

# Bayesian Inference on Uncertain Kinetic Parameters for the Pyrolysis of Composite Ablators

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**Garden Grove, California, USA**

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# Motivation: atmospheric entry

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## Thermal protection systems (TPS)



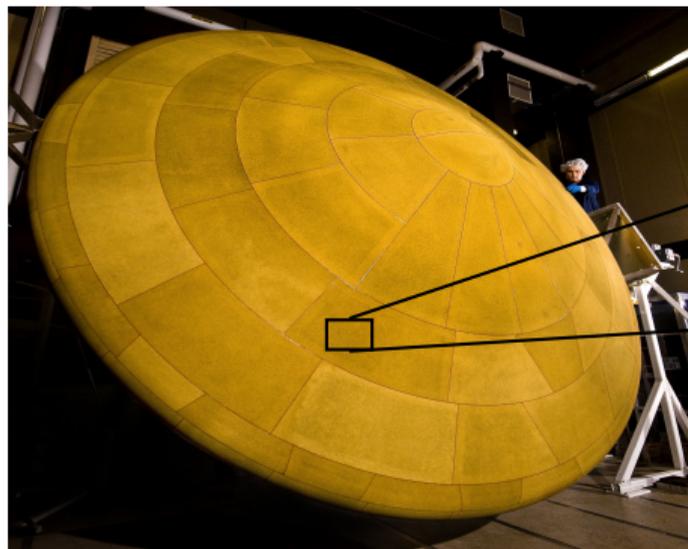
Dragon capsule (Space X)

## Space debris



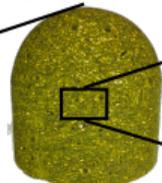
ATV-1 *Jules Verne* (ESA)

# Ablative materials for thermal protection systems

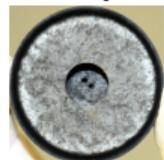


Mars Science Laboratory thermal protection system (NASA)

Porous thermal protection material



[Helber, 2016]



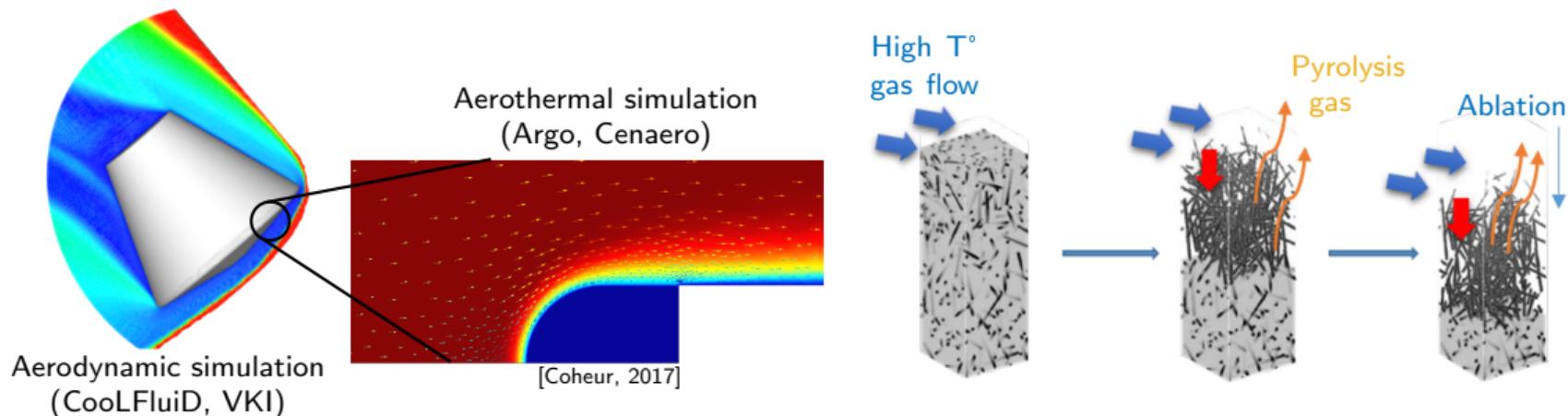
Bottom view  
(after burn)

Fibers  
+ resin



Lawson et al.

# Modeling the pyrolysis of ablative thermal protection materials



- ▶ CFD codes requires accurate model for ablation [Lachaud, 2014; Schrooyen, 2016], e.g. conservation of mass species

$$\frac{\partial \epsilon_g \langle \rho_i \rangle_g}{\partial t} + \nabla \cdot (\epsilon_g \langle \rho_i \rangle_g \langle u \rangle_g) = \nabla \cdot \langle J_i \rangle + \langle \dot{\omega}_i^{\text{pyro}} \rangle$$

- ▶ The governing laws for the pyrolysis gas production  $\langle \dot{\omega}_i^{\text{pyro}} \rangle$  must be derived from recent experimental studies and their associated uncertainties must be quantified

