

# Supporting Information: Low-Energy Phases of Bi Monolayer Predicted by Structure Search in Two Dimensions

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This Supporting Information (SI) contains details of the optimized crystal structures for all the six studied Bi monolayers; atomic orbital resolved electronic band structure; the spin-texture calculated at constant energy surfaces for the  $\gamma$  phase; electronic band structure calculated without spin-orbit coupling (SOC) effects; topological characterization and edge states; strain-stabilized phonon spectrum for the  $\gamma$ -phase; a list of potential substrates for the  $\alpha$ -phase; and density of states (DOS). It also contains the electronic band structure of *b*-hexa, *p*-mono, and flat hexagonal Bi monolayer (*f*-hexa) calculated considering Bi  $5d^{10}$  electrons in the core.

## 1 Role of semicore Bi- $5d$ electrons on the electronic band structure

Figure 1 shows the electronic band structure of *b*-hexa and *p*-mono phases calculated with SOC, considering Bi  $5d^{10}$  electrons in the core. The optimized lattice parameters using this pseudopotential are 4.38 Å and 6.71 Å for *b*-hexa and *p*-mono monolayers, respectively. The resulting Bi-Bi bond length is:  $d = 3.07$  Å for *b*-hexa phase, and  $d_{in-plane} = 3.10$  Å,  $d_{out-of-plane} = 3.10$  Å for *p*-mono monolayer structure. The inclusion of  $5d^{10}$  valence electrons is crucial to produce accurate results for Bi: it does not change the shape of the electronic bands near the Fermi level, but it does affect the overall bandgap very strongly. It also considerably changes the optimized lattice parameters of Bi (see the main text for more details).

To ensure that the observed  $5d$ -electron potential does not suffer from additional pathologies (e.g. from missing, even deeper, semi-core states) we check the VASP calculations against all electron LAPW results, using the ELK code. We use standard input parameters and atomic configuration files, include spin orbit coupling, and employ the same geometry relaxed using the plane wave code. Figure 2 shows the electronic band structure within LAPW with SOC for the *f*-hexa structure.

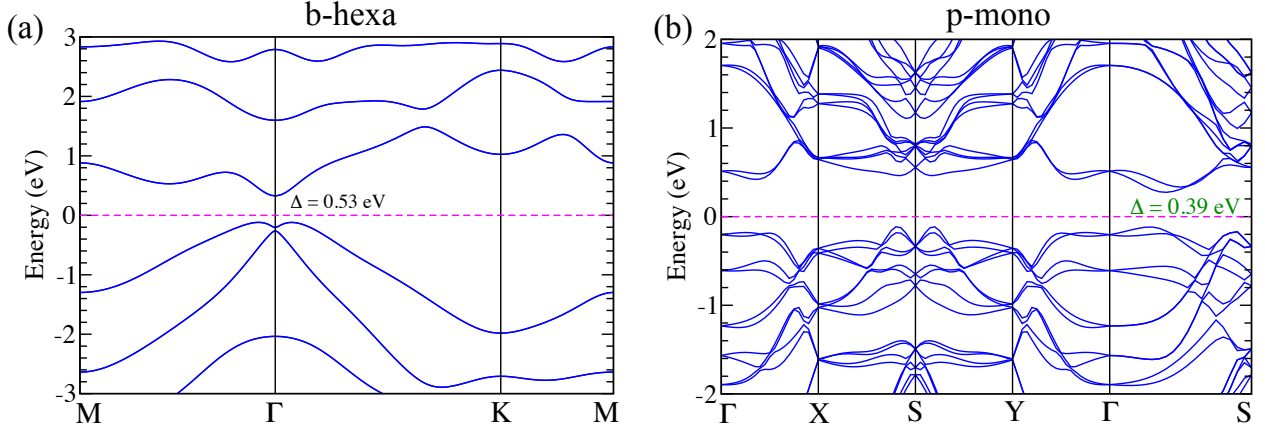


Figure 1: Electronic band structure (PBE+SOC) calculated considering Bi  $5d^{10}$  electrons in the core for (a) *b*-hexa, and (b) *p*-mono monolayers.

We also compare to the flat Bi monolayer structure used by Reis et al. [Science 2017, 357, 287–290]. Note that no substrate included here, and lattice spacing is  $5.35 \text{ \AA}$ . Upon flattening, the gap shifts to K and opens significantly, to a value coincidentally close to that at  $\Gamma$  in the *b*-hexa phase, without *d*-electrons in the valence. We notice negligible changes in the electronic band structure of flat Bi monolayer with inclusion of  $5d$  electrons in valence, as shown in Figure 3. This confirms that the Bi  $5d$  electrons play a crucial role in determining the electronic band structure of Bi monolayer through the crystal field effects, which are negligible in the *f*-hexa phase.

