

22 **Abstract**

23 The main challenge in tissue engineering consists in understanding and controlling the growth
24 process of *in vitro* cultured neotissues towards obtaining functional tissues. Computational
25 models can provide crucial information on appropriate bioreactor and scaffold design but also on
26 the bioprocess environment and culture conditions. In this study the development of a 3D model
27 using the level set method to capture the growth of a microporous neotissue domain in a dynamic
28 culture environment (perfusion bioreactor) was pursued. In our model, neotissue growth velocity
29 was influenced by scaffold geometry as well as by flow induced shear stresses. The neotissue was
30 modeled as a homogenous porous medium with a given permeability and the Brinkman equation
31 was used to calculate the flow profile in both neotissue and void space. Neotissue growth was
32 modeled until the scaffold void volume was filled thus capturing already established
33 experimental observations, in particular the differences between scaffold filling under different
34 flow regimes. This tool is envisaged as a scaffold shape and bioprocess optimization tool with
35 predictive capacities. It will allow control fluid flow during long-term culture, whereby neotissue
36 growth alters flow patterns, in order to provide shear stress profiles and magnitudes across the
37 whole scaffold volume influencing, in turn, the neotissue growth.

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43 **Introduction**

44 Skeletal Tissue Engineering (TE) strategies hold a great promise for the regeneration of bone and
45 cartilage based on the combination of bioreactors, 3D biomaterials and mesenchymal stem cells
46 (MSCs). MSCs are progenitor cells crucial for skeletal TE applications due to their ability to
47 undergo osteogenic and chondrogenic differentiation under the influence of various biochemical,
48 biophysical and, importantly, biomechanical cues. Perfusion bioreactors have been extensively
49 employed for the expansion and differentiation of MSCs providing sufficient mass transport for
50 cell growth and differentiation (Sikavitsas, Bancroft et al. 2005; Grayson, Marolt et al. 2011;
51 Sonnaert, Papantoniou et al. 2014). Furthermore, shear stress can determine early stem cell
52 lineage commitment (Song, Dean et al. 2013) but also promote terminal osteogenic
53 differentiation of bone marrow, periosteum and adipose derived MSCs and enhance extracellular
54 matrix (ECM) deposition (McCoy and O'Brien 2010; Rauh, Milan et al. 2011; Papantoniou, Chai
55 et al. 2013). There is a substantial body of literature illustrating the osteogenic effect of
56 mechanical stimulation either due to fluid flow or mechanical compression (or stretching) on the
57 differentiation of MSCs when cultured in dynamic environments seeded on 3D scaffolds in vitro
58 (Wang and Chen 2013; Delaine-Smith and Reilly 2011).

59 In scaffold-based perfusion bioreactor culture, 3D cell growth and neotissue formation has been
60 observed to begin with 2D cell proliferation on the scaffold strut surface. Subsequently, cells
61 bridge scaffold struts and start growing towards the pore void followed by ECM deposition.
62 Eventually, 3D cell growth will result in scaffold void filling, something that has been studied
63 recently using computed tomography imaging techniques (Voronov, VanGordon et al. 2013;
64 Papantoniou, Sonnaert et al. 2014). The calculation of wall shear stress values in empty scaffold
65 geometries are therefore indicative of the shear stress experienced by cells during early culture

66 time (Truscello, Schrooten et al. 2011). A current challenge to further advance and utilize
67 computational modeling strategies in the TE field is to comprise a domain composed of cells and
68 ECM (a growing permeable neotissue) on real 3D scaffold geometries. There are intriguing 2D
69 studies investigating this (Sacco, Causin et al. 2011; Hossain, Bergstrom et al. 2014) however it
70 was recently shown that without the third dimension, model parameters were overestimated
71 losing accuracy in the representation of neotissue growth (Nava, Raimondi et al. 2013).

72 Computational fluid dynamics (CFD) modeling has been extensively used in the field of TE (for
73 review see (Hutmacher and Singh 2008) (Hossain, Chen et al. 2012) (Patrachari, Podichetty et al.
74 2012)). The quantification of flow-associated shear stresses as well as their spatial distribution
75 within various 3D scaffold geometries has been thoroughly investigated in perfused bioreactor
76 setups (Raimondi, Boschetti et al. 2004; Porter, Zauel et al. 2005; Boschetti, Raimondi et al.
77 2006; Cioffi, Boschetti et al. 2006; Jungreuthmayer, Donahue et al. 2009; Maes, Ransbeeck et al.
78 2009; Voronov, VanGordon et al. 2010). In previous studies, local shear stresses were defined as
79 a function of flow rate of the culture medium, bioreactor configuration, porosity and porous
80 scaffold micro-architecture (Voronov, VanGordon et al. 2010; Pham, Voronov et al. 2012). Most
81 3D CFD studies to date only use empty scaffold geometries to calculate shear stress magnitude
82 and distribution across the empty scaffold surface i.e. wall shear stress. The aforementioned
83 studies do not take into account the transient nature of the 3D neotissue domain as a result of
84 neotissue growth which has been observed experimentally (Papantoniou, Sonnaert et al. 2013).
85 There have been few attempts to capture 3D neotissue growth on scaffolds, limited to gradually
86 increasing scaffold struts (Lesman, Blinder et al. 2010), representing the neotissue layer as an
87 impermeable domain. A recent interesting study describes the growth of a ‘biomass’ domain

88 (similar to what we term neotissue in this study) whose growth was coupled to oxygen
89 concentration and shear stress (Nava, Raimondi et al. 2013).

90 In addition to flow associated shear stress, neotissue growth kinetics have been also linked to the
91 geometric features of scaffolds. Pore size (Zeltinger et al. 2001b), pore shape (Knychala et al.
92 2013), and more specific features such as local curvature (Rumpler et al. 2008; Gamsjager et al.
93 2013) have been demonstrated to control cell fate both for in vitro (Rumpler et al. 2008; Guyot et
94 al. 2014) and in vivo applications (Bidan et al. 2012).

95 Building on our previous studies where we investigated the effect of local curvature on in vitro
96 3D growth (Guyot et al., 2014) and where we used the evolving neotissue growth in a 3D
97 scaffold to accurately determine the shear stress values in both the void space and the neotissue
98 during the bioreactor culture process (Guyot et al., 2015), this study demonstrates the added value
99 of using also the shear stress as a parameter that influences the neotissue growth. Hereto this
100 study shows the development of a level-set based computational tool able to capture the
101 difference in growth of a microporous neotissue domain in a dynamic culture environment
102 (perfusion bioreactor) under different flow rates. Capturing this difference is only possible by the
103 explicit incorporation of flow-induced shear stresses as a parameter in the calculation of the
104 neotissue growth velocity (alongside the scaffold geometry which was already present (Guyot et
105 al., 2014)). This tool provides the ability to steer fluid flow during long term culture in order to
106 provide given shear stress profiles and magnitudes across the whole scaffold volume.

107
108 **Methods**

109 In this section, the model describing the neotissue growth, the calculation of the flow-induced
110 shear stresses and the influence of the latter on the former will be explained.

