



“Seuls les sages, même réduits à l’extrême mendicité, sont riches.”

Cicéron, 1er siècle avant J.-C.

An introduction to optical/IR interferometry

Brief summary of main results obtained during the past lectures:

$$\rho = R / z$$

$$T_{\text{eff}} = (F/\sigma)^{1/4} = (f / \sigma \rho^2)^{1/4}$$

$$E = A(z) \exp[i2\pi\nu t]$$

$$E = A(z, t) \exp[i2\pi\nu t]$$

$$\tau = 1 / \Delta\nu \quad \lambda_{\text{eff}} = \lambda^2 / \Delta\lambda$$

$$I = A A^* = |A|^2 = a^2 .$$



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If $\Delta \geq \lambda / (2B)$, fringe disappearance!

$$I_q = I + I + 2I |\gamma_{12}(0)| \cos(\beta_{12} - 2\pi\nu\tau)$$

$$\gamma_{12}(\tau) = \langle V_1^*(t) V_2(t - \tau) \rangle / I$$

Fringe visibility:
$$v = \left(\frac{I_{\max} - I_{\min}}{I_{\max} + I_{\min}} \right) = |\gamma_{12}(0)|$$



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$$V = \left| \gamma_{12}(0, u, v) \right| = \left| \iint_S I'(\xi, \eta) \exp\{-i2\Pi(u\xi + v\eta)\} d\xi d\eta \right|$$

$$I'(\xi, \eta) = \iint \gamma_{12}(0, u, v) \exp\{i2\Pi(\xi u + \eta v)\} d(u)d(v)$$

- For the case of a 1D uniformly brightening star whose angular diameter is $\phi = b/z'$, we found that the visibility of the fringes is zero when $\lambda/B = b/z' = \phi$ where B is the baseline of the interferometer
- For the case of a double star with an angular separation $\phi = b/z'$, we found that the visibility of the fringes is zero when $\lambda/2B = b/z' = \phi$

An introduction to optical/IR interferometry

8.1 The fundamental theorem

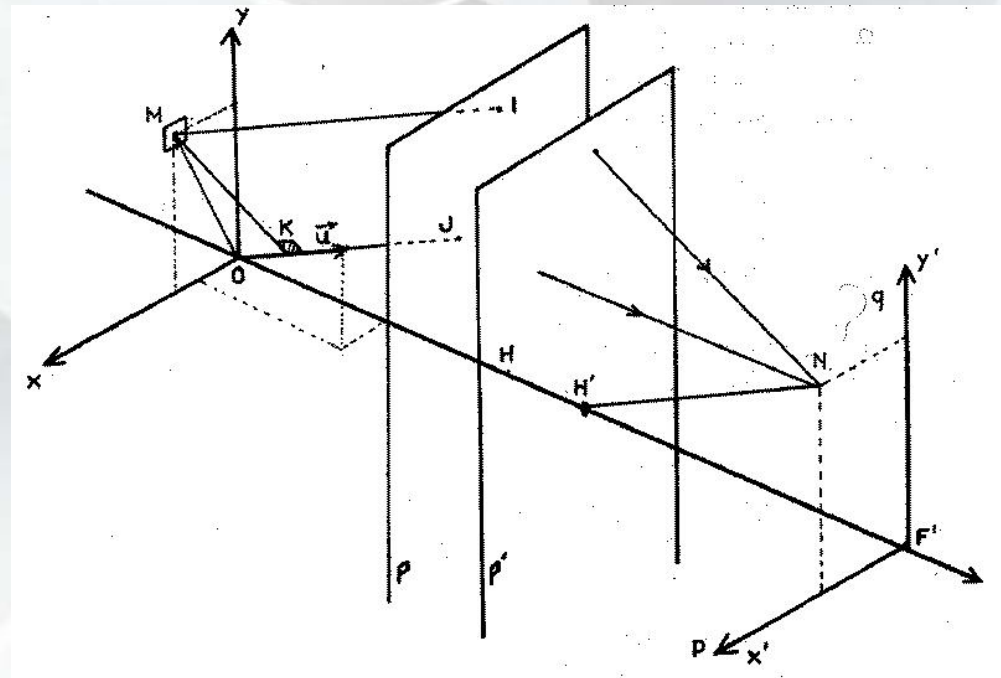
$$a(p, q) = \text{TF}_-(A(x, y))(p, q),$$

$$a(p, q) = \int_{R^2} A(x, y) \exp[-i2\pi(px + qy)] dx dy,$$

with

$$p = x' / (\lambda f)$$

$$q = y' / (\lambda f)$$



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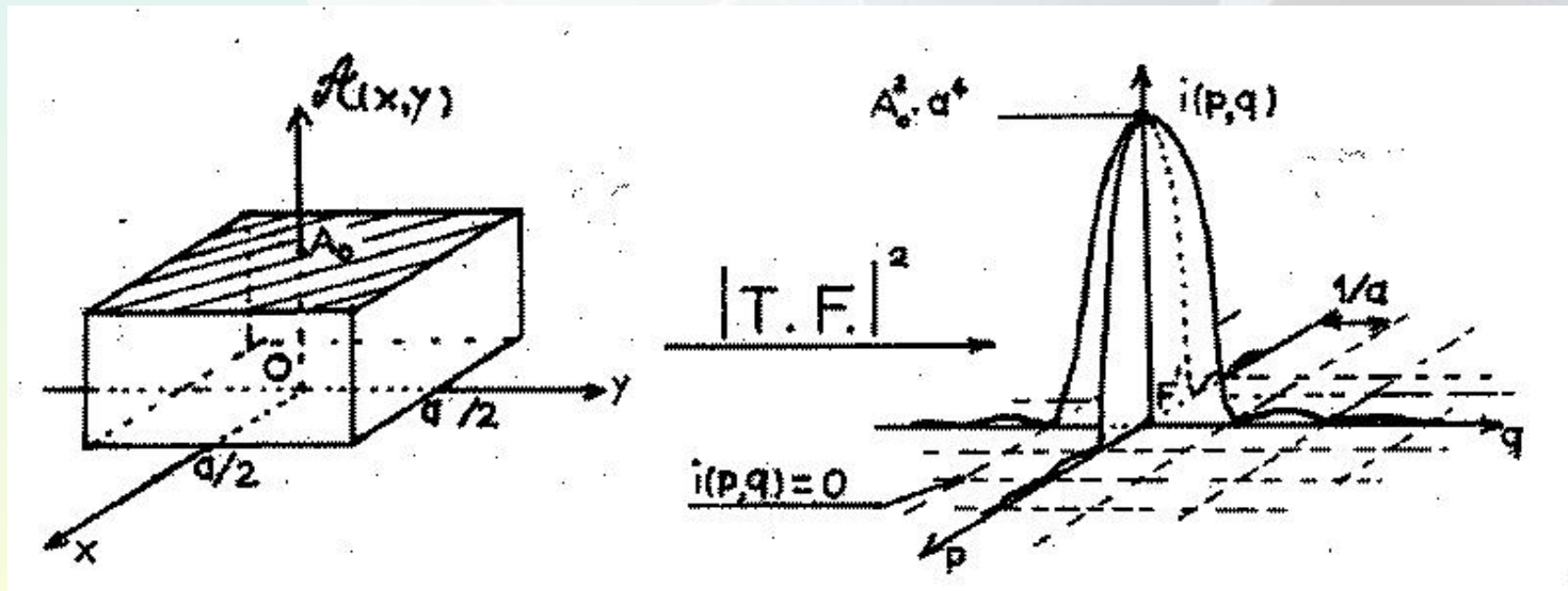
8.1 The fundamental theorem

The distribution of the complex amplitude $a(p,q)$ in the focal plane is given by the Fourier transform of the distribution of the complex amplitude $A(x,y)$ in the entrance pupil plane.