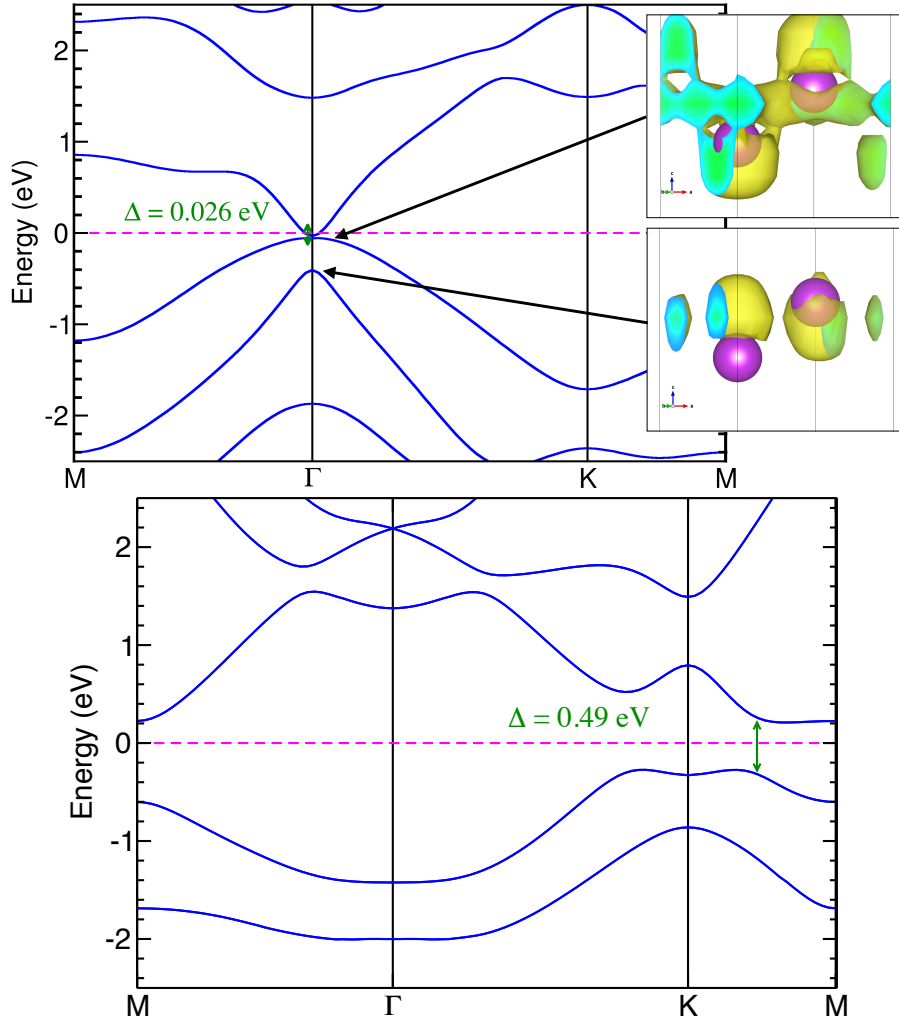


Figure 2: Electronic band structure (with SOC) calculated with the all-electron LAPW method and the ELK code. Top: *b*-hexa, with isodensities of the last two valence states at  $\Gamma$ . These show clear dangling bonds inwards and outwards from the buckled layer, which are sensitive to the presence of semi-core d electrons. Bottom: Flattened (graphene analog) *f*-hexa phase monolayer.



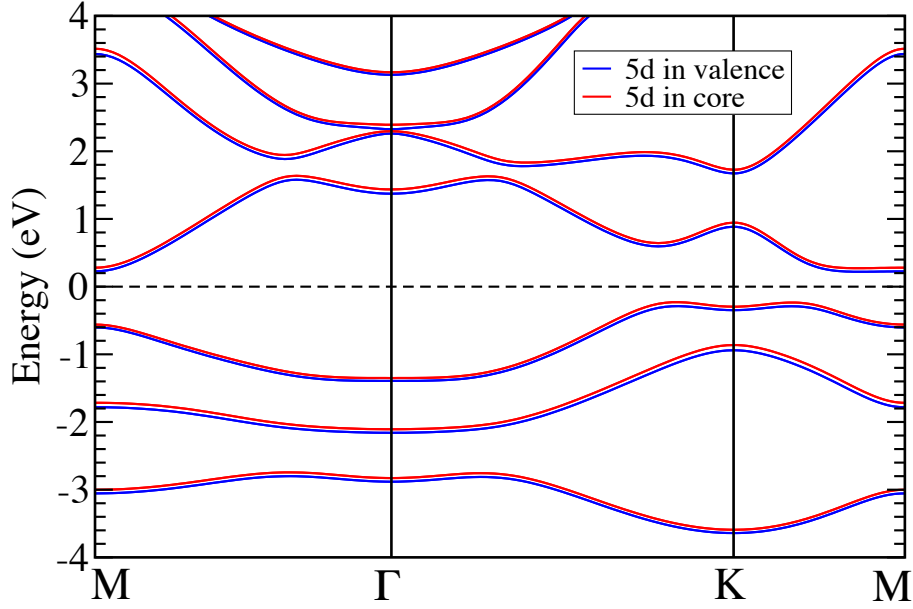


Figure 3: Role of Bi  $5d$  electrons on the electronic band structure of flat Bi monolayer ( $f$ -hexa) calculated with SOC.

