Alfvénic Acceleration Sustain Ganymede’s Footprint Tail Aurora

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Key Points
- First in-situ particles and fields measurements connected to Ganymede’s auroral tail are reported.
- Alfvén wave activity is observed with Poynting fluxes of ~100 mW/m² capable of accelerating electrons into the atmosphere.
- Ganymede’s footprint tail contains electron populations consistent with Alfvénic acceleration, and precipitating energy fluxes of ~11 mW/m².

Abstract
Integrating simultaneous in-situ measurements of magnetic field fluctuations, precipitating electrons, and ultraviolet auroral emissions, we find that Alfvénic acceleration mechanisms are responsible for Ganymede’s auroral footprint tail. Magnetic field perturbations exhibit enhanced Alfvénic activity with Poynting fluxes of ~100 mW/m². These perturbations are capable of accelerating the observed broadband electrons with precipitating fluxes of ~11 mW/m², such that Alfvénic power is transferred to electron acceleration with ~10% efficiency. The UV emissions are consistent with in-situ electron measurements, indicating 13 ± 3 mW/m² of precipitating electron flux. Juno crosses flux tubes with both upward and downward currents connected to the auroral tail exhibiting small-scale structure. We identify an upward electron conic in the downward current region, possibly due to acceleration by inertial Alfvén waves near the Jovian ionosphere. In concert with in-situ observations at Io’s footprint tail, these results suggest that Alfvénic acceleration processes are universally applicable to magnetosphere-satellite interactions.

Plain Language Summary
Jupiter’s moon Ganymede interacts with the planet’s rapidly rotating magnetic field, which generates an aurora in the Jovian upper atmosphere. The Juno spacecraft crossed magnetic field lines connected to this aurora. We found that a specific type of wave, similar to a wave produced when a string is plucked, is responsible for accelerating the electrons sustaining this aurora. This type of interaction between a moon and the planet it orbits is likely a common process occurring at other exoplanetary systems.

1. Introduction
Jupiter’s auroral features are diverse and complex, consisting of separate emission regions linked to different locations in the magnetosphere and many different physical mechanisms [e.g. Grodent et al. 2015]. While many of these auroral features map to broad areas in the magnetosphere, the aurorae generated by the interaction of Jupiter’s rotating magnetosphere with the Galilean moons allow for a more precise determination of the auroral source regions and physical processes sustaining their aurora.

Ganymede’s auroral signature was discovered in ultraviolet (UV) images of Jupiter’s aurora [Clarke et al. 2002] and is observed to have multiple auroral spots downstream of its main spot, whose separation varies with Ganymede’s System III longitude [Bonfond et al. 2013]. A more continuous auroral tail has also been observed for Ganymede, and remote sensing observations suggest this tail is sustained by reflections of Alfvén waves generated by Ganymede’s interaction with Jupiter’s magnetosphere [Bonfond et al. 2017a]. The morphology and generation of Ganymede’s auroral signatures are expected to be similar to Io’s tail [Bonfond et al. 2017b], such that it has a Main Alfvén Wing (MAW) spot, a Transhemispheric Electron Beam [Bonfond et al. 2008], and a footprint tail aurora. Similar to Io, high-resolution infrared
images have revealed Ganymede’s tail to be highly structured [Mura et al. 2018]. In-situ observations by the Juno spacecraft have bolstered the case that Alfvénic acceleration is sustaining Io’s footprint tail aurora [Bonfond et al. 2009] from measurements of broadband precipitating electron fluxes [Szalay et al. 2018] and Alfvénic magnetic signatures [Gershman et al. 2019]. Since our understanding of Io’s auroral interaction is coalescing around a dominantly Alfvénic acceleration mechanism sustaining both the MAW and tail, and similar acceleration processes are operating at Saturn’s moon Enceladus [e.g. Sulaiman et al. 2018], identifying the acceleration mechanism at another Jovian moon allows us to test the assertion that this process is more broadly applicable to moon-magnetosphere interactions as a whole [Bonfond et al. 2017a; 2017b].