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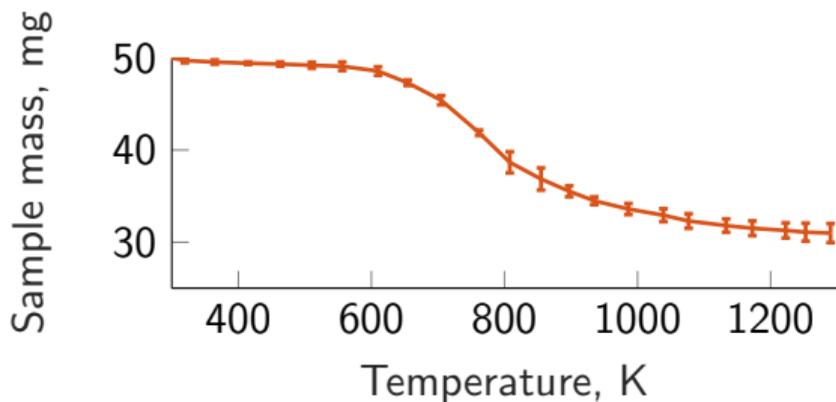
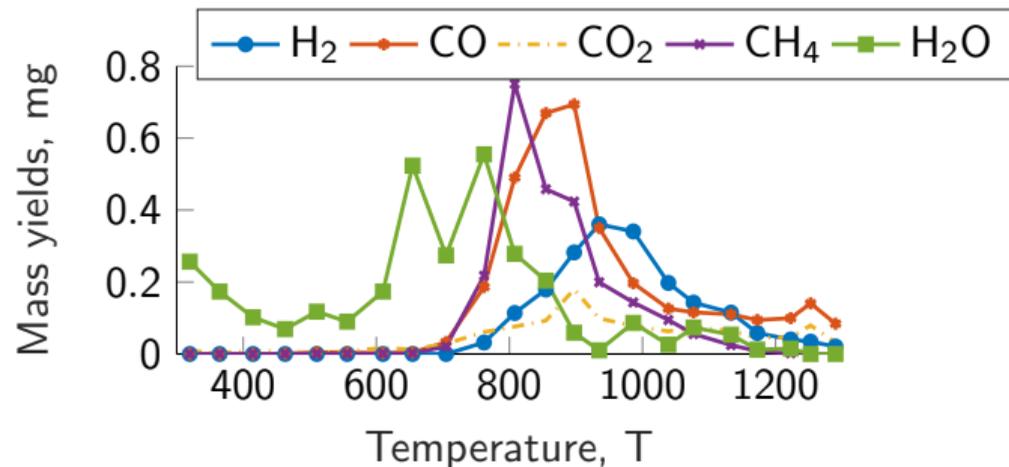
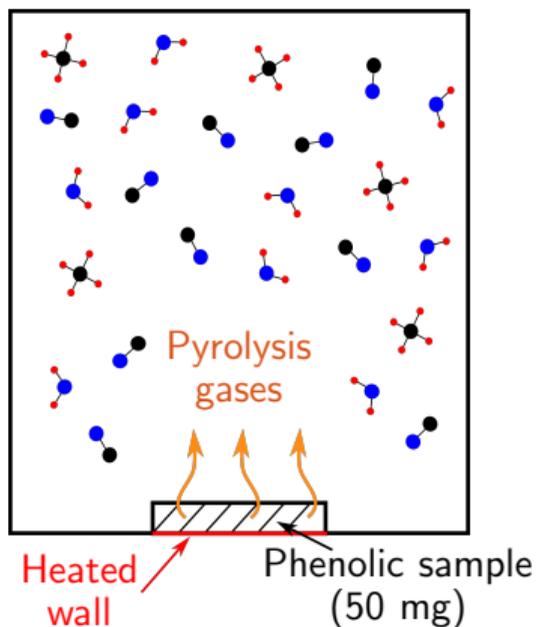
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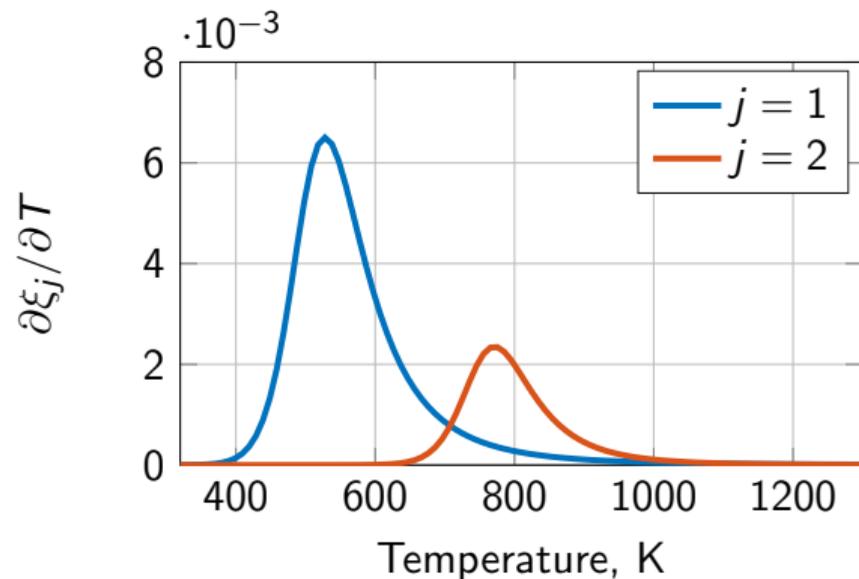
Conclusions

# Pyrolysis experiments [Wong et al., 2015]



# Phenomenological laws for pyrolysis mass loss and gas species production

- ▶ Previous pyrolysis experiments showed that the pyrolysis decomposition of ablative materials follows successive reaction rates [Goldstein, 1969; Trick, 1997]

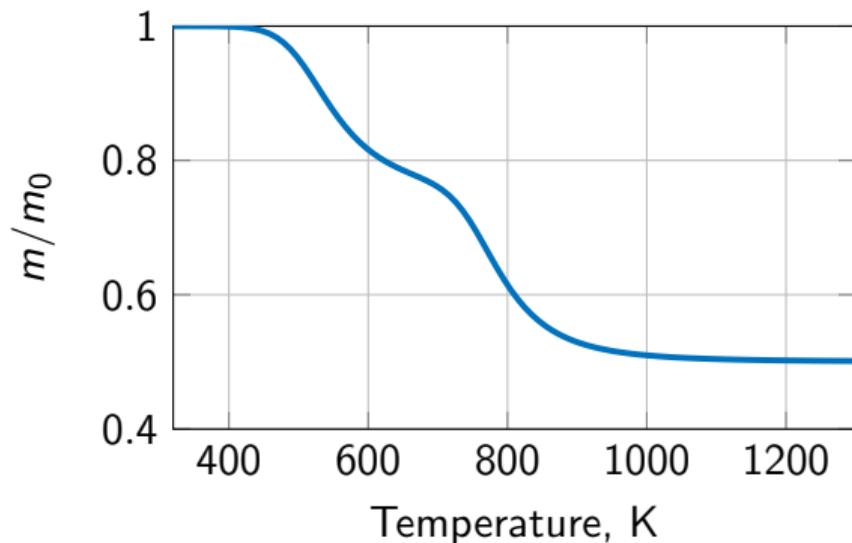


$$\langle \dot{\omega}_i^{\text{pyro}} \rangle = \sum_j^{N_p} F_{ij} \frac{\partial \xi_j}{\partial T} \tau m_0$$
$$\frac{\partial \xi_j}{\partial T} = (1 - \xi_j)^{n_j} \frac{A_j}{\tau} \exp\left(-\frac{E_j}{\mathcal{R}T}\right)$$