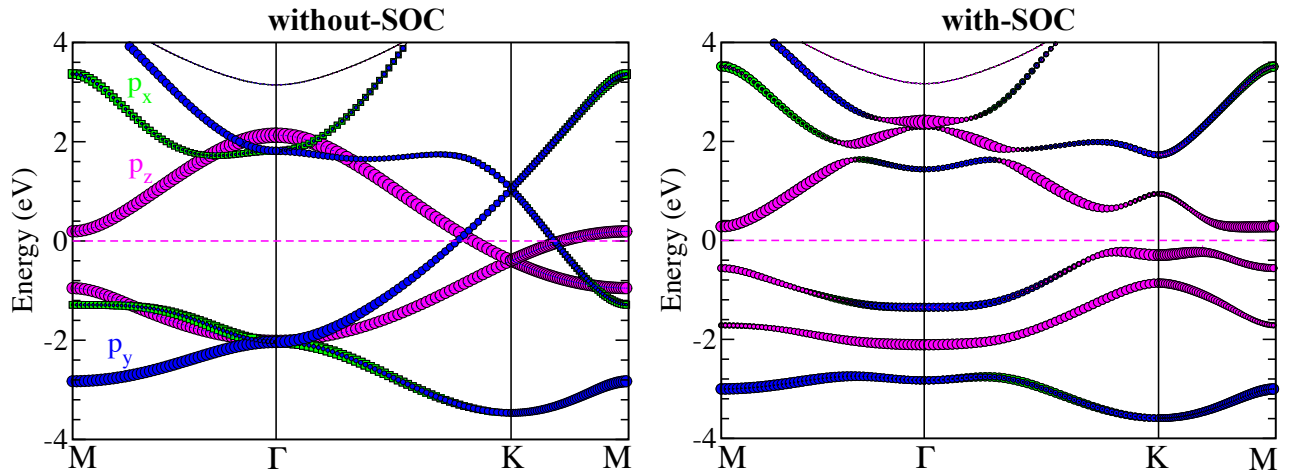


Figure 4: Orbital resolved electronic band structure of flat Bi monolayer ( $f$ -hexa) calculated without (left) and with SOC (right) considering  $5d$  electrons in core.

## 2 Optimized lattice parameters and fractional coordinates of atoms

Table 1: Puckered monoclinic structure

Lattice vectors ( $\text{\AA}$ )	x	y	z
<b>a</b>	6.722	-0.077	0.000
<b>b</b>	-0.393	6.710	0.000
<b>c</b>	0.000	0.000	15.981
Fractional coordinates	x	y	z
Bi	0.9679	0.4643	0.3728
Bi	0.5076	0.9239	0.5650
Bi	0.9677	0.9642	0.5920
Bi	0.4677	0.4642	0.5920
Bi	0.0076	0.4239	0.5650
Bi	0.5078	0.4240	0.3997
Bi	0.4679	0.9643	0.3728
Bi	0.0078	0.9240	0.3998

Table 2: Buckled hexagonal structure

Lattice vectors ( $\text{\AA}$ )	x	y	z
<b>a</b>	4.598	0.000	0.000
<b>b</b>	-2.299	3.982	0.000
<b>c</b>	0.000	0.000	20.000
Fractional coordinates	x	y	z
Bi	0.6667	0.3333	0.0591
Bi	0.0000	0.0000	0.1408

Table 3:  $\alpha$ -phase

Lattice vectors ( $\text{\AA}$ )	x	y	z
<b>a</b>	3.561	0.000	0.000
<b>b</b>	0.334	5.768	0.000
<b>c</b>	0.163	-0.222	16.497

Fractional coordinates	x	y	z
Bi	0.0776	0.5872	0.3966
Bi	0.5759	0.0873	0.4036
Bi	0.0810	0.5776	0.2050
Bi	0.5823	0.0781	0.2120

Table 4:  $\beta$ -phase

Lattice vectors ( $\text{\AA}$ )	x	y	z
<b>a</b>	7.955	0.000	0.000
<b>b</b>	3.232	7.258	0.000
<b>c</b>	0.000	0.000	16.000

Fractional coordinates	x	y	z
Bi	0.7665	0.5962	0.2457
Bi	0.5352	0.8655	0.3793
Bi	0.8430	0.0196	0.3793
Bi	0.3430	0.5196	0.3793
Bi	0.0352	0.3655	0.3793
Bi	0.6121	0.2883	0.2457
Bi	0.2665	0.0962	0.2457
Bi	0.1121	0.7883	0.2457

Table 5:  $\gamma$ -phase

Lattice vectors ( $\text{\AA}$ )	x	y	z
<b>a</b>	7.613	0.000	0.000
<b>b</b>	2.086	7.321	0.000
<b>c</b>	0.000	0.000	16.000

Fractional coordinates	x	y	z
Bi	0.8142	0.5676	0.2729
Bi	0.5631	0.8177	0.4030
Bi	0.8142	0.0676	0.3205
Bi	0.3142	0.5676	0.3199
Bi	0.0652	0.3167	0.4029
Bi	0.5651	0.3169	0.1897
Bi	0.3142	0.0676	0.2721
Bi	0.0632	0.8175	0.1897

Table 6: Flat hexagonal (*f*-hexa) phase

Lattice vectors ( $\text{\AA}$ )	x	y	z
<b>a</b>	5.329	0.000	0.000
<b>b</b>	-2.664	4.615	0.000
<b>c</b>	0.000	0.000	20.767