While Io’s auroral signature is the best catalogued of the Galilean moons due to it being the brightest satellite interaction, Ganymede’s auroral footprint characteristics have also been well-documented. Ganymede’s footprint aurora varies on three different timescales, the largest corresponding to Jupiter’s orbital period of ~10 hrs, where the footprint emission is strongest when Ganymede is exposed to the largest plasma sheet densities [Grodent et al. 2009]. The average emitted power of Ganymede’s auroral footprint is estimated to range from 0.2-1.5 GW [Grodent et al. 2009] and 0.8-6.3 GW [Bonfond et al. 2017b]. The total theoretically calculated Poynting fluxes generated by Ganymede’s interaction are estimated to lie between 1-150 GW depending on Ganymede’s position in Jupiter’s plasma sheet [Saur et al. 2013]. The main footprint emission feature covers an area of 5 x 10^9 km^2 at Jupiter’s atmosphere [Grodent et al. 2009], corresponding to a region 8-20 R_G wide in Ganymede’s orbital plane (R_G = 2634 km is Ganymede’s radius). This region is significantly larger than the diameter of Ganymede due to Ganymede’s own magnetosphere providing a factor of ~10 larger obstacle to the plasma flow and enhancing the size of its interaction region with Jupiter’s magnetospheric plasma [Jia et al. 2008; Paty et al. 2008; Saur et al. 2013]. Accounting for the footprint area, the average emitted power over the footprint is 0.4-13 mW/m^2. Assuming a 10% conversion efficiency from incident electron power to auroral emission power [Gérard & Singh, 1982; Grodent et al. 2001; Gustin et al. 2012], this corresponds to an average incident electron energy flux of 4-130 mW/m^2 for Ganymede’s main auroral spot. Similar incident energy fluxes are inferred from a decade of Hubble Space Telescope observations of Ganymede’s auroral footprint [Wannawichian et al., 2010]. Its footprint tail fluxes are expected to follow an exponential decay in power as a function of angular separation from the main spot, with an e-folding distance of 11° [Bonfond et al. 2017a].

The Juno mission [Bolton et al. 2017] provides an unparalleled opportunity to investigate the acceleration processes driving Ganymede’s complex auroral phenomena as the spacecraft directly transits flux tubes connected to this aurora. Magnetic field perturbations and the characteristics of the electrons responsible for causing emission in the Jovian atmosphere hold the clues to how this aurora is sustained. One auroral signature of particle acceleration/heating is the production of conics, where charged particle populations exhibit peak fluxes at oblique pitch angles [Klumpar et al. 1979]. Observations of conics give insight into how electrons are accelerated on auroral field lines. Electron conics were first observed with the DE-1 spacecraft [Menietti and Burch, 1985], and have been associated with various auroral phenomena [Lysak 1993] such as parallel or perpendicular heating [e.g. Menietti and Burch, 1985; Temerin and Cravens, 1990] and low-frequency parallel electric fields [e.g. Lundlin et al. 1987; Andre and Eliasson, 1992; Thompson and Lysak, 1996]. They have been observed frequently at Jupiter with Juno, likely driving the cyclotron maser instability, and are suggested to be due to inertial Alfvén waves [Louarn et al. 2018].
We focus on data taken when Juno transited magnetic field lines connected to Ganymede’s auroral footprint tail, reporting on a variety of auroral phenomena such as broadband precipitating electrons, Alfvénic perturbations to the magnetic field, and up-going electron conics. In Section 2, we discuss the observation geometry and describe the correlation between remote UV and in-situ fields and particle measurements necessary to link these measurements to the Ganymede footprint tail. We discuss the specific details of the in-situ measurements in Section 3 and conclude with a discussion in Section 4.

2. Correlating Remote and In-Situ measurements

Ganymede orbits Jupiter at 15 Jovian radii (R_J), twice per orbit crossing Jupiter’s current sheet that distends the magnetospheric magnetic field and makes it difficult to precisely trace magnetic field lines from its location to the Jovian atmosphere. Instead of correlating features in Juno plasma data to their equatorially mapped regions as done previously [e.g. Allegrini et al. 2017; Szalay et al. 2017], we link in-situ measurements of particles and fields with near-simultaneous UV images of the Jovian auroral emissions [e.g. Ebert et al. 2019; Gérard et al. 2019].

The top panel of Figure 1 shows a false-color UV image taken by Juno’s Ultraviolet Spectrograph (UVS) [Gladstone et al. 2017] accumulated just prior to Juno’s 20th perijove over 10 minutes on 2019-149 (May 29) from 7:35:02 to 7:45:02 UTC (during which Ganymede has moved less than 1°). These colors can be qualitatively compared to precipitating electron energies observed in-situ.

