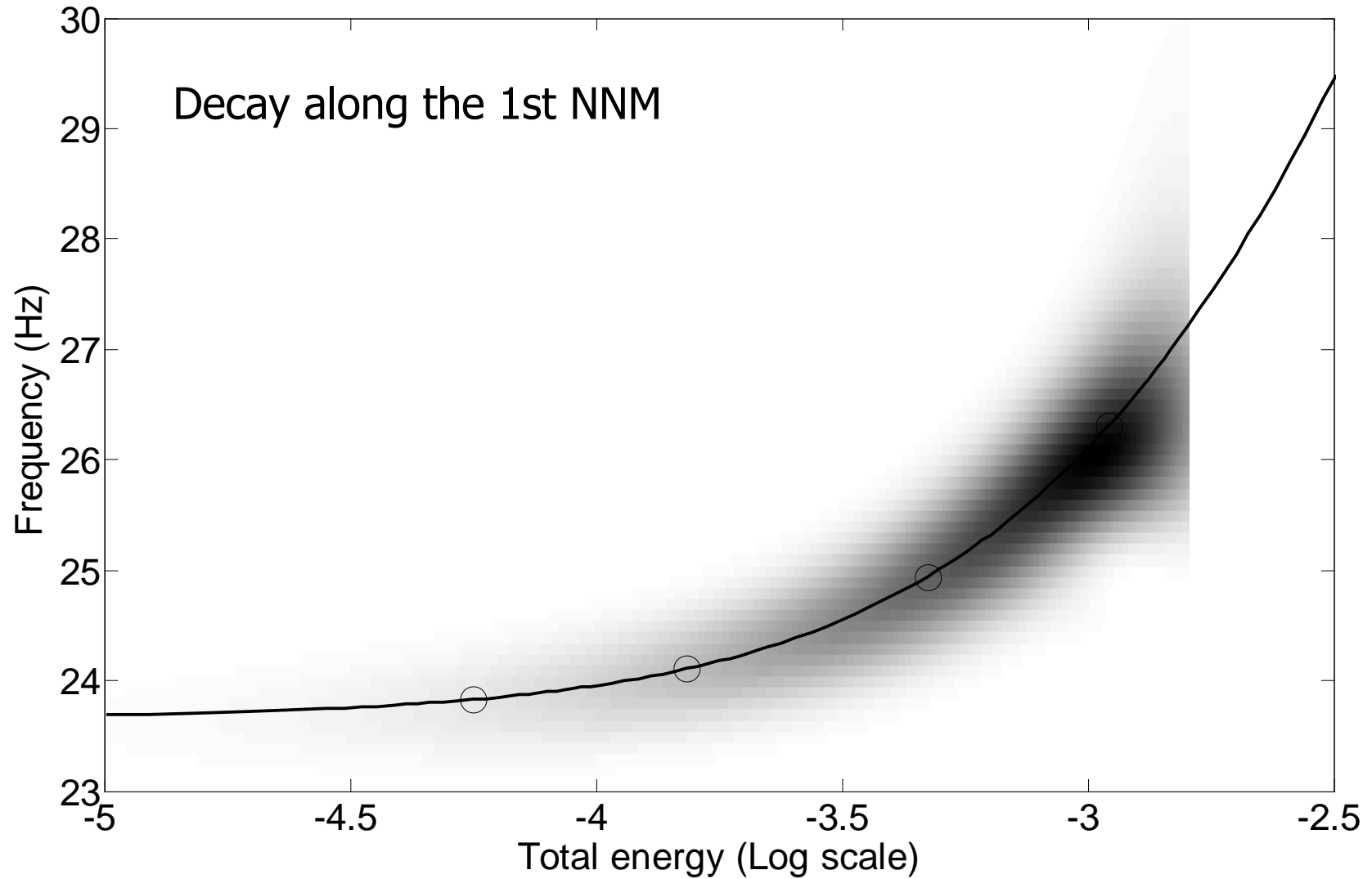


Turn Off the Shaker



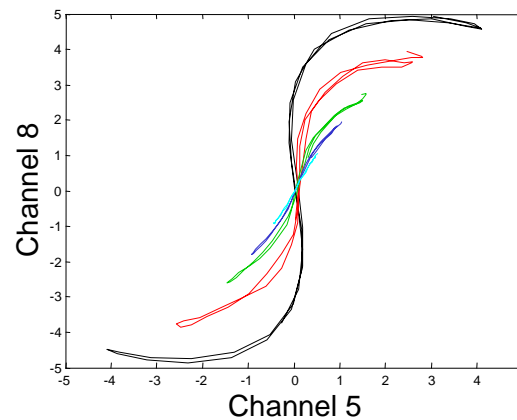
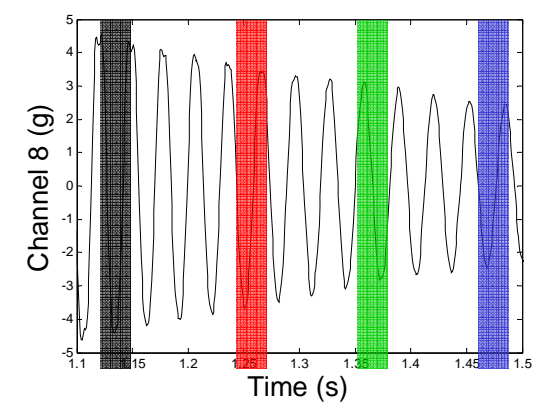
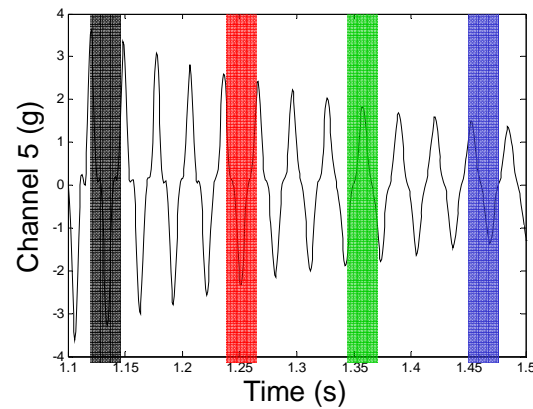
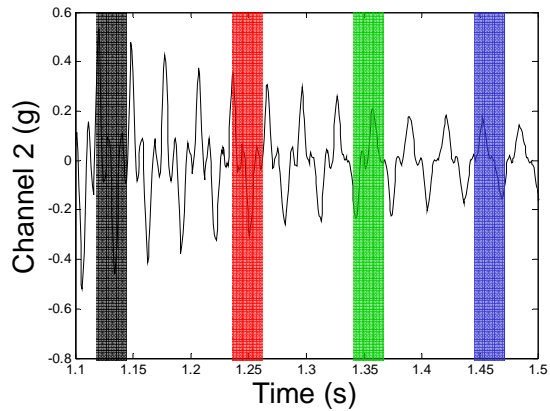
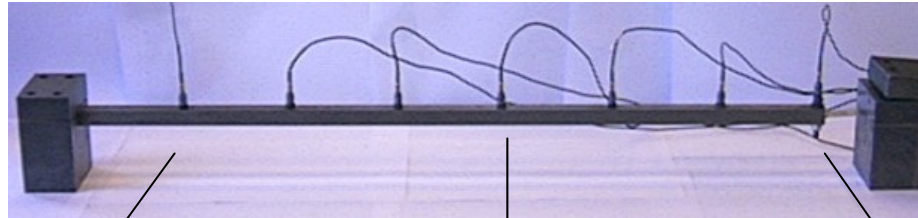
Burst sine