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8.1 The fundamental theorem

Application: Point Spread Function determination



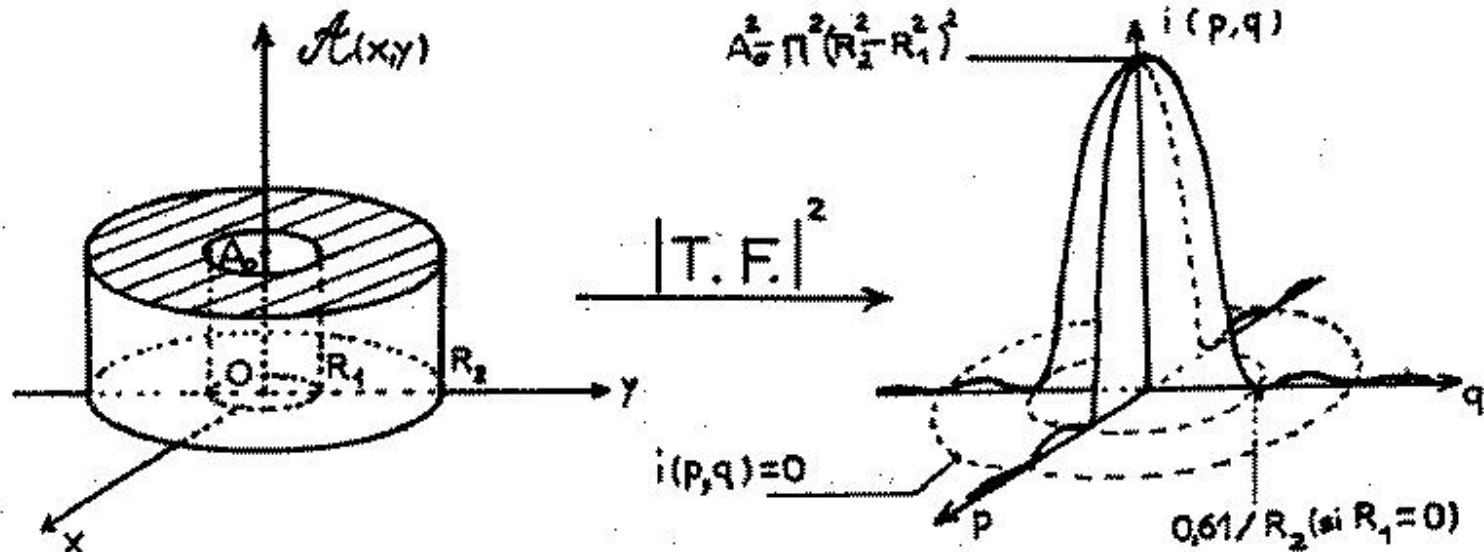
$$\Delta p = \Delta x' / (\lambda f); \Delta q = \Delta y' / (\lambda f) = 2/a \rightarrow \Delta \phi_x = \Delta \phi_y = 2\lambda/a \quad (8.1.7)$$

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8.1 The fundamental theorem

Application: Point Spread Function dete

$$h(p,q) = \text{TF}_-(P(x,y))(p,q)$$



$$i(\rho') = |a(\rho')|^2 = (A_0 \pi)^2 [R_2^2 \frac{2 J_1(Z_2)}{Z_2} - R_1^2 \frac{2 J_1(Z_1)}{Z_1}]^2, \quad (8.1.8)$$

$$\text{with } Z_2 = 2\pi R_2 \rho' / (\lambda f) \text{ and } Z_1 = 2\pi R_1 \rho' / (\lambda f). \quad (8.1.9)$$

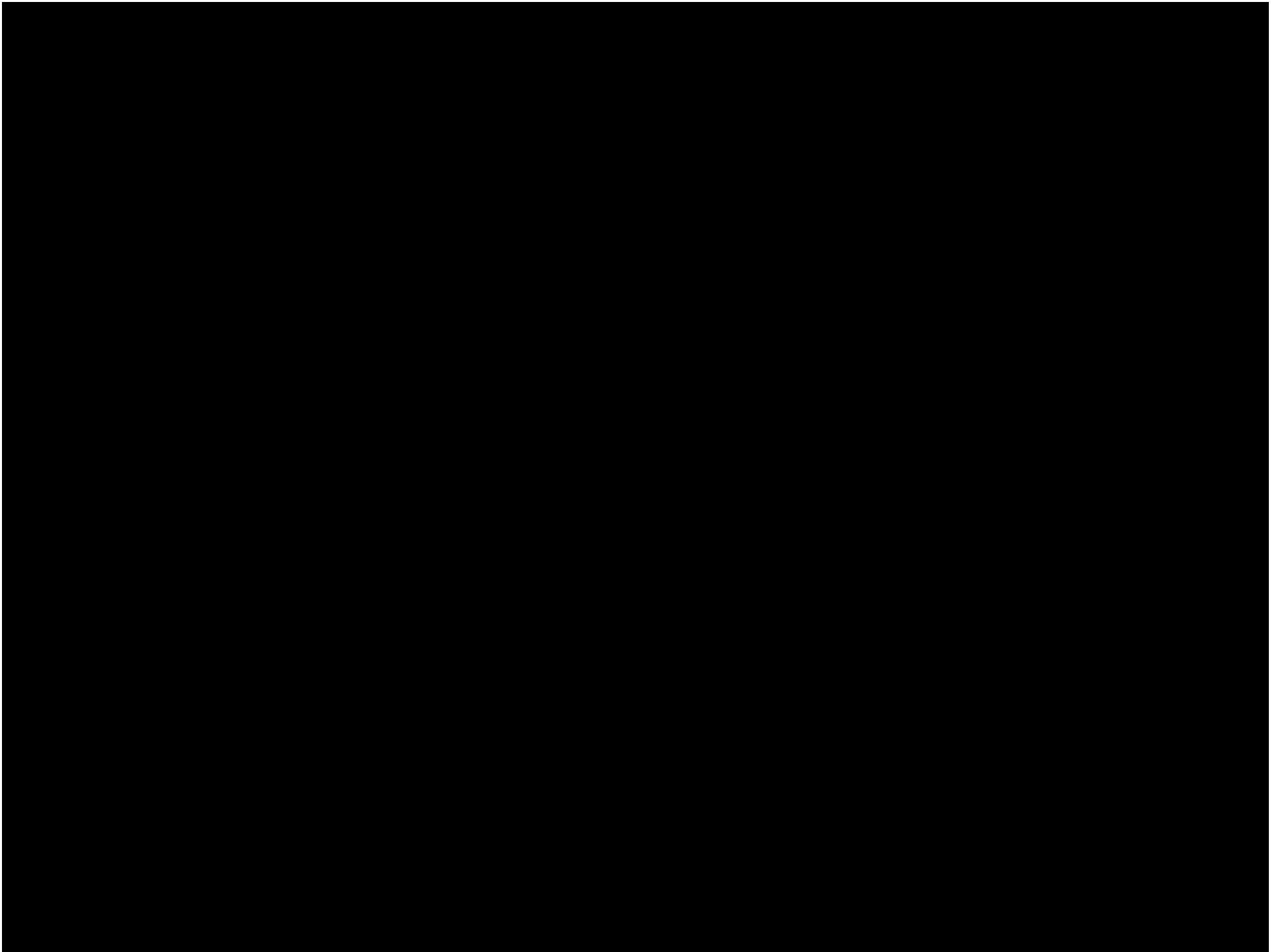
From the previous result, i.e. for the case of a circular aperture with a radius R , the distribution of the complex amplitude in the focal plane is given by the expression:

$$a(\rho') = (A_0 \pi) [R^2 2 J_1(Z) / Z],$$

where $Z = 2\pi R \rho' / (\lambda f)$

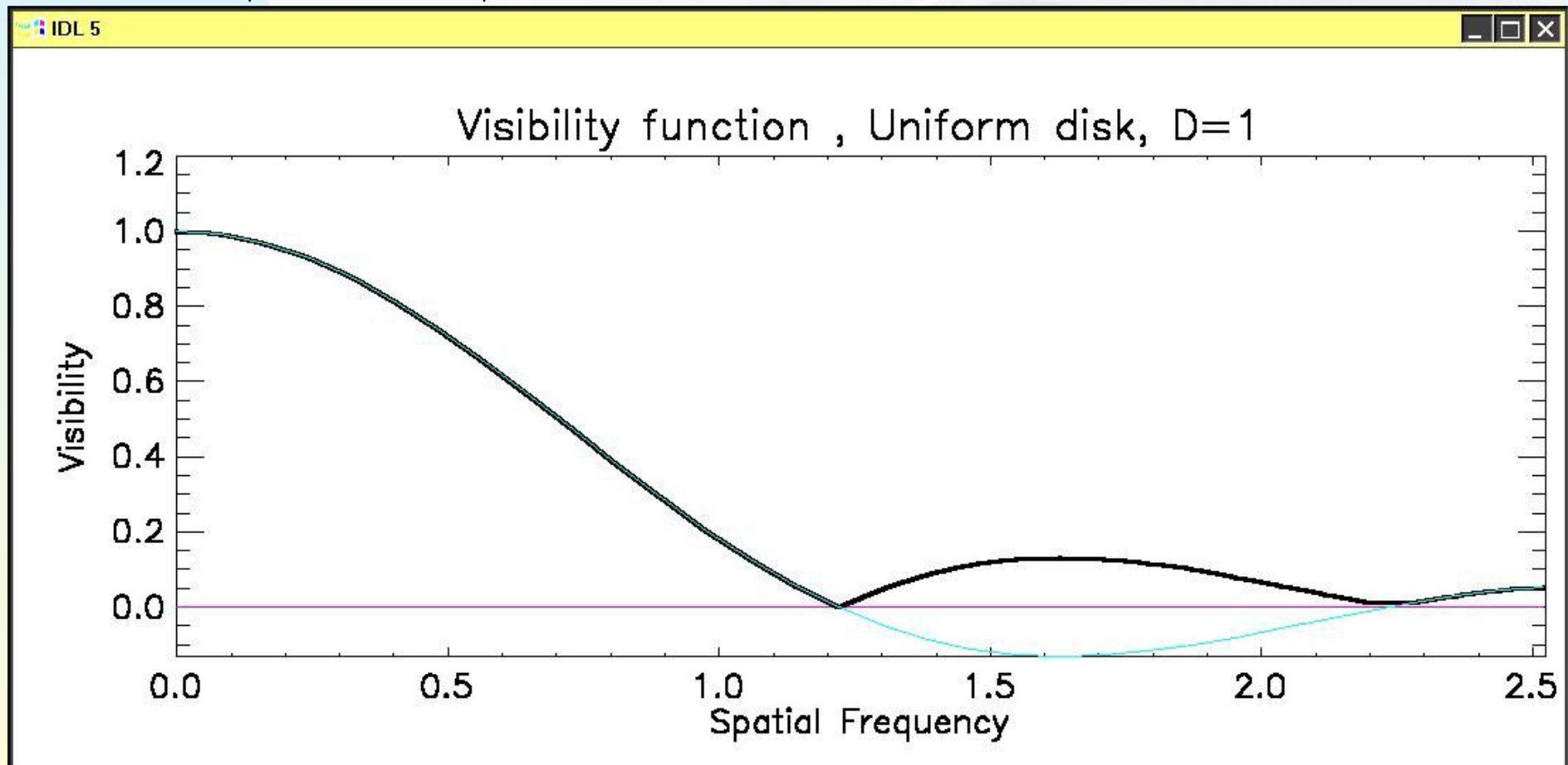
one should be able to demonstrate the next result, i.e., the visibility V of the fringes observed for the case of a uniformly bright circular disk source with an angular diameter θ_{UD} by means of an interferometer with a baseline B is given by:

$$v = \left(\frac{I_{\max} - I_{\min}}{I_{\max} + I_{\min}} \right) = |\gamma_{12}(0)| = TF(I') = \frac{2J_1(\pi\theta_{UD}B / \lambda)}{\pi\theta_{UD}B / \lambda}$$