Fractional coordinates	x	y	z
Bi	0.3333	0.6667	0.0000
Bi	0.6667	0.3333	0.0000

### 3 Electronic band structures calculated without SOC

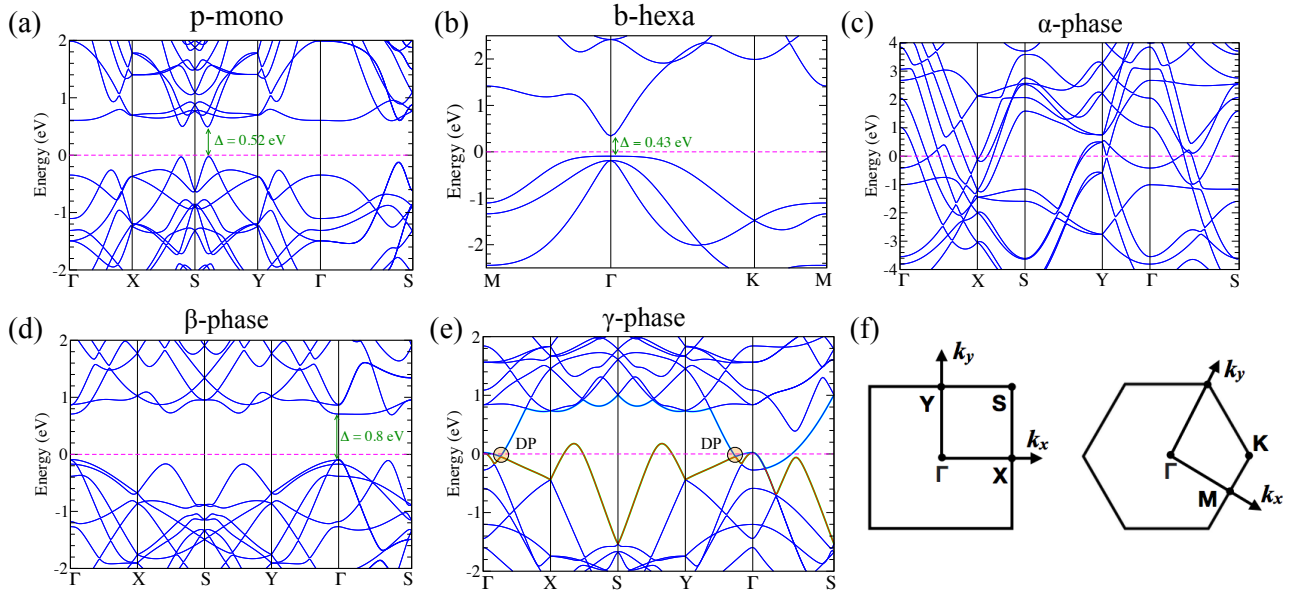


Figure 5: Electronic band structures calculated without considering SOC effects for (a) *p*-monoclinic, (b) *b*-hexagonal, (c)  $\alpha$ -, (d)  $\beta$ -, and (e)  $\gamma$ -phases. (f) Brillouin zones marked with high-symmetry points.

## 4 Strain-stabilized phonon spectrum for $\gamma$ -phase

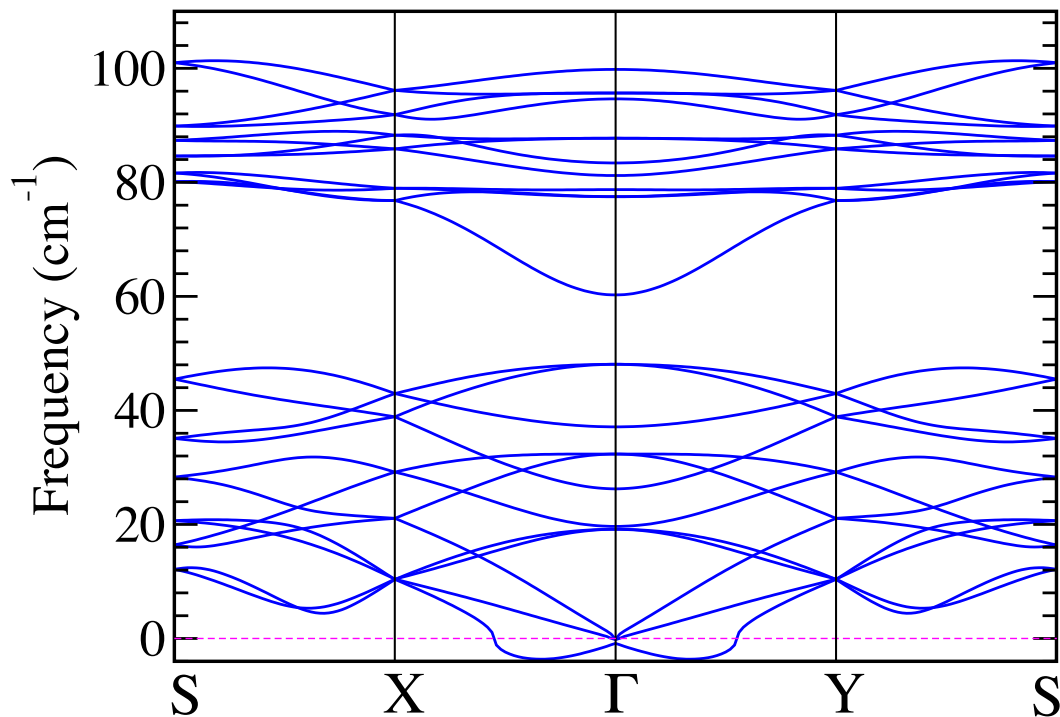


Figure 6: 8.0% tensile biaxial strain stabilized phonon spectrum of the  $\gamma$ -phase. The elastic instabilities near the  $\Gamma$  points are due to poor numerical convergence at  $k \rightarrow 0$  in 2D systems: the expected quadratic 2D mode is difficult to interpolate using Fourier methods, though all calculated frequencies are positive. The inner-coordinates were optimized after applying biaxial strain, and phonon spectrum was calculated using a supercell of size  $2 \times 2 \times 1$ .