Experimental Demonstration

NNM Extraction



1st NNM of the beam at
different energy levels



1. Introduction

2. Theoretical Modal Analysis of Nonlinear Systems

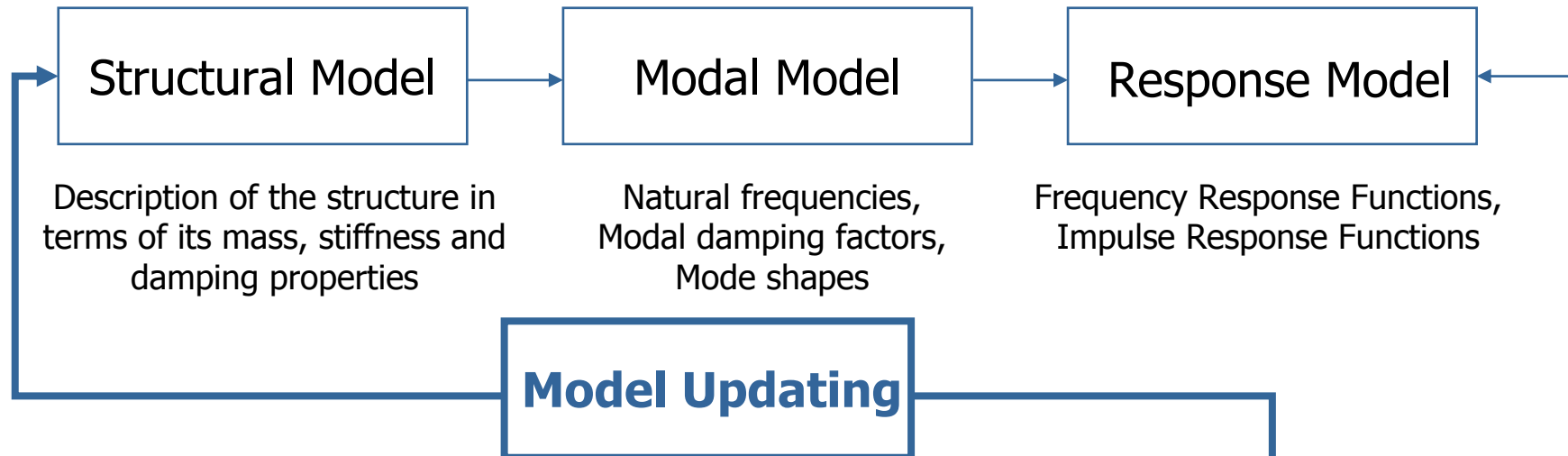
3. Nonlinear Experimental Modal Analysis

4. Model Parameter Estimation Techniques

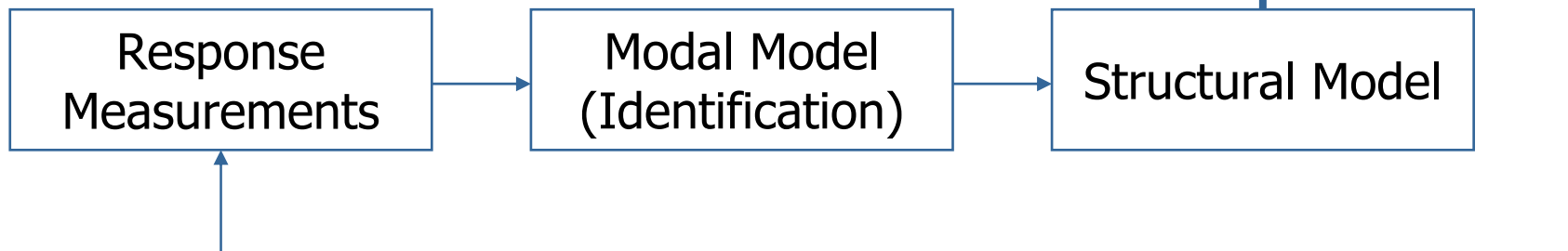
- **Parameter Estimation Using POD**
- **Parameter Estimation Using Nonlinear EMA**

5. Concluding Remarks

Theoretical Approach – Direct Problem



Experimental Approach – Inverse Problem





Assumption

- The linear counterpart of the structure is known (updated).

Methodology

- Estimation of nonlinear parameters only (which will be based on FE updating techniques).

Parameters for Model Updating (Crucial step!)

The number of parameters :

- should be kept small to avoid problems of ill-conditioning,
- should be chosen with the aim of correcting recognised features in the model.

→ requires physical insight → leads to **knowledge-based** models.

Mathematical Background

Consider the general equation governing the dynamics of a structure

$$\mathbf{M} \ddot{\mathbf{x}} + \mathbf{C} \dot{\mathbf{x}} + \mathbf{K} \mathbf{x} + \mathbf{f}_{NL}(\mathbf{x}, \dot{\mathbf{x}}) = \mathbf{g}(t)$$



Vector of nonlinear forces

Step 1: definition of a penalty function involving modal features of the system (residual between analytical and measured dynamic behaviour)

The measured quantities may be assembled into a measurement vector \mathbf{z} .

The vector of modal features \mathbf{z} depends on parameters \mathbf{p}

$$\mathbf{z} = \mathbf{z}(\mathbf{p})$$

The choice of parameters is a crucial step in model updating. For nonlinear identification purposes, we will assume that a *knowledge-based* model exists (the physically meaningful model and the associated parameters are supposed to be known).

Penalty function methods are based on the Taylor series expansion of the modal data in terms of the unknown parameters

$$\mathbf{z} = \mathbf{z}(\mathbf{p}_0) + \left[\frac{\partial \mathbf{z}}{\partial \mathbf{p}} \right]_{\mathbf{p}=\mathbf{p}_0} (\mathbf{p} - \mathbf{p}_0) + O(\mathbf{p}^2)$$

↓
↳

Sensitivity matrix
Initial estimation of the parameters

This expansion is often limited to the first two terms.


Define the weighted penalty function

$$J = \boldsymbol{\varepsilon}^T \mathbf{W} \boldsymbol{\varepsilon}$$

 **Positive definite
weighting matrix**

where $\boldsymbol{\varepsilon} = \Delta \mathbf{z} - \mathbf{S} \Delta \mathbf{p}$ is the error in the predicted measurements.

$$\mathbf{S} = \left[\frac{\partial \mathbf{z}}{\partial \mathbf{p}} \right]$$

 is the sensitivity matrix.

