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The fracture studies of polycrystalline silicon based MEMS

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University of Liège Shantanu S. Mulay, G. Becker, L. Noels

Université Catholique de Louvain Renaud Vayrette, Jean-Pierre Raskin, Thomas Pardoen

> Université Libre de Bruxelles Montserrat Galceran, Stéphane Godet



Aerospace & Mechanical engineering



- Introduction
- Numerical fracture framework for polycrystalline silicon
 - Discontinuous Galerkin (DG) method
 - Hybrid DG/Extrinsic cohesive law (ECL)
 - Orthotropic plane-stress Hooke's law for core of grains
 - Intra-granular fracture
 - Thickness effect
 - Preliminary results
 - Observations

• Future work

- Characterize inter-granular strength
- Compare with experiments
- Apply to robust design

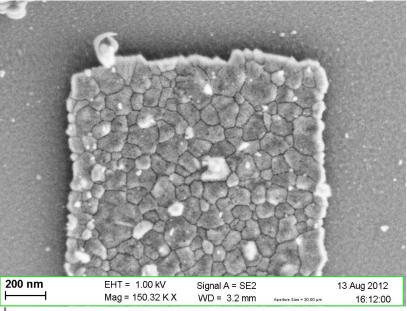




Introduction

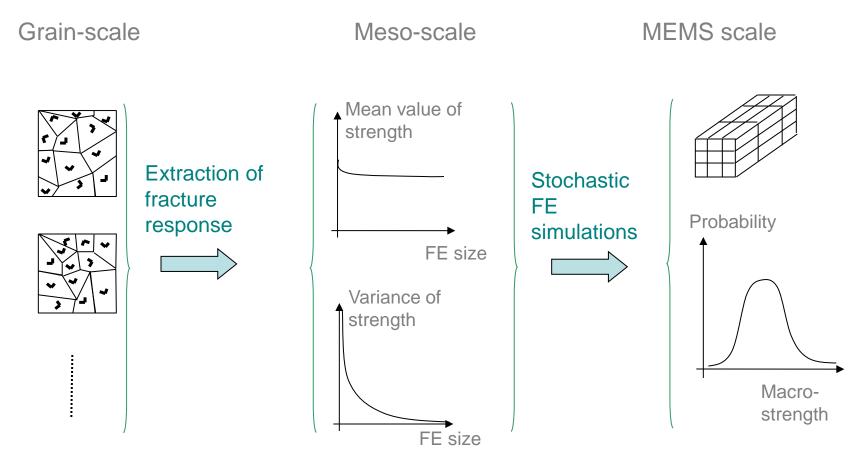
Purpose

- Develop a numerical method to predict MEMS fracture
- Difficulties
 - Grains sizes are no longer negligible compared to the structure size
 - Silicon is anisotropic
 - Inter/intra granular fractures
 - Dimensions are not perfectly controlled
 - Two MEMS will have
 - Different grains orientations/sizes
 - Different dimensions/surface profiles
- The numerical method should thus be probabilistic
 - But impossible to perform many direct numerical simulations with grain size resolutions





 Objective is to develop a robust design procedure of MEMS based on numerical stochastic 3-scale approaches



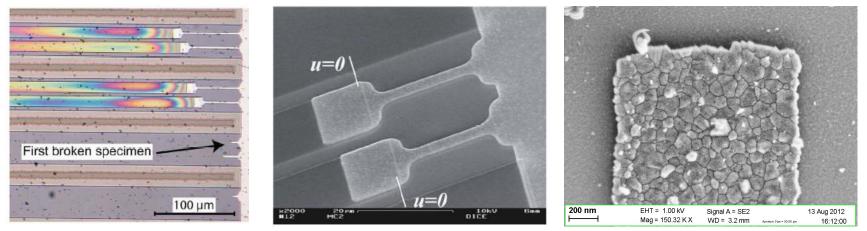
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Introduction

Methodology

- Develop a numerical fracture framework for polycrystalline structures (ULg)
- Validate tool with on-ship testing (UcL)