## 5 Topological characterization

### 5.1 Inversion parity eigenvalues

Table 7: The  $Z_2$  topological invariant computed from the product of the inversion parity eigenvalues of all occupied bands at four TRIM points  $\Gamma$  (0.0, 0.0),  $M_1$  (0.5, 0.0),  $M_2$  (0.0, 0.5), and  $M_3$  (0.5, 0.5)

	Product of the parity eigenvalues				$Z_2$
	$\Gamma$	$M_1$	$M_2$	$M_3$	
<i>b</i> -hexa	-	+	+	+	1
<i>f</i> -hexa	+	+	+	+	0
$\alpha$ -phase	-	+	+	+	1

### 5.2 Edge states

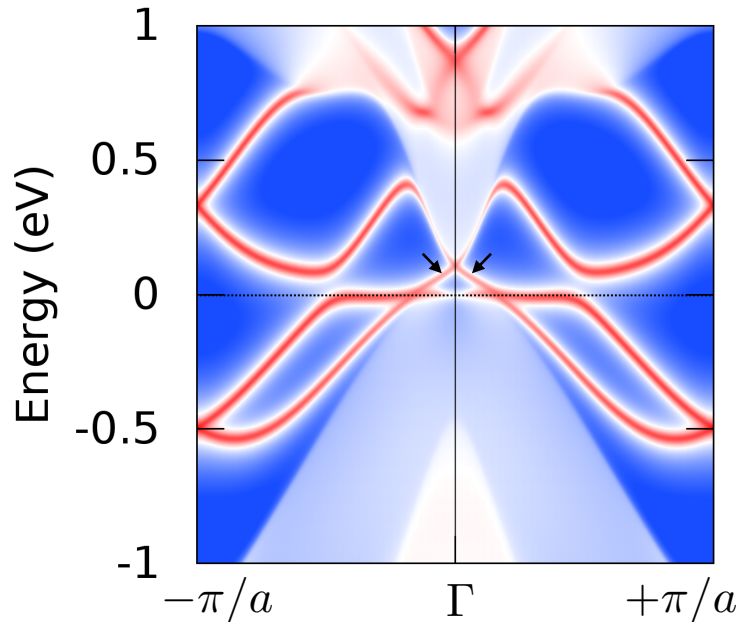


Figure 7: *b*-hexa phase: The edge state spectrum showing gapless dispersion of edge modes in the *b*-hexa phase. Arrows mark the topological nontrivial edge states. The flat edge states at the Fermi level are distinctive of the van Hove singularity. Red/Blue color represents the spectral density. Red color denotes the projection of edge states while white/blue color denotes the projection of monolayer states.

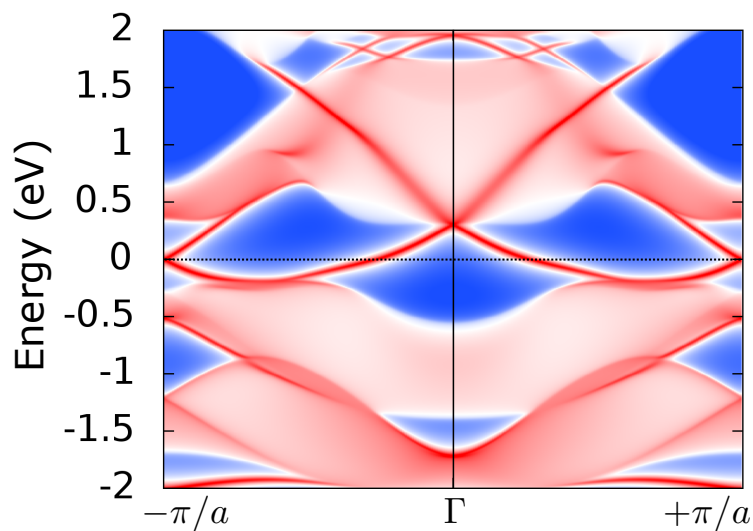


Figure 8: *f*-hexa phase: The edge state spectrum of *f*-hexa phase. Even number of edge states crossing the Fermi level along  $\Gamma \rightarrow \pi/a$  kpath reveal the trivial topological nature of *f*-hexa phase, as predicted by calculated  $Z_2$  topological invariant. Colors are as defined in Fig. 7. As noted by Reis et al. [Science 2017, 357, 287–290], the topology becomes nontrivial once SiC substrate is included and planar mirror symmetry is broken.