Juno’s magnetic footprint is shown, traced from Juno’s position to 400 km above the 1 bar level using the JRM09 field model [Connerney et al. 2018], which incorporates a current sheet model [Connerney et al. 1981]. Juno’s footprint moves from left to right in this image and the time between circles is 1 minute. A nominal location for the main auroral oval is denoted with the solid red and dashed white line in the left half of the image. Ganymede’s main auroral spot is the pink feature circled under the “Ganymede Footprint” label and arrow, with the tail extending up and left such that Juno’s mapped footprint crossed it 8° down-tail. Beginning on the left side of the image, Juno is transiting across Jupiter’s polar region, where there is limited auroral emission coverage by UVS. Juno then encounters magnetic field lines connected to main auroral oval, where the polar edge has a higher average energy incident to the atmosphere (a), shown in red, immediately followed by an auroral region with lower incident energy (b), shown in pink/maroon. Following this, it transits a region with very little auroral emission for about a minute (c). It then encounters Ganymede’s footprint tail (d), which extends from the Ganymede MAW upward in this image. After this, it encounters a small region with low emission (e), followed by a region with high incident energy for about a minute and a half (f) and then encounters another region with low emission (g).

Uncertainties in the magnetic field mapping can lead to time shifts when comparing the emission thought to be underfoot to in-situ particle and fields data. These effects can be more drastic when mapping from the northern hemisphere, due to the localized magnetic anomaly in this hemisphere [Grodent et al. 2008]. Additionally, the latitude of Ganymede’s auroral footprint can vary by 2° depending on current sheet properties [Grodent et al. 2009], adding to the complexities in determining its exact auroral mapping. Therefore, we do not compare Juno’s location with field lines expected to map to Ganymede’s orbit in the equatorial plane. Instead, we directly correlate emission features along Juno’s mapped footprint to precipitating electron fluxes measured by the Jovian Auroral Distributions Experiment (JADE) [McComas
et al. 2017], which exhibit similar temporal and energy structure. The bottom panel of Figure 1 shows JADE electron differential energy fluxes (DEF) for precipitating electrons able to produce auroral emissions (observed in the loss cone) during this time. Pitch angles are calculated using the broadcast magnetic field onboard Juno [Connerney et al. 2017] at a cadence of 1 Hz. Overlaid on this spectrogram is the total downward precipitating energy flux (EF), described in more detail in Section 3.

In region (a), JADE observes high energy electrons, initially extending above JADE’s energy range, then peaking in DEF at ~20 keV towards the end of (a). Subsequently, it measures a downward flux of electrons at lower average DEF energies of 2-20 keV in region (b). Region (c) is initially mostly void of downward electron fluxes, then transitions to a smaller low energy component extending below 0.1 keV. In region (d), we observe a highly structured enhancement in downward energy flux, which is the focus of this study and will be shown in more detail in Section 3. Region (e) returns to lower energy fluxes, similar to the 2nd half of region (c). In region (f), JADE observed an enhancement in high-energy downward fluxes extending above ~10 keV. In region (g), JADE observes very low fluxes without appreciable structure.

Comparison of the temporal order of features seen as Juno transits this region by both UV and in-situ electron measurements provides a high degree of correlation, both in average incident energy and feature duration. Given the high fidelity of correlation between the auroral emission image and downward energy flux, we attribute the enhancement observed in region (d) by JADE to the Ganymede footprint tail. Following this association, we describe the characteristics of auroral fields and particles connected to Ganymede’s footprint tail in Section 3.

### 3. In-Situ Measurements

Figure 2 shows in-situ measurements zoomed into region (d) from Figure 1, taken by Juno during 2019-149 from 7:36:44 to 7:38:02 UTC. The Ganymede footprint tail transit is identified as the period with enhanced fluxes in the center of this figure, from 7:37:14 to 7:37:42 and marked with the small color strip running along the top of the figure, where the colors indicate time from 7:37:22 and are used to identify spectra in Figure 3. Figure 2a shows the downward, precipitating electron differential energy flux, taken for measurements in the loss cone of 32˚ during this time (JRM09 model). Overlaid on this panel in white is the downward precipitating energy flux, which peaks at 11 mW/m² in the enhancement from 7:37:19 to 7:37:24. The downward energy flux is calculated as $\mathcal{E}_F = \pi \sum_i \text{DEF}_i \cdot \Delta E_i$, summing occurs over the JADE energies (subscript $i$), where $\pi$ is the area-projection weighted size of the loss cone above Jupiter’s atmosphere and $\Delta E_i$ is the width of each energy bin following previous analyses [e.g. Mauk et al. 2017a]. We calculate energy flux by summing from 50 eV to 40 keV. Since the fluxes above 40 keV are very near the noise floor, their omission does not substantially affect the analysis in this study. Figure 2c shows the same type of information as the first panel, but for the fluxes moving away from Jupiter in the upward loss cone.