Minimising J with respect to $\Delta \mathbf{p}$ leads to

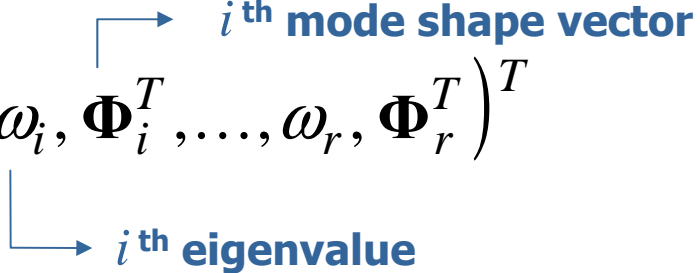
$$\Delta \mathbf{p} = \left(\mathbf{S}^T \mathbf{W} \mathbf{S} \right)^{-1} \mathbf{S}^T \mathbf{W} \Delta \mathbf{z}$$

With the assumption that the number of measurements is larger than the number of parameters, the matrix $\mathbf{S}^T \mathbf{W} \mathbf{S}$ is square and hopefully full rank.

Definition of the measurement vector \mathbf{z} containing the modal features.

- In the case of **linear** systems

$$\mathbf{z}^T = \left(\omega_1, \Phi_1^T, \dots, \omega_i, \Phi_i^T, \dots, \omega_r, \Phi_r^T \right)^T$$



- In the case of **nonlinear** systems
 - Proper Orthogonal Decomposition
 - Nonlinear Modal Analysis

Linear Systems

$$\mathbf{M} \ddot{\mathbf{x}} + \mathbf{C} \dot{\mathbf{x}} + \mathbf{K} \mathbf{x} = \mathbf{p}(t)$$

Deterministic approach

Eigenvalue problem:

$$(\mathbf{K} - \omega^2 \mathbf{M}) \Phi = \mathbf{0}$$

Response:

$$\mathbf{x}(t) = \sum_{i=1}^n \eta_i(t) \Phi_{(i)}$$

Spatial information

Natural frequencies

$$\eta_i = A_i \cos(\omega_i t) + B_i \sin(\omega_i t)$$

Nonlinear Systems

$$\mathbf{M} \ddot{\mathbf{x}} + \mathbf{C} \dot{\mathbf{x}} + \mathbf{K} \mathbf{x} + \mathbf{f}_{NL}(\mathbf{x}, \dot{\mathbf{x}}) = \mathbf{p}(t)$$

Statistical approach

Proper Orthogonal Decomposition:

$$\mathbf{X} = \mathbf{U} \Sigma \mathbf{V}^T$$

Response:

$$\mathbf{x}(t) = \sum_{j=1}^n a_j(t) \mathbf{u}_{(j)}$$

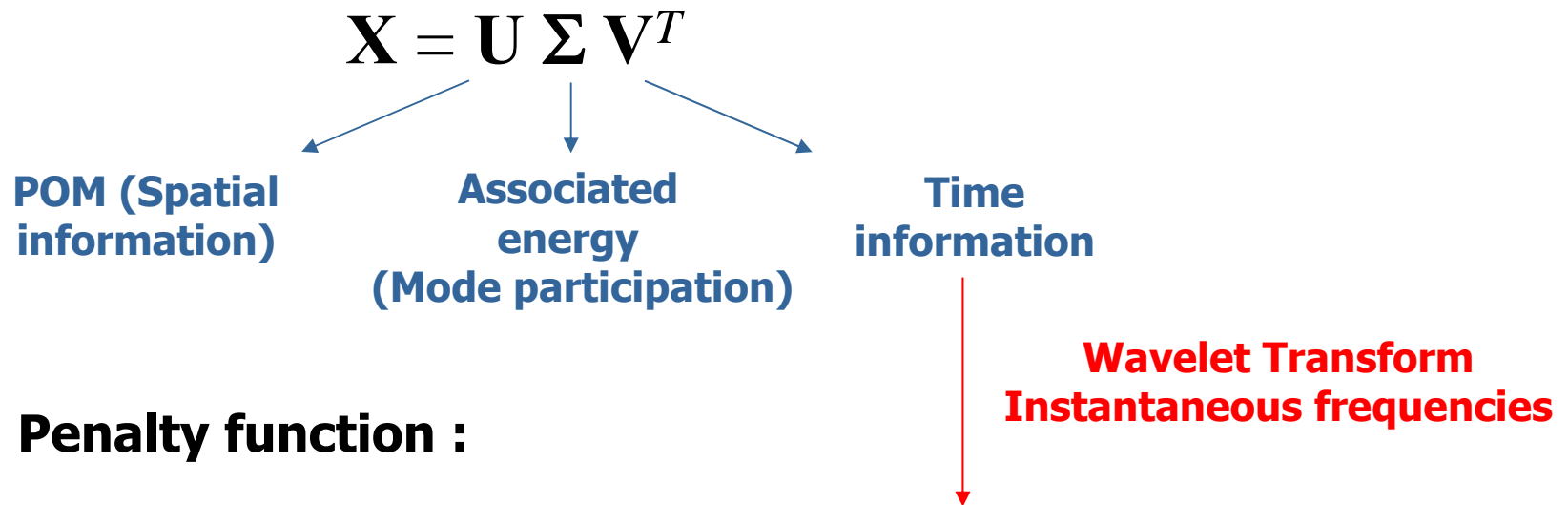
Spatial information

Time information
→ Instantaneous frequencies

Parameter Estimation Using POD

Principle of the method :

Minimise the residuals between the bi-orthogonal decompositions of the measured and simulated data.



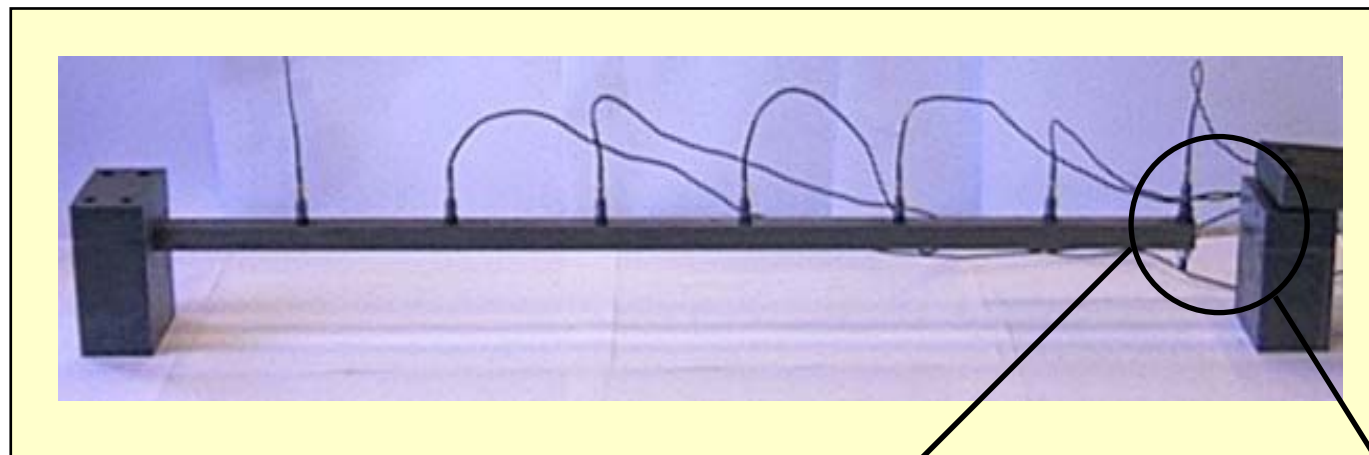
Penalty function :

$$J = \sum_i \sum_j (\Delta U_{ij})^2 + \sum_j (\Delta \Sigma_{jj})^2 + \sum_j \sum_k (\Delta V_{jk})^2$$