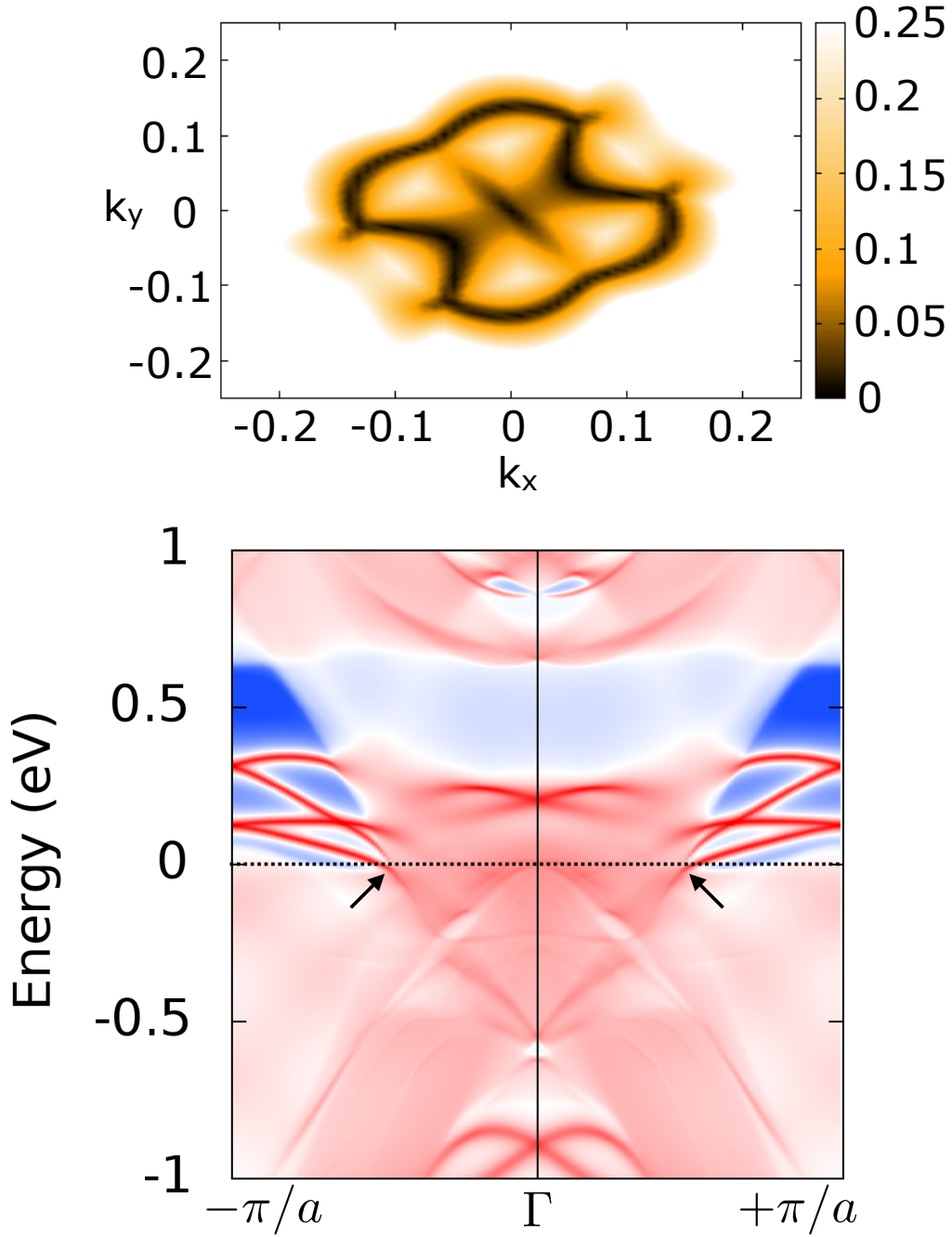


Figure 9:  $\alpha$ -phase: Top panel shows the energy bandgap plotted in color scale (eV units) in a  $k_x - k_y$  plane centered at the  $\Gamma$  point. A gapless Dirac nodal line can be noticed. Bottom panel depicts the edge state spectrum showing gapless nontrivial edge modes in the  $\alpha$ -phase. Arrows mark the topological nontrivial edge states originating from the Dirac point.