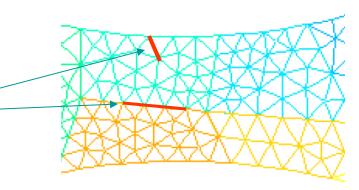
[[]Gravier et al., JMEMS 2009]

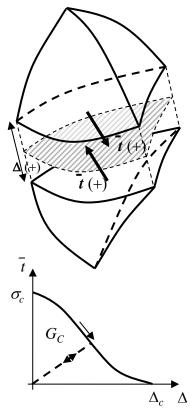
 Exploit numerical fracture framework in the 3-scale stochastic method (future work)



- Fracture challenges
 - Fracture can be
 - Inter-granular
 - Intra-granular -
 - Grains are anisotropic
 - Initially there is no crack
- Numerical approach
 - Cohesive elements inserted between two bulk elements
 - They integrate the cohesive Traction Separation Law
 - Characterized by
 - Strength σ_c &
 - Critical energy release rate G_C
 - Can be tailored for
 - Intra/inter granular failure
 - Different orientations

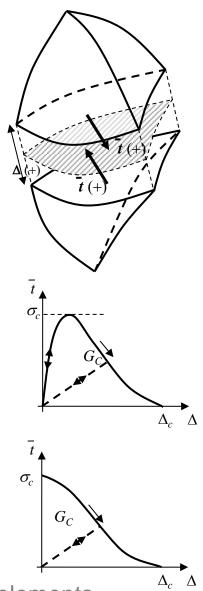






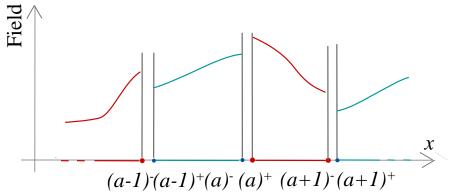


- Problems with cohesive elements
 - Intrinsic Cohesive Law (ICL)
 - Cohesive elements inserted from the beginning
 - Drawbacks:
 - Efficient if a priori knowledge of the crack path
 - Mesh dependency [Xu & Needelman, 1994]
 - Initial slope modifies the effective elastic modulus
 - This slope should tend to infinity [Klein et al. 2001]:
 - » Alteration of a wave propagation
 - » Critical time step is reduced
 - Extrinsic Cohesive Law (ECL)
 - Cohesive elements inserted on the fly when failure criterion is verified [Ortiz & Pandolfi 1999]
 - Drawback
 - Complex implementation in 3D (parallelization)
- Solution
 - Use discontinuous Galerkin method embedding interface elements



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- Discontinuous Galerkin (DG) methods
 - Finite-element discretization
 - Same discontinuous polynomial approximations for the
 - **Test** functions φ_h and
 - Trial functions $\delta \varphi$



- Definition of operators on the interface trace:
 - Jump operator: $\llbracket \bullet \rrbracket = \bullet^+ \bullet^-$
 - Mean operator: $\langle \bullet \rangle = \frac{\bullet^+ + \bullet^-}{2}$
- Continuity is weakly entorced, such that the method
 - Is consistent
 - Is stable
 - Has the optimal convergence rate

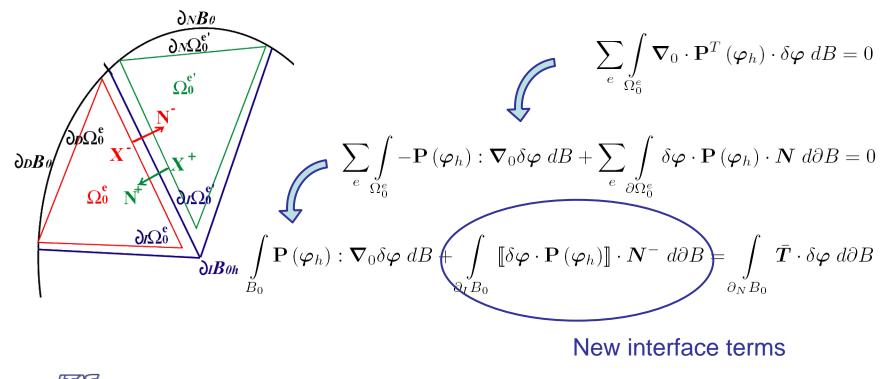




- Discontinuous Galerkin (DG) methods (2)
 - Formulation in terms of first Piola-Kirchhoff stress tensor P