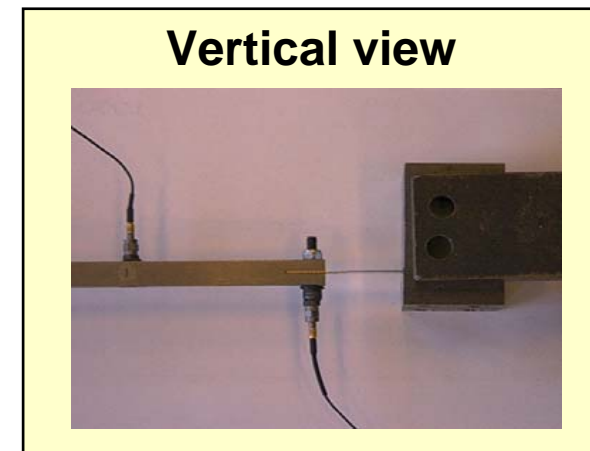
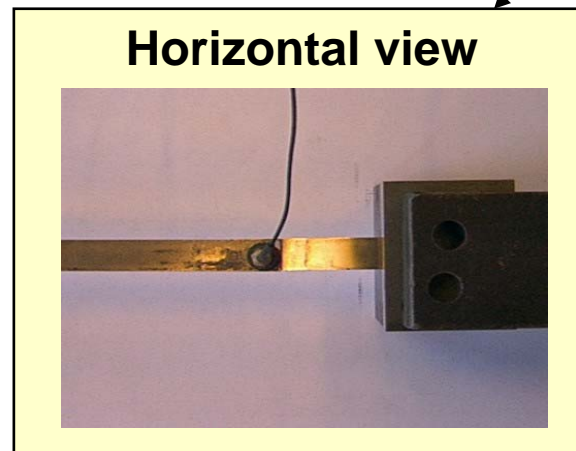
→ Selection of the POMs with the highest POV

Parameter Estimation Using POD

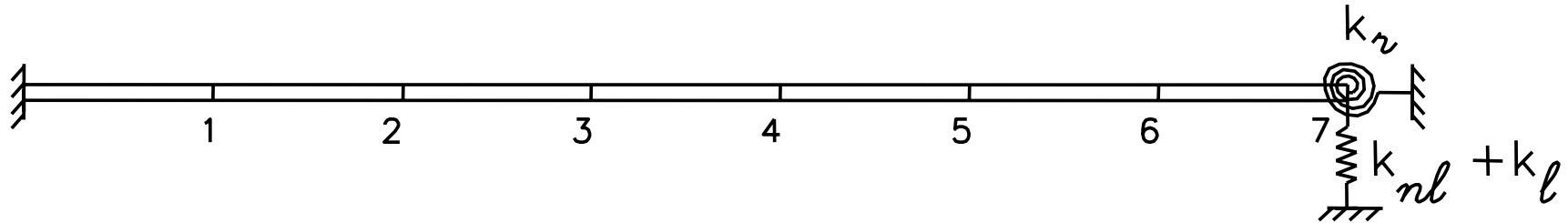
Benchmark of the European COST Action F3 « Structural Dynamics »



Experimental
set-up



Finite Element model



The nonlinear stiffening effect of the thin beam is modelled by a nonlinear function in displacement of the form:

$$f_{nl} = A |x|^\alpha \text{sign}(x)$$

where A and α are nonlinear parameters to be identified.

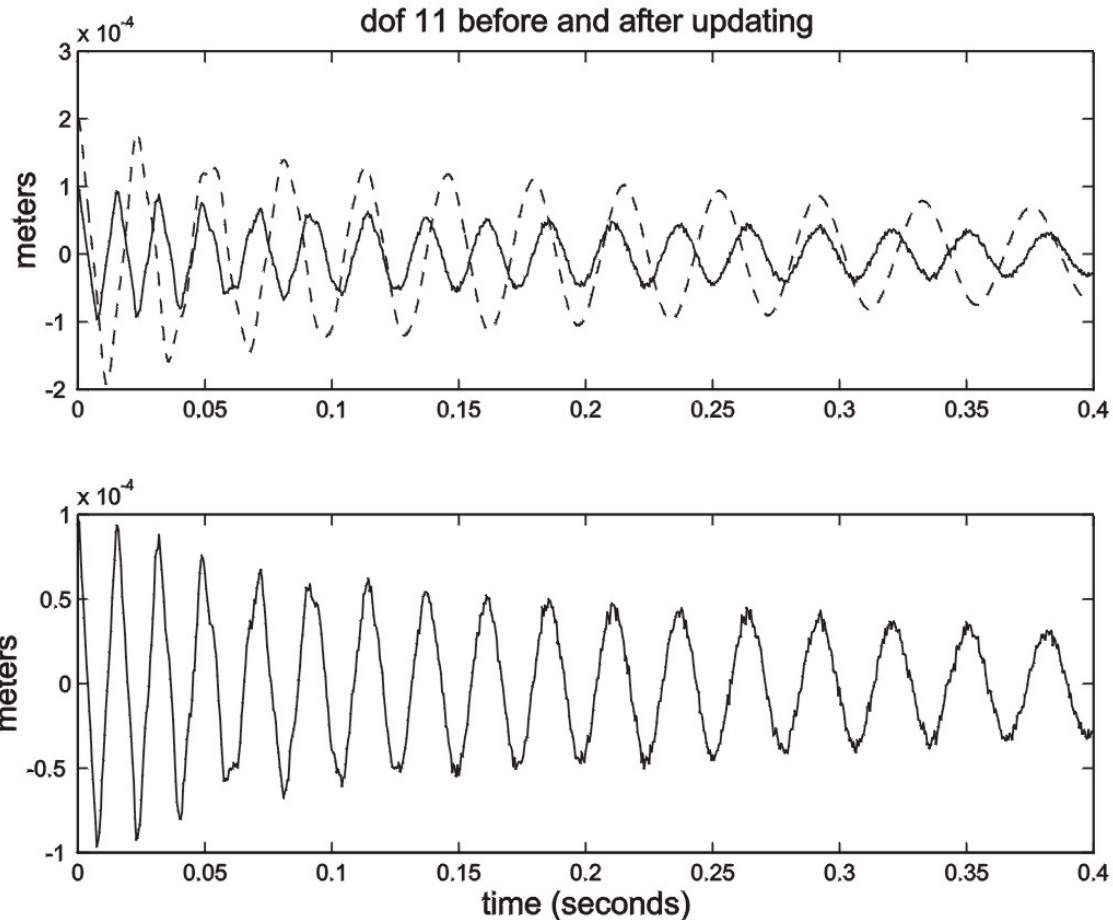
Simulated results

Identification of linear and nonlinear parameters

- 2 parameters : nonlinear stiffness + Young's modulus
- Penalty function in terms of the first POM
- Simulation time = 0.4 sec
- Gaussian white noise of 1 %
- Nonlinear parameter correction < 10 %
- Linear parameter correction < 50 %