- ▶  $\xi_j$ : advancement of reaction of the fictitious resin component  $j$

# Phenomenological laws for pyrolysis mass loss and gas species production

- ▶ Previous pyrolysis experiments showed that the pyrolysis decomposition of ablative materials follows successive reaction rates [Goldstein, 1969; Trick, 1997]



$$m = m_0 - \sum_j^{N_p} F_j \xi_j m_0$$
$$\frac{\partial \xi_j}{\partial T} = (1 - \xi_j)^{n_j} \frac{A_j}{\tau} \exp\left(-\frac{E_j}{\mathcal{R}T}\right)$$

- ▶  $\xi_j$ : advancement of reaction of the fictitious resin component  $j$

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# Mathematical description

- ▶ Forward problem:  $d_i = \eta(\mathbf{x}_i, \mathbf{p})$

code output ←  $\eta$  → model

- ▶ Inverse problem: Find  $\mathbf{p}$  knowing the observations  $d_i^{\text{obs}}$  associated with errors  $\varepsilon_i$

$$d_i^{\text{obs}} = \eta(\mathbf{x}_i, \mathbf{p}) + \varepsilon_i$$

Deterministic inverse problem

*What is the best estimate of  $\mathbf{p}$ ?*

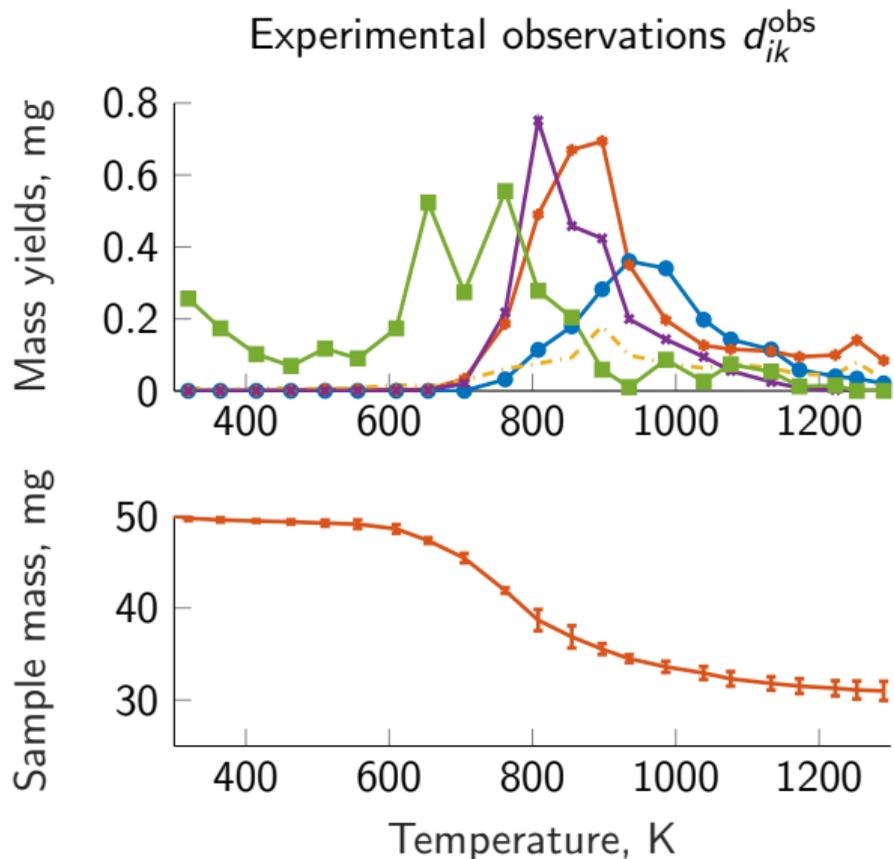
$$\tilde{\mathbf{p}} = \arg \min_{\mathbf{p} \in \mathcal{P}_{\text{ad}}} \|\mathbf{d}^{\text{obs}} - \mathbf{d}\|$$

Statistical inverse problem

*What is the plausibility about  $\mathbf{p}$ ?*

$$\pi(\mathbf{p}|\mathbf{d}^{\text{obs}}) = \frac{\pi(\mathbf{d}^{\text{obs}}|\mathbf{p})\pi_0(\mathbf{p})}{\int_{\mathbb{R}^p} \pi(\mathbf{d}^{\text{obs}}|\mathbf{p})\pi_0(\mathbf{p})d\mathbf{p}}$$

# Pyrolysis experiments: the forward problem



## Code outputs

- ▶ Collected mass of species  $i$

$$d_{ik} = \int_{t_{k-1}}^{t_k} \langle \dot{\omega}_i^{pyro} \rangle |_{T=T_k} dt$$
$$= \sum_j^{N_p} F_{ij} m_0 \left( \xi_j^{(t_f)} - \xi_j^{(t_i)} \right)$$

- ▶ Mass of the sample

$$m_k = m_0 - \sum_{i=1}^{n_s} \sum_{l=1}^k d_{il}$$

## Bayesian inference: the inverse problem

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$$\pi(\mathbf{p}|\mathbf{d}) = \frac{\pi(\mathbf{d}|\mathbf{p})\pi_0(\mathbf{p})}{\int_{\mathbb{R}^p} \pi(\mathbf{d}|\mathbf{p})\pi_0(\mathbf{p})d\mathbf{p}} \quad (\text{Bayes' theorem})$$