$$\boldsymbol{\nabla}_{0} \cdot \mathbf{P}^{T} = 0 \text{ in } \boldsymbol{\Omega} \quad \boldsymbol{\&} \quad \left\{ \begin{array}{l} \mathbf{P} \cdot \boldsymbol{N} = \bar{\boldsymbol{T}} \text{ on } \partial_{N} \boldsymbol{\Omega} \\ \boldsymbol{\varphi}_{h} = \bar{\boldsymbol{\varphi}}_{h} \text{ on } \partial_{D} B \end{array} \right.$$

– Weak formulation obtained by integration by parts on each element Ω^e





- Discontinuous Galerkin (DG) methods (3)
 - Interface terms rewritten as the sum of 3 terms
 - Introduction of the numerical flux h

$$\int_{\partial_{I}B_{0}} \left[\!\left[\delta\varphi \cdot \mathbf{P}\left(\varphi_{h}\right)\right]\!\right] \cdot \mathbf{N}^{-} \, d\partial B \to \int_{\partial_{I}B_{0}} \left[\!\left[\delta\varphi\right]\!\right] \cdot h\left(\mathbf{P}^{+}, \mathbf{P}^{-}, \mathbf{N}^{-}\right) \, d\partial B$$

$$\overset{\bullet}{\rightarrow} \text{Has to be consistent:} \begin{cases} h\left(\mathbf{P}^{+}, \mathbf{P}^{-}, \mathbf{N}^{-}\right) = -h\left(\mathbf{P}^{-}, \mathbf{P}^{+}, \mathbf{N}^{+}\right) \\ h\left(\mathbf{P}_{\text{exact}}, \mathbf{P}_{\text{exact}}, \mathbf{N}^{-}\right) = \mathbf{P}_{\text{exact}} \cdot \mathbf{N}^{-} \end{cases}$$

$$\overset{\bullet}{\rightarrow} \text{One possible choice:} \quad h\left(\mathbf{P}^{+}, \mathbf{P}^{-}, \mathbf{N}^{-}\right) = \langle \mathbf{P} \rangle \cdot \mathbf{N}^{-}$$

- Weak enforcement of the compatibility

$$\int_{\partial_I B_0} \left[\!\!\left[\boldsymbol{\varphi}_h\right]\!\!\right] \cdot \left\langle \frac{\partial \mathbf{P}}{\partial \mathbf{F}} : \boldsymbol{\nabla}_0 \delta \boldsymbol{\varphi} \right\rangle \cdot \boldsymbol{N}^- \ d\partial B$$

- Stabilization controlled by parameter β , for all mesh sizes h^s

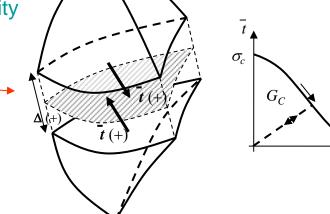
$$\int_{\partial_I B_0} \llbracket \boldsymbol{\varphi}_h \rrbracket \otimes \boldsymbol{N}^- : \left\langle \frac{\beta}{h^s} \frac{\partial \mathbf{P}}{\partial \mathbf{F}} \right\rangle : \llbracket \delta \boldsymbol{\varphi} \rrbracket \otimes \boldsymbol{N}^- \ d\partial B :$$

- Can also be explicitly derived from a variational form [Noels & Radovitzky, IJNME 2006 & JAM 2006]

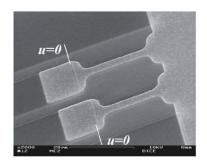


Hybrid DG/ECL

- Interface terms exist at the beginning
 - DG method ensures consitency/stability [Seagraves, Jerusalem, Radovitzky, Noels, CMAME 2012]



