A radio-map of the colliding winds in the very massive binary system HD 93129A

P. Benaglia\textsuperscript{1,*}, B. Marcote\textsuperscript{2}, J. Moldón\textsuperscript{3}, E. Nelan\textsuperscript{4}, M. De Becker\textsuperscript{5}, S. M. Dougherty\textsuperscript{6}, and B. S. Koribalski\textsuperscript{7}

\textsuperscript{1} Instituto Argentino de Radioastronomía, CCT-La Plata, CONICET, Argentina; e-mail: paula@iar-conicet.gov.ar
\textsuperscript{2} Departament d’Astronomia i Meteorologia, Institut de Ciències del Cosmos, Universitat de Barcelona, IEEC-UB, Martí i Franquès 1, E08028 Barcelona, Spain
\textsuperscript{3} ASTRON Netherlands Institute for Radio Astronomy, Postbus 2, 7990AA Dwingeloo, The Netherlands
\textsuperscript{4} Space Telescope Science Institute, USA
\textsuperscript{5} Department of Astrophysics, Geophysics and Oceanography, University of Liège, 17 Allée du 6 Août B5c, B-4000 Sart Tilman, Belgium
\textsuperscript{6} NRC Herzberg Astronomy and Astrophysics, DRAO, PO Box 248, Penticton B.C., Canada V0H 1K0
\textsuperscript{7} Australia Telescope National Facility, CSIRO Astronomy & Space Science P O Box 76, Epping, NSW 1710, Australia

Received: accepted

ABSTRACT

Context. Radio observations are an effective tool to discover particle acceleration regions in colliding-wind binaries, through detection of synchrotron radiation; these regions are natural laboratories for the study of relativistic particles. Wind-collision region (WCR) models can reproduce the radio continuum spectra of massive binaries that contain both thermal and non-thermal radio emission; however, key constraints for models come from high-resolution mapping. Only five WCRs have been resolved to date at radio frequencies at milliarcsecond (mas) angular scales. The source HD 93129A, prototype of the very few known O2 I stars, is a promising target for study: recently, a second massive, early-type star about 50 mas away was discovered, and a non-thermal radio source detected in the region. Preliminary long-baseline array data suggest that a significant fraction of the radio emission from the system comes from a putative WCR.

Aims. We sought evidence that HD 93129A is a massive binary system with colliding stellar winds that produce non-thermal radiation, through spatially resolved images of the radio emitting regions.

Methods. We completed observations with the Australian Long Baseline Array (LBA) to resolve the system at mas angular resolutions and reduced archival Australia Telescope Compact Array (ATCA) data to derive the total radio emission. We also compiled optical astrometric data of the system in a homogeneous way. We reduced historical Hubble Space Telescope data and obtained absolute and relative astrometry with milliarcsec accuracy.

Results. The astrometric analysis leads us to conclude that the two stars in HD 93129A form a gravitationally bound system. The LBA data reveal an extended arc-shaped non-thermal source between the two stars, indicative of a WCR. The wind momentum-rate ratio of the two stellar winds is estimated. The ATCA data show a point source with a change in flux level between 2003-4 and 2008-9, that is modeled with a non-thermal power-law spectrum with spectral indices of $-1.03 \pm 0.09$ and $-1.21 \pm 0.03$ respectively. The mass-loss rates derived from the deduced thermal radio emission and from the characteristics of the WCR are consistent with estimates derived by other authors.

Key words. Stars: massive – binaries: general – Radio continuum: stars – Radiation mechanisms: non-thermal – stars: individual: HD 93129A

1. Introduction

Wind-collision regions (WCR) in binary systems with two hot massive stars constitute natural laboratories to study relativistic particles \citep{Eichler1993, Benaglia2003, Pittard2006, DeBecker2007}, as do supernova remnants but with somewhat different physical and geometrical parameters. A WCR is the site of particle acceleration that generates a population of relativistic electrons, as now demonstrated in more than 40 massive binary systems. The particle-accelerating colliding-wind binary (PACWB) status of these objects is in most cases revealed by radio observations through the detection of synchrotron emission \citep{DeBecker2013}. WCR models are able to reproduce the radio continuum spectra of these massive binaries \citep{Pittard2006}. The radiation from these systems is the superposition of thermal emission from the stellar winds of each binary component and the non-thermal emission from the WCR. A key ingredient for constraining models of the non-thermal component is expected to come from high-resolution radio imaging. The non-thermal emission region associated with colliding winds has been mapped at radio frequencies for WR 140 \citep{Dougherty2005}, WR 146 \citep{O'Connor2005}, WR 147 \citep{Williams1997}, Cyg OB2 #5 \citep{Contreras1997}, Cyg OB2 #7 \citep{Ortiz-Leon2011} and Cyg OB2 #9 \citep{Dougherty2006}. In each case, the WCR is resolved by observations with angular resolutions of a few milliarcseconds (mas).

* Fellow of CONICET
HD 93129A is the brightest source in the most crowded part of Tr 14 in the Carina nebula and has been studied for several decades. The record of its special spectral characteristics can be traced back to Payne (1927) (see also Table 1 for information on other studies). It now is recognised as hosting the prototype O2 If+ (Walborn et al. 2002). Radio emission from HD 93129A has been detected between 1.4 and 24.5 GHz (e.g. Benaglia et al. 2006), showing a flux decreasing with frequency indicative of synchrotron radiation, most probably produced in a WCR between stars in a colliding-wind binary system. The Australian Long Baseline Array (LBA) observation presented by Benaglia et al. (2010) suggested that a significant fraction of the radio emission comes from a region compatible with a WCR.