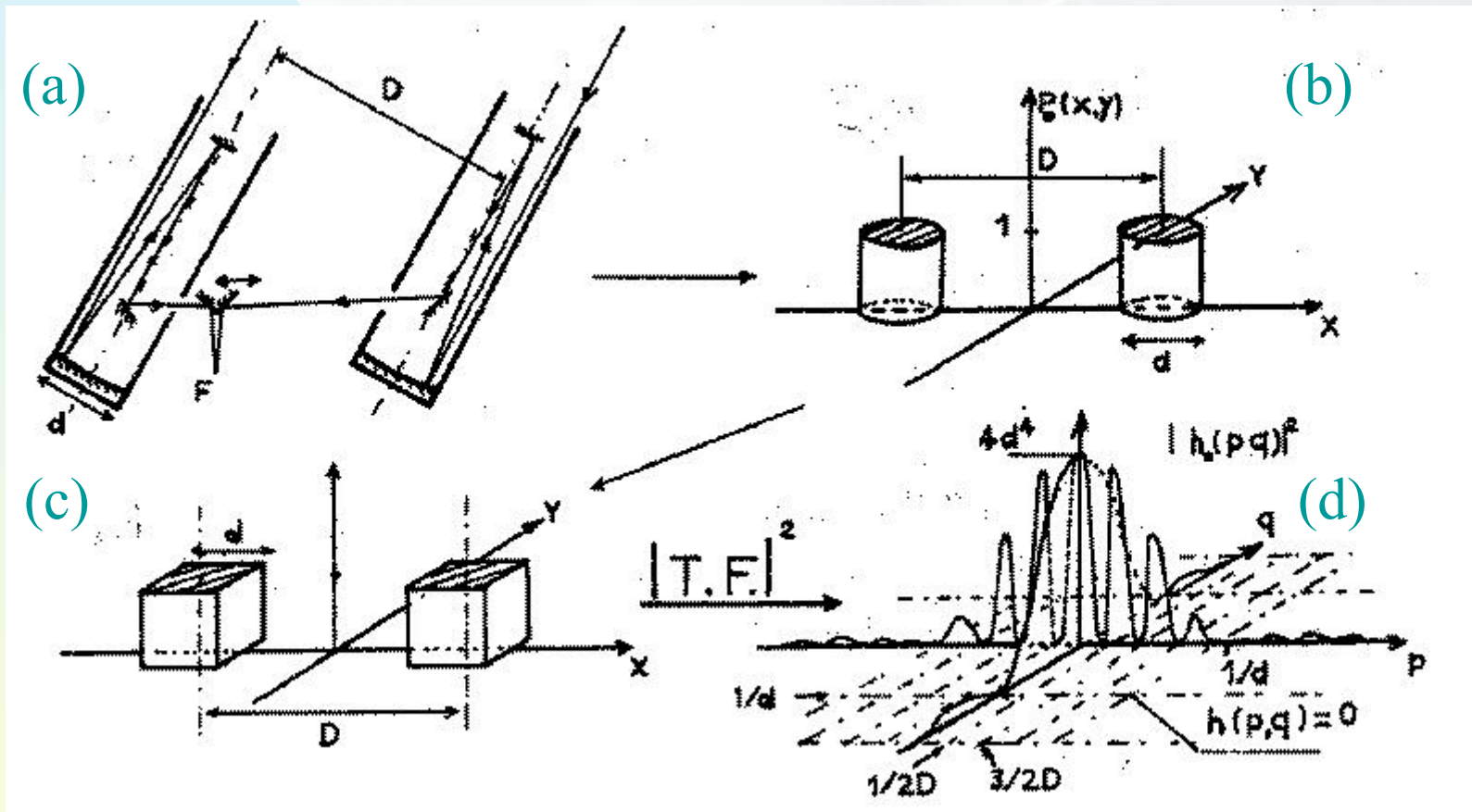
If the source is characterized by a uniform disk light distribution, the corresponding visibility function is given by

$$v = \left(\frac{I_{\max} - I_{\min}}{I_{\max} + I_{\min}} \right) = |\gamma_{12}(0)| = TF(I') = \frac{2J_1(\pi\theta_{UD}B/\lambda)}{\pi\theta_{UD}B/\lambda}$$



An introduction to optical/IR interferometry

8.1 The fundamental theorem: 2 telescope interferometer



Two coupled optical telescopes: simplified optical scheme (a). Distribution of the complex amplitude for the case of two circular (b) or square (c) apertures and corresponding impulse response (d).

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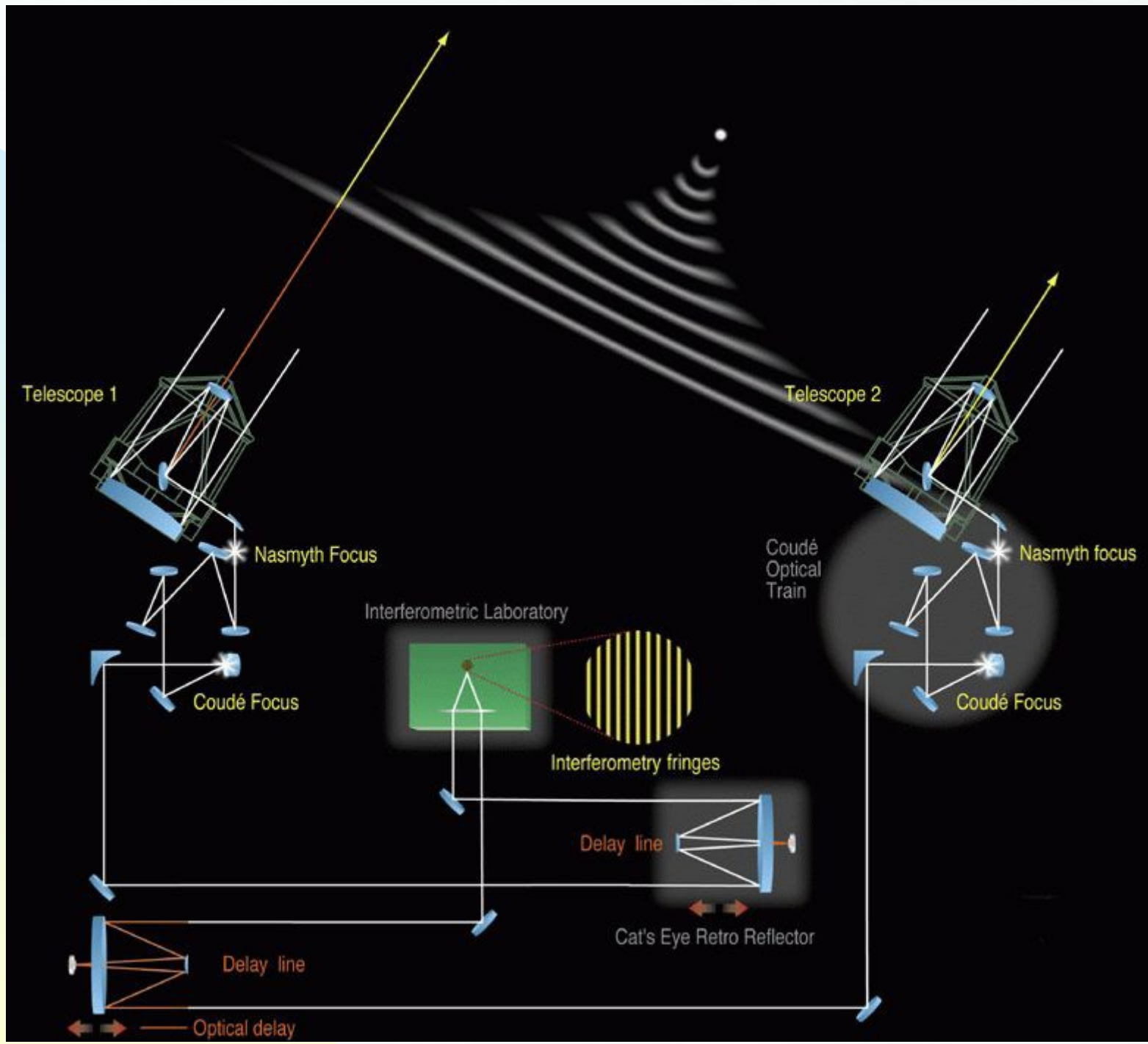
8.1 The fundamental theorem: 2 telescope interferometer

$$h(p, q) = TF(P(x, y))(p, q) = \int_{R^2} P(x, y) \exp[-i2\pi(px + qy)] dx dy \quad (8.1.10)$$

$$\begin{aligned} h(p, q) &= TF(P_0(x + D/2) + P_0(x - D/2))(p, q) = \\ &TF(P_0(x + D/2))(p, q) + TF(P_0(x - D/2))(p, q) = \\ &\exp(i\pi D) TF(P_0(x))(p, q) + \exp(-i\pi D) TF(P_0(x))(p, q) = \\ &(\exp(i\pi D) + \exp(-i\pi D)) TF(P_0(x))(p, q) = \\ &2 \cos(\pi D) TF(P_0(x))(p, q) \end{aligned} \quad (8.1.11)$$

