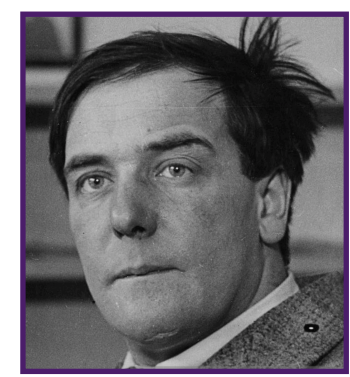
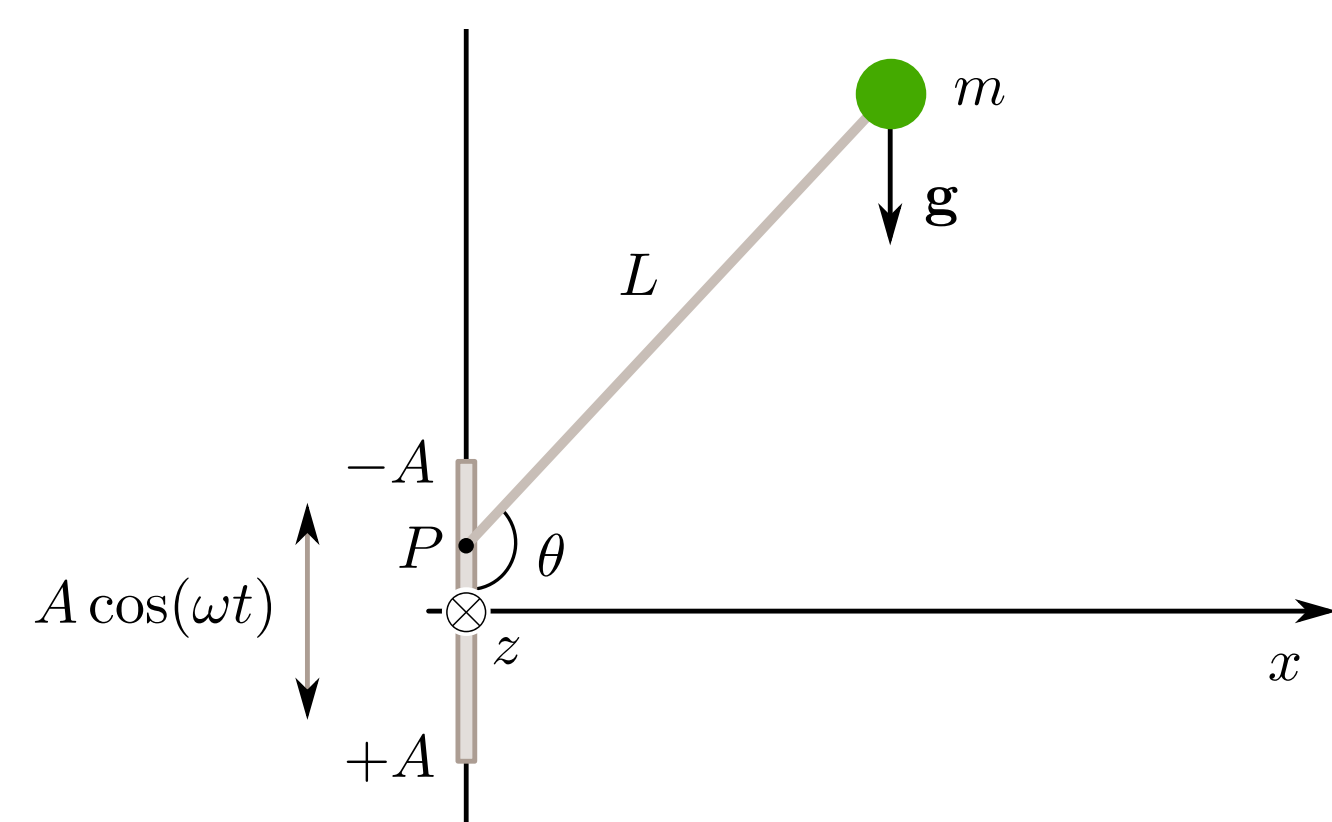


Inverted driven pendulum

Pendulum submitted to an oscillating force of frequency ω with zero average value:



Piotr Kapitza (1894-1984)