Figure 2b shows DEF electron pitch angles for the same energy range. The JADE sensors provide 240˚ of azimuthal coverage (perpendicular to Juno’s spin axis) and can electrostatically deflect up to 35˚ above/below Juno’s spin plane to track the magnetic field direction. The missing 120˚ of azimuthal coverage leads to gaps in pitch angle, shown as the triangular regions.
with no observed fluxes in the second panel. At the onset of the identified Ganymede footprint tail transit (7:37:14), a short-duration electron conic is observed, peaking at around 30°. There is also a similar feature before the transit at 7:37:00. From 7:37:15 to 7:37:25, strong downward fluxes are observed, peaking in flux within the loss cone at 7:37:17 and 7:37:21. From 7:37:25 to 7:37:27, Juno encounters an upward beam that transitions to an electron conic from 7:37:27 to 7:37:29 with increasing peak pitch angle. The conic diminishes for a single time step (1 s) and an intense upgoing electron flux is observed at 7:37:30 spread more broadly in pitch angle, but also dominantly contained in the upgoing loss cone.

Overlaid on Figure 2b are background-subtracted magnetic field fluctuations in the azimuthal component ($B_\phi$ in System III coordinates), smoothed over 0.3 s to remove high-frequency fluctuations. The fluctuations are extracted from Juno magnetic field vector observations after removing the planetary magnetic field by subtracting the JRM09 model [Kotsiaros et al. 2019]. Fluctuations reaching magnitudes up to ~200 nT are readily apparent in the transverse components (perpendicular to the magnetic field direction) and the absence of analogous compressive fluctuations (Supporting Information Figure 1) indicate that the currents are mostly field aligned. The magnetic field structure in $\delta B_\phi$ is typical of a northern Alfvén wing crossing with negative $\delta B_\phi$ within the main wing and positive $\delta B_\phi$ outside of the main wing [e.g. Neubauer 1980; Saur et al. 1999; Jacobsen et al. 2007]. The lack of positive $\delta B_\phi$ around 7:37:15 could be due to the detailed geometry of the tail substructures with respect to Juno’s crossing.

The directionality of the currents can be determined by their azimuthal component. Specifically, the sharp decrease in $\delta B_\phi$ around 7:37:20 followed by a gradual increase at 7:37:24 and a sharp increase around 7:37:30 suggests Juno crossed a localized upward current (downward electron beam) followed by a downward current (upward electron beam) typical of the northern sub-Jovian Alfvén wing flank [Saur et al. 1999, Jacobsen et al. 2007]. Figure 2d shows the power spectral densities of the transverse magnetic field fluctuations. Following previous analysis techniques [Gershman et al. 2019] by determining the RMS fluctuation power over the given interval and assuming the Alfvén speed is the speed of light in such a strong magnetic field, the total Poynting flux is calculated to be ~100 mW/m². The Waves instrument onboard Juno [Kurth et al. 2017] also observed enhancements in electric and magnetic field fluctuations coincident with the tail transit time identified here and support these directionalities (Supporting Information Figures 2–4). The power in the magnetic field fluctuations residing in the non-compressional (transverse) component up to 5 Hz, together with the absence of compressive magnetic fluctuations, is consistent with the presence of a non-dispersive Alfvénic mode transmitting field-aligned currents.

Figure 3a shows the intensity (differential number flux, the DEF divided by proton energy for each energy bin), along with characteristic energy for the precipitating electrons. The characteristic energy is calculated via $E_{\text{char}} = \sum_i I_i E_i \Delta E_i / \sum_i I_i \Delta E_i$ [e.g. Clark et al. 2018], where $I$ is the intensity and the sum is performed over an energy range from 50 eV to 40 keV. The characteristic energy varies from ~0.1 keV to 0.5 keV in the peak downward flux region, with a dip at ~0.1 keV at the center of the feature. Figure 3b shows intensity spectra for all timesteps in the top panel, color coded by the bar at the top of the figure. During the times around the peak downward energy flux (blue/red), the electron spectra exhibit a broad energization above ~0.5 keV. There is no clear evidence for a peak in the intensity profile indicative of a quasi-static parallel potential structure. The “sawtooth” pattern in some of the spectra is an instrumental effect due to JADE observing rapidly evolving electron fluxes, that
vary on timescales shorter than the 1 s it takes for JADE to sweep every other energy step through all available energies.