- Onset of fracture
 - When interface traction reaches σ_{c}
 - The cohesive law substitutes for the DG terms
- Advantages
 - Consistent
 - Easy to implement
 - Highly parallelizable
- In this work 2D plane-stress structures are studied







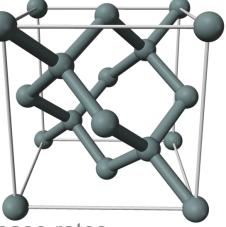
 $\Delta_c \quad \Delta$

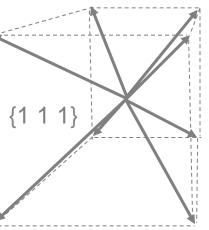
Silicon crystal

- Diamond-cubic crystal
- Has symmetry-equivalent surfaces
- Orthotropic material (at least two orthogonal planes of symmetry)
- Different fracture strengths and critical strain energy release rates along crystal lattice planes
 - 6 {1 0 0}-directions, 12 {1 1 0}-directions, 8 {1 1 1}-directions $\{1 \ 0 \ 0\}$ $\{1 \ 1 \ 1\}$ $\{1\ 1\ 0\}$

 $\sigma_{100} = 1.53 \text{ GPa}, \sigma_{110} = 1.21 \text{ GPa}, \sigma_{111} = 0.868 \text{ GPa}$





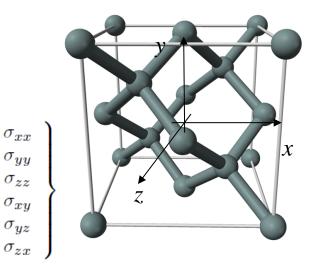




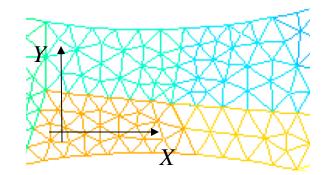
Bulk law

- In the referential (x, y, z) of the crystal
 - 9 constants (actually 3 ≠)

$$\left\{ \begin{array}{c} \epsilon_{xx} \\ \epsilon_{yy} \\ \epsilon_{zz} \\ \epsilon_{xy} \\ \epsilon_{yz} \\ \epsilon_{zx} \end{array} \right\} = \left[\begin{array}{ccccc} \frac{1}{E_x} & \frac{-\nu_{yx}}{E_y} & \frac{-\nu_{zx}}{E_z} & 0 & 0 & 0 \\ \frac{-\nu_{xy}}{E_x} & \frac{1}{E_y} & \frac{-\nu_{zy}}{E_z} & 0 & 0 & 0 \\ \frac{-\nu_{xz}}{E_x} & \frac{-\nu_{yz}}{E_y} & \frac{1}{E_z} & 0 & 0 & 0 \\ 0 & 0 & 0 & \frac{1}{2G_{xy}} & 0 & 0 \\ 0 & 0 & 0 & 0 & \frac{1}{2G_{yz}} & 0 \\ 0 & 0 & 0 & 0 & 0 & \frac{1}{2G_{zx}} \end{array} \right]$$



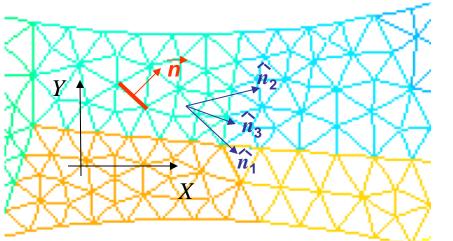
- Is rotated in the referential axes (X, Y, Z)
 - Different angles for different grains
 - Plane stress state $\sigma_{ZZ} = 0$

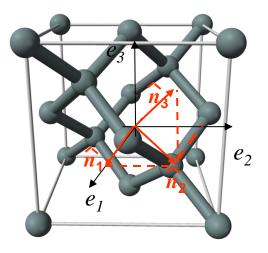


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- Intra-granular fracture
 - Different fracture strengths along crystal lattice planes
 - 6 {1 0 0}-directions \hat{n}_1 , 12 {1 1 0}-directions \hat{n}_2 , 8 {1 1 1}-directions \hat{n}_3
- - Mesh-interfaces are not along a fracture direction