Natural oscillation frequency: $\omega_0 = \sqrt{\frac{g}{L}}$
 Reduced acceleration: $\varepsilon = \frac{A\omega^2}{g} = \frac{A}{L} \left(\frac{\omega}{\omega_0}\right)^2$
 Moment of inertia: $I = mL^2$
 Gravitational potential energy: $U_L = mgL = I\omega_0^2$
 Hamiltonian: $H = \frac{p_\theta^2}{2I} - U_L \cos \theta (1 + \varepsilon \cos(\omega t))$

Effective potential and stabilization

→ timescale separation when the pivot point vibrates with a small amplitude ($A \ll L$) and with a large frequency ($\omega \gg \omega_0$):

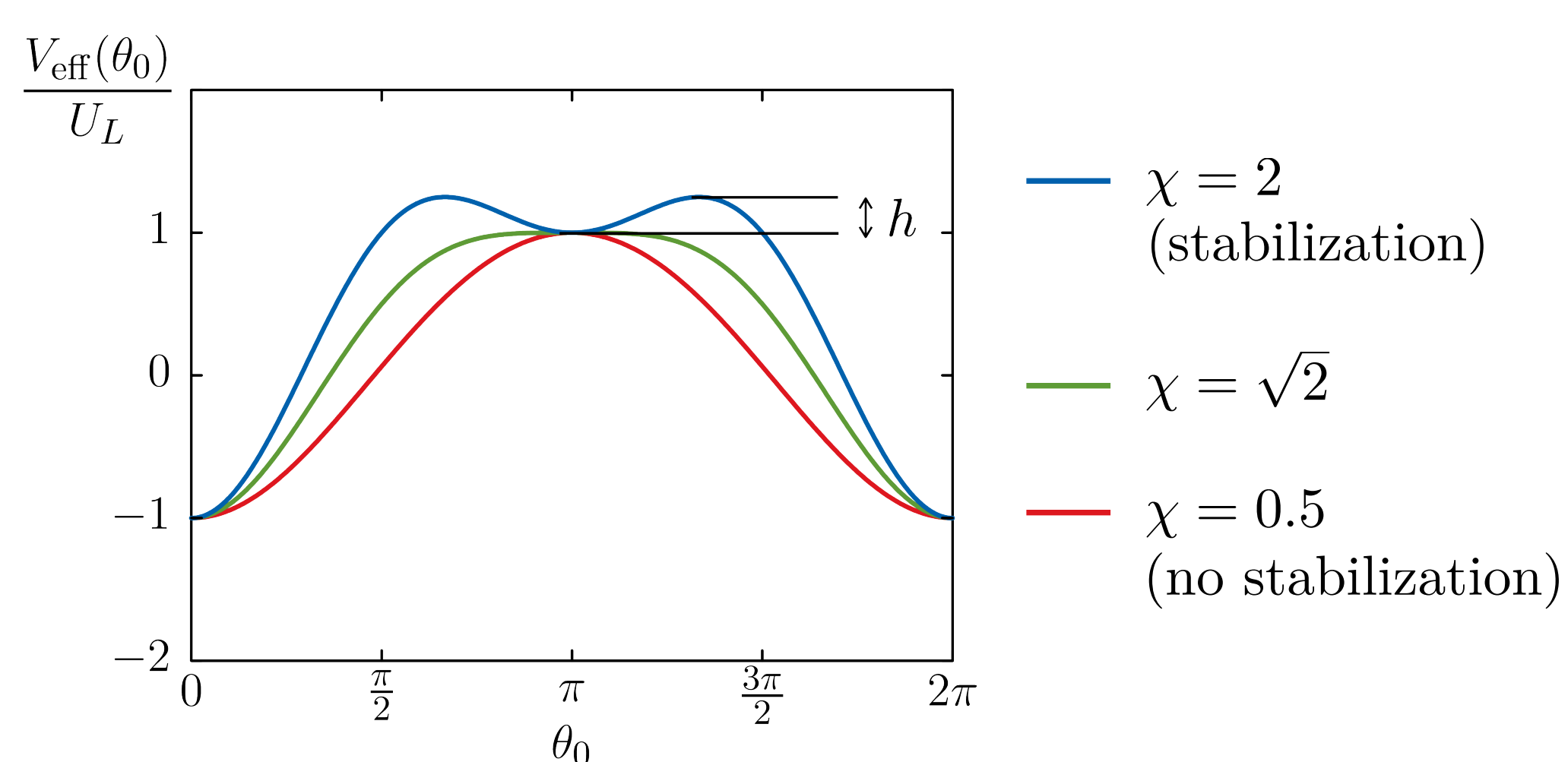
$$\theta(t) = \theta_0(t) + \xi(t) \quad \text{with} \quad |\xi(t)| \ll 2\pi$$