- ▶ Choice for the prior  $\pi_0(\mathbf{p})$ : uniform pdf (bounded support)
- ▶ Choice for the likelihood  $\pi(\mathbf{d}^{\text{obs}}|\mathbf{p})$

$$\pi(\mathbf{d}^{\text{obs}}|\mathbf{p}) = \frac{1}{\prod_i (2\pi\sigma_i^2)^{n/2}} \exp\left(-\sum_{i=1}^{n_s} \sum_{k=1}^n \frac{[d_{ik}^{\text{obs}} - \eta_i(\mathbf{x}_k, \mathbf{p})]^2}{2\sigma_i^2}\right)$$

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## Metropolis-Hastings algorithm

1. Initialize  $\mathbf{p}_0$
2. Iteration  $i$  :
  - a) Draw a new value  $\mathbf{p}^*$  from proposal distribution  $J(\cdot|\mathbf{p}_{i-1})$ .  $J$  specifies  $\mathbf{p}^*$  based on previous value  $\mathbf{p}_{i-1}$ , and  $J(\cdot|\mathbf{p}_{i-1})$  should not be interpreted as a conditional density.
  - b) Construct acceptance probability

$$\alpha(\mathbf{p}^*|\mathbf{p}_{i-1}) = \min \left( 1, \frac{\pi(\mathbf{p}^*|\mathbf{d})J(\cdot|\mathbf{p}_{i-1})}{\pi(\mathbf{p}_{i-1}|\mathbf{d})J(\mathbf{p}_{i-1}|\cdot)} \right) \quad (1)$$

- c) Sample a Uniform(0,1) random variable  $U$ . Set

$$\mathbf{p}^j = \begin{cases} \mathbf{p}^*, & \text{if } U \leq \alpha(\mathbf{p}^*|\mathbf{p}_{i-1}) \\ \mathbf{p}^{j-1}, & \text{otherwise} \end{cases} \quad (2)$$

- ▶ Choice for the proposal  $J(\mathbf{p}^* | \mathbf{p}_{i-1}) \rightarrow$  random walk metropolis

$$J(\mathbf{p}^* | \mathbf{p}_{i-1}) = \frac{1}{\sqrt{(2\pi)^m \det \Sigma}} \exp \left( -\frac{1}{2} (\mathbf{p}^* - \mathbf{p}_{i-1})^T \Sigma^{-1} (\mathbf{p}^* - \mathbf{p}_{i-1}) \right)$$
$$\Sigma = \begin{pmatrix} \sigma_1^2 & \rho_{12}\sigma_1\sigma_2 & \cdots & \rho_{1m}\sigma_1\sigma_m \\ \rho_{21}\sigma_2\sigma_1 & \sigma_2^2 & \rho_{2m}\sigma_2\sigma_m & \\ \vdots & \vdots & \ddots & \vdots \\ \rho_{m1}\sigma_m\sigma_1 & \rho_{m2}\sigma_m\sigma_2 & \cdots & \sigma_m^2 \end{pmatrix}$$

- ▶ We assume  $\Sigma$  to be a diagonal matrix
- ▶ Using  $\log(A)$  and  $\log(E)$  values increases the mixing

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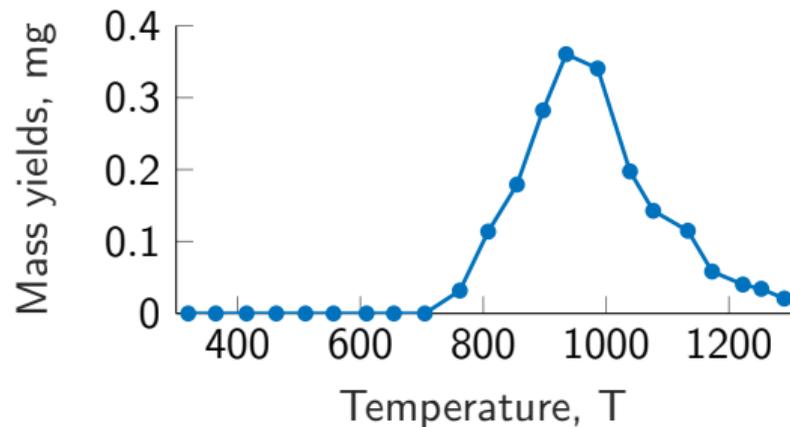
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## Numerical set-up: one reactant model with one species