Simulated results

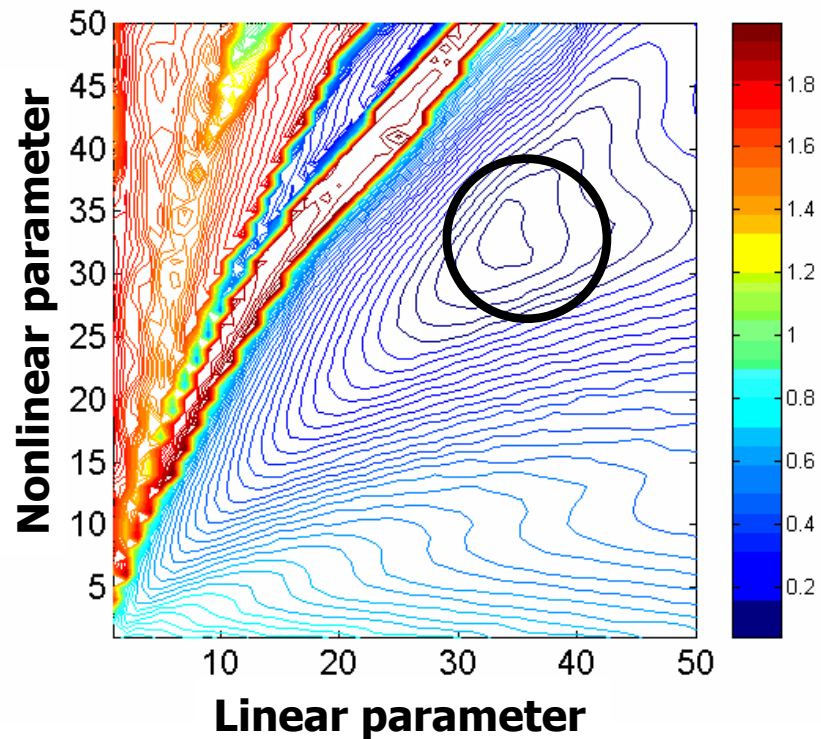
Comparison between the original (—) and the reconstructed (--) signals



Simulated results

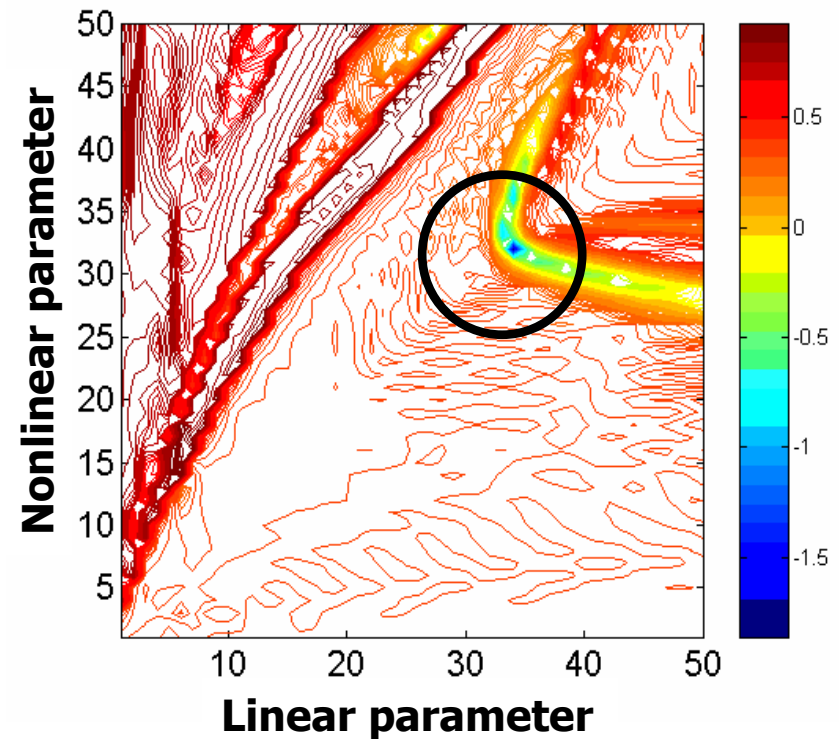
Contour Plot

Penalty Function (use of WT)



Well-conditioning

Penalty Function (no WT)



Ill-conditioning

Experimental results (Vertical set-up)

Model of the nonlinear stiffness

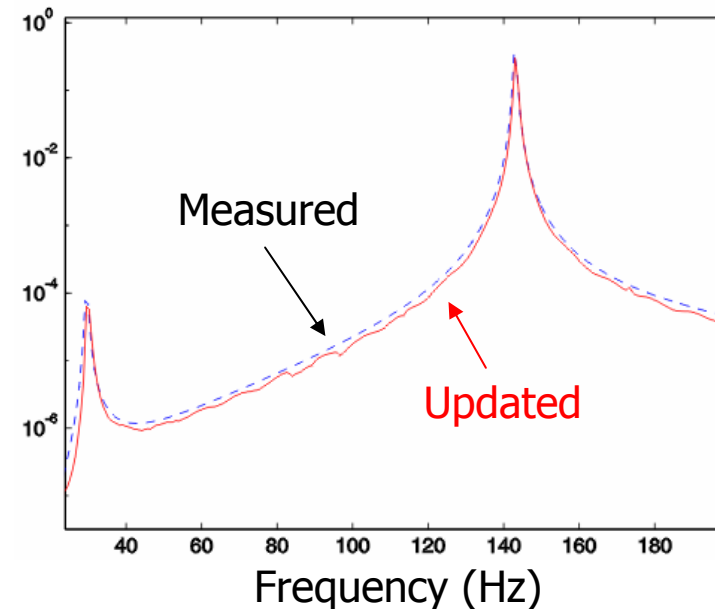
$$f_{nl}(x) = A|x|^{\alpha} \text{sign}(x)$$

Results of the identification of the nonlinear parameters based on the model updating method:

$$\alpha = 2.8$$

$$A = 1.65 \cdot 10^9 \text{ N/m}^{2.8}$$

PSD of the time evolution of the 1st POM

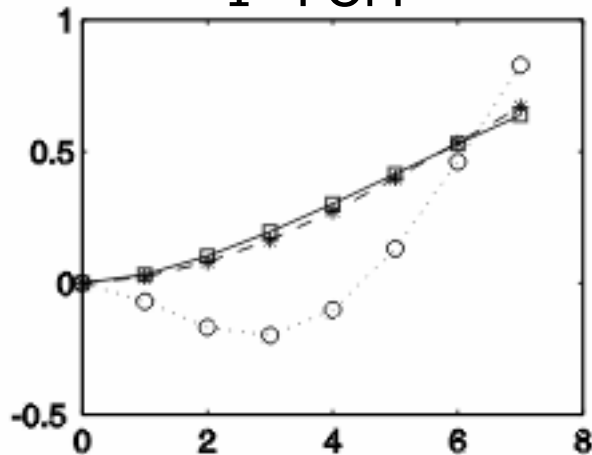


Parameter Estimation Using POD

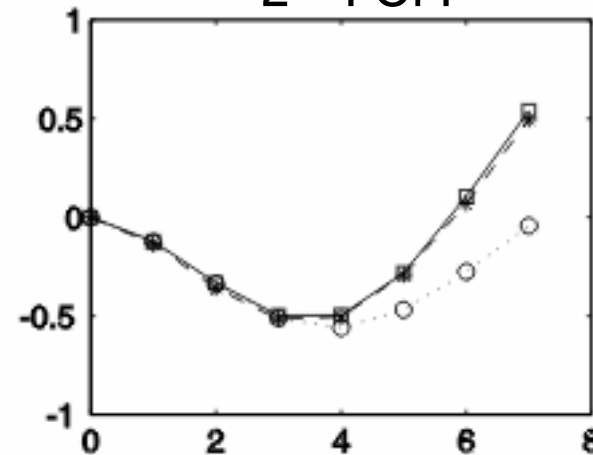
Experimental results (Vertical set-up)