$\theta_0(t)$: slowly varying component

$\xi(t) = \frac{A}{L} \sin \theta_0(t) \cos(\omega t)$: rapidly varying component

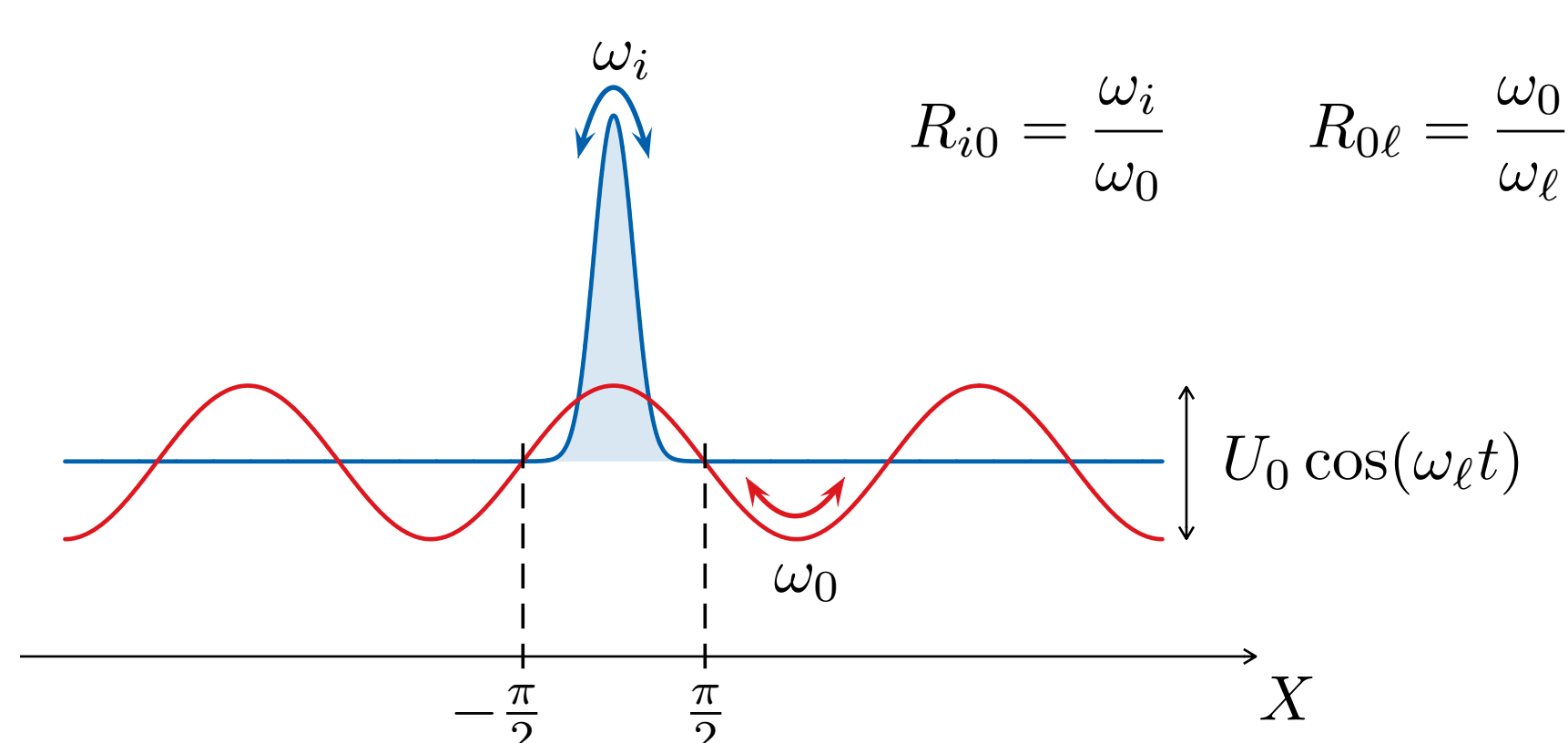
→ Euler-Lagrange equation for the slow component with the effective potential:

$$V_{\text{eff}}(\theta_0) = \frac{F^2}{2m\omega^2} = -U_L \left(\cos \theta_0 - \frac{\chi^2}{4} \sin^2 \theta_0 \right) \quad \text{with} \quad \chi = \frac{A}{L} \frac{\omega}{\omega_0} = \varepsilon \frac{\omega_0}{\omega}$$