In this paper, we offer an in-depth discussion of a more recent LBA observation of HD 93129A, aiming to provide compelling evidence of the WCR origin for the bulk of the radio emission, in agreement with the so-called standard scenario for PACWBs. These results allow us to spatially resolve for the first time the synchrotron emission region and to derive specific information about the non-thermal radio contribution. The two component nature of HD 93129A is presented, including the most recent astrometric measurements that show the relative position of the components over time (Sect. 2), and a new analysis of Hubble Space Telescope/Fine Guidance Sensor (HST/FGS) data that allow us to correlate the optical positions of the two components with the position of the radio emission (Sect. 3). The LBA observations are presented in Sect. 4. The main results and general discussion are in Sections 5 and 6, with a final statement of our conclusions in Sect. 7.

2. The HD 93129A stellar system

The object HD 93129A was first classified as an O3 star (Walborn 1982), though later was reclassified as the prototype O2 I star (Walborn et al. 2002). Mason et al. (1998) had suggested it was a speckle binary, and was confirmed as having two components, Aa and Ab, from HST/FGS observations (Nelan et al. 2004); these authors provide the separation at the epoch of the observations, the difference in apparent visual magnitude between the primary and the secondary and the position angle (PA) (see also Nelan et al. 2010). Maiz-Apellaniz et al. (2008) detected proper motion along the radius vector between the components after analysis of data from 1996, 2002 and 2004: the binary nature is supported by the relative velocity of about 2 mas yr\(^{-1}\), higher than the expected relative velocities in the field. The authors suggested a highly elliptical or inclined orbit, or both, although they did not determine the orbital parameters. The companion is assumed coeval with the primary and of the same spectral type as various other OB stars nearby (O3.5 V) (see Table 1). Through detailed ultra-violet and H\(\alpha\) line modeling, Taresch et al. (1997) and Repolust et al. (2004) derived the main parameters of HD 93129A (see Table 1).

The relative motion of the two stars can be traced by archived HST/FGS and HST/ACS observations from 1996 to 2009, along with recent VLT/NACO and VLTI/PIONIER observations of Sana et al. (2014). Figure 1 shows the relative astrometry from these measurements over time. The position angles for 1996 and 2002 have large uncertainties because the binary components were resolved along only one FGS axis, due to the H\(\text{ST}\) roll angle resulting in a small projected separation along the unresolved axis. For later epochs, the approximate position angle was known, so the observations were planned to avoid this problem.

HD 93129A was included in the recent Galactic O star Spectroscopic Survey of Sota et al. (2014). The coordinates from the Tycho catalogue (Hog et al. 2000) assume a single star since the presence of a companion star was unknown at the time. Hipparcos did not resolve the system. These coordinates have errors of about 0.1', an order of magnitude larger than the expected angular resolution of the LBA data of ~ 10 mas.

Table 1. Adopted parameters of HD 93129A, relevant here.

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Value</th>
<th>Unit</th>
<th>Ref.</th>
</tr>
</thead>
<tbody>
<tr>
<td>Aa Spectral type</td>
<td>O2 If+</td>
<td></td>
<td>Wal02</td>
</tr>
<tr>
<td>Ab Spectral type</td>
<td>O3.5 V</td>
<td></td>
<td>Wal06</td>
</tr>
<tr>
<td>(\Delta m)</td>
<td>0.9±0.05</td>
<td>mag</td>
<td>Ne110</td>
</tr>
<tr>
<td>Wind terminal velocity(\dagger)</td>
<td>3200±200</td>
<td>km s(^{-1})</td>
<td>Tar97</td>
</tr>
<tr>
<td>System mass</td>
<td>200±45</td>
<td>M(_{\odot})</td>
<td>Mai08</td>
</tr>
<tr>
<td>Effective temperature(\dagger)</td>
<td>42500</td>
<td>K</td>
<td>Rep04</td>
</tr>
<tr>
<td>System luminosity (log)</td>
<td>6.2</td>
<td>L/L(_{\odot})</td>
<td>Rep04</td>
</tr>
<tr>
<td>X-ray luminosity(\dagger)</td>
<td>1.3 x 10(^{-7})</td>
<td>L(_{\odot})</td>
<td>Coh11</td>
</tr>
<tr>
<td>Distance</td>
<td>2.5 kpc</td>
<td></td>
<td>Wal95</td>
</tr>
</tbody>
</table>

\(\dagger\): Derived from a single star. \(\dagger\): Derived from the energy band (0.5 – 8) keV.

References. Wal02: Walborn et al. (2002); Wal06: Walborn, priv. communication; Ne110: Nelan et al. (2010); Tar97: Taresch et al. (1997); Mai08: Maiz-Apellaniz et al. (2008); Rep04: Repolust et al. (2004); Coh11: Cohen et al. (2011); Wal95: Walborn (1995).

3. Hubble Space Telescope/FGS data analysis

3.1. Astrometry

To improve upon the coordinates of HD 93129A from the Tycho catalogue, we re-analyzed HST/FGS measurements. The three Fine Guidance Sensors on the H\(\text{ST}\) are 2-channel (x,y) white-light shearing interferometers that provide line-of-sight positional measurements of guide stars to enable accurate pointing and stabilization of the H\(\text{ST}\). The FGS is capable of resolving sources down to ~15 mas when operated in its high angular TRANS (or Transfer) mode and sub-milliarcsecond astrometry when operated in POS (or Position) mode. A description of the FGS and its operation can be found in Nelan (2014).