Comparison of the POM

1st POM

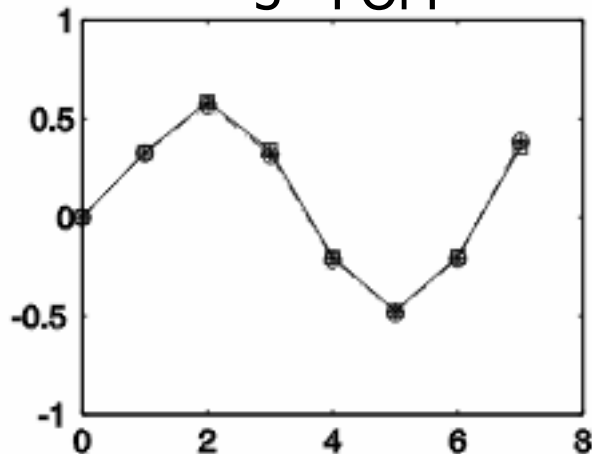


2nd POM

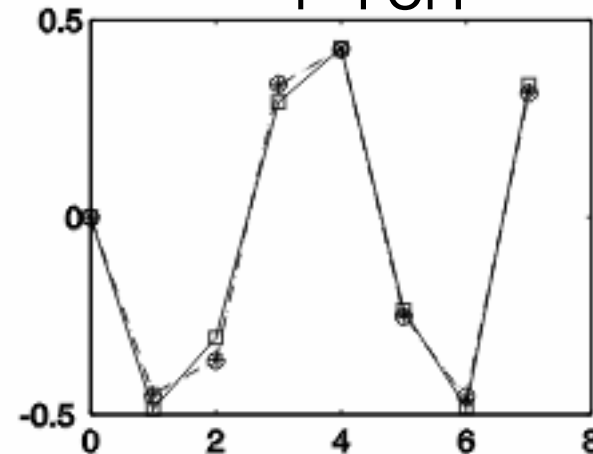


- experimental
- * nonlinear model (after updating)
- linear model (before updating)

3rd POM



4th POM



Nonlinear MDOF systems

$$\mathbf{M} \ddot{\mathbf{x}} + \mathbf{K} \mathbf{x} + \mathbf{f}_{NL}(\mathbf{x}, \dot{\mathbf{x}}) = 0$$

The concept of Nonlinear Normal Modes (NNMs) is a rigorous extension of the concept of eigenmodes to nonlinear systems.

Caution: the solution is energy-dependent !

Parameter Estimation Using Nonlinear EMA

Vector of modal features: $\mathbf{z}^T = (\mathbf{z}_1^T, \dots, \mathbf{z}_i^T, \dots, \mathbf{z}_r^T)$

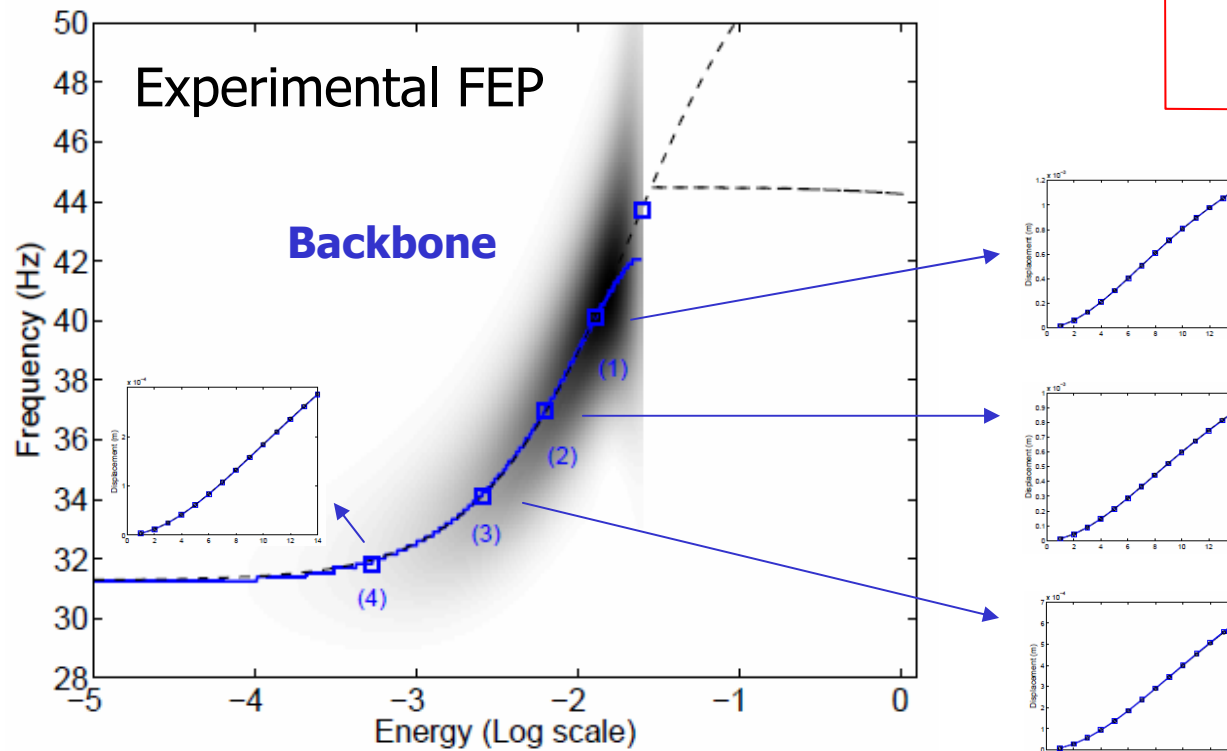
i^{th} backbone

energy level

$$\mathbf{z}_i^T = (\omega_i^{(1)}, \Phi_i^{(1)T}, \omega_i^{(2)}, \Phi_i^{(2)T}, \omega_i^{(3)}, \Phi_i^{(3)T}, \omega_i^{(4)}, \Phi_i^{(4)T}, \dots)^T$$