BEC in oscillating optical lattice

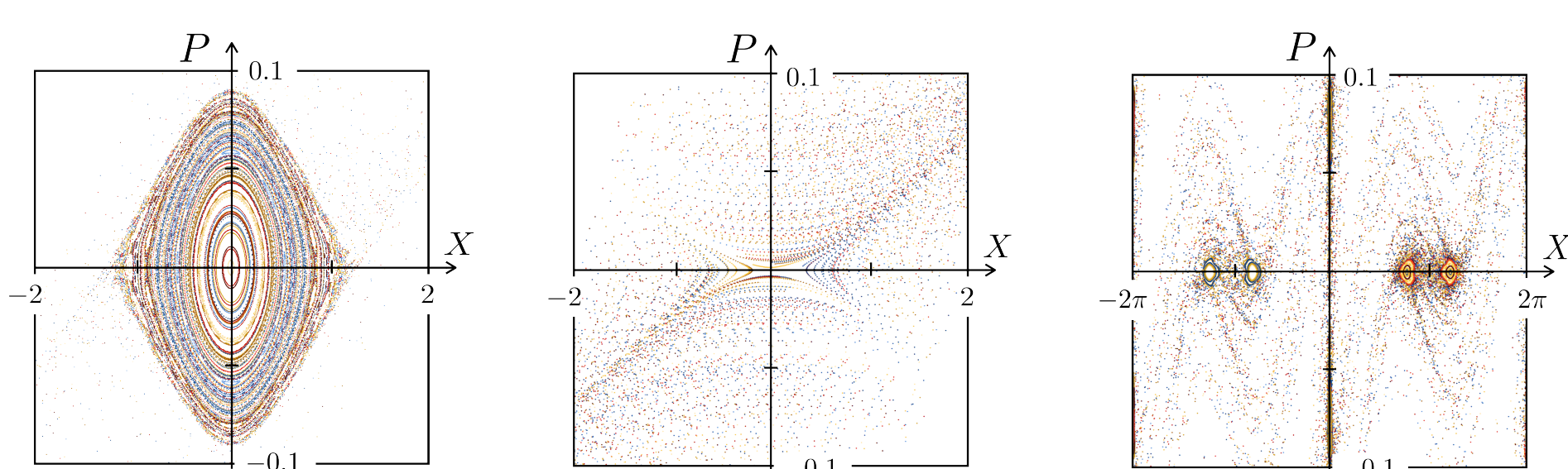
Idea: repulsive atom-atom interactions give rise to a mean-field potential with an unstable fixed point that can be transformed into a stable fixed point by subjecting the BEC to an optical lattice with oscillating amplitude (oscillating force)



Classical dynamics

$$\text{Hamiltonian: } \frac{H}{U_0} = 8P^2 - R_{i0}^2 X^2 + \cos X \cos(2\pi T)$$

→ inverted harmonic potential
→ driving



Poincaré sections for the frequency ratios $R_{i0} = 0.075$ and $R_{0l} = 0.45$ (left), $R_{i0} = 0.15$ and $R_{0l} = 0.2$ (middle), and $R_{i0} = 0.02$ and $R_{0l} = 0.7$ (right).

Kapitza stabilization of a repulsive BEC

Quantum evolution with GPE

In reduced variables, Gross-Pitaevskii equation reads

$$i\partial_T \Psi = 4\pi R_{0l} \left(-\hbar_{\text{eff}} \partial_X^2 + \frac{\cos(2\pi T) \cos X}{8\hbar_{\text{eff}}} + \frac{\bar{g}|\Psi|^2}{2} \right) \Psi$$