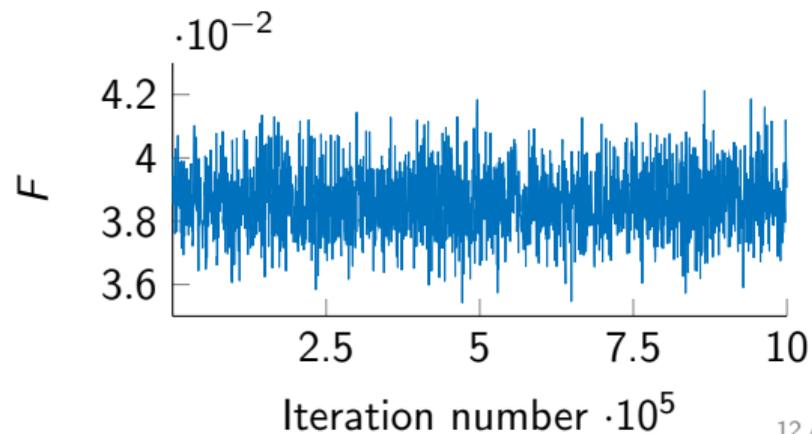
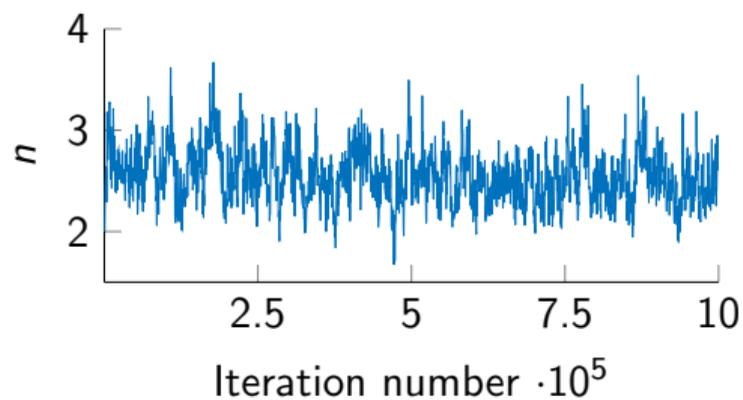
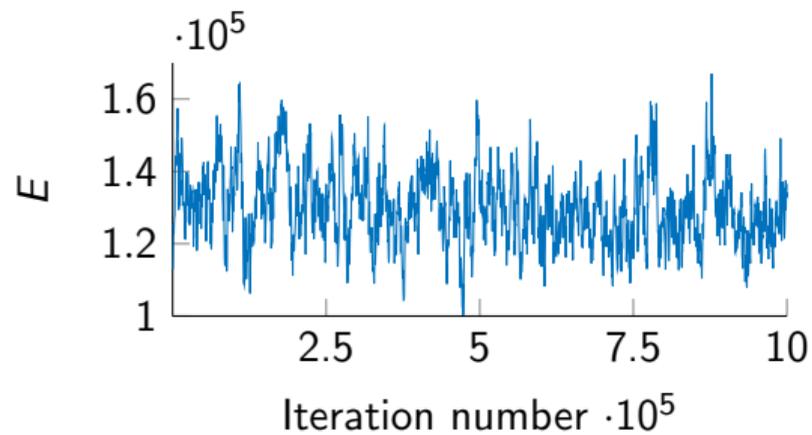
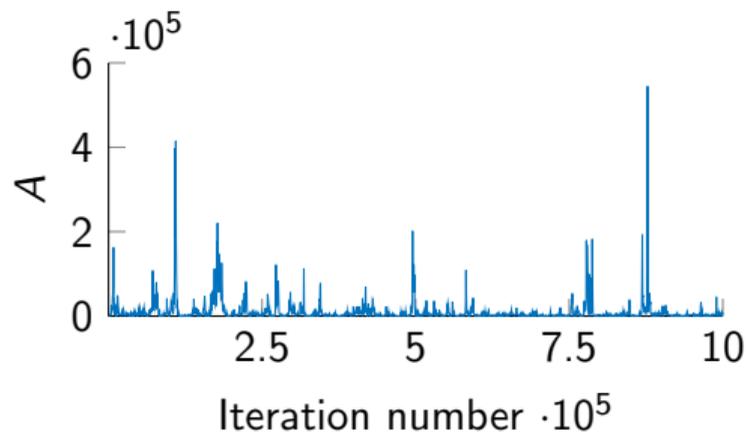
- ▶ Observable (data)  $d_{\text{H}_2k}^{\text{obs}}$ : mass yields of species  $\text{H}_2$  at each temperature.  $N_p = 1$ .

$$d_{\text{H}_2k} = F_{\text{H}_2} m_0 \left( \xi^{(k)} - \xi^{(k-1)} \right),$$
$$\frac{\partial \xi^{(k)}}{\partial t} = \left( 1 - \xi^{(k)} \right)^n A \exp \left( -\frac{E}{\mathcal{R} T_k} \right),$$
$$\xi^{(0)} = 0.$$

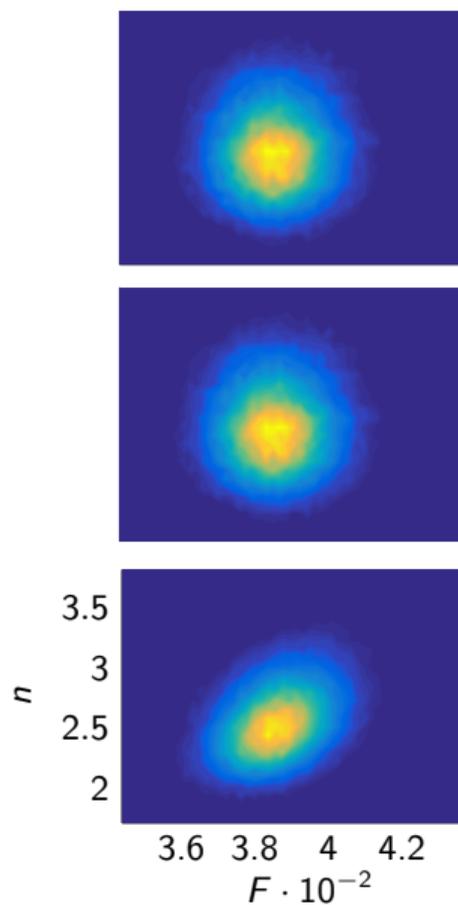
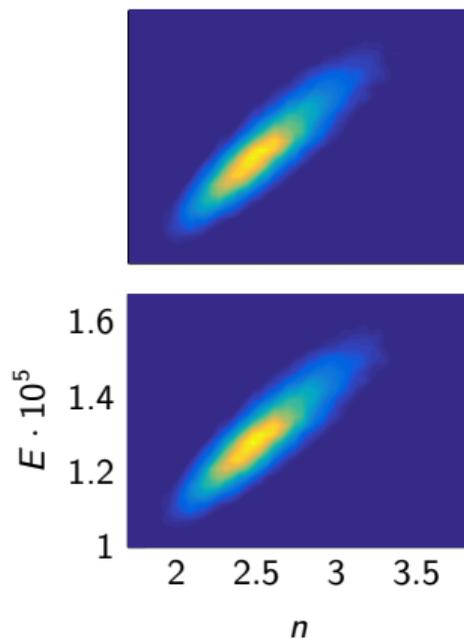
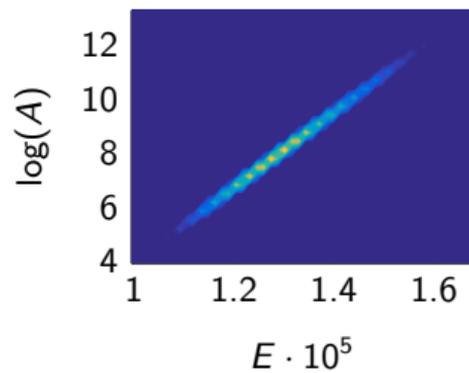