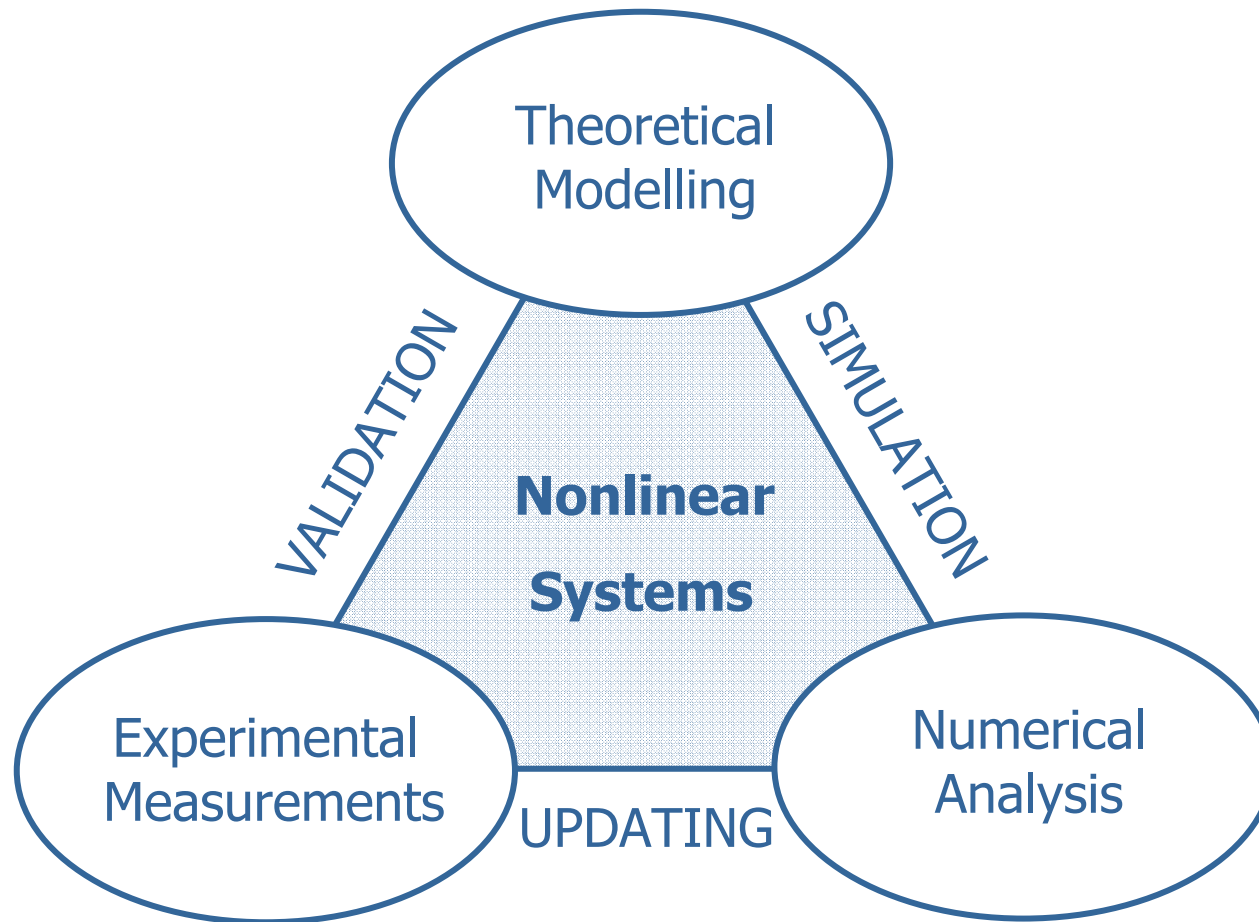
modal shape

frequency

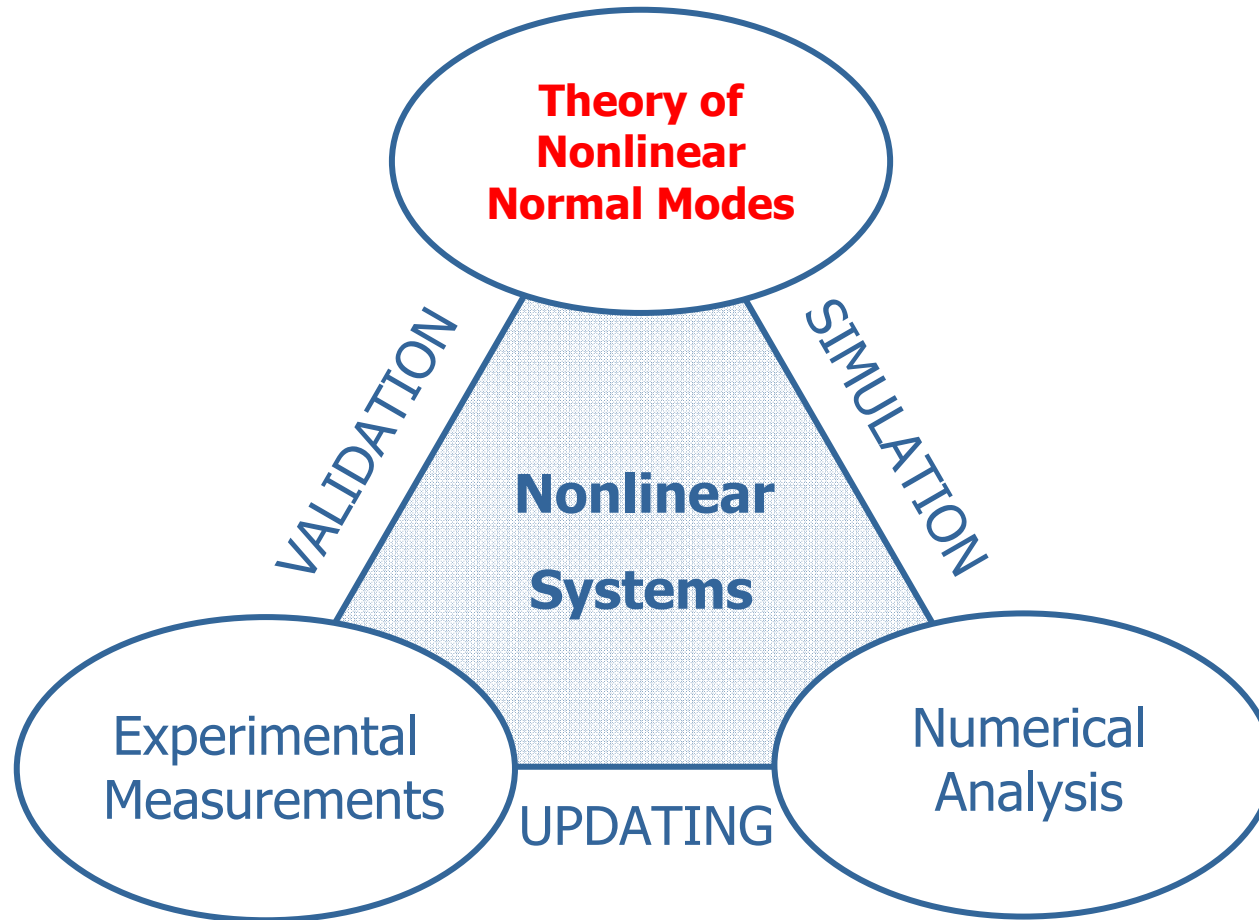




The Structural Dynamicist's Toolkit

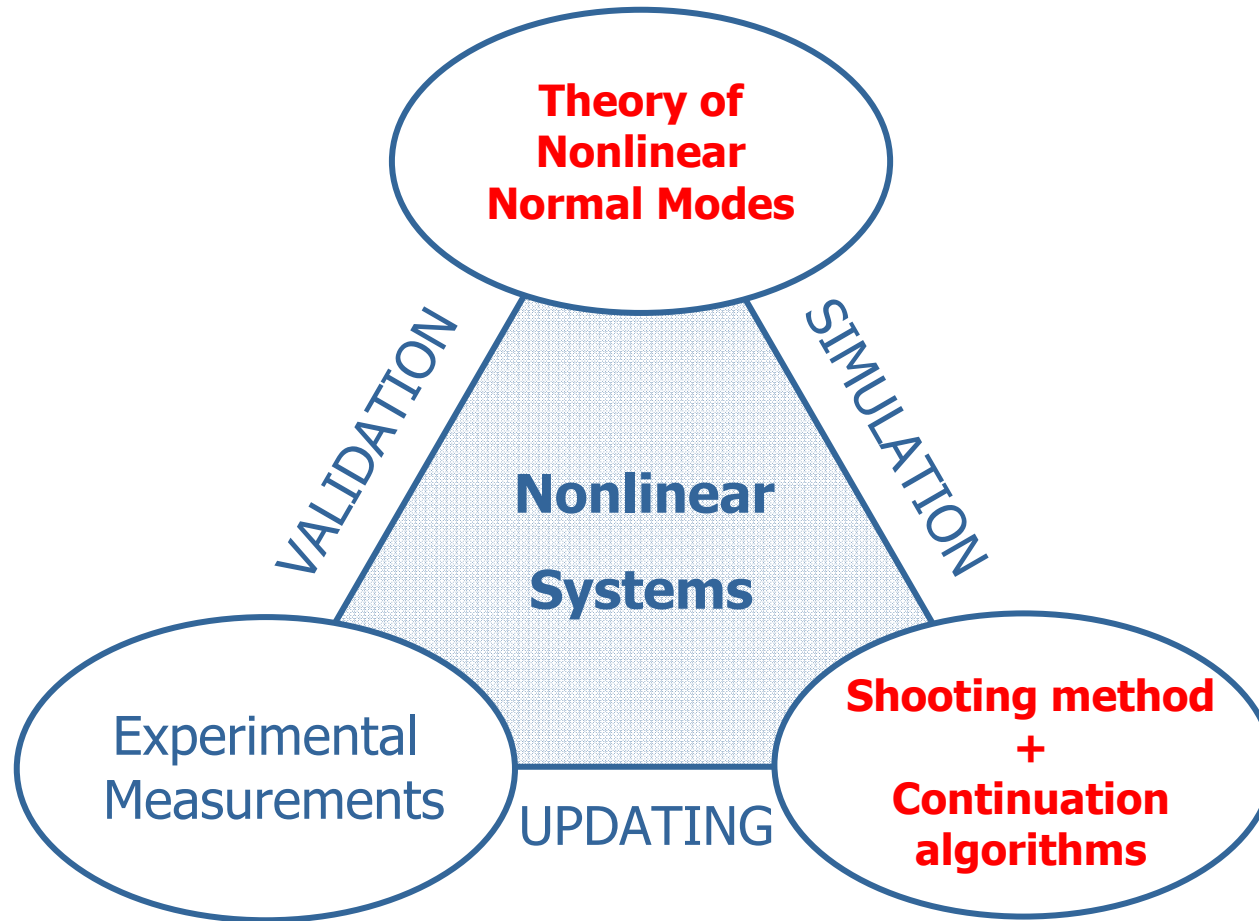


The Structural Dynamicist's Toolkit