X : position in units of the lattice spacing

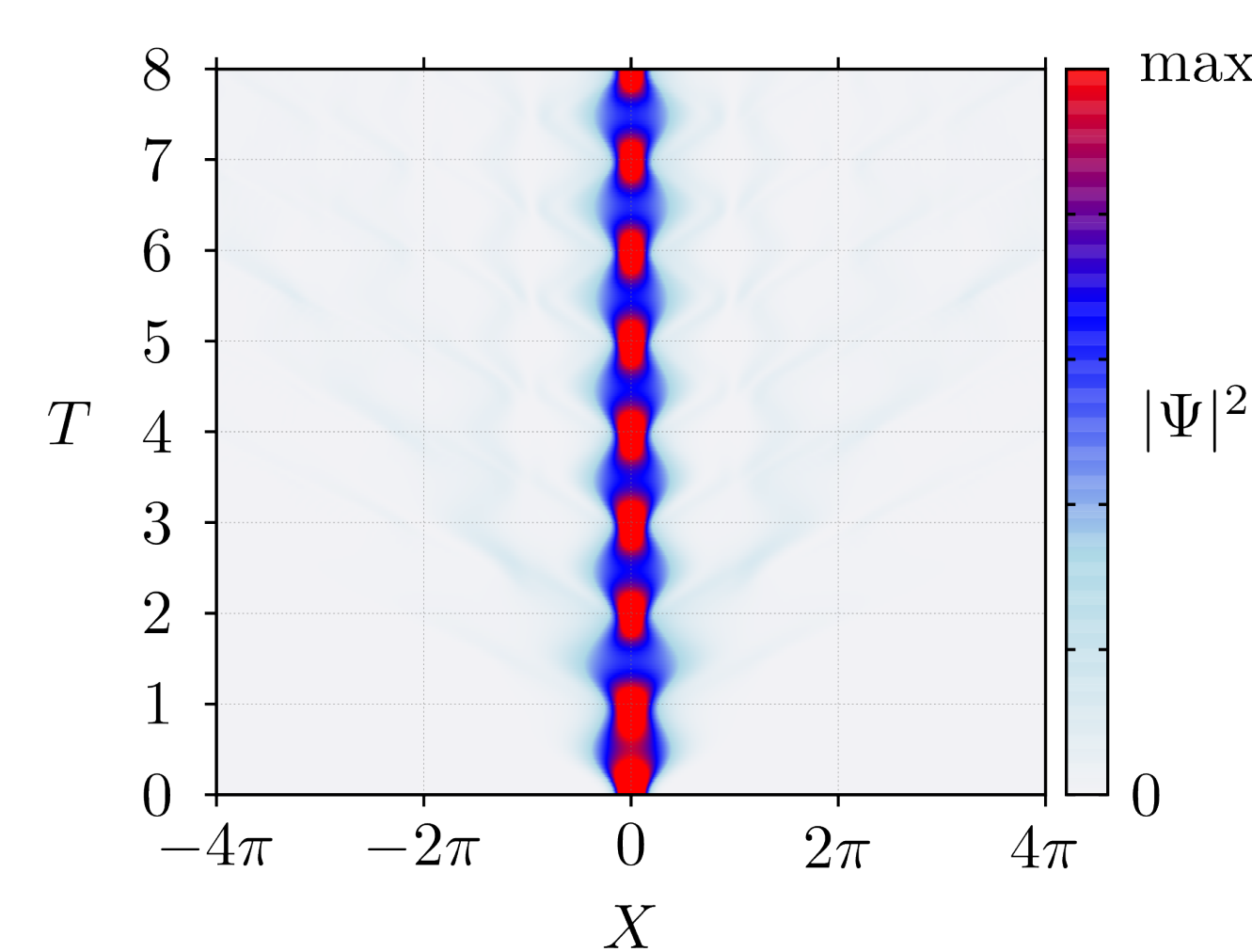
T : time in units of the driving period

$R_{0l} = \omega_0/\omega_l$: frequency ratio