111 *Neotissue growth via the level set method*

112 The growth induced changes in the neotissue topology during the culture process can be seen as a
113 moving interface between two different domains (Sethian 1999): in this study, one domain
114 represents the neotissue volume Ω_{nt} , and the other one is the void Ω_v , separated by an interface
115 Γ , with a normal \mathbf{n}_Γ (see Fig 1B). The level set method (LMS) is a technique that has been
116 developed to deal with this kind of moving interfaces and it is used in the context of this study to
117 simulate the evolution of the frontline Γ between neotissue and void space in a mesh-free
118 manner. The principle of the LSM consists in defining a signed distance function φ in $\Omega =$
119 $\Omega_{nt} \cup \Omega_v$ with the following properties:

$$\begin{cases} \varphi > 0 \text{ in } \Omega_{nt} \\ \varphi < 0 \text{ in } \Omega_v \\ \varphi = 0 \text{ on } \Gamma \end{cases} \quad (1)$$

120 In order to capture the evolution of the moving interface Γ , the following transient advection
121 equation is solved with a given growth velocity \mathbf{v}_G and homogenous Neumann boundary
122 condition ($\partial\varphi/\partial\mathbf{n} = 0$, with \mathbf{n} being the normal to the computational domain Ω).

$$\frac{\partial\varphi}{\partial t} + \mathbf{v}_G \cdot \nabla\varphi = 0 \text{ in } \Omega. \quad (2)$$

123 With $\mathbf{v}_G = V_G \cdot \mathbf{n}_\Gamma$, and $\mathbf{n}_\Gamma = \frac{\nabla\varphi}{|\nabla\varphi|}$. The expression of the growth velocity magnitude V_G will be
124 described later. The initial configuration of the distance function φ corresponds to a homogenous
125 single cell layer over the scaffold struts with a thickness equal to 20 μm (Darling and Guilak

126 2008). The time discretization of equation (2) was done using the backward Euler method and the
 127 advection term was treated with the method of characteristics. For more details about the
 128 implementation, please refer to (Guyot, Papantoniou et al. 2014).

129 *Shear stress distribution*

130 During dynamic culture, the neotissue grows, eventually filling up the whole scaffold void. These
 131 changes affect the flow patterns developed, depending on the presence (or not) of neotissue and
 132 so the flow profile has to be treated differently in Ω_{nt} and Ω_v (as described in detail in Guyot et
 133 al., 2015). In Ω_v , due to a low Reynolds number ($Re < 1$), the flow profile was approximated via
 134 the Stokes equation (3). In Ω_{nt} , the complex structure of the neotissue can be seen as a
 135 homogenous porous medium with a given permeability K_0 , leading to the flow profile being
 136 calculated with the Darcy equation (4). According to the definition of the LSM, the interface Γ is
 137 never conforming to the computational mesh, resulting in difficulties to couple equation (3) and
 138 (4) together with suitable boundary conditions at the interface. To overcome this, the Brinkman
 139 equation (5) was used to calculate the flow profile in the whole domain Ω (see Fig 1C). A no-slip
 140 boundary condition was applied on the scaffold surface as well as on the chamber walls. The
 141 scaffold was placed at a sufficient distance from the inlet avoiding refluxes occurring at the
 142 entrance of the scaffold (Papantoniou, Guyot et al. 2014), so the boundary condition for the
 143 inflow was set to a Poiseuille profile with a velocity u corresponding to the given flow rate Q .

$$\begin{cases} -\mu\Delta\mathbf{u} + \nabla p = 0 & \text{in } \Omega_v \\ \nabla \cdot \mathbf{u} = 0 & \text{in } \Omega_v \end{cases} \quad (3)$$

$$\begin{cases} \frac{\mu}{K_0}\mathbf{u} + \nabla p = 0 & \text{in } \Omega_{nt} \\ \nabla \cdot \mathbf{u} = 0 & \text{in } \Omega_{nt} \end{cases} \quad (4)$$

$$\begin{cases} -\mu\Delta\mathbf{u} + \nabla p + \frac{\mu}{K}\mathbf{u} = 0 & \text{in } \Omega \\ \nabla \cdot \mathbf{u} = 0 & \text{in } \Omega \end{cases} \quad (5)$$

144 This technique can be seen as a penalization method, indeed, when K is equal to the neotissue
 145 permeability ($K \ll 1$) in Ω_{nt} , it makes the Stokes term in (5) negligible. At the contrary, when
 146 K is set to a very high value ($K \gg 1$) in Ω_v , it leads to the Darcy term to be close to zero. To
 147 avoid numerical problems, this switch between the two values was implemented using a space
 148 and time dependent smeared out Heaviside function H (6), and K was updated according to this
 149 function (7) ($\varepsilon = 1.5h$, with h the mesh size).

150

$$H(\varphi) = \begin{cases} 0, & \varphi < -\varepsilon \\ \frac{1}{2} + \frac{\varphi}{2\varepsilon} + \frac{1}{2\pi} \sin\left(\frac{\pi\varphi}{\varepsilon}\right), & -\varepsilon < \varphi < \varepsilon \\ 1, & \varphi > \varepsilon \end{cases} \quad (6)$$

$$K(\mathbf{x}) = 10^{30}(1 - H) + K_0H \quad (7)$$

151 The parameter K_0 was estimated using the random fibers theory approximation. In (Nabovati A.
 152 2009), the authors provide an approximation of the permeability of a porous media made of a
 153 random fiber web:

$$K_0 = 0.491\delta^2 \sqrt{\frac{1-\psi_c}{1-\psi}}^{2.31} \quad (8)$$

154

155 In this equation, ψ represents the porosity of the media (neotissue, assumed to be equal to 90% in
 156 this study, which is in the range of porosities of soft tissues and hydrogels) and ψ_c corresponds to
 157 a percolation threshold or a threshold porosity where flow is permitted (set to zero in this study).

158 Finally, δ represents the micro-pore size of the neotissue and was set to 50 μm , assuming the
 159 pore size equal to half to one third of the size of a fully spread cell of the type used in this study
 160 (Eyckmans, Lin et al. 2012).

161 In this study, we distinguished two different wall shear stresses acting on cells depending on their
 162 location. The first one (SS_{surf}) is the shear stress acting on the interface Γ due to the different
 163 flow profiles from either side and is calculated with the usual definition:

$$SS_{surf} = \sqrt{\tau_{12}^2 + \tau_{23}^2 + \tau_{13}^2}, \quad \text{with} \quad \tau_{ij} = \mu \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right). \quad (9)$$

164 The second one (SS_{in}) is the shear stress acting within the neotissue, it is associated with the
 165 interstitial flow through the micro-porous neotissue and is approximated following the method
 166 presented in (Whittaker, Booth et al. 2009). This method is based on the fact that the micro-
 167 porous neotissue is assumed to be composed of cylindrical ducts of diameter δ . Since the Darcy
 168 equation just gives the average (Darcy) velocity u , the interstitial velocity magnitude u_{in} is
 169 calculated from results of equation (5) and is turned into a Poiseuille velocity profile (u_p) in a
 170 cylindrical channel in order to have an analytical expression of local wall shear stress (on the
 171 wall of the cylindrical ducts, representing here the micro-pores), allowing for the estimation of
 172 SS_{in} .

$$u_{in} = |\mathbf{u}|/\psi \quad (10)$$

$$u_p \approx 2u_{in} \left(1 - \left(\frac{2r}{\delta} \right)^2 \right) \quad (11)$$

$$SS_{in} = \mu \left| \frac{\partial u_p}{\partial r} \right|_{r=\delta/2} \approx \frac{8\mu u_{in}}{\delta} \approx \frac{8\mu |\mathbf{u}|}{\delta\psi} \quad (12)$$