- ▶ Simple point-mass model (0D),  $k = 1, \dots, n_T$
- ▶  $\mathbf{p} = \{A, E, n, F_{\text{H}_2}\}$

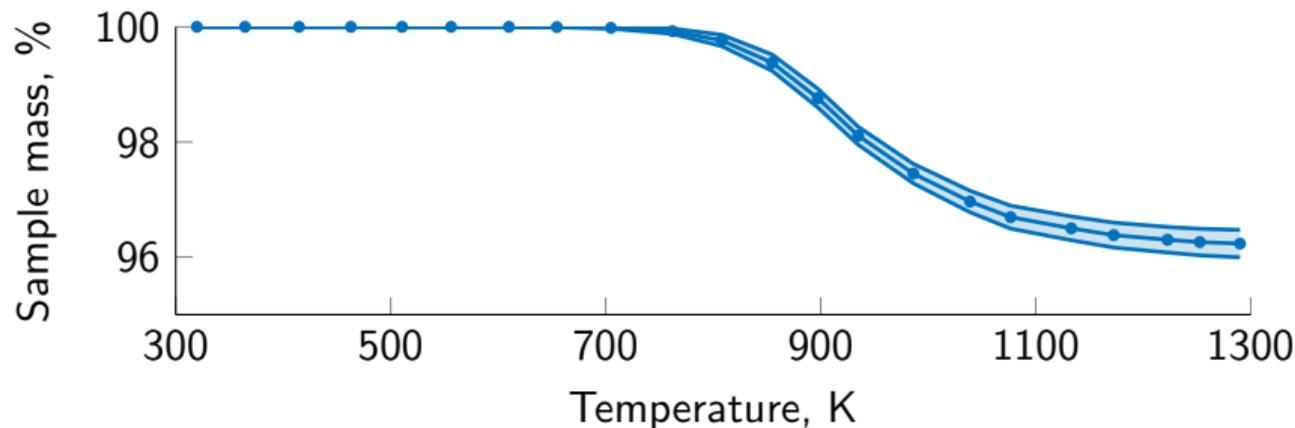
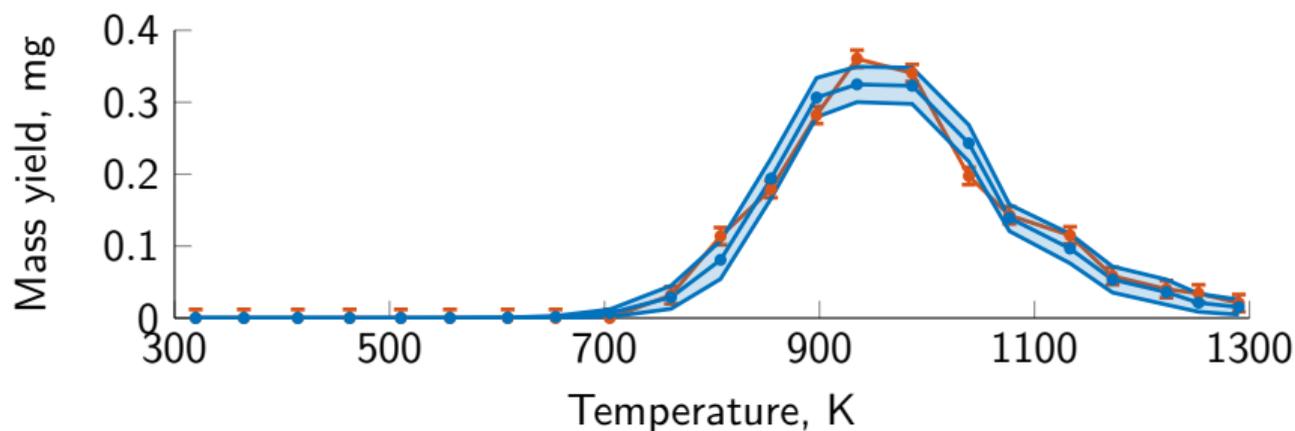
# Application to pyrolysis experiments: Markov-Chain



# Bivariate posterior PDFs



## Posterior predictive checks assuming 1 reactant (4 parameters)



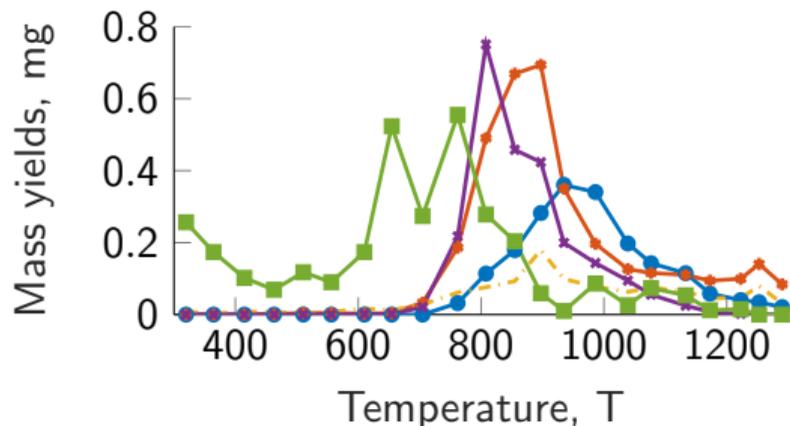
## Numerical set-up: 3 reactants model with 5 species