$\hbar_{\text{eff}} \propto U_0^{-1/2}$: effective Planck constant

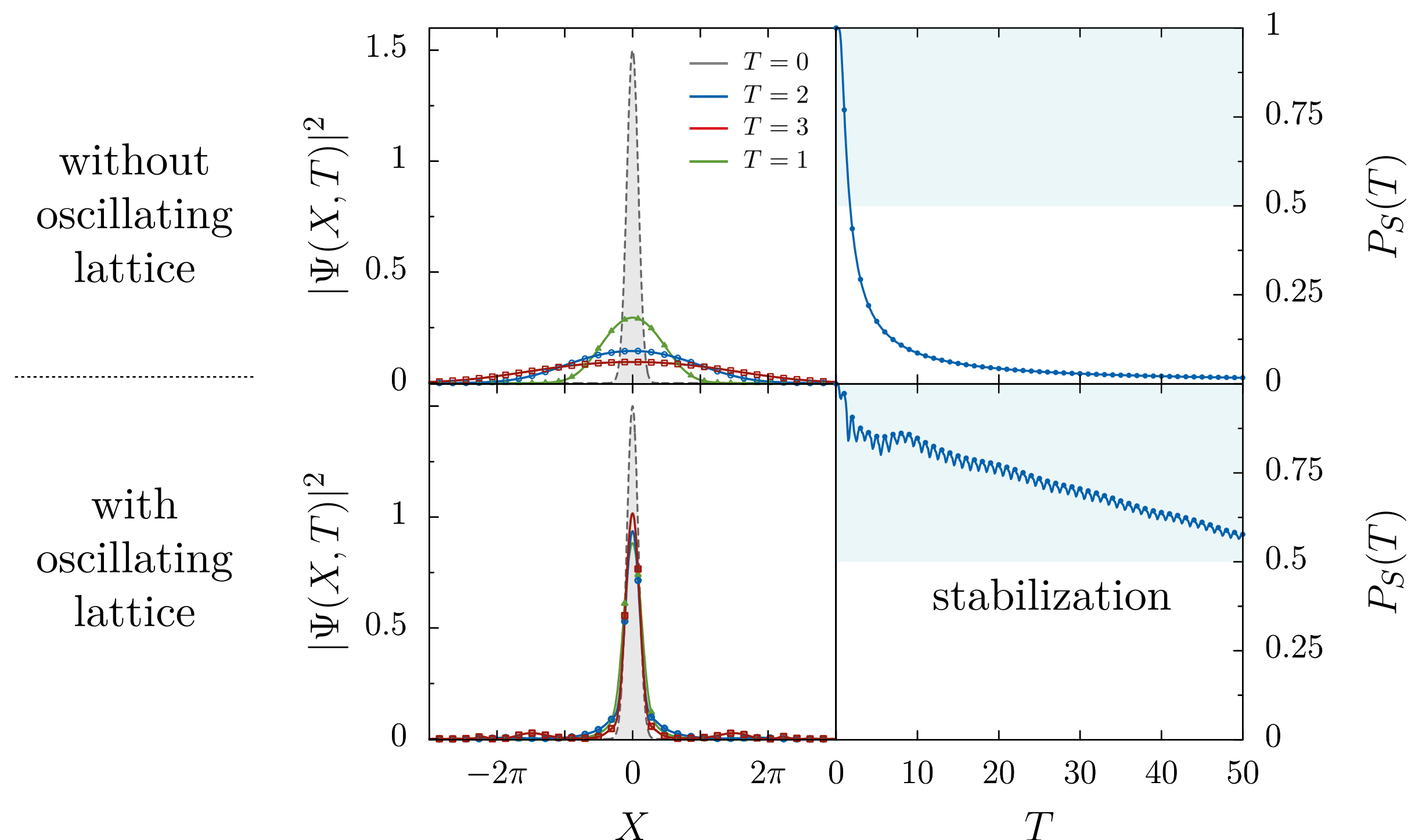
\bar{g} : atom-atom interaction strength