## 6 Atomic orbital resolved band structure of $\gamma$ phase

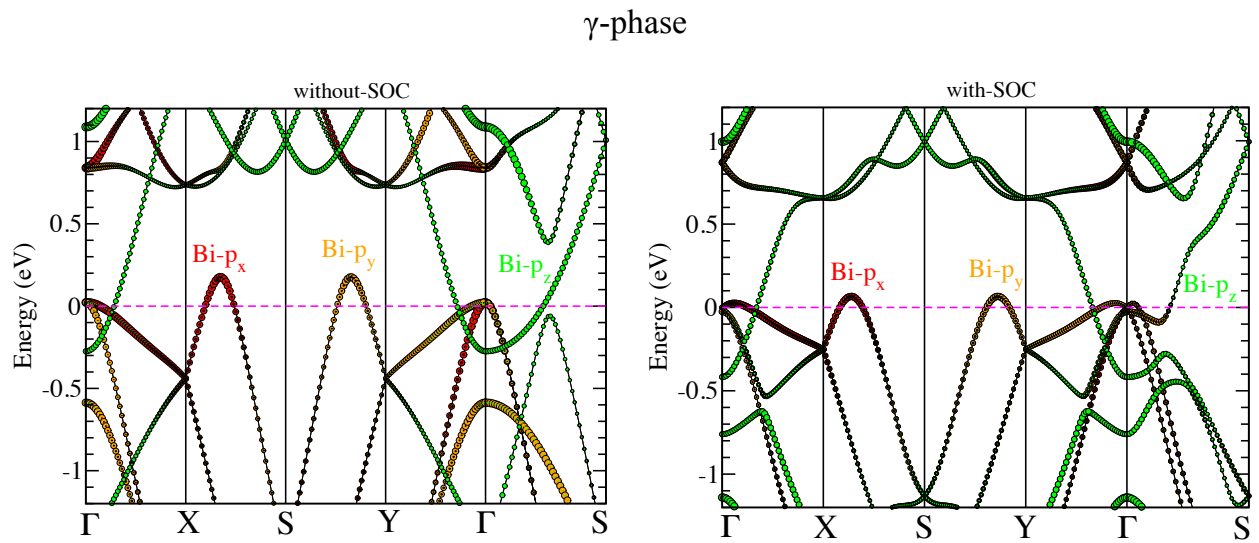


Figure 10: Atomic orbital resolved electronic band structure calculated without (left) and with (right) inclusion of SOC.

## 7 Spin-texture in $\gamma$ phase

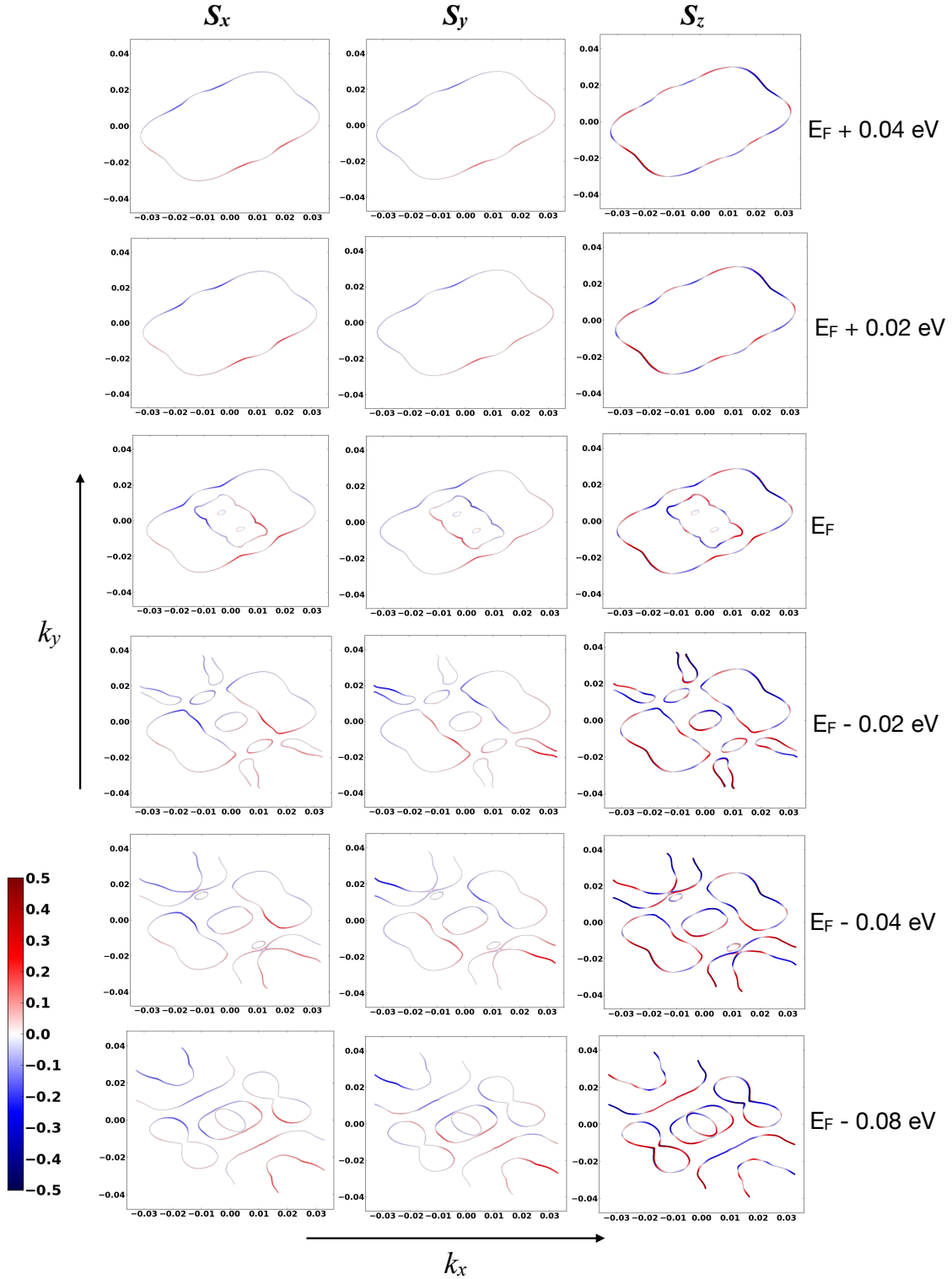


Figure 11:  $\gamma$  phase : Spin-texture calculated at several different constant energy surfaces near the Fermi level in a  $k_x - k_y$  plane centered at the  $\Gamma$  point of Brillouin zone.