173

174 *Neotissue growth velocity*

175 A key parameter of the study is the local growth velocity of the neotissue. In (Guyot,
176 Papantoniou et al. 2014), this space dependent velocity was only depending on the local mean
177 curvature of the interface as it has been shown in (Bidan, Kommareddy et al. 2012) or (Rumpler,
178 Woesz et al. 2008). Another important growth influencing factor is the local shear stress on cells.
179 In (Nava, Raimondi et al. 2013), authors present a biomass growth model for cartilage describing
180 an interface moving in time in function of the fluid-induced shear stress and in (Chapman,
181 Shipley et al. 2014) the authors introduce a growth model for cell aggregates in hollow fiber
182 bioreactors where cell population growth increased or decreased depending on wall shear stresses
183 experienced by cells. In this study, a similar approach is used and the local neotissue growth
184 velocity V_G is described as a function of both the mean curvature and the flow-induced shear
185 stress:

$$V_G = A \cdot f(SS_{surf}) \cdot g(\kappa). \quad (13)$$

186 The neotissue growth velocity parameter A was estimated from the experimental data obtained
187 for the low flow rate results obtained in (Papantoniou, Sonnaert et al. 2014) using a trial and error
188 approach and was set equal to $4 \cdot 10^{-14}$ m²/s. As described extensively in (Guyot et al, 2014),
189 the basis of the influence of curvature on the neotissue growth comes from the observation that
190 neotissue grows faster where the curvature is higher and that it does not grow if the curvature is
191 negative or equal to zero (Bidan, Kommareddy, et al. 2013). The mean curvature influence
192 function $g(\kappa)$ can therefore be expressed mathematically as follows:

$$g(k) = \begin{cases} -\kappa & \text{if } \kappa > 0 \\ 0 & \text{if } \kappa \leq 0 \end{cases} \quad (14)$$

193 where κ is the local mean curvature ($\kappa = \nabla \cdot \mathbf{n}_\Gamma$) and the second row of equation (14) depicts that
 194 there is no growth when the curvature is null or negative. The negative sign in equation (14)
 195 comes from the fact that according to our definition of φ , the normal \mathbf{n}_Γ points toward neotissue,
 196 so growth has to be towards the opposite of $\nabla\varphi$. The surface shear stress influence function
 197 $f(SS_{surf})$ (unit less) was inspired by (Nava, Raimondi et al. 2013) (Fig 1D) and defined as a
 198 continuous function

$$f(SS_{surf}) = \begin{cases} 0.5 + \frac{0.5 \cdot SS_{surf}}{a_1} & 0 \leq SS_{surf} < a_1 \\ 1 & a_1 \leq SS_{surf} < a_2 \\ \frac{SS_{surf} - a_3}{a_2 - a_3} & a_2 \leq SS_{surf} < a_3 \\ 0 & a_3 \leq SS_{surf} \end{cases} \quad (15)$$

199 This function was established in order to have an optimal shear stress influence that enhances the
 200 growth ($a_1 < SS_{surf} < a_2$), and a critical threshold ($SS_{surf} > a_3$) above which shear stress
 201 inhibits the cells and growth cannot occur anymore, in this study a_1 , a_2 , a_3 were respectively set
 202 to 0.05, 0.15 and 0.2 Pa (Chapman, Shipley et al. 2014). Table 1 summarizes all the parameters
 203 used in the calculation of the neotissue growth velocity v_G .

204

205 Table 1: model parameters used for the calculation of the neotissue growth rate v_G .

Parameter	Value	References
Neotissue porosity, ψ	90 %	(Guyot, Papantoniou et al,

		2015)
Neotissue micro-pore size, δ	50 μm	(Guyot, Papantoniou et al, 2015)
Neotissue growth rate, A	$4 \cdot 10^{-14} \text{ m}^2/\text{s}$	Determined from (Papantoniou, Sonnaert et al. 2014)
Minimal shear stress value enhancing neotissue growth, a_1	0.01	(Chapman, Shipley et al. 2014).
Maximal shear stress value enhancing neotissue growth, a_2	0.03	(Chapman, Shipley et al. 2014).
Critical shear stress value, a_3	0.05	(Chapman, Shipley et al. 2014).

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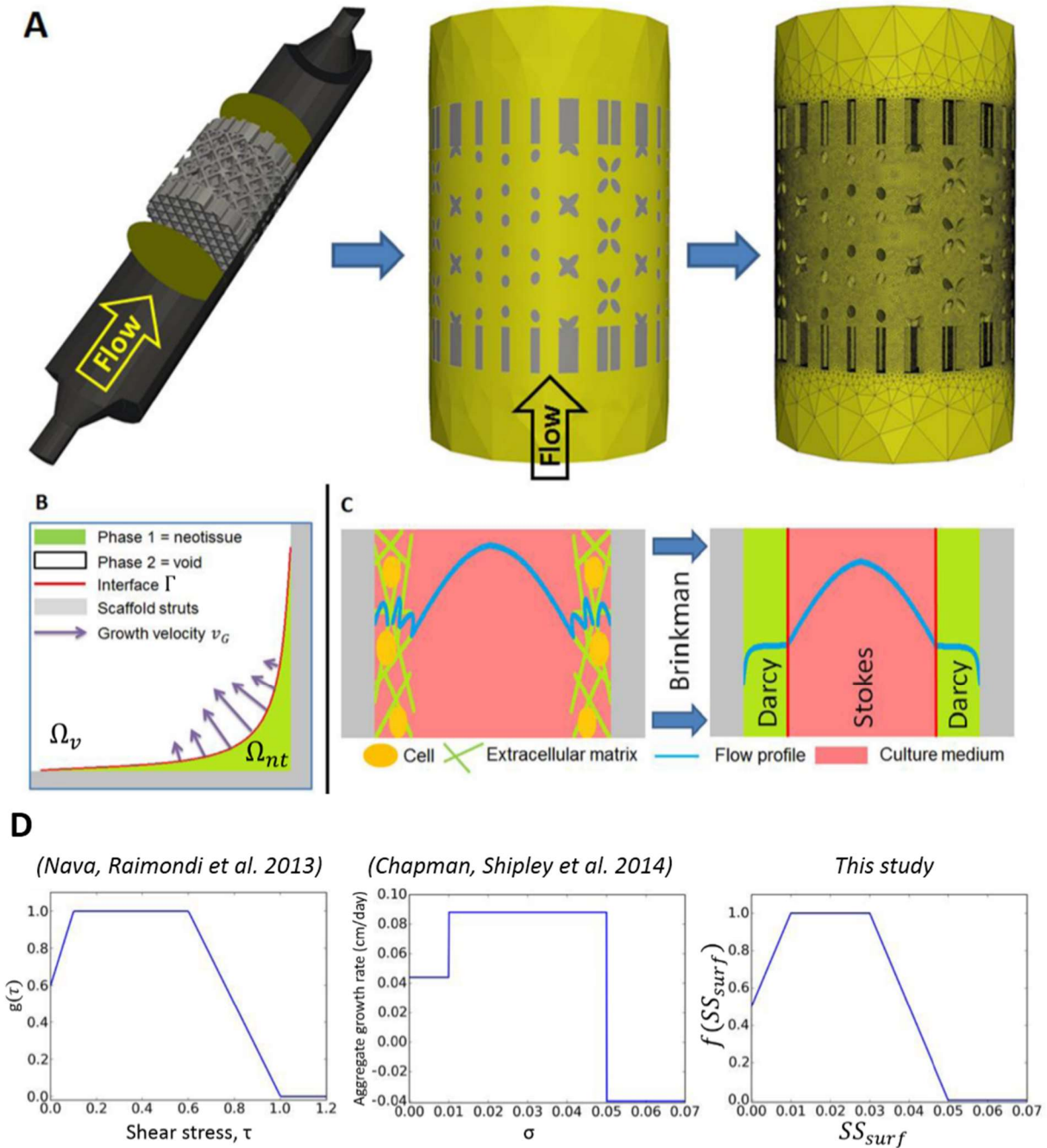
207 *Implementation*