Probability density evolution

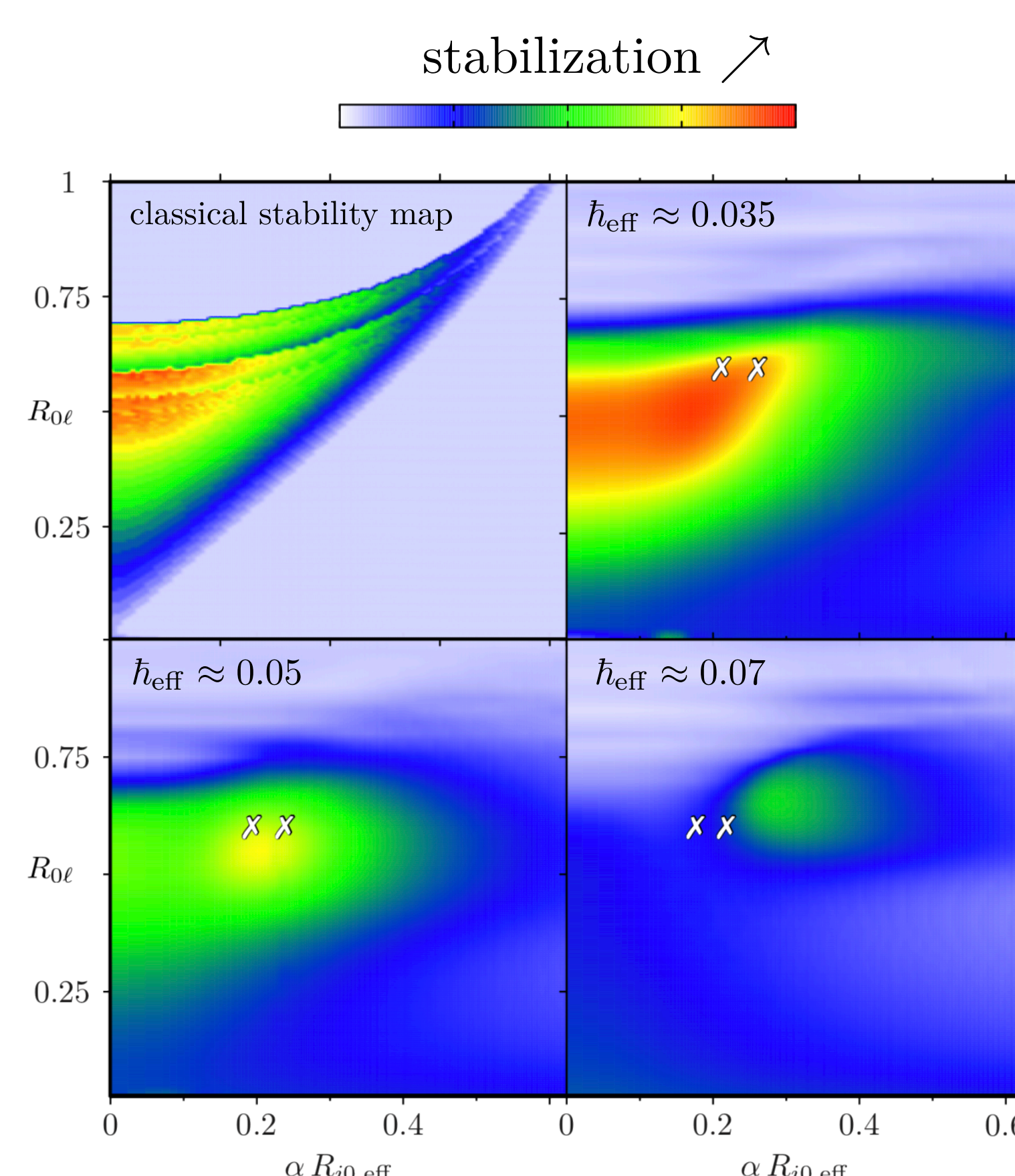


Probability density as a function of position and time for $R_{0l} = 0.6$, $\bar{g} = 0.6$, $\hbar_{\text{eff}} \approx 0.035$ and an initial Gaussian wave packet with $R_{i0,\text{eff}} \approx 1.34$.

$$\text{Probability inside the initial "well": } P_S(T) = \int_{-\pi/2}^{+\pi/2} |\Psi(X, T)|^2 dX$$



Stability diagram



Stabilized fraction of the condensate as quantified by $\overline{P_S} = \frac{1}{10} \int_{3T_{\text{sp}}}^{3T_{\text{sp}}+10} P_S(T) dT$ where T_{sp} is a characteristic spreading time of the initial wave packet in the absence of driving ($U_0 = 0$) determined by $P_S^{=0}(T_{\text{sp}}) = 0.75$, as a function of R_{0l} and $R_{i0,\text{eff}} = \frac{\omega_{i,\text{eff}}}{\omega_0}$ (multiplied by a scaling coefficient $\alpha = 0.2$) with $\frac{\omega_{i,\text{eff}}}{\omega_0} \propto \sqrt{\frac{\hbar_{\text{eff}} \bar{g}}{\sigma^3}}$ for an initial Gaussian wave packet of width σ .

References

See J. Martin, B. Georgeot, D. Guéry-Odelin and D. L. Shepelyansky, Phys. Rev. A **97**, 023607 (2018) & arXiv:1709.07792 and references therein