## 8 List of potential substrates for $\alpha$ phase

Table 8: List of potential substrates for the  $\alpha$  phase: This table shows a list of potential substrates that might be suitable for the synthesis of the  $\alpha$  phase. For the substrate search, we employed the available algebraic algorithm provided by the MPInterfaces package, to search for geometrical matches with structures available in a material databases. We calculate the work-function of the matched substrates, and use it as a criterion to define a good substrate. The maximum area mismatch between the substrate candidates and Bi monolayer was set to 40 Å, and cutoff for the maximum lattice mismatch was set to be 0.05 Å with a maximum cell angle difference of 1°.

Compound	space group number
LiS <sub>2</sub> Ti	164
PtS <sub>2</sub>	164
Mg <sub>2</sub> Tb <sub>4</sub>	164
Fe <sub>6</sub> Se <sub>8</sub>	12
SnS <sub>2</sub>	164
Mg <sub>2</sub> Y <sub>6</sub>	44
As <sub>2</sub> Cr <sub>2</sub>	186
MgB <sub>2</sub>	191
PbO	187
Mn <sub>4</sub> Te <sub>4</sub>	62
InTe	123
As <sub>3</sub> Co <sub>3</sub> Fe <sub>3</sub>	189
PbS	166
Fe <sub>2</sub> Se <sub>2</sub>	194
Tl <sub>12</sub> O <sub>6</sub>	166
Ni <sub>2</sub> Tb <sub>2</sub> C <sub>4</sub>	38
Mn <sub>6</sub> P <sub>3</sub>	150
MgCl <sub>2</sub>	164
As <sub>4</sub> Rh <sub>4</sub>	62
TiSe <sub>2</sub>	164

## 9 Density of States

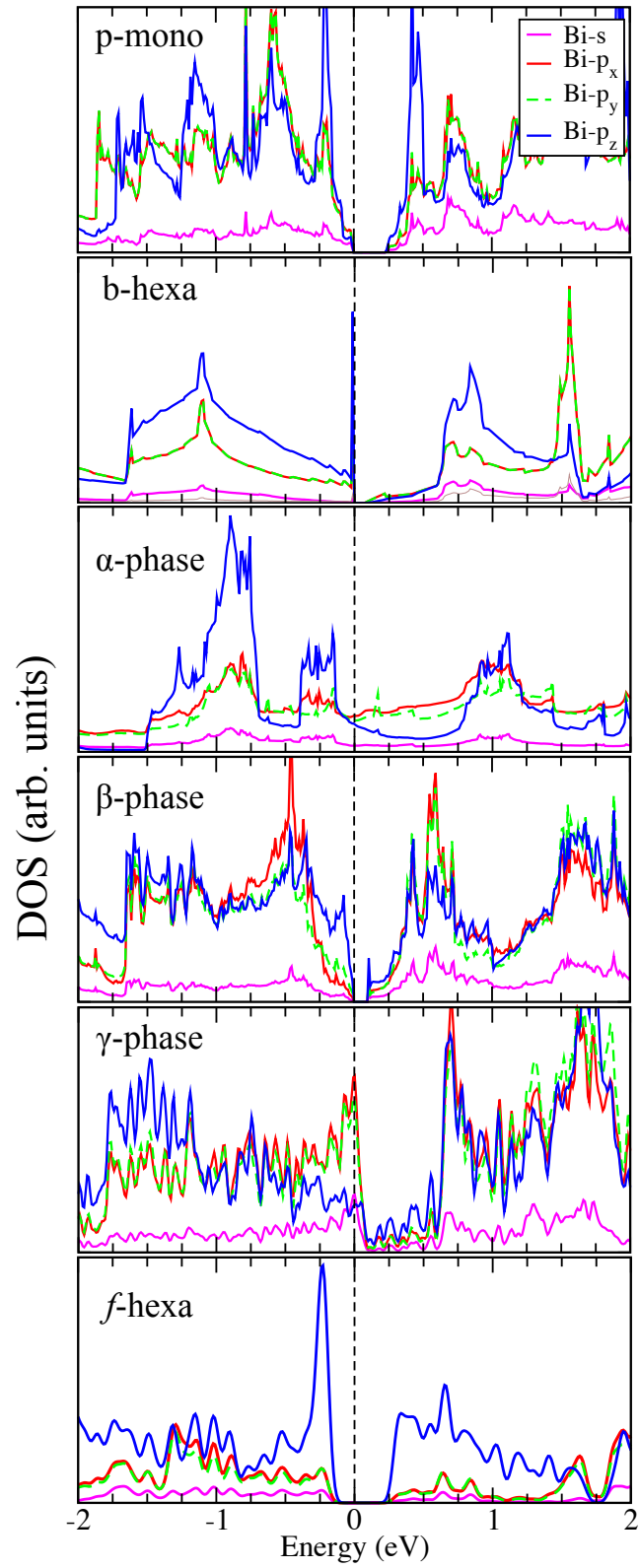


Figure 12: DOS calculated with SOC using a  $k$ -mesh of size  $21 \times 21 \times 1$  for  $p$ -mono,  $b$ -hexa,  $\alpha$ -,  $\beta$ -,  $\gamma$ , and  $f$ -hexa phases of monolayer Bi (arranged in a top-to-bottom order, respectively).