4. Discussion and Conclusions

Juno has provided the first in situ measurements of particles and fields in the Ganymede footprint tail. Magnetic field measurements indicate strong Alfvénic activity, with Poynting fluxes ~100 mW/m² that provide sufficient energy to power the Ganymede footprint tail aurora. The observed Alfvén waves can accelerate electrons into the Jovian atmosphere, where JADE observed precipitating electron energy fluxes peaking at 11 mW/m², ~10% of the available Poynting flux energy. The precipitating electrons observed by JADE are sufficient to generate the instantaneously measured auroral emission from the Ganymede footprint tail. UVIS observed the main auroral spot (Main Alfvén Wing – MAW) and near-spot auroral tail at this time to have an average H₂ photon flux of 130 ± 30 kR in the region associated with the Ganymede footprint (within the circle pointed to with the arrow in Figure 1). This value was calculated by summing the brightness within the 155–162 nm spectral range and then multiplying by a factor of 8.1 to scale it to the entire H₂ Lyman and Werner bands UV spectrum [Bonfond et al., 2017c], where the error bar is due to systematic uncertainties in the radiometric calibration [Gérard et al. 2019]. The UV intensities correspond to a precipitating electron energy flux of 13 ± 3 mW/m² [Grodent et al. 2001; Gustin et al. 2012]. These observed values are also fully consistent with the expected values at Δλ₆ = 8˚ angular separation from the MAW of 2-63 mW/m², assuming the fluxes are ~50% of the value at the MAW with an e-folding distance of 11˚ [Bonfond et al. 2017a]. Figure 4 summarizes these findings.

Before and after the Ganymede footprint tail transit, the downward electron intensity spectra were power-law-like as a function of energy, monotonically decreasing with energy resembling a power-law function and without substantial substructure in energy. During the period of largest downward energy flux, the intensities increased by roughly an order of magnitude, with broad enhancements in the range of 0.5 to 40 keV. These downward electron spectra are characteristic of broadband acceleration, similar to those observed throughout Jupiter’s main auroral [Mauk et al. 2017b] and polar region [Ebert et al. 2017] emissions, and do not exhibit peaked features associated with discrete acceleration. While the shape of these spectra is not easily approximated by any of the canonical analytic distributions, to compare with similar previous fits, we fit the spectra to a Kappa distribution [e.g. 3.12 of Livadiotis and McComas, 2013], which yielded Kappa values of K = 1.5 – 2.8. A Kappa value of 2.4 was found to best fit the precipitating electron fluxes at Io [Bonfond et al. 2009] and is within the range of values reported in this study for Ganymede, suggesting similar acceleration processes are involved.

At 8˚ separation from the MAW, the emissions map to a distance of ~60 R₉ downstream in Ganymede’s equatorial plane. During this event, Ganymede is at 280˚ System III longitude, such that it is within the northern portion Jupiter’s plasma sheet and near the peak emission longitude of ~310˚ [Grodent et al. 2009]. The broad structure of the feature is also consistent with Ganymede’s MAW spot size. Mapping the width of the feature discussed here down to an auroral emission altitude of 400 km above the 1-bar level, it spans a distance of 660 km in Jupiter’s atmosphere (Supporting Information Table 1), consistent with high-resolution observations showing the MAW spot to be ~700 km wide [Mura et al. 2018]. The auroral current system is also highly structured, where the sawtooth/checkerboard pattern observed by JADE suggests the electron fluxes are varying on local scales smaller than 50 km (the distance Juno transits locally in JADE’s 1 s sampling period).
During the transit, Juno crossed a current system with separate upward and downward current regions adjacent to each other and exhibiting small-scale structures. In addition to these small-scale variations, electron beams and conics are observed in the upward current region (Figure 2b). Since there is no corresponding downward beam during this time, and more importantly since this population is in the loss cone, the conic is not a reflected loss-cone distribution and instead indicates heating and/or acceleration below the spacecraft’s altitude of 0.5 R$_J$

The up-going electron beam at 7:37:25 indicates parallel acceleration occurs below the spacecraft’s altitude of 0.5 R$_J$. If this beam and adjacent conic are generated by similar acceleration mechanisms, the conic could also have been driven by parallel acceleration. Low frequency parallel electric field fluctuations, such as those from Alfvén waves, have been found to be able to generate electron conics [André and Eliasson, 1992]. There is a trend of increasing peak pitch angle and decreasing characteristic energy and total energy flux in this conic. If the beam and conic are all formed due to similar parallel electric field fluctuations, this could indicate the acceleration region is getting further from Jupiter (closer to Juno) and decreasing in strength as Juno moves across the downward current region.