- Observable (data)  $d_{ik}^{\text{obs}}$  with  $i = \{\text{H}_2, \text{CO}, \text{CO}_2, \text{CH}_4, \text{H}_2\text{O}\}$ : mass yields at each temperature iteration  $k$ .  $N_p = 3$ .

$$d_{ik} = \sum_j^{N_p} F_{ij} m_0 \left( \xi_j^{(k)} - \xi_j^{(k-1)} \right),$$

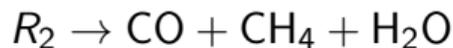
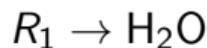
$$\frac{\partial \xi_j^{(k)}}{\partial t} = \left( 1 - \xi_j^{(k)} \right)^{n_j} A_j \exp \left( -\frac{E_j}{RT_k} \right),$$

$$\xi_j^{(0)} = 0.$$

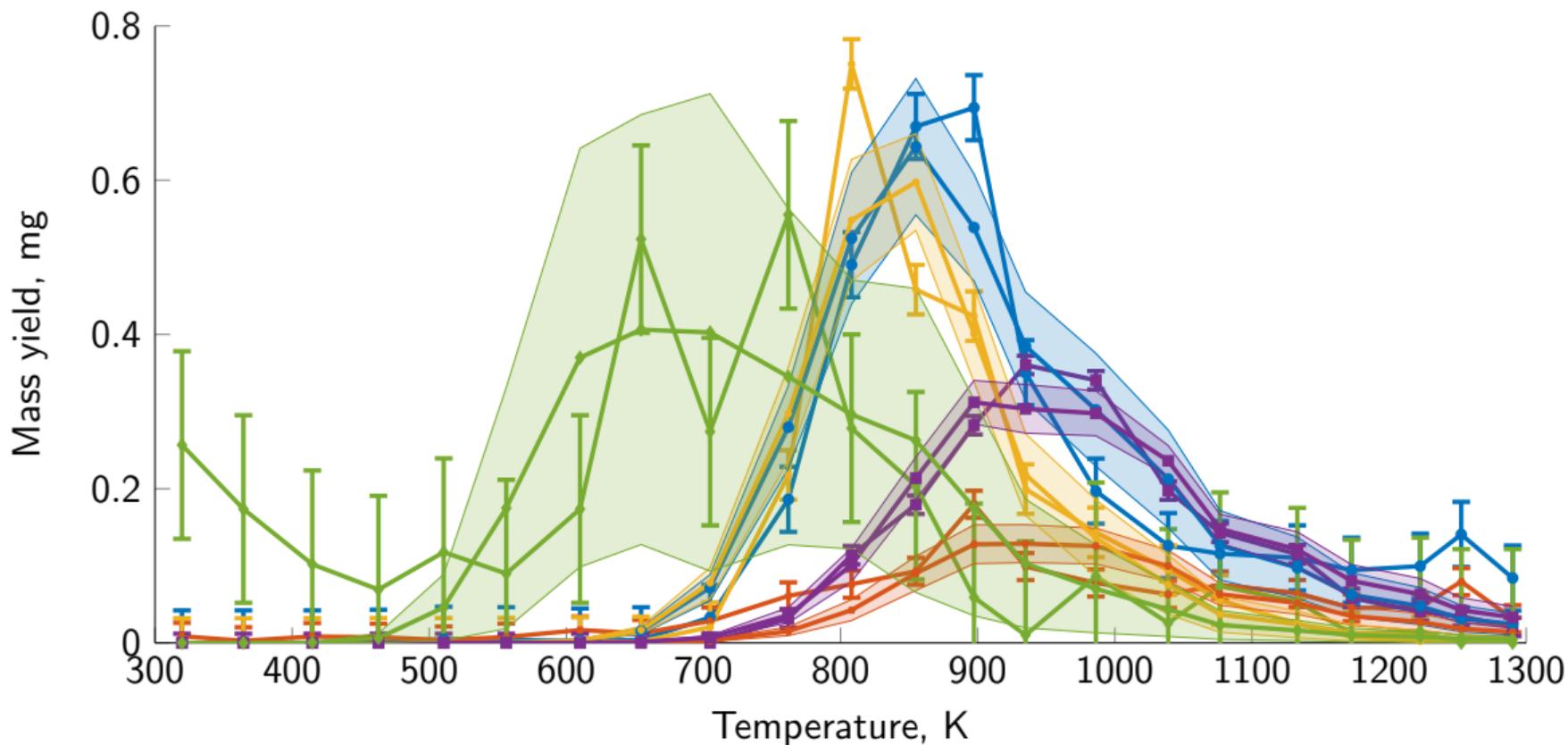


- Simple point-mass model (0D),  $k = 1, \dots, n_T$

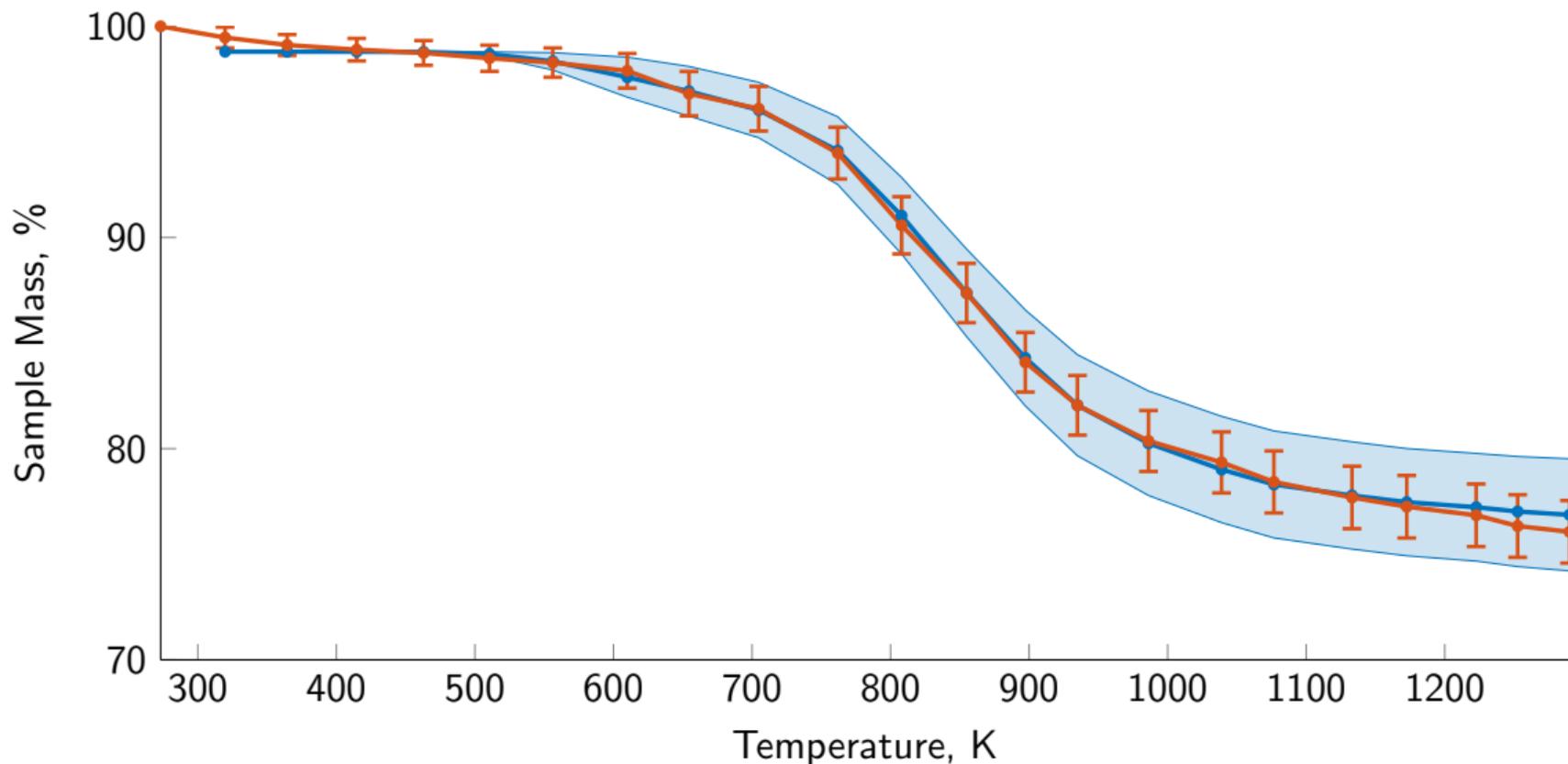
$$\mathbf{p} = \left\{ \begin{array}{l} \{A_3, E_3, n_3, F_{\text{H}_2\text{O},3}\} \\ \{A_2, E_2, n_2, F_{\text{CO},2}, F_{\text{CH}_4,2}, F_{\text{H}_2\text{O},2}\} \\ \{A_1, E_1, n_1, F_{\text{CO},1}, F_{\text{CO}_2,1}, F_{\text{H}_2,1}\} \end{array} \right\}$$



## Posterior predictive checks: 3 reactants model with 5 species



## Posterior predictive checks: 3 equations, 5 species



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## Conclusions and further work

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### Conclusions:

- ▶ Application of a simple model to simulate pyrolysis experiments
- ▶ Inference on parameter laws using Bayesian approach
- ▶ Efficient method to find kinetic parameters and characterize their uncertainties
- ▶ Simulation of mass yields with error bars using the simulated posterior samples. Good agreement with experimental curves.