In eight epochs between 2006 and 2009, FGS observed HD 93129A in TRANS mode, along with “wide angle” POS mode observations of the binary itself and seven reference field stars within one arcminute (see Fig 2). The TRANS mode observations resolve HD 93129Aab, providing the component separation, position angle, and magnitude difference at each epoch. The POS mode observations supply the relative positions of the field stars and composite HD 93129A system. Combining data from these two observing modes gives the position of Aa and Ab relative to these field stars at each epoch. Combining all epochs yields a local FGS catalog of the relative positions and proper motions of reference stars and binary components.

Figure 3 shows the interference fringes of HD 93129A for the Jan 27, 2009 observations, the best fitting binary star model and, for comparison, the fringes of an unresolved point source.
Table 2. PPMXL and FGS astrometry of the HD 93129A field.

<table>
<thead>
<tr>
<th>Star</th>
<th>RA (h)</th>
<th>Dec (deg)</th>
<th>μ_α</th>
<th>μ_δ</th>
<th>ξ</th>
<th>η</th>
<th>μ_α</th>
<th>μ_δ</th>
</tr>
</thead>
<tbody>
<tr>
<td>R1</td>
<td>10.7325347 (13)</td>
<td>-59.535576 (19)</td>
<td>-6.6 (2.5)</td>
<td>-0.6 (2.5)</td>
<td>-2.6531 (0.9)</td>
<td>43.2955 (0.5)</td>
<td>-5.78 (0.29)</td>
<td>1.54 (0.31)</td>
</tr>
<tr>
<td>R2</td>
<td>10.7322835 (14)</td>
<td>-59.541006 (16)</td>
<td>——</td>
<td>——</td>
<td>-9.5351 (1.0)</td>
<td>23.4780 (0.5)</td>
<td>-5.78 (0.53)</td>
<td>1.71 (0.50)</td>
</tr>
<tr>
<td>R3</td>
<td>10.7340882 (13)</td>
<td>-59.541208 (19)</td>
<td>-4.8 (3.2)</td>
<td>0.6 (3.2)</td>
<td>39.9116 (1.2)</td>
<td>23.1199 (0.5)</td>
<td>-6.73 (0.32)</td>
<td>1.38 (0.36)</td>
</tr>
<tr>
<td>R4</td>
<td>10.7344007 (49)</td>
<td>-59.552781 (74)</td>
<td>24.6 (17)</td>
<td>60.5 (17)</td>
<td>48.0711 (1.2)</td>
<td>-19.2493 (0.5)</td>
<td>-5.34 (0.44)</td>
<td>2.22 (0.45)</td>
</tr>
<tr>
<td>R6</td>
<td>10.7326563 (13)</td>
<td>-59.560735 (19)</td>
<td>-4.8 (2.1)</td>
<td>3.0 (2.1)</td>
<td>0.8902 (1.1)</td>
<td>-47.2647 (0.4)</td>
<td>-6.29 (0.32)</td>
<td>1.64 (0.34)</td>
</tr>
<tr>
<td>R7</td>
<td>10.7315615 (13)</td>
<td>-59.557925 (19)</td>
<td>-6.6 (2.5)</td>
<td>1.6 (2.5)</td>
<td>-29.0924 (1.1)</td>
<td>-37.1969 (0.2)</td>
<td>-5.46 (0.29)</td>
<td>1.60 (0.30)</td>
</tr>
<tr>
<td>R8</td>
<td>10.7296359 (15)</td>
<td>-59.548569 (15)</td>
<td>-6.4 (2.0)</td>
<td>2.5 (2.0)</td>
<td>-81.8936 (1.6)</td>
<td>-3.6759 (0.3)</td>
<td>-5.53 (0.32)</td>
<td>2.12 (0.30)</td>
</tr>
</tbody>
</table>

"Star": Reference star identification, same as in Figure 2. All errors are shown within parenthesis. PPMXL coordinates (J2000): RA in hours, Dec in degrees, with errors in mas. The J2000 FGS tangential plane coordinates, ξ, η, are given in arcsec, relative to the point mid-way between the binary components; errors are in mas. Proper motions (μ) in mas yr^{-1}, with errors in the same unit. *: Reference star 2 is not in the PPMXL catalog, so coordinates from UCAC4 are quoted.

The binary model takes into account the light contributed by HD 93129B which, 2.73” away, accounts for about 5 percent of the light of HD 93129A (FGS optics collimates a source onto a circular beam with a full width-half maximum intensity diameter of about 3”). Figure 2 also shows the location on the fringes that the POS mode observations of HD 93129A “lock” on to. The model fit to the TRANS mode observations provides the offset of each component from this point, which thereby allows us to create pseudo POS mode exposures for both HD 93129Aa and HD 93129Ab, and hence accurate positions for each star relative to the reference stars. The bright O star HD 202124 was used as the point-source calibrator since it is nearly the same B-V color as HD 93129A, and it was observed by FGS only eight days after HD 93129A. The orientation of HST at the time of the observation caused the projected separation of HD 93129Aa and B along the FGS y-axis to be only 0.233”, as is evident in the figure.

To compare the FGS coordinates of HD 93129Aab to the radio source, it is necessary to convert the relative FGS positions to absolute International Celestial Reference System (ICRS) coordinates. The FGS coordinate frame lies in a tangential plane, with the zero point arbitrarily set to the point mid way between the binary components (at the zero point the FGS Cartesian plane is tangential to the celestial sphere). We used the PPMXL catalog4 for the coordinates and proper motions of the reference stars as input to a model that finds the RA,Dec of this tangent point. This is a two step process. The FGS observa-

---

4 PPMXL is a catalog of positions, proper motions, 2MASS and optical photometry of 900 million stars and galaxies.
for directly without any external constraints. Table 2 lists the reference star identifier, PPMXL coordinates and proper motions with the catalog errors, and the resulting FGS tangent plane coordinates and proper motions and errors. The positional errors for RA, Dec are in mas, while the proper motion errors are in mas yr$^{-1}$. The FGS $\xi, \eta$ coordinates are in arcsecs, and their errors are in mas. The FGS proper motions are now absolute rather than relative. We note that PPMXL does not contain reference star R2. For this star we used the UCAC4 (Fourth US Naval Observatory CCD Astrograph Catalog) coordinates but solved for the proper motion using only FGS data.