208 The full model was implemented with the free partial differential equation solver *FreeFem++*
209 (Hecht 2012). The computational domain (see Fig 1A) chosen consisted of the part of the
210 perfusion bioreactor chamber comprising the entire scaffold including 2mm of the chamber at
211 each scaffold side. This was done in order to avoid numerical reflexes and was meshed with
212 approximately 3 million tetrahedrons. The numerical problem was solved in parallel using a 12
213 core facility. In order to avoid unnecessary computational effort, the Brinkman equation was

214 solved on only one quarter of the full mesh cut along the flow axis and with respect to
215 symmetrical boundary condition.

216 *Summary of Experimental set-up*

217 In this study, results presented in (Papantoniou, Sonnaert et al. 2014) were used in a first
218 calibration and validation set-up. The set-up of the experiment is briefly repeated below. Three-
219 dimensional additive manufactured open porous Ti6Al4V scaffolds ($\emptyset = 6$ mm, $h = 6$ mm,
220 porosity = $73\% \pm 1\%$, strut diameter = 245 ± 2 μ m and pore size = 755 ± 3 μ m), produced on an
221 in-house developed selective laser melting (SLM) machine (Van Bael, Kerckhofs et al. 2011)
222 were used. Human Periosteal Derived Stem Cells (hPDCs, one of the known sources of MSCs)
223 were isolated from periosteal biopsies of different donors as described previously (Eyckmans and
224 Luyten 2006) and expanded in culture flasks. When a sufficient amount of these cells was
225 produced in this 2D culture (number of population doublings = 15), cells were harvested by
226 trypsinization with Tryple Express (Invitrogen) and seeded on the scaffolds, which marked the
227 start of the bioreactor experiment. The TE constructs were cultured in an in-house developed
228 bioreactor for 14, 21 and 28 days under dynamic culture conditions ($n = 9$ per flow rate –
229 triplicates per time point). Two different perfusion flow rates were used: 4 mL/min (Q_1 , high) and
230 0.04 mL/min (Q_2 , low). After culture, three constructs for each time point were prepared for
231 contrast-enhanced nanofocus Computed Tomography (NanoCT) imaging (Kerckhofs 2013)
232 allowing for visualization and quantification of the neotissue volume formed inside the
233 constructs.



234

235 *Fig 1: Model setup. (A) Representation of the scaffold and the bioreactor chamber, the region of*
 236 *interest is delimited by two yellow circles. The scaffold is then removed from a cylinder*
 237 *representing the delimited area. Finally, a mesh is created in the computational domain. (B)*
 238 *Schematic representation of curvature driven neotissue growth using the level set method. (C)*

239 *Schematic representation of the Brinkman approximation used in this study. (D) The influence of*
240 *shear stress on the neotissue growth velocity in this study (right) is a combination of the*
241 *continuous shape proposed by (Nava, Raimondi et al. 2013) and the values by (Chapman, Shipley*
242 *et al. 2014) (middle).*

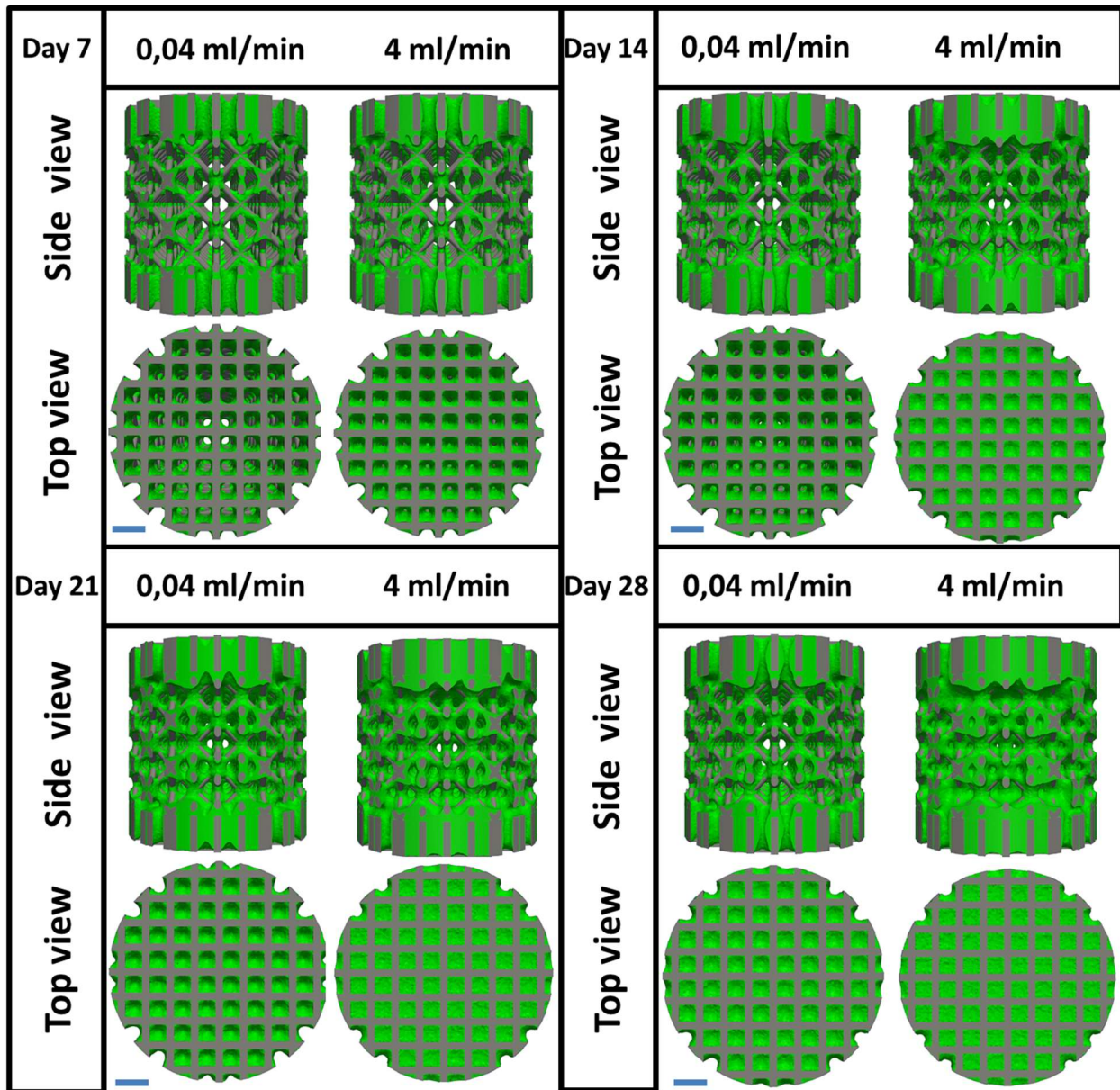
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244 **Results**