### Further work:

- ▶ Improve algorithms to allow the use of more elaborated models. Adaptive MCMC has a high rejection rate. Use Hamiltonian Monte Carlo?
- ▶ Inference on reaction mechanisms together with parameter inference [Galagali and Marzouk, 2016]
- ▶ Results interpretation and their used in CFD codes

# Acknowledgments

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## Additional Information

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Bibliography

Mathematical model

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## Simple model for simulating pyrolysis experiments

For point-mass material, conduction is assumed to be instantaneous and  $T$  is uniform inside the material (lumped capacitance model,  $Bi = \frac{L_c h}{k} \ll 1$ ).

- ▶  $T = T_i = \text{cte}$  in  $\Delta t = [0, t_f]$

$$\xi(t) = 1 - \left[ (1-n) \left( C_i - A \exp\left(-\frac{E}{RT_i}\right) t \right) \right]^{1/(n-1)}, \quad \text{with } C_i = \frac{(1-\xi(0))^{1-n}}{1-n}$$

- ▶  $T_{\text{out}} = \tau t$ ,  $\tau$  heating rate

$$\xi(T) = 1 - \left\{ (1-n) \left[ -\frac{A}{\tau} T \exp\left(\frac{-E}{RT}\right) - \frac{A E}{\tau R} \text{Ei}\left(-\frac{E}{RT}\right) + C \right] \right\}^{\frac{1}{1-n}}$$
$$C = \frac{(1-\xi_0)^{1-n}}{1-n} + \frac{A}{\tau} T_0 \exp\left(\frac{-E}{RT_0}\right) + \text{Ei}\left(\frac{-E}{RT_0}\right) \frac{EA}{\tau}, \quad \text{where } \text{Ei}(x) \equiv \int_{-\infty}^{-x} \frac{\exp(v)}{v} dv$$

For more complex temperature evolution, integration should be performed numerically. When  $Bi > 1$  ( $L_c, h \nearrow$ , or  $k \searrow$ ) more complex model should be used (Argo).