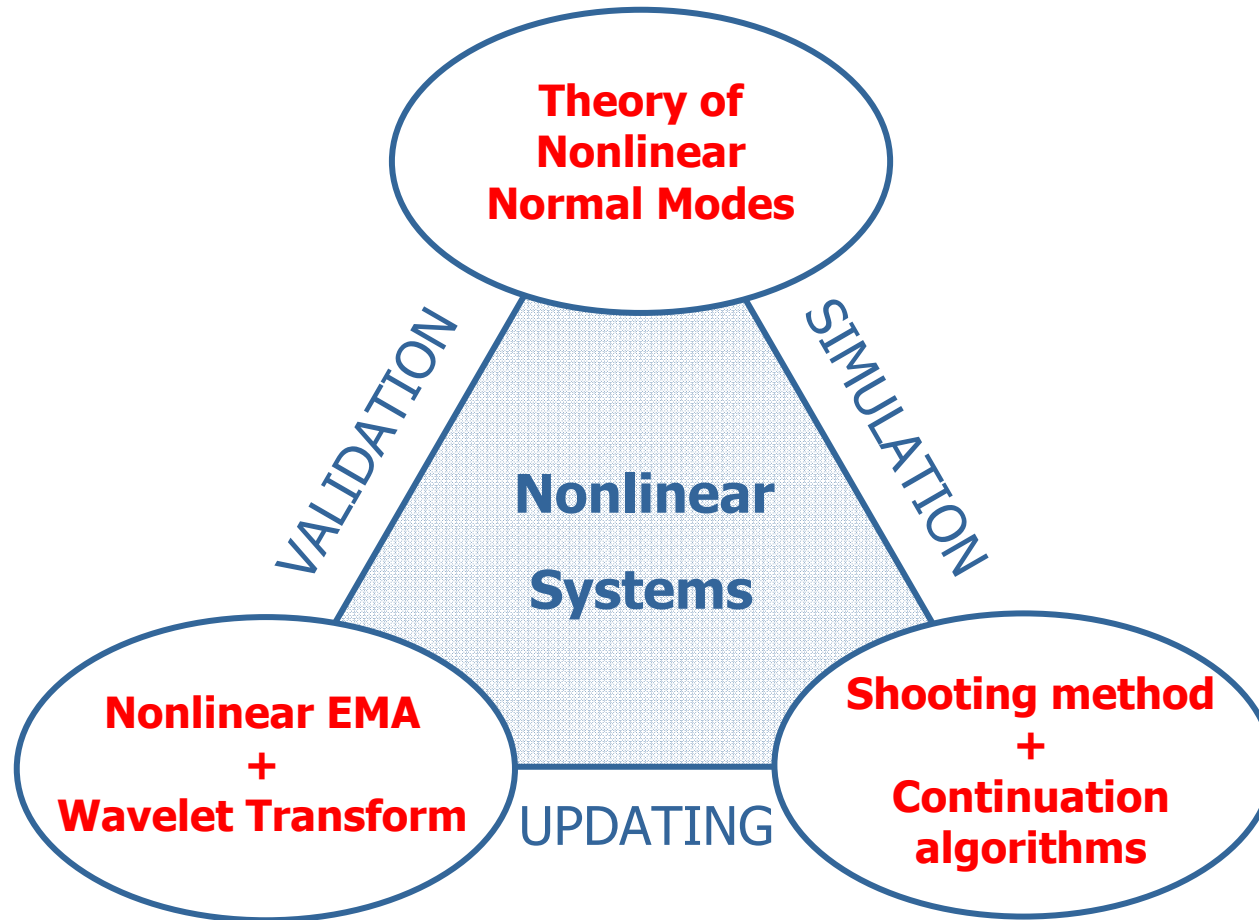


The Structural Dynamicist's Toolkit





The Structural Dynamicist's Toolkit



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Thank you for your attention.