245 The model was run for different flow rates according to the experimental set up described in
246 (Papantoniou, Sonnaert et al. 2014). Simulations show a significant difference between the two
247 conditions regarding the total amount of neotissue that has been produced at different time points
248 (Fig 2 and Fig 3). Indeed, under the high flow rate , the local shear stress acting on the neotissue
249 interface is higher than for the lowest flow rate, resulting in an acceleration of neotissue growth.
250 Although differences can be observed between the simulations and experiments, the simulations
251 are capable of capturing the experimentally observed differences between the two flow rates in
252 terms of volume filling (Fig 3B). The model was also able to compute quantitative data (shown in
253 Fig 4) regarding important culture variables during neotissue growth such as pressure drop across
254 the scaffold along the flow axis, average surface and inner shear stresses, giving an overview of
255 different mechanical stimuli acting on cells over culture time.

256 The pressure drop (Fig 4A) across the scaffold varied from almost 4 Pa in the first days of culture
257 to 11 Pa at day 28 for the high flow rate Q_1 , while it ranged from 0.03 Pa to 0.08 Pa for the low
258 flow rate Q_2 in an equivalent period of time. The surface shear stresses (Fig 4B) and inside shear
259 stresses (Fig 4C) differed by a factor of 100 between the two different flow rates. For both flow

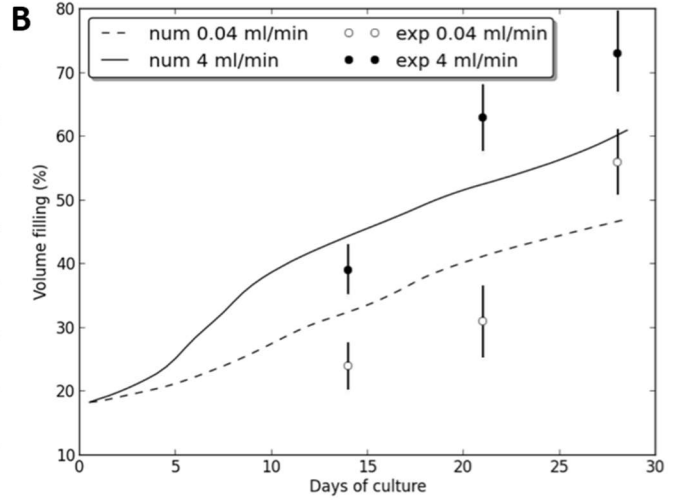
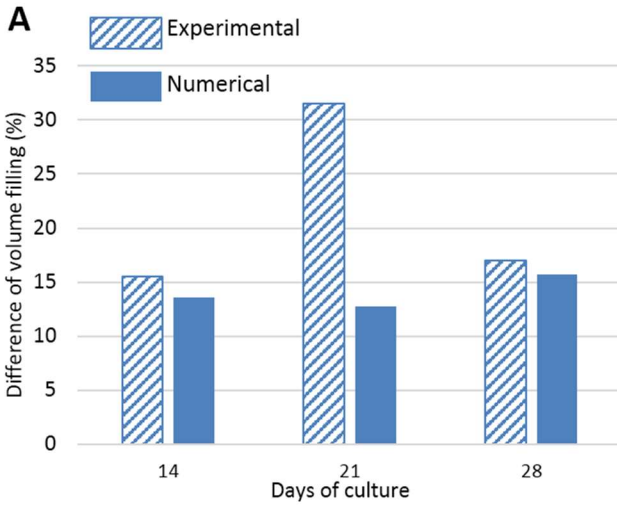
260 rates, the inner shear stress, the stress acting on cells embedded into the neotissue, can be around
261 3 (for early time points) to 20 (for late time points) times bigger than the interface shear stress
262 acting on cells at interface in contact with the free flow. Fig 5 shows how the flow regime
263 changes from when the scaffold is almost empty (Fig 5A) to where it is partially filled (Fig 5B).
264 Fig 6 depicts the local growth velocity v_G showing the influence of the distributed shear stress as
265 well as the local mean curvature of the surface. Maximal neotissue growth velocities of 3 and 7
266 $\mu\text{m}/\text{day}$ were obtained for the low and high flow rates respectively.



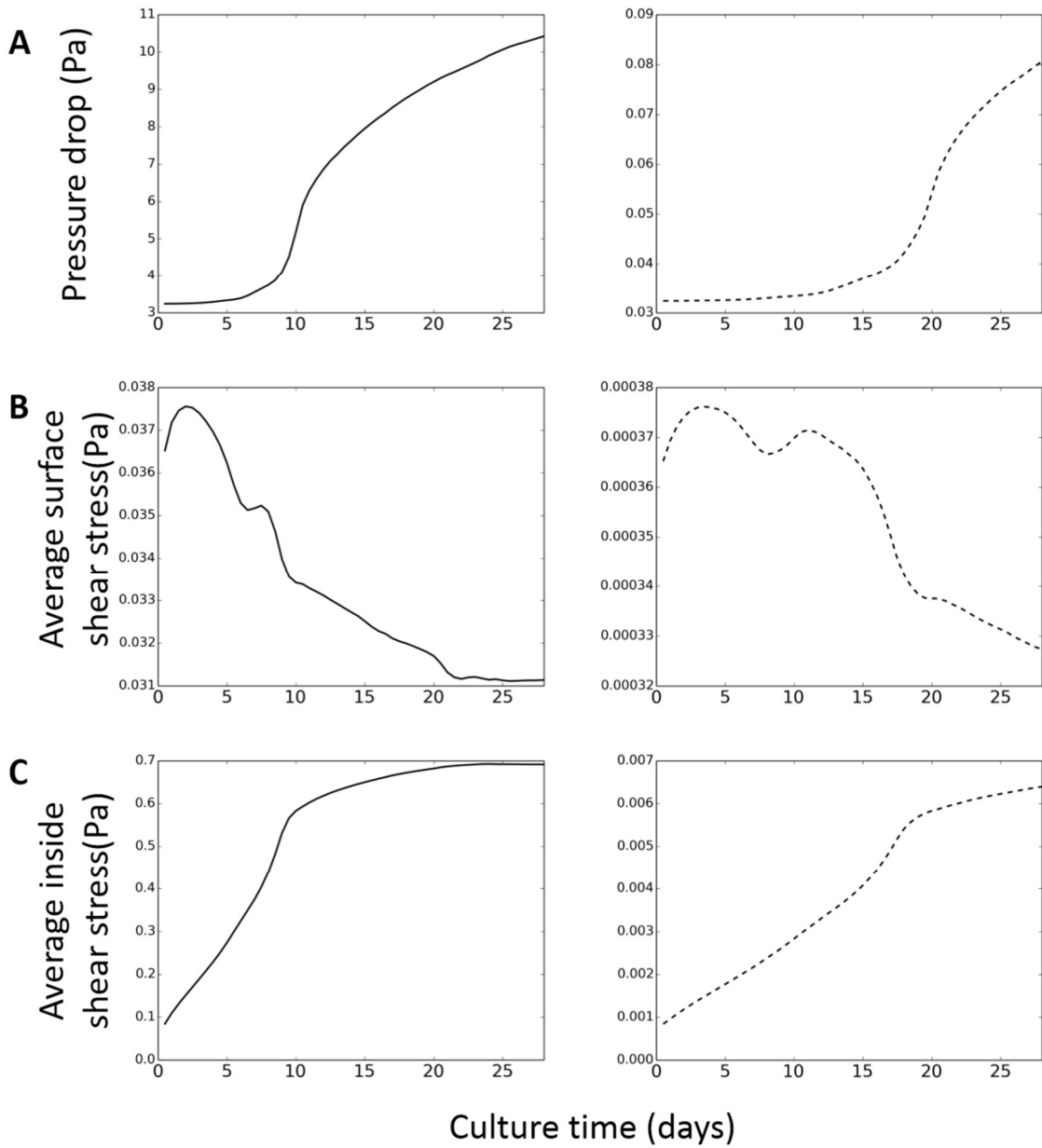
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268 *Figure 2: Neotissue growth (green) on scaffold (grey) at different time points for two different*
 269 *flow rates and two different views (direction of flow rate is from bottom to top in side view).*
 270 *Gradual scaffold pore closure can be observed. Scale bar represents 1 mm.*