Next step is to find the RA,Dec of the zero point of the FGS tangential plane. The accurate FGS proper motions are used to update the PPMXL coordinates to epoch 2009, which is the epoch of the FGS master plate. Using the algorithms outlined in Smart [1960], (1) initial values for the RA,Dec of the FGS tangent point are given, for which the Guide Star Catalog v2.3 coordinates of HD 93129A are used, (2) the PPXML (RA,Dec) of each reference star are converted to Cartesian coordinates in this tangential plane, which are treated as “observations with errors”, (3) the PPXML Cartesian coordinates are used to set the scale and orientation of the FGS frame to (East, North). (4) An iterative process is applied to find a global solution that provides the best-fit of the observed FGS and PPXML coordinates in the tangential plane. The solution creates a hybrid FGS-PPMXL catalog and computes the RA,Dec of the tangent point. Once convergence is achieved, it is a straightforward process to compute the (RA,Dec) of any point in the FGS frame. We used the GAUSSFIT [Jefferys et al. 1987] least-squares program to build the models to carry out these calculations.

We compute the RA,Dec of the HD 93129/Aa and Ab components, adjusted for proper motion to the Aug 6, 2008 epoch of the radio observations, equinox J2000:

$$\alpha, \delta(\text{Aa}) = 10^h43^m57.455^s, -59^\circ32'51.36''$$
$$\alpha, \delta(\text{Ab}) = 10^h43^m57.456^s, -59^\circ32'51.33''$$

The uncertainty in each coordinate is ±27 mas, dominated by the PPXML catalog errors in the position and proper motion of the seven FGS reference stars surrounding HD 93129A. However the position of Aa relative to Ab is accurate to about 1 mas. From these positions, the separation and position angle of Aa in Aug 2008 was 36 ± 1 mas and $PA = 12 \pm 1^\circ$.

The proper motions derived from the FGS data are $\mu_{\alpha}, \mu_{\delta}(\text{Aa}) = -8.4, 2.6$ mas yr$^{-1}$ and $\mu_{\alpha}, \mu_{\delta}(\text{Ab}) = -9.0, 0.0$ mas yr$^{-1}$. Note that the proper motion of the binary’s components include any putative orbital motion along with the system motion. The difference from the reference star proper motions (Table 2), particularly in RA, could reflect internal motion within the core of Tr 14, while the reference stars lie outside the core.

3.2. Orbit estimation

The relative motion of the components Aa and Ab hints at HD 93129A being a gravitationally bound system. However, the relative positions shown in Figure 1 cover a rather small part of the orbit and thus are inadequate to perform a standard determination of orbital elements. No definitive radial velocity variation has been reported, which is not surprising given the wide separation and apparent near linear trajectory of the stars. Nevertheless, a first approximation of an orbital fit was attempted. An algorithm that minimizes the weighted square distance between the measured data points and the fitted orbit (Casco & Vila, private communication) was used to determine a preliminary set
of orbital parameters. Our results suggest an orbital period of the order of 200 years, an eccentricity larger than 0.9, and semi-major axes of about 93 and 37 mas respectively for components a and b. Our best-fit solution points also to a periapsis argument of about 220 degrees and an inclination of about 103 degrees. An attempt to fix the inclination value close to zero yielded a much poorer fit (larger weighted root mean square rms). According to this preliminary solution, the next periastron passage is expected to take place in 2024. It should be emphasized that such a preliminary solution was calculated to give a rough idea of the orbit, even though the present astrometry does not allow for a convincing characterization of the orbit accompanied by relevant error bars on the estimated parameters. This calculation has however the merit to provide a preliminary basis to organize future observation campaigns dedicated to HD 93129A.

4. Radio observations


We observed HD 93129A on June 22, 2007 (MJD 54273) between 03:00 and 13:00 UTC at 2.3 GHz as part of eVLBI experiment vt11D3 that used Parkes, Mopra, and Australia Telescope Compact Array (ATCA) to search for suitable phase calibrators near the target (Benaglia et al. 2010). The sources PKS 0637–752 (J0635–7516) and J1047–6217 were used as flux and phase calibrators, respectively. The total time observing HD 93129A was 3 hours.

Standard VLBI calibration procedures were used (see also Sect. 4.2) and the source imaged with a synthesized beam of 0.2"×0.05". The source has a measured flux density of ~3 mJy. See more details in Benaglia et al. (2010) and Fig. 5.

4.2. Observations of August 6, 2008

Further observations were conducted on Aug 6, 2008 (MJD 54684) from 01:00 to 12:00 UTC with the LBA at an observing frequency of 2.3 GHz, and using five antennas: ATCA (in tied-array mode, i.e., working as a single station), Parkes, Mopra, Ceduna, and Hobart (experiment V191B). The baseline lengths range from 100 to 1700 km, providing an angular resolution of ~15 mas at that frequency. The data were recorded at 256 Mbps provided by 4 sub-bands of 16-MHz bandwidth at each of the two circular polarizations. The data were correlated with the Distributed FX (DiFX) software correlator located at Curtin University of Technology (WA, Australia), using 128 channels per sub-band, each band 125 kHz wide, and using 2-second integrations. The data from the ATCA antennas were also correlated as an independent interferometer to obtain the total flux density from a low-resolution observation. The time on HD 93129A amounted to 3.2 hours. The $uv$ coverage is shown in Fig. 6.
5. Results