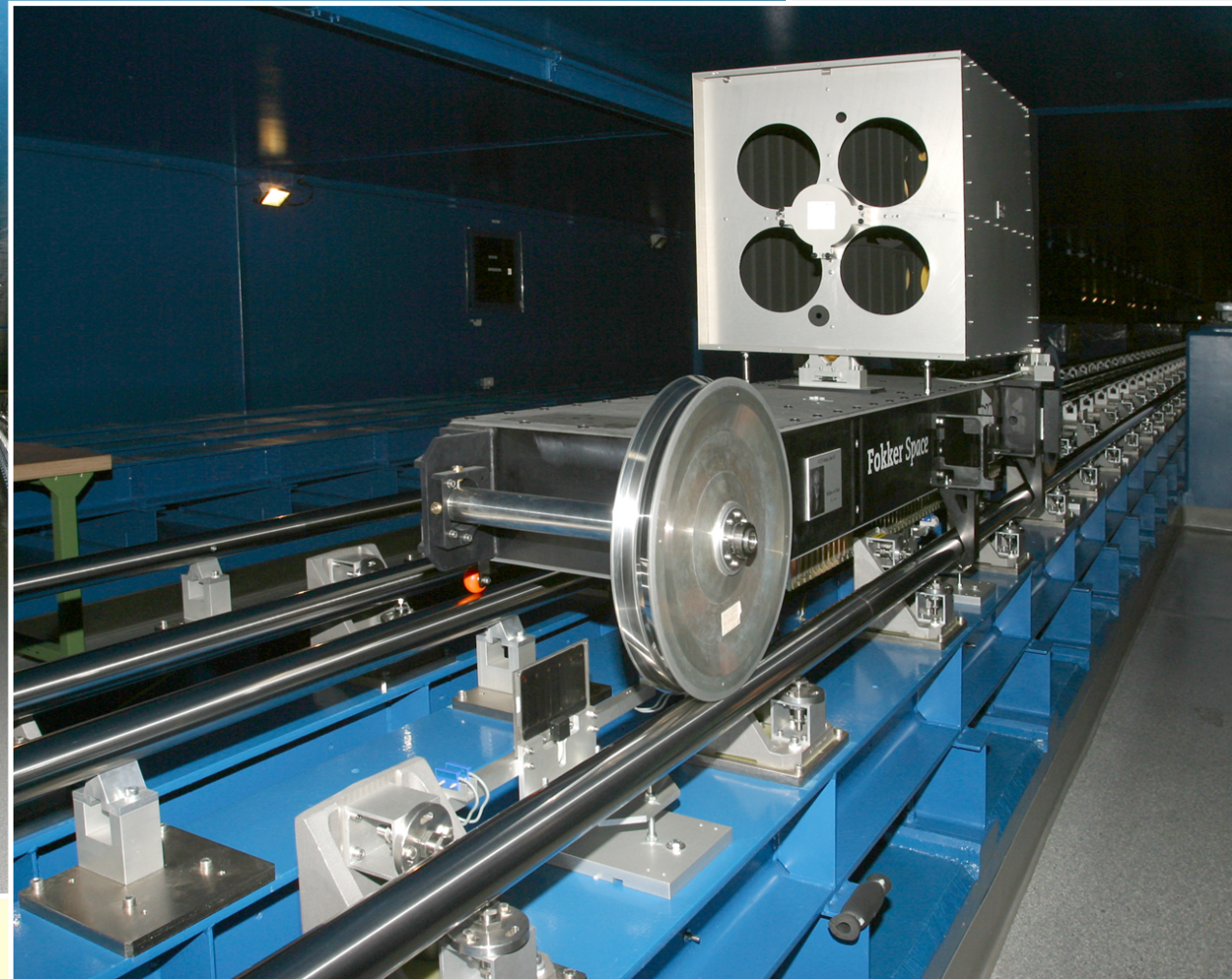
For the particular case of two square apertures:

$$i(p, q) = |h(p, q)|^2 = 4 \cos^2(\pi p D) d^4 \left(\frac{\sin(\pi q d)}{\pi q d} \right)^2 \left(\frac{\sin(\pi p d)}{\pi p d} \right)^2 \quad (8.1.12)$$



Delay lines at the VLT

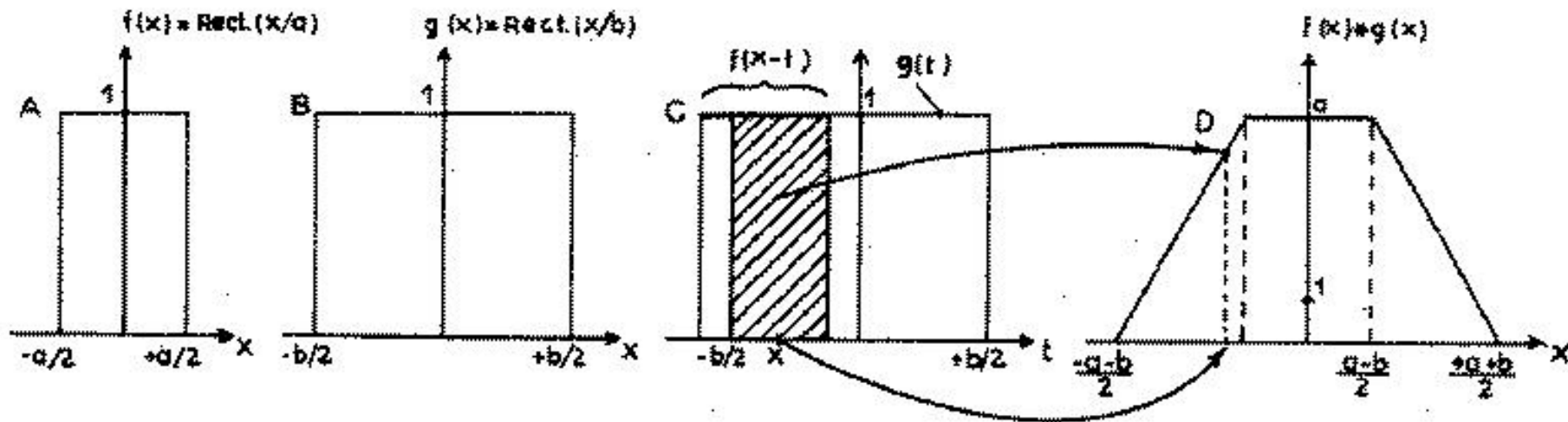
VLT (Chile)