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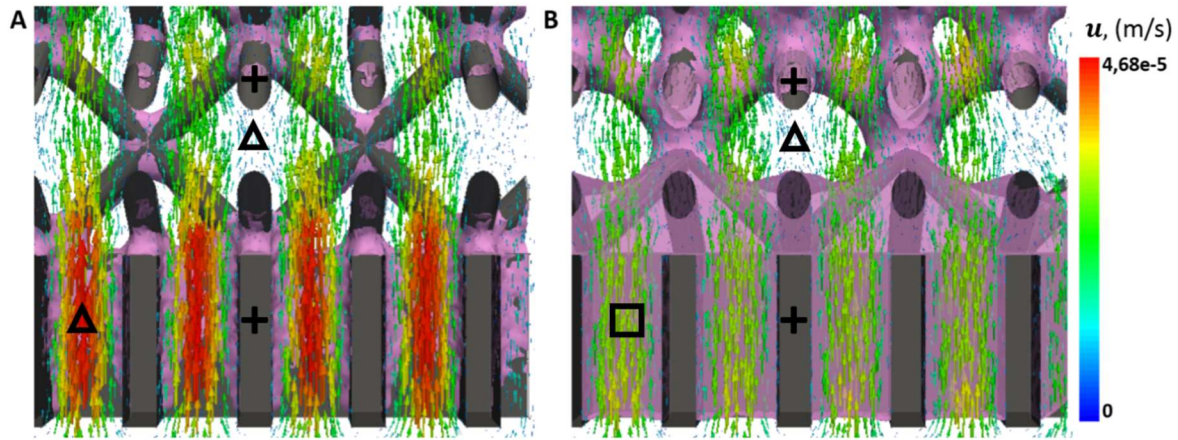


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 273 *Fig 3: Comparison between simulations and experimental results (Papantoniou, Sonnaert et al.*
 274 *2014). (B) Evolution of volume filling over culture time for the two different flow rates used in*
 275 *this study. Experimental results are presented through the mean and standard deviation (n=3).*
 276 *(A) Quantitative differences in volume filling between the different flow rates, indicating that the*
 277 *simulated differences are similar to those experimentally observed.*



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279 Fig 4: Mechanical characterization of the predicted tissue growth for flow rates Q_2
 280 (0.04ml/min , left) and Q_1 (4ml/min , right). (A) Pressure drop between the entrance and the exit of
 281 the scaffold. (B) Average neotissue surface shear stress. (C) Average inside neotissue shear
 282 stress. Notice the difference in scale on the vertical axes between left and right figures.



283

284 *Fig 5: Close-up view on a section of the scaffold illustrating the different flow profiles (u) at*
 285 *early (A) and late (B) time points (neotissue volume is indicated in pink). In image (A), the*
 286 *scaffold is almost empty, so the flow is mostly ruled by the Stokes equation leading to Poiseuille*
 287 *flow in the bottom part. In (B), the bottom part of the scaffold is completely filled with neotissue,*
 288 *so the flow profile in this area is ruled by the Darcy equation. +=scaffold, □=neotissue, Δ=void.*

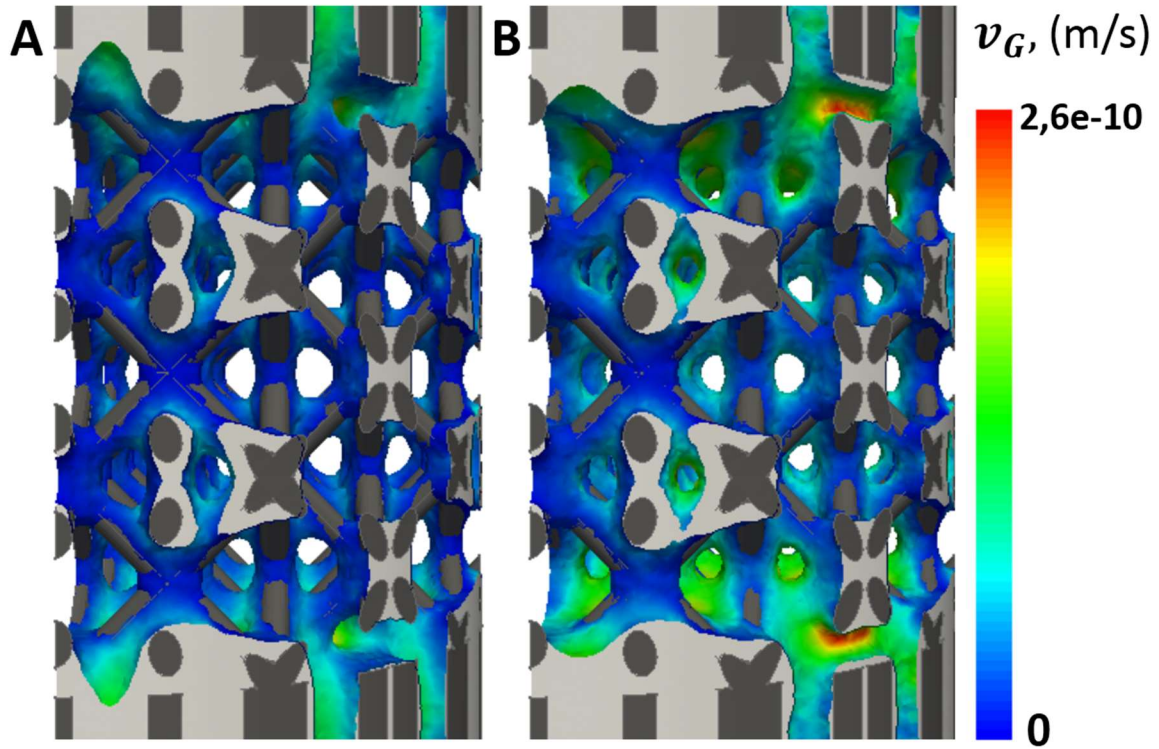
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295 *Fig 6: Local distribution of the neotissue growth velocity v_g (m/s) at the neotissue-void interface*
 296 *at 28 days for flow rates Q_2 (0.04ml/min, left) and Q_1 (4ml/min, right). the neotissue is indicated*
 297 *in light grey).*

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304 **Discussion**

305 Understanding spatiotemporal cell growth in dynamic cultures constitutes a challenge in TE, in
306 particular the quantification, distribution and interplay between scaffold geometry, shear stress
307 and neotissue growth. The present model provides a foundation that will allow the quantitative
308 investigation of the effect of shear stresses on cell growth and ECM production, which today is
309 still not well understood. The most important contribution of this study is the extension of a
310 previously developed model of local curvature-dependent neotissue growth to incorporate the
311 influence of the local shear stresses on the neotissue growth, on complete scaffolds in 3D. The
312 importance of 3D models in comparison to 2D has been shown by (Nava, Raimondi et al. 2013),
313 the latter suffering from a lack of accurate representation which leads to loss of information and
314 an incorrect parameter estimation.