5.1. Radio flux density

In Fig. [7] we show the radio image from LBA 2008 observations of HD 93129A at 2.3 GHz. Two final images are presented: one at maximum angular resolution (Fig. [7] left: high-resolution image), and a second one using tapering that increases the sensitivity to larger scale emission (Fig. [7] right: low-resolution image). The high-resolution image has an rms of 0.21 mJy beam$^{-1}$, and shows a “comma-shaped” extended emission with a flux density of 1.5 ± 0.5 mJy. The value quoted for the error is the difference in the integrated fluxes above 2 and 3$\sigma$, where $\sigma$ is the measured rms of the image. In the low-resolution image the rms noise is lower and more flux is recovered from the source, as expected. The rms of the image is 0.13 mJy beam$^{-1}$, and a gaussian fit to the source gives an integrated flux density detected by the LBA of 2.9 ± 0.5 mJy, with a peak intensity of 1.8 ± 0.2 mJy beam$^{-1}$.

We note that the amplitude calibration of the data, which is based on the system temperature of the antennas is limited by the uncertainty on the antennas’ response and gain curves. The flux density measured for the phase calibrator is ∼ 30% lower than its unresolved flux density measured with ATCA between Feb 2007 and Nov 2008: 1.46 ± 0.01 Jy to 1.42 ± 0.02 Jy. This indicates that calibrator is partially resolved by the LBA or it is variable, or both. We also note that the LBA-measured flux of HD 93129A is ∼ 35% of the total flux measured in simultaneous Aug 2008 ATCA data at 2.3 GHz. We conclude that the rest of the flux from HD 93129A is resolved out by the LBA.

5.2. Radio morphology and astrometry

From the 2008 LBA image the angle subtended by the emitting source as seen from the secondary star (so-called opening angle $\theta_{\text{w}}$) can be estimated, leading to half-opening angle $\theta_{\text{w}} \approx 65^\circ$. This is related to the wind-momentum rate ratio $\eta$ by $\theta_{\text{w}} \approx 180^\circ \times \eta/(1 + \eta)$, thus implying $\eta \approx 0.6$.

The radio emission at maximum angular resolution presents two maxima. The position of the centroid of the two fitted gaussians is $\alpha_c = 10:43:57.456, \delta_c = -59:32:51.34, J2000$. A systematic uncertainty of 1.1 mas in Right Ascension and 2.8 mas in Declination comes from the position error of the phase reference calibrator in the ICRF2 (International Celestial Reference Frame, second realization). Therefore, the absolute uncertainty of the position of the radio emission is about 3 mas in the ICRF. In Fig. [7] the crosses mark the positions of HD 93129 Aa and Ab, with the size of the crosses indicating the absolute uncertainty, although the relative separation of the components has an uncertainty of ∼ 1 mas (see Sect. [3]).

A systematic uncertainty of 1.1 mas in Right Ascension and 2.8 mas in Declination comes from the position error of the phase reference calibrator in the ICRF2 (International Celestial Reference Frame, second realization). Therefore, the absolute uncertainty of the position of the radio emission is about 3 mas in the ICRF. In Fig. [7] the crosses mark the positions of HD 93129 Aa and Ab, with the size of the crosses indicating the absolute uncertainty, although the relative separation of the components has an uncertainty of ∼ 1 mas (see Sect. [3]).

6. Discussion

6.1. Wind-momentum rates

The radio source appears in the LBA observations as a bow-shaped region, slightly curved around the component Ab, and reminiscent of the WCRs resolved in WR 140 (Dougherty et al. 2005), WR 146 (O’Connor et al. 2005), Cyg OB2 #5 (Ortiz-León et al. 2011), as examples. If we assume that the radio source is centered on the stagnation point of the WCR, the wind momentum rates ratio $\eta$ can be expressed as:

$$\eta \approx \left( \frac{R_b}{R_a} \right)^2 \frac{M_b v_b}{M_a v_a},$$

Table 3. Radio flux density of HD 93129A at different frequencies obtained with ATCA and LBA data. We include the project code and the flux density calibrator name with its flux value.

<table>
<thead>
<tr>
<th>Array/project and Date</th>
<th>$\nu$ [GHz]</th>
<th>$S_\text{cal}$ [mJy]</th>
<th>Flux cal name</th>
<th>$S_\text{real}$ [Jy]</th>
</tr>
</thead>
<tbody>
<tr>
<td>ATCA/C678</td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>2003-01-28</td>
<td>4.8</td>
<td>4.1 ± 0.4$^a$</td>
<td>1934−638</td>
<td>2.84</td>
</tr>
<tr>
<td>2003-01-28</td>
<td>8.6</td>
<td>2.0 ± 0.2$^a$</td>
<td>1934−638</td>
<td>5.83</td>
</tr>
<tr>
<td>2003-12-20</td>
<td>1.4</td>
<td>9.4 ± 0.9$^a$</td>
<td>1934−638</td>
<td>14.98</td>
</tr>
<tr>
<td>2003-12-20</td>
<td>2.4</td>
<td>7.8 ± 0.4$^a$</td>
<td>1934−638</td>
<td>11.59</td>
</tr>
<tr>
<td>2004-05-05</td>
<td>2.4</td>
<td>1.8 ± 0.15$^a$</td>
<td>Mars</td>
<td></td>
</tr>
<tr>
<td>2004-05-05</td>
<td>14.5</td>
<td>1.5 ± 0.35$^a$</td>
<td>Mars</td>
<td></td>
</tr>
<tr>
<td>ATCA/V191B</td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>2008-08-06</td>
<td>2.3</td>
<td>7.5±0.11</td>
<td>0637−752</td>
<td>5.32</td>
</tr>
<tr>
<td>ATCA/C1726</td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>2009-01-18</td>
<td>4.8</td>
<td>5.6±0.3</td>
<td>1934−638</td>
<td>5.83</td>
</tr>
<tr>
<td>2009-01-18</td>
<td>8.6</td>
<td>2.9±0.3</td>
<td>1934−638</td>
<td>2.86</td>
</tr>
<tr>
<td>LBA/V191B</td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>2008-08-06</td>
<td>2.3</td>
<td>2.9±0.51$^b$</td>
<td>0637−752</td>
<td></td>
</tr>
</tbody>
</table>

References. a: Benaglia et al. 2006; b: Derived from image portrayed in Fig. [7] right.

