Progressive Models of Current Sources for the Magnetic Vector Potential Finite Element Formulation

*Patrick Dular,* †Mauricio V. Ferreira da Luz, †Patrick Kuo-Peng, †João Pedro Assumpção Bastos and †Laurent Krähenbühl

*University of Liege, ACE, Belgium* †Universidade Federal de Santa Catarina, GRUCAD, Brazil

E-mail: Patrick.Dular@ulg.ac.be

**Abstract**—Progressive refinements of the current sources for the magnetic vector potential finite element formulation are done with a subproblem method. The sources are first considered via perfect magnetomotive force or Biot-Savart models up to their volume finite element models, from statics to dynamics. Conversions of the common volume sources to surface sources are first done to lighten the computational resources. Accuracy improvements are then efficiently obtained for local currents and fields, and global quantities, i.e. inductances, resistances, Joule losses and forces.

**Index Terms**—Finite element method, inductor, model refinement, subproblems.

I. INTRODUCTION – OBJECTIVES

The current sources in finite element (FE) magnetodynamic problems can generally be considered via Biot-Savart (BS) models, possibly giving the conductors some simplified geometries, with, e.g., filament, circular or rectangular cross sections [1]. The portions of conductors that are slotted into magnetic core regions can first be omitted by considering their interfaces with the cores as perfect magnetic walls, thus neglecting the slot leakage flux. The associated sources are then magnetomotive forces (MMF) [2].

MMF sources are inherently associated with surfaces whereas BS models define source fields (SFs) that are originally volume sources (VSs). It is here proposed to convert the BS SFs into surface sources (SSs) as well, to lighten the computational resources. The developments are performed in the frame of the magnetic vector potential \( \mathbf{a} \) formulation. Accuracy improvements up to volume FE representations of the conductors, that improve the local field distributions, and from static to dynamic excitations, that accurately render skin and proximity effects, can be done at a second step via the subproblem (SP) method (SPM) [3], [4], that defines a general frame for the whole modeling procedure.

II. PROGRESSIVE MODELS – METHODOLOGY

**MMF model**— Each slot boundary \( \Gamma_{slot} \), possibly including air gap boundary portions (Fig. 1a), is given an essential boundary condition (BC) that fixes a zero normal trace of the magnetic flux density \( \mathbf{b} \). The key is to express this BC with the tangential trace \( \mathbf{n} \times \mathbf{a}\Gamma_{slot} \) (with \( \mathbf{b} = \text{curl} \mathbf{a}, \mathbf{n} \) is the unit normal) replaced by the gradient of a surface scalar potential \( u \), i.e., \( \mathbf{n} \times \text{grad} u\Gamma_{slot} \), or in 2-D via a floating \( \mathbf{a} \) on \( \Gamma_{slot} \) (constant but unknown) (Fig. 1b) [2]. MMF global sources are weakly associated with all \( \Gamma_{slot} \) as it will be shown in the full paper.

**BS model**— With the SPM, the BS SF evaluations can be limited to the material regions \( \Omega_m \) [3], [4], instead of the whole domain with the common method [1]. Such a support reduction already allows to lighten the BS calculations. Then, for accurate combinations with the reaction fields, the BS SFs gain at being projected onto similar function spaces (edge FE). Also, instead of volume projections of the SFs in the mesh of \( \Omega_m \), the SFs gain at being calculated there via a FE problem with their boundary values as BCs on \( \partial \Omega_m \), thus already limiting the BS evaluations to surfaces.

To go one step further, such a preliminary FE problem can be avoided through its inclusion in the main SP. The key is to think of two successive SPs-a-b actually solved together. SP-a first prevents the field to enter \( \Omega_m \), thus with a reaction field in \( \Omega_m \) opposing the BS field, keeping unchanged the outer field to \( \Omega_m \) (this is simply expressed via both normal and tangential field trace discontinuities through \( \partial \Omega_m \) as SSs; the result is an exactly zero field in \( \Omega_m \) with no need of volume calculation). Then, SP-b considers the actual physical properties in \( \Omega_m \), with no more VSs, which is a great advantage. Combining SPs-a-b thus gives a single SP that only requires BS evaluations on \( \partial \Omega_m \) for SSs (Figs. 1c-d).

Couplings between MMF and BS models can be defined in 3-D at the interfaces between slots and end winding regions. Additional SPs for volume FE models of the conductors can follow (Figs. 1g-h) [3], [4]. The advantages of the proposed methodology will be shown to be numerous and significant.

![Fig. 1. Current source in a slot with air gap: field lines for (a) full model solution, (b) flux wall solution, (c) BS SF with its projection limited to the core boundary, (d) total solution with BS SF, (e)-(f) window zooms of (a)-(d), (g) volume correction of coil and its surrounding, (h) with its elevation (z-component of \( \mathbf{a} \)) pointing out the field trace discontinuities.](image)

**REFERENCES**