315 Recent experimental observations have revealed that hPDCs (the MSC source used in this study)
316 may grow in the third dimension by bridging scaffold struts, leading to complete filling of
317 scaffold pores when cultured in perfusion bioreactors in glucose based growth medium (Sonnaert,
318 Papantoniou et al. 2014) and in growth factor containing osteogenic medium (Papantoniou Ir,
319 Chai et al. 2013). NanoCT imaging has been employed to confirm that complete pore closure
320 could be obtained depending on bioreactor operating conditions (Papantoniou, Sonnaert et al.
321 2014). Similar 3D cell growth behavior upon seeding on scaffolds has been also observed for
322 other cell types such as human bone marrow MSCs (Zhao, Pathi et al. 2005; Li, Tang et al. 2009),
323 primary fibroblasts (Joly, Duda et al. 2013) and cell lines such as MC3T3-E1 pre-osteoblast cells
324 (Kommareddy, Lange et al. 2010). This opens the possibility to create 3D in vitro tissue-like
325 structures where cells may grow independently of the initial scaffold surface allowing for the
326 study of the interaction between cells and their own ECM and the culture environment.

327 Additionally, this means that when simulating the experimentally observed void filling due to 3D
328 neotissue growth, models should take into account the specificity of the experimentally used cell
329 type (Chapman, Shipley et al. 2014). In this work we attempt to investigate this for adult stem
330 cell types in particular for hPDCs.

331 Neotissue growth kinetics in 3D dynamic and scaffold-based culture set ups, has been linked with
332 oxygen level (Zhao, Pathi et al. 2005), shear stress (Nava, Raimondi et al. 2013) and scaffold
333 geometry - in particular curvature (Bidan, Kommareddy et al. 2013; Knychala, Bouropoulos et al.
334 2013). In our study, oxygen was not investigated since for the flow rate and scaffold dimensions
335 used in this work, taking into account the the oxygen consumption rate of hPDCs (Lambrechts,
336 Papantoniou et al. 2014), there is no significant oxygen drop detected from inlet to outlet.
337 Therefore the assumption that the whole cell population is exposed to the same level of dissolved
338 oxygen tension (atmospheric) seems to be justified. When other flow regimes and/or other cell
339 types will be studied, this assumption has to be re-evaluated. An additional assumption was the
340 existence of an initial single cell layer at the start of the simulations. This layer was assumed to
341 be 20 μm based on 2D suspended MSC diameter values (Darling and Guilak 2008), leading to an
342 initial filling percentage of 18 %. Provided that a high number of cells is used for cell seeding
343 (Chen, Bloemen et al. 2011), a homogeneous distribution can indeed be achieved across the
344 scaffold. This step has been studied in detail for regular geometry scaffolds (Melchels, Tonnarelli
345 et al. 2011). The initial filling percentage is likely an overestimation of the real filling and is (in
346 its current way of calculating) dependent on the available surface of the scaffold under study.
347 With the currently used measurement technique (NanoCT) the experimental value for the initial
348 filling cannot be obtained due to the absence of matrix produced by the cells immediately after
349 seeding. This parameter therefore merits further study in future work. Neotissue shrinkage/loss

350 due to cell death was not incorporated in the presented model, however, there is a possibility that
351 this shrinkage/loss occurs for those cases where excessively harsh culture conditions are developed
352 i.e. very low dissolved oxygen tension (anoxia) and glucose concentration or high lactate
353 concentration (Flaibani, Magrofuoco et al. 2010). Further efforts are required in order to
354 incorporate the aforementioned physico-chemical components in our model. However this is out
355 of the scope of the present study, where geometry and fluid dynamics are the main regulators of
356 neotissue growth.

357 In this study neotissue growth velocity was coupled to the local scaffold topography (curvature)
358 (Guyot, Papantoniou et al. 2014) and local fluid flow induced shear stress. Shear stress was
359 modeled to enhance increasingly neotissue growth from 0 to a_1 reaching a maximum value
360 between a_1 and a_2 (equal to 0.01 and 0.03 Pa respectively). The lower amount of neotissue
361 observed at the periphery of the scaffolds is due to the low values of fluid flow induced shear
362 stresses at that location. This in turn is due to the flow demonstrating a Poiseuille profile when
363 entering the scaffold, thereby positively influencing to a greater extent the neotissue formation at
364 the center of the scaffold. When the shear stress reaches threshold a_3 equal to 0.05 Pa, its effect
365 on neotissue growth was modeled to change into a prohibitive one with decreasing neotissue
366 growth for increasing shear stresses beyond a_3 . Experimental observations of 0.015 Pa as an
367 optimum shear stress level for human bone marrow MSCs cultured in perfusion bioreactors
368 confirmed this hypothesis (Li, Tang et al. 2009). A similar range of shear stresses has been
369 reported for human bone marrow MSCs by (Grayson, Bhumiratana et al. 2010) although these
370 values referred to empty scaffold geometries. We have recently experimentally observed the
371 adverse effects of excessive shear stress on neotissue growth by hPDCs in perfusion bioreactors
372 resulting in inhomogeneous tissue engineered constructs (Papantoniou, Guyot et al. 2014). This
373 was observed for shear stress values in excess of 0.05 Pa which is the threshold value used in this

374 study. Interestingly (McCoy, Jungreuthmayer et al. 2012) estimated a critical threshold for human
375 bone MSC detachment from irregular scaffolds at 0.088 Pa closely matching the one used here,
376 taking into account the differences in scaffold architecture and cell type.

377 The growth velocity term A was estimated in a trial and error fashion based on in-house
378 experimental observations (Papantoniou, Sonnaert et al. 2014) measuring neotissue growth on the
379 simulated scaffold for the low flow rate Q_2 . The prediction obtained for the high flow rate Q_1
380 shows an increase in volume filling, similar to the experimental data (Fig 3B). The discrepancy
381 observed between experimental and computationally-derived neotissue growth kinetics seen in
382 Fig 3 could be attributed to numerical and physical factors. For the former (numerical factor), we
383 did not go through an objective and rigorous optimization process to find the value for A because
384 of the insufficient quantity of experimental data available. In a follow-up study, dedicated
385 experimental data will be generated to parametrize the model. As to the latter (physical factor),
386 neotissue growth kinetics do not solely depend upon scaffold geometry and shear stress
387 magnitude. The introduction in the model of additional physico-chemical parameters that are well
388 known for their influence on neotissue growth such as dissolved oxygen tension (Grayson, Zhao
389 et al. 2007; Dos Santos, Andrade et al. 2010), glucose concentration (Saki, Jalalifar et al. 2013)
390 and lactate concentrations (Schop, Janssen et al. 2009) in the medium, could help to improve the
391 agreement between experimentally and computationally determined values. Even though for the
392 experiments used in this study (Papantoniou, Sonnaert et al. 2014), no global changes in oxygen
393 tension were observed between the inlet and outlet for the flow rate and cell source used, the
394 model allows to quantify these variable throughout the entire scaffold, potentially showing
395 pockets of decreased oxygen tension (and neotissue growth) coming from improper perfusion due
396 to local geometrical particularities. Additionally, the possibility that higher flow rates might lead

397 to the secretion of more ECM could also explain to a certain extent the differences in volume
398 filling patters between simulations and experiments.