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8.2 The convolution theorem

$$f(x) * g(x) = (f * g)(x) = \int_{R^n} f(x-t)g(t)dt$$

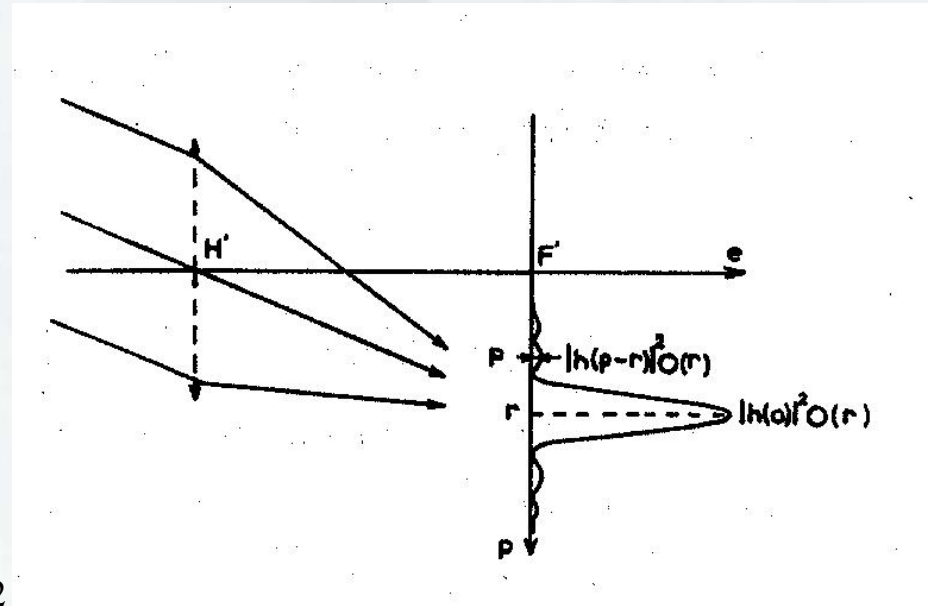


Convolution product of two 1D rectangle functions. A) $f(x)$, B) $g(x)$, C) $g(t)$ and $f(x-t)$; the dashed area represents the integral of the product of $f(x-t)$ and $g(t)$ for the given x offset, D) $f(x)*g(x) = (f*g)(x)$ represents the previous integral as a function of x .

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8.2 The convolution theorem

$$e(p,q) = O(p,q) * |h(p,q)|^2,$$



$$e(p,q) = \int_{R^2} O(r,s) |h(p-r, q-s)|^2 dr ds$$

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8.2 The convolution theorem

For the case of a point-like source:

$$O(p,q) = E \delta(p,q), \quad (8.2.1)$$

$$[\delta(x) = 0 \text{ if } x \neq 0, \delta(x) = \infty \text{ if } x = 0] \text{ and} \quad (8.2.2)$$

$$e(p,q) = O(p,q) * |h(p,q)|^2 = E \delta(p,q) * |h(p,q)|^2 = E |h(p,q)|^2 \quad (8.2.3)$$

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8.2 The convolution theorem

More generally, since

$$\text{TF}_-(f * g) = \text{TF}_-(f) \text{TF}_-(g). \quad (8.2.4)$$

We find, because

$$e(p,q) = O(p,q) * |h(p,q)|^2 \quad (8.2.5)$$

that:

$$\text{TF}_-(e(p,q)) = \text{TF}_-(O(p,q)) \text{TF}_-(|h(p,q)|^2), \quad (8.2.6)$$

and, finally,

$$O(p,q) = \text{TF}^{-1} [\text{TF}_-(e(p,q)) / \text{TF}_-(|h(p,q)|^2)]. \quad (8.2.7)$$

An introduction to optical/IR interferometry

8.2 The convolution theorem

$$O(p,q) = (\lambda^2 E / \phi^2) \Pi(p \lambda / \phi) \Pi(q \lambda / \phi). \quad (8.2.8)$$

$$e(p,q) = O(p,q) * |h_0(p,q)|^2.$$

$$e(p) = O(p) * |h_0(p)|^2, \quad (8.2.9)$$

$$e(p) = 2 d^2 (\lambda / \phi) \sqrt{E} \int_{p-\phi/2\lambda}^{p+\phi/2\lambda} \left(\frac{\sin(\pi r d)}{\pi r d} \right)^2 \cos^2(\pi r D) dr \quad (8.2.10)$$

$$\left(\frac{\sin(\pi r d)}{\pi r d} \right)^2 \approx \text{Cte sur } [p-\phi/2\lambda, p+\phi/2\lambda], \quad \text{and} \quad (8.2.11)$$

$$e(p) = 2 d^2 (\lambda / \phi) \sqrt{E} \left(\frac{\sin(\pi p d)}{\pi p d} \right)^2 \int_{p-\phi/2\lambda}^{p+\phi/2\lambda} \cos^2(\pi r D) dr. \quad (8.2.12)$$

An introduction to optical/IR interferometry

8.2 The convolution theorem

$$e(p) = 2d^2 \left(\frac{\sin(\pi p d)}{\pi p d} \right)^2 \left[O(p) * \cos^2(\pi p D) \right], \quad (8.2.13)$$

$$e(p) = 2d^2 \left(\frac{\sin(\pi p d)}{\pi p d} \right)^2 \left[\frac{1}{2} \int_R O(p) dp + \frac{1}{2} O(p) * \cos(2\pi p D) \right] \quad (8.2.14)$$

$$e(p) = A \left[B + \frac{1}{2} \operatorname{Re}(O(p) * \exp(i2\pi p D)) \right], \quad (8.2.15)$$

$$A = 2d^2 \left(\frac{\sin(\pi p d)}{\pi p d} \right)^2 \quad \text{et} \quad B = \frac{1}{2} \int_R O(p) dp, \quad (8.2.16)$$

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8.2 The convolution theorem

$$e(p) = A \left[B + \frac{1}{2} \operatorname{Re} \left(\int_R O(r) \exp(i2\pi(p-r)D) dr \right) \right], \quad (8.2.17)$$

$$e(p) = A \left[B + \frac{1}{2} \cos(2\pi pD) \operatorname{TF}_-(O(r))(D) \right], \quad (8.2.18)$$

$$\gamma(D) = (e_{\max} - e_{\min}) / (e_{\max} + e_{\min}), \quad (8.2.19)$$

$$\gamma(D) = \operatorname{TF}_-(O(r))(D) / (2B) = \operatorname{TF}_-(O(r))(D) / \int O(p) dp. \quad (8.2.20)$$

An introduction to optical/IR interferometry

8.3 The Wiener-Khintchin theorem

In our case, this theorem merely states that the Fourier transform of the PSF (see Eq. (8.2.7)) is the auto-correlation function of the distribution of the complex amplitude in the pupil plane:

$$TF(|h(a,b)|^2) = \iint P^*(x,y)P(x-a,y-b)dx dy$$

Démonstration (1/2):

Let us evaluate: $\iint P^*(x, y) P(x - a, y - b) dx dy$

Let us also remind:

$$h(p, q) = \iint P(x, y) \exp(-i2\pi(px + qy)) dx dy$$

And thus (Fourier inverse transform),

$$P(x, y) = \iint h(p, q) \exp(i2\pi(px + qy)) dp dq, \text{ and also}$$

$$P^*(x, y) = \iint h^*(p', q') \exp(-i2\pi(p'x + q'y)) dp' dq'$$

We then find that

$$\iint P^*(x, y) P(x - a, y - b) dx dy =$$

$$\iint \iint h^*(p', q') h(p, q) \exp(-i2\pi(ap + bq)) dp dq$$

$$\iint \exp(-i2\pi((p' - p)x + (q' - q)y)) dx dy dp' dq'$$

Démonstration (2/2):