The data reduction was performed using the Astronomical Image Processing System (AIPS) software package. Bad data were flagged using telescope a-priori flag files and information provided in the observing logs. The amplitude calibration was performed using the antenna system temperatures and phase solutions on the calibrators were obtained using the AIPS task FRING. The compact source J1047−6217 was used as the phase calibrator. The fringe finders 0637−790, J1051−6646 and the phase calibrator V191B were used for the bandpass calibration. First, an accurate model of the calibrator was obtained after several self-calibration and imaging cycles. The calibrator, slightly resolved on the longest baselines, had a flux density of 0.99 ± 0.02 Jy. Finally, the phase solutions and the amplitude scale were transferred to the target source, and images produced by deconvolving the visibility data.

The data from ATCA were correlated in standard interferometer mode. The resulting fluxes for the phase calibrator J1047−6217 and HD 93129A were 1.300±0.002 Jy and 7.5±0.1 mJy respectively. The flux and uncertainty for the phase calibrator was derived by fitting a point source to the image, using all baselines.

4.3. ATCA observations

The Australia Telescope Online Archive (ATOA) was searched for observations of HD 93129A. In addition to project C678 (Benaglia et al. 2006) and the observations associated with the LBA projects described in the previous Section, project C1726 targeted HD 93129A for 3.3 h at 4.8 and 8.6 GHz. Table [3] summarizes the dates and observing frequencies for these observations. In all observations with the ATCA, HD 93129A is not resolved by ATCA and appears as a point source.

http://www.aips.nrao.edu
Fig. 7. The LBA observations of HD 93129A at 2.3 GHz. a) Left: At highest angular resolution with an image rms of 0.21 mJy beam$^{-1}$ and contour levels shown for -0.6, 0.6 (3σ), 0.7, 0.8 and 0.9 mJy beam$^{-1}$. b) Right: a tapered image that increases the weight of the shorter baselines and increases sensitivity to larger scale emission; the image rms is 0.13 mJy beam$^{-1}$ and the contour levels are -0.4, 0.4 (3σ), 0.6, 1, 1.4 and 1.8 mJy beam$^{-1}$. The synthesized beams are 15mas × 11mas and 31mas × 23mas respectively, shown in the bottom right-hand corner of the images. The crosses mark the positions of the system components Aa (south) and Ab (north), at the epoch of the radio observations based on the position derived in Sec. 2. North is up and East is to the left.

where $R_i$ is the distance between the WCR and star $i$ (Usov 1992). Note that the value of $\eta$ is independent of orbit inclination.

From the relative position of Aa, Ab and the centroid of the radio emission in Figure 7, we estimate $R_b/R_a \approx 0.7$. The stellar winds of the two stars are assumed to have reached their terminal velocities before colliding, considering the wide orbit, and we further assume that the terminal velocities of both components are similar (see the values quoted for instance in Table 1 of Muijres et al. 2012), then $\eta$ reduces to a mass-loss rate ratio, leading to $\dot{M}_b \approx 0.5 \dot{M}_a$.

In this work, we hesitate to use the shape of the WCR to estimate $\eta$. As Pittard & Dougherty (2006) (section 4.3.3) show for WR 140, the emission in the WCR detected in VLBI is close to the stagnation point and the opening angle has not reached the asymptotic value. They also noted the challenge of identifying the stagnation point relative to the stellar components due to opacity effects. In the best case, we can perhaps assume the shape of the contact discontinuity (CD) between the two stellar winds (Cantó et al. 1996) is represented by the radio emission. In Figure 8, the positions of the CD corresponding to different values of $\eta$ are plotted relative to the emission detected in the LBA observation. The distance and angle between stars are known, so the only free parameters are the offsets with respect to the radio image and $\eta$. The $\eta$ values that best match the radio source are between ~ 0.4 and ~ 0.6.

6.2. On the binary nature of HD 93129A

Do HD 93129Aa and Ab form a gravitationally bound binary system? The system is being monitored in the OWN Survey project (Barbá et al. 2010). A radial velocity (RV) curve from preliminary data shows slight variations though no period has yet been determined. The putative low amplitude RV variations may be due to a number of factors, including insufficient data over a potential orbital period, a very low mass companion or unfavorable orbital inclination (Barbá & Gamen, private communication).

An argument in support of HD 93129Aab being a bound system is the small angular separation and the change in the rate at which the two stars are approaching one another. In the simplest scenario, if they are in the plane of the sky, with a separation of 26.5 mas (epoch 2013.09 from Sana et al. 2014) and at a distance of 2.5 kpc, they have a physical separation of only 66 AU, which implies a bound system for such massive stars. However, this is strongly dependent on the inclination being in the plane of the sky. More compelling is the increase in the rate at which the two stars are approaching one another. Based upon FGS and ACS
data, between 1996 and 2009 the separation had been closing at approximately 2.4 mas yr\(^{-1}\), but VLTI data shows this rate to have accelerated to 4.2 mas yr\(^{-1}\) between 2011.19 and 2013.09 \citep{Sana2014}, consistent with a binary system approaching periastron. The relative positions in Fig. 1 suggest either an orbital inclination close to 90 degrees, or a nearly face-on orbit of high eccentricity. The results of the very preliminary orbit fits given in Sect. 3 favor the former hypothesis.