399 In a very interesting recent study, neotissue was modeled in 3D by (Nava, Raimondi et al. 2013)
400 using the arbitrary Lagrangian–Eulerian (ALE) method to implement mesh movement. However,
401 the flow was only modeled in the void space, which is acceptable for the early growth phase, but
402 is not suitable to simulate complete filling of the scaffold. In this study, we employed the level set
403 method to represent neotissue growth as described previously in (Guyot, Papantoniou et al. 2014
404 and Guyot et al. 2015). This method, separating the movement of the interface from the definition
405 of the mesh allows tracking the neotissue kinetics until full scaffold pore filling is reached. A
406 maximal neotissue growth velocity of 3 and 7 $\mu\text{m}/\text{day}$ was obtained for the low and high flow
407 rates respectively as seen in Fig 6. The average growth velocities from the present study are the
408 same order of magnitude than the ones estimated in (Nava, Raimondi et al. 2013) which were in
409 the order of 1 $\mu\text{m}/\text{day}$ and higher than the values adopted in previous tissue growth models
410 (Sacco, Causin et al. 2011). This could be due to the fact that chondrocytes modeled in that work
411 are slow proliferating cells compared to hPDCs.

412 Numerical predictions obtained in this study (Fig 4 and Fig 5) are in general agreement with
413 published literature with regard to the range of shear stresses calculated at the neotissue
414 interphase during neotissue growth (Boschetti, Raimondi et al. 2006; Cioffi, Boschetti et al. 2006;
415 Lesman, Blinder et al. 2010; Nava, Raimondi et al. 2013). However shear stresses developed
416 within the neotissue (SS_{in}) were seen to be an order of magnitude higher than the surface shear
417 stresses. It is interesting to note that the average values obtained for inner shear stress for the high
418 flow rate in this work, reaching 0.7 Pa, compared closely to the ones determined for native bone
419 (1–3 Pa) (Zeng, Cowin et al. 1994), closer than what has been reported to date for tissue
420 engineered constructs. In this study it was assumed that the pores within the neotissue had a size

421 of 50 μm , this was justified, for our case study, for large hPDC cells whose size has been shown
422 to range between 100-150 μm when fully spread (Eyckmans, Lin et al. 2012). This value could
423 become smaller upon 3D confluency when cells ECM secretion might lead to fuller neotissue
424 structures. Interestingly, the calculated inner neotissue (microporous) shear stress magnitude for
425 the selected neotissue pore size matched closely to those determined via single cell simulations
426 (Jungreuthmayer, Donahue et al. 2009; Verbruggen, Vaughan et al. 2014; Zhao, Vaughan et al.
427 2014). Empty scaffold simulations may be useful for early time points and have been successfully
428 used to determine stem cell lineage commitment (Song, Dean et al. 2013) however for later
429 culture time points these models seem to be insufficient.

430 Using the whole scaffold geometry it was also possible to determine changing pressure drop
431 values during neotissue growth and subsequent scaffold pore closure. The maximum pressure
432 drop values obtained for the high flow rate, for completely full scaffolds, equaled to 11 Pa for the
433 parameters of this study (cylindrical scaffold dimensions $d = 6\text{ mm}$, $h = 6\text{ mm}$; neotissue porosity
434 90%, pore size = 50 μm). This value is close to pressure drop measured computationally in
435 microporous polyester scaffolds in flow through perfusion bioreactors ($\Delta P = 46\text{ Pa}$, scaffold $h =$
436 20 mm , $d = 100\text{ mm}$, flow rate 10 ml/min (Podichetty, Bhaskar et al. 2014)) but also with
437 experimentally measured pressure drop values for chitosan-gelatin scaffolds (with 80-92 %
438 porosity and pore size $\sim 100\text{ }\mu\text{m}$, $d = 100\text{ mm}$, $h = 2\text{ mm}$ (Podichetty, Dhane et al. 2012)).
439 Pressure drop is a parameter that can be measured online during bioreactor culture and, when
440 linked to this whole-scaffold model, could provide a non-invasive readout to monitor neotissue
441 growth in perfusion bioreactors. The continuous secretion of ECM, in particular upon reaching
442 3D ‘confluency’, is expected to lead a decrease in neotissue permeability. This will occur mostly
443 towards the later stages of perfusion culture. To achieve this in the present model a time-
444 dependent parameter could be coupled to the porosity parameter making it decrease over time,

445 reflecting ECM deposition. This could be experimentally validated by measuring pressure drop
446 across the tissue engineered construct over time, for fixed flow rates.

447 Scaffold design in perfusion bioreactors can affect neotissue growth in two ways. First, due to
448 their geometry, the scaffolds provide topographies that will enhance 3D cell growth and neotissue
449 formation (Melchels, Barradas et al. 2010; Van Bael, Chai et al. 2012). Second, their design will
450 define fluid flow patterns throughout the entirety of the scaffold affecting the mechanical stimuli
451 exerted on the cells (Hutmacher and Singh 2008) and the resulting growth kinetics. The
452 computational tool developed in this study could be used to assess biomechanical regimes that
453 will develop in a particular scaffold during neotissue growth but also to evaluate the effect of
454 specific geometries using scaffold CAD designs on these regimes. This model could also suggest
455 ‘ideal’ geometries where shear stress variation exerted on cells across the scaffold could be
456 minimized, resulting thus in a more homogeneous cell population phenotype. The spatiotemporal
457 mapping of shear stress levels will allow to more accurately link phenotypic responses in
458 bioreactors (Gomes, Sikavitsas et al. 2003; Yu, Botchwey et al. 2004; Grayson, Marolt et al.
459 2011) with the experienced biomechanical microenvironment. Moreover, already observed
460 phenomena such as shear stress dependent ECM secretion and mineralization (Gomes, Sikavitsas
461 et al. 2003; Sikavitsas, Bancroft et al. 2005; Papantoniou, Chai et al. 2013) could be also linked
462 to the experienced microenvironment. For their validation such models will require 3D tools of
463 high resolution such as NanoCT as presented here or synchrotron X ray microCT as reported
464 elsewhere (Albertini, Giuliani et al. 2009; Voronov, VanGordon et al. 2013).

465

466 **Conclusion**

467 In this study, a 3D model of microporous neotissue growth in a dynamic culture environment was
468 presented in which the neotissue growth velocity depends on scaffold geometry and fluid flow

469 induced shear stress. The obtained simulation results showed a correspondence with established
470 experimental observations. Although the model can be extended to include additional
471 determinants of the growth process, in its current state it is already able to act as a scaffold shape
472 and bioprocess optimization tool, allowing for a control of the flow-induced mechanical
473 stimulation and growth of the neotissue.

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