Since

$$\iint \exp(-i2\pi((p'-p)x + (q'-q)y)) dx dy = \delta(p-p')\delta(q-q'), \text{ we then find}$$

$$\iint P^*(x,y) P(x-a,y-b) dx dy =$$

$$\iint \iint h^*(p',q') h(p,q) \delta(p-p')\delta(q-q') \exp(-i2\pi(ap+bq)) dp dq dp' dq'$$

and finally

$$\iint P^*(x,y) P(x-a,y-b) dx dy =$$

$$\iint |h(p,q)|^2 \exp(-i2\pi(ap+bq)) dp dq =$$

$$= TF(|h(p,q)|^2)(a,b).$$

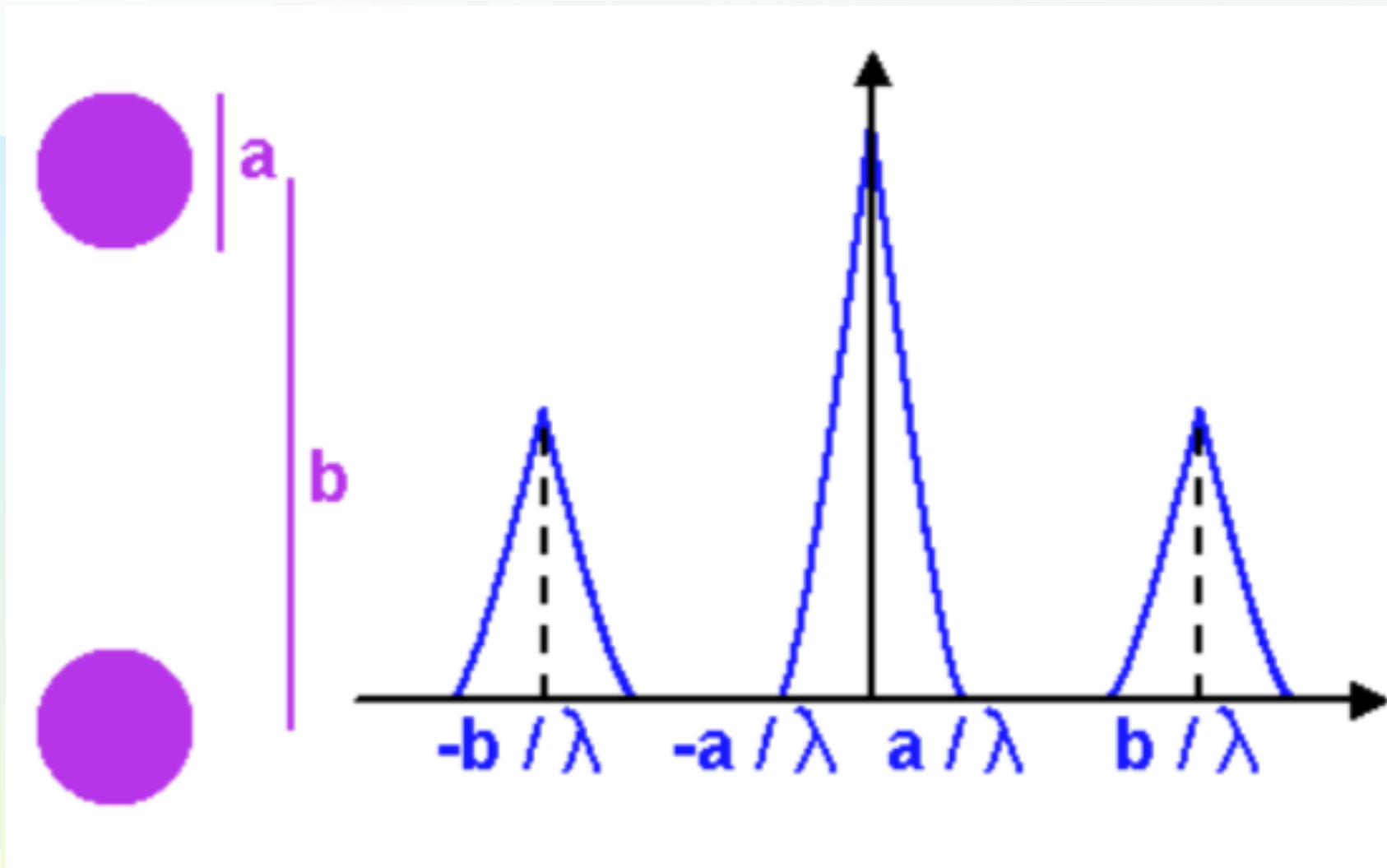
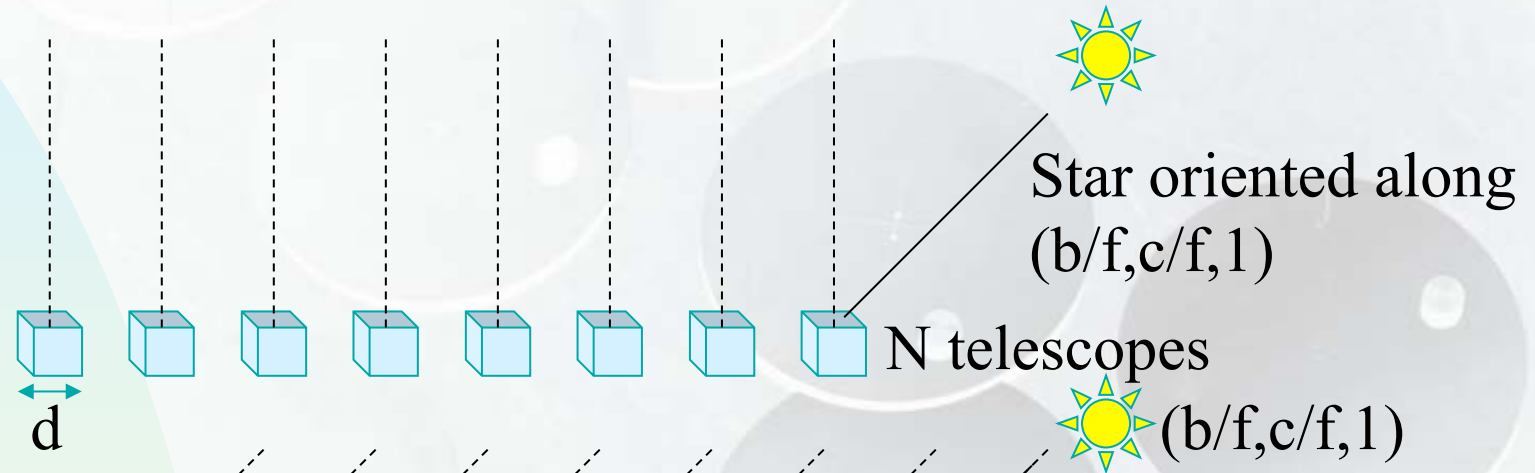