We observe a broad distribution of upgoing pitch angles at the edge of the downward current region, which suggests that acceleration occurs at a more dispersed range of altitudes near this boundary. The field-aligned upward electron beams in the loss cone in the downward current region could also generate conjugate auroral emissions in the southern hemisphere. However, if a similar acceleration process is occurring in the southern hemisphere, the lack of a down-going conic in the upward current regions suggests that the downward electron flux may be more locally accelerated, and not due to a conjugate acceleration near the southern ionosphere. The one-way transit time to travel between hemispheres is ~16 min for 50 eV and ~0.5 min for 40 keV electrons and Ganymede has moved less than 1˚ during the transit. If such a conjugate population exists, they could have been pitch-angle scattered during this transit. Ganymede’s auroral tail features are also observed to be patchy and highly structured [Mura et al. 2018], therefore it is possible that the features discussed in this study are a consequence of where Juno has crossed this patchy tail. Additionally, there is no evidence for a bifurcated tail like that intermittently observed for Io [Mura et al. 2018; Szalay et al. 2018].

While this study is focused on the electrons responsible for generating the auroral tail, there is some indication of proton acceleration (Supporting Information Figure 2), similar to that observed during the Io footprint tail transits [Szalay et al. 2019]. Unfortunately, JADE does not have complete ion pitch angle coverage through this observation. Future studies could compare these accelerated protons to those accelerated in the Io footprint tail aurora, particularly on how the source of available protons and the Alfvén reflection scheme through the plasma sheet differs compared to the Io torus. Additionally, comparisons could be made between Io’s and Ganymede’s signatures in the high energy component, where energetic proton depletions [Paranicas et al. 2019] have been observed connected to Io’s footprint tail aurora.

By simultaneously comparing observations of the Alfvén waves responsible for electron acceleration, the precipitating electrons they accelerate, and the auroral emissions those electrons produce, we confirm that Ganymede’s footprint tail is sustained primarily by Alfvénic acceleration processes. Together with the finding that Io’s footprint tail is also sustained primarily by similar processes [Bonfond et al. 2009; Hess et al. 2010; Szalay et al. 2019; Gershman et al. 2019; Damiano et al. 2019], Alfvénic acceleration appears to be a characteristic feature in Jupiter’s moon-magnetosphere interactions and further bolsters the conclusion that broadband acceleration processes are dominant in Jupiter’s aurorae [Mauk et al. 2017b; Saur
et al. 2018]. Saturn’s moon Enceladus also exhibits similar auroral acceleration [Sulaiman et al. 2018], hence, similar processes are operating at multiple outer planets. By extension, Alfvénic acceleration is likely a more universal process applicable to similar interactions at other exoplanet-moon systems [Saur et al. 2013; Bonfond et al. 2017a].

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References


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Figure 1. (Top) UV false color of the auroral emissions near Juno’s mapped magnetic footprint. Grey areas mark regions UVS either did not observe or registered no photons while viewing. Qualitatively, red features correspond to emissions deep in the atmosphere due to higher energy precipitating electrons, while pink/white features correspond to higher altitude emissions due to lower energy incident electrons. (Bottom) The color spectrogram shows JADE downward precipitating differential energy flux (DEF). The overlaid line shows energy flux (EF) within the loss cone, with the separate EF axis on the right of the spectrogram.
Figure 2. JADE and MAG data during the Ganymede footprint tail transit. (a) and (c) show the downward and upward electron differential energy flux (DEF) within the loss cone respectively. Precipitating energy flux is overlaid on these panels with its separate axes on the right of the spectrograms. (b) shows electron pitch angles with $\delta B_{\parallel}$ overlaid. (d) shows the transverse B-field power spectral densities. Red/blue bars indicate approximate up/downward current regions inferred from MAG data.
Figure 3. (a) Precipitating electron intensity spectrogram, with the characteristic energy from 0.05-40 keV overlaid. (b) Intensity for individual spectra, color-coded by the bar at the top of (a).
Figure 4. Summary schematic of the auroral processes involved in sustaining Ganymede’s footprint tail aurora 8° downtail. Juno observations are consistent with a 10% efficiency coupling Alfvénic power to electron acceleration.