### 6.3. On the mass-loss rates of the HD 93129A system

HD 93129A has been the target of different studies aimed at deriving the mass-loss rate \( \dot{M} \). \cite{Sana2006} estimated \( \dot{M} \) of HD 93129A assuming it was a single star, by fitting the H\( \alpha \) profile and also using a more complete spectral synthesis approach in the ultraviolet. These two approaches led to consistent values of \( 1.8 \times 10^{-5} \) and \( 2.1 \times 10^{-5} M_\odot \text{yr}^{-1} \), respectively. The authors considered that these values should be reduced by a factor \( 2^{1/4} \), a reduction of about 40\%, if HD 93129A consisted of two very similar components. Using H\( \alpha \) line diagnostics exclusively, \cite{Repolust2004} derived a value of \( 2.4 \times 10^{-5} M_\odot \text{yr}^{-1} \) for the mass-loss rate, noting that the results could strongly suffer from contamination by a companion star.

\cite{Benaglia2004} presented flux density measurements of HD 93129A system from ATCA data at two frequencies, 4.8 and 8.6 GHz, and derived a spectral index \( \alpha = -1.2 \pm 0.3 \). Their interpretation was that both thermal and non-thermal emission were contributing to the radio emission.

\cite{Wright1975} showed that the mass-loss rate of a star with a thermal wind can be expressed in terms of the radio flux density. Using the \cite{Wright1975} formula and if \( f_\text{T} \) is the fraction of thermal emission at a given frequency, it is possible to express the mass-loss rate in terms of \( f_\text{T} \) as

\[
\dot{M} = f_\text{T} M_\odot \text{yr}^{-1} \quad \text{at} \quad v = 8.64 \text{GHz} \quad \text{for} \quad \text{stellar distance} \quad 2.5 \text{ kpc.}
\]

This leads to a mass-loss rate of \( 5 \times 10^{-5} M_\odot \text{yr}^{-1} \) if \( f_\text{T} = 0.5 \).

At the epoch of the observations presented by \cite{Benaglia2004}, the stars Aa and Ab were \( \sim 50 \) mas apart in the plane of the sky (see Sect. 2). The angular resolution of the ATCA data at \( \sim 1 \) arcsec precludes separation of the contributions to the total flux of both stellar winds and a colliding-wind region. An estimate of the thermal emission from the two stellar winds can be obtained from the radio observations in \cite{Benaglia2006} from ATCA at 17.8 and 24.5 GHz. At these frequencies, non-thermal emission is negligible and can be disregarded to zeroth order. The total flux is then the sum of the thermal emission from the two stellar winds, and characterized with a spectral index of \( +0.6 \) \citep{Wright1975}. Consequently, the thermal contribution to the flux at lower frequencies can be derived by extrapolation and assuming no variations between 2003 (8.64 GHz data) and 2004 (17.8 GHz data). This leads to a thermal contribution at 8.64 GHz of 1.2 mJy, compared with a total flux of 2.0 mJy at 2.36 GHz (see Table 3), an \( f_\text{T} \approx 0.6 \), and \( M = 4.9 \times 10^{-5} M_\odot \text{yr}^{-1} \).

The main conclusion here is that mass-loss rate determinations are challenging when dealing with a binary system. Considering the high fraction of binaries among O-type stars (see e.g. \cite{Sana2014} and references therein), many observational determinations of \( \dot{M} \) should certainly be viewed with caution, including those of HD 93129A.

### 6.4. Flux density variability of HD 93129A

The flux density measurements of HD 93129A derived from ATCA data (Table 3) cover frequencies from 1.4 to 24.5 GHz. Observations at 2.4, 4.8 and 8.6 GHz were obtained twice, at different epochs. The two measurements at 2.4 GHz were observed in 2003 and 2008, and the fluxes agree within the uncertainties.

The fluxes at 4.8 GHz (from 2003 and 2009) differ by \( \sim 37\% \), increasing with time. Similarly, the 8.6 GHz fluxes at the same epochs differ by \( \sim 45\% \). In all the archival observations, relative errors are 10\% or less. We checked for calibration errors in the individual datasets. For all of them, the flux calibrator was 1934-638, from which the flux scale of phase calibrators and source were bootstrapped.

The data indicate that the flux increased by a significant fraction compared to the uncertainties between 2003 and 2009. Such a trend is indeed anticipated if the system is approaching periastron passage c.f. WR 140 \citep{Dougherty2005}.

A synchrotron component that dominates over the thermal emission from the stellar winds of the two stars can explain the measured flux densities. We have modeled the data with a power-law spectrum with a low-frequency exponential cutoff. We have computed two possible models. The first model fits the data from January 2003 - December 2003, and the 2004 data at the higher frequencies (Fig. 9: solid dots). It is assumed that the flux does not vary significantly within this range of time, given the separation of Aa to Ab in that period. These data are well-fit by a power-law, with a turnover around 1.4 GHz. However, the data from 2008 and 2009 are modeled by a similar power-law spectrum but with a turnover at a slightly higher frequency. An additional thermal component is required to fit the data at 20 GHz. The spectral indices from the fits to the 2003-4 and 2008-9 data are, respectively, \( -1.03 \pm 0.09 \) and \( -1.21 \pm 0.03 \).