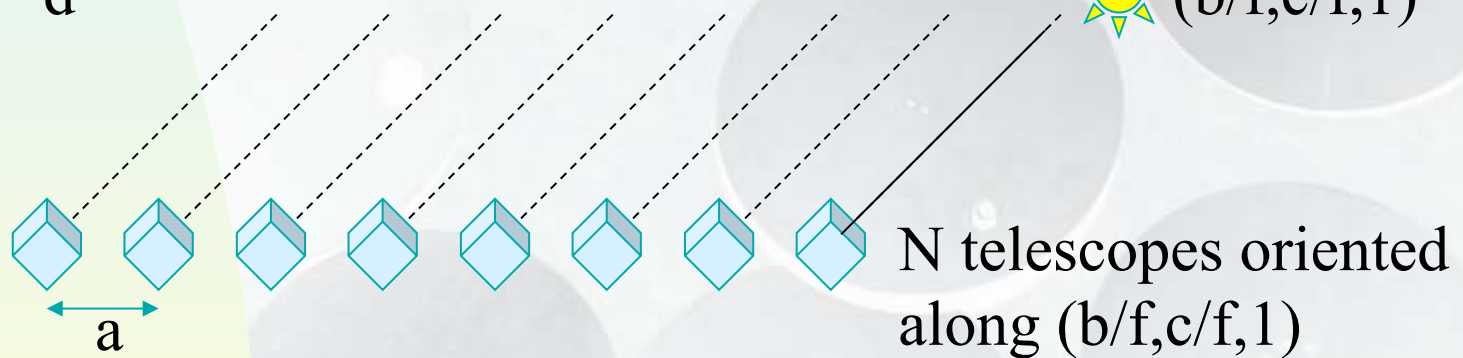
Diagram illustrating the autocorrelation as a function of the space frequency, for the case of an interferometer composed of 2 telescopes, with a diameter a , separated by the baseline b . The autocorrelation of the pupil gives access to high space frequencies.

- **Exercices:** Calculate the response functions (PSFs) for the cases described in the figures below!

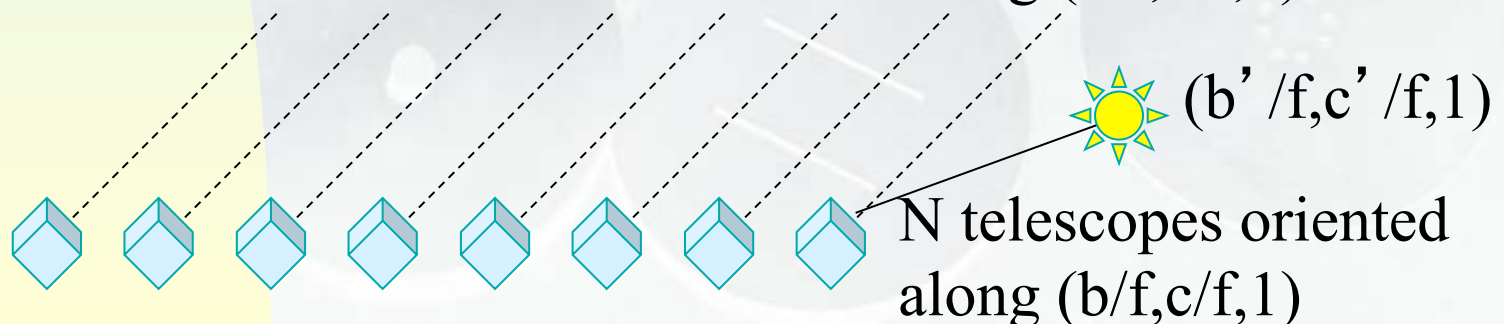
A)



B)

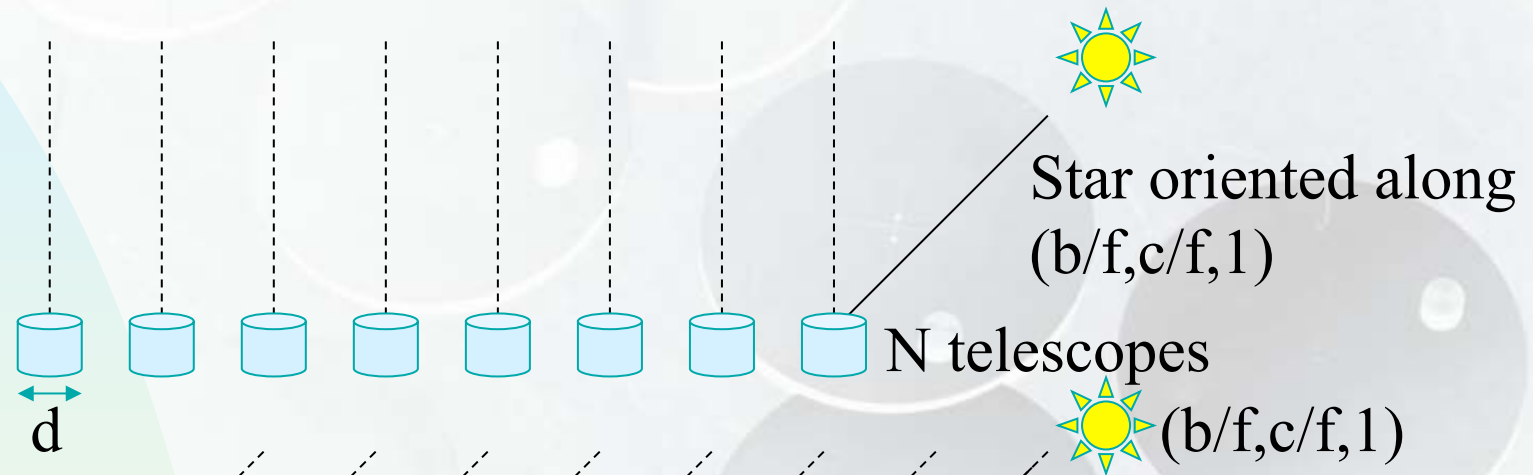


C)

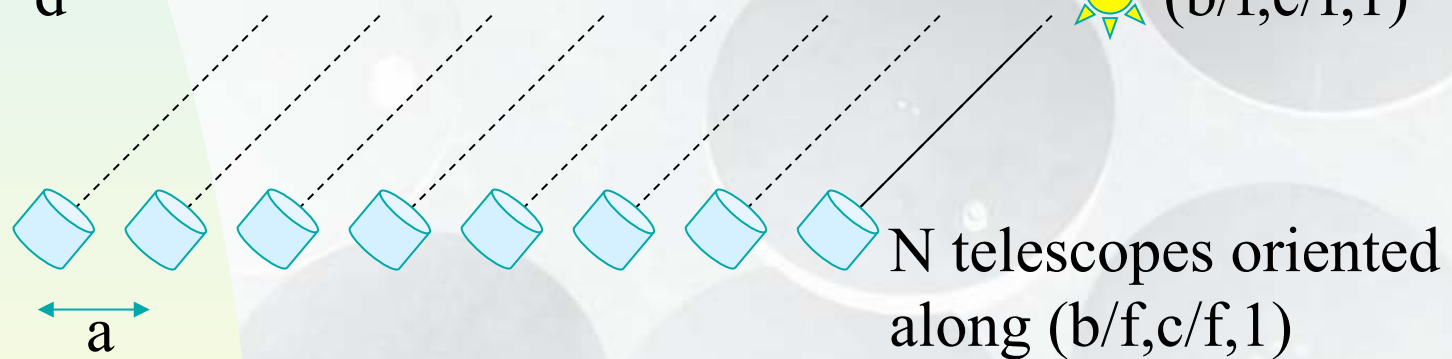


- **Exercices:** Calculate the response functions (PSFs) for the cases described in the figures below!

A)



B)



C)