Up to now, the results from observations at high-energy spectral bands remain inconclusive in detecting emission from a WCR. In soft X-rays (below 10 keV), the work by \cite{Cohen2011} concluded that the X-ray spectrum is dominated by emission from the individual stellar winds in the system. On the other hand, \cite{Gagne2011} reported a ratio between the X-ray and bolometric luminosities \((L_X/L_{bol})\) of \( 1.3 \times 10^{-7} \). This value suggests that only a small contribution of the X-rays come from the colliding-wind region. It is not unexpected to see colliding-wind binaries made of very early-type stars with no bright X-ray emission in excess of the canonical \( L_X/L_{bol} \) value of \( \sim 10^{-7} \). For very early-type stars with strong and dense winds such a the members of HD 93129A, the individual X-ray emission should be somewhat fainter than suggested by the canonical ratio \citep{Owocki2013,DeBecker2013}. A low emission measure of a WCR due to a large stellar separation on top of such an underluminous:

<table>
<thead>
<tr>
<th>Reference</th>
<th>Method</th>
<th>Value \times 10^3</th>
</tr>
</thead>
<tbody>
<tr>
<td>Tar97</td>
<td>H( \alpha ) profile</td>
<td>1.8</td>
</tr>
<tr>
<td>Tar97</td>
<td>Ultraviolet lines</td>
<td>2.08</td>
</tr>
<tr>
<td>Rep04</td>
<td>H( \alpha ) profile</td>
<td>2.36</td>
</tr>
<tr>
<td>BK04</td>
<td>8.64GHz continuum flux</td>
<td>(3/4) \times 7.2</td>
</tr>
</tbody>
</table>

This work: Separating T and NT fluxes 4.9††

**References.**
- \cite{Tar97}:
- \cite{Taresch1997}:
- \cite{Repolust2004}:
- \cite{BK04}:

---

\footnotetext[6]{For an overview of these mass-loss rate determination techniques, including a comprehensive discussion of their advantages and limitations, we refer to the review by \cite{Puls2008}.}
wind emission leads therefore to a not so high overall $L_x/L_{bol}$ ratio.

HD 93129A is rather close to the Fermi source 1FGL J1045.2–5942 (in the 0.1–100 GeV energy range). This source is associated with η Car, the only PACWB identified on the basis of γ-ray emission. It seems very unlikely that HD 93129A contributes to the emission detected by Fermi, as clarified by Abdo et al. (2010). At even higher energies, no TeV source has been found at the position of HD 93129A in the TevCat catalogue (http://tevcat.uchicago.edu/), even when considering newly announced sources and source candidates. We conclude that there is no known γ-ray source associated with HD 93129A. Observations with the forthcoming Cherenkov Telescope Array, with unprecedented sensitivity and source location techniques, will surely provide useful constraints to better understanding of the high energy phenomena in these objects.

7. Summary and conclusions

An LBA observation of the massive binary HD 93129A detected an extended and arguably curved radio emission component with a flux density of 2.9 mJy at 2.3 GHz between the two components Aa and Ab. Following a detailed analysis of recent high angular resolution astrometry of the system, we provide compelling evidence that the radio emission component is coincident with the expected position of a wind interaction region between components Aa and Ab, and suggest a wind-momentum ratio of $\sim 0.5$.

Historic ATCA observations of the HD 93129A system, from 1 to 25 GHz, re-reduced in a uniform way, show that the flux increased between 2003-4 and 2008-9, both epochs being well fit by a power-law spectrum with a steep spectral index. Similar increases in flux have been seen in other systems as they approach periastron. Thus, the results presented in this paper lend support to the idea that wind-collision regions are the sites where relativistic particles are accelerated in particle-accelerating colliding-wind binaries.

Very Long Baseline Interferometer (VLBI) observations of particle-accelerating colliding-wind binaries help to quantify specifically the properties of the non-thermal radio emission. Such measurements are crucial for models aimed at reproducing the particle acceleration and non-thermal physics at work in these objects. VLBI observations of the HD 93129A system across a wide frequency range will allow more accurate properties of the spectrum to be derived, and hence the properties of the relativistic electron populations involved in the synchrotron emission process, and potentially reveal associations between flux variations and changes in the properties of the wind-collision region.

In parallel, new ATCA observations, preferentially simultaneous across all the observing frequencies are necessary to obtain accurate total fluxes. Repeated observations in the leading up to the putative periastron passage will reveal if the flux increases as observed in other PACWB systems at this phase of the orbit.

Acknowledgements. The authors thank an anonymous referee for detailed comments, C. Phillips and A. Tzoumis for their work in a preliminary stage of the study, A. López-Sánchez for providing some of the ATCA observations, Rodolfo Barba, Roberto Gamen, Ed Fomalont, Jamie Stevens, Jesús Maíz-Apellániz, and very especially G. Vila and N. Casco for help with the orbits. P.B. is supported by the ANPCyT PICT-2012/00878 and UNLP G11/115 project. B.M. acknowledges support by the Spanish Ministerio de Economía y Competitividad (MINECO) under grants AYA2013-47447-C3-1-P and BES-2011-049886. This study was partially supported by NASA through grant GO-10898 from the Space Telescope Science Institute, which is operated by AURA, Inc., under NASA contract NAS 5-26555. The research has made use of the Washington Double Star Catalog maintained at the U.S. Naval Observatory.

References


Benaglia, P., Koribalski, B., & Albacete Colombo, J. F. 2006, PASA, 23, 50


De Becker, M. 2007, A&ARv, 14, 171

De Becker, M. 2013, NewA, 25, 7


Walborn, N. R. 1995, RMxAA, 2, 51