- Assumption: FE mesh > silicon crystal cell size (5.43 Å)
 - Compute effective fracture strength on any required plane
- But: \hat{n}_1 , \hat{n}_2 & \hat{n}_3 do not form an orthonormal basis
 - Consider the dual basis \hat{n}^1 , \hat{n}^2 & \hat{n}^3

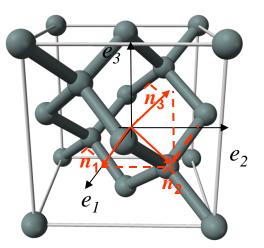




- Intra-granular fracture (2)
 - Surface normals of (1 0 0), (1 1 0), (1 1 1) known
 - \hat{n}_1 , \hat{n}_2 & \hat{n}_3 do not form an orthonormal basis
 - Consider the dual basis \hat{n}^1 , \hat{n}^2 & \hat{n}^3

$$\hat{n}_{1} = \hat{e}_{1}
\hat{n}_{2} = (1/\sqrt{2})(\hat{e}_{1} + \hat{e}_{2})
\hat{n}_{3} = (1/\sqrt{3})(\hat{e}_{1} + \hat{e}_{2} + \hat{e}_{3})$$

$$\hat{n}^{1} = \hat{e}_{1} - \hat{e}_{2}
\hat{n}^{2} = \sqrt{2}(\hat{e}_{2} - \hat{e}_{3})
\hat{n}^{3} = \sqrt{3} \hat{e}_{3}$$



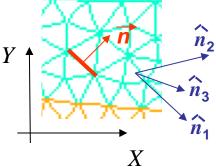
• Extract component of surface normal in the dual basis

$$\begin{cases} n^{100} = \vec{n} \cdot \hat{n}^1 \\ n^{110} = \vec{n} \cdot \hat{n}^2 \\ n^{111} = \vec{n} \cdot \hat{n}^3 \end{cases}$$

• Interpolate strength from strength along {1 0 0}, {1 1 0} and {1 1 1}

$$\vec{\sigma}_{eff} = \left[\sigma_{100} \ n^{100} + \frac{\sigma_{110} \ n^{110}}{\sqrt{2}} + \frac{\sigma_{111} \ n^{111}}{\sqrt{3}}\right]\hat{e}_1 + \left[\frac{\sigma_{110} \ n^{110}}{\sqrt{2}} + \frac{\sigma_{111} \ n^{111}}{\sqrt{3}}\right]\hat{e}_2 + \left[\frac{\sigma_{111} \ n^{111}}{\sqrt{3}}\right]\hat{e}_3$$





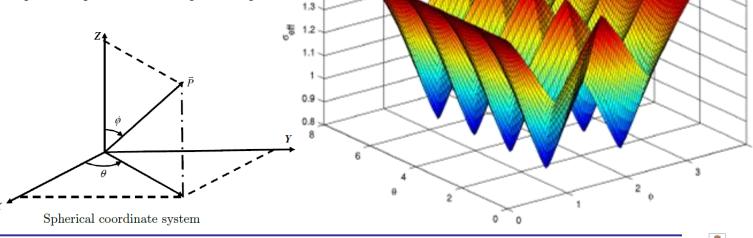
- Intra-granular fracture (3)
 - At the end of the day
 - $\sigma_{100} = 1.53 \text{ GPa}, \sigma_{110} = 1.21 \text{ GPa}, \sigma_{111} = 0.868 \text{ GPa}$

•
$$\|\vec{\sigma}_{eff}\| = \sqrt{\left(\sigma_{100} n^{100} + \frac{\sigma_{110} n^{110}}{\sqrt{2}} + \frac{\sigma_{111} n^{111}}{\sqrt{3}}\right)^2 + \left(\frac{\sigma_{110} n^{110}}{\sqrt{2}} + \frac{\sigma_{111} n^{111}}{\sqrt{3}}\right)^2 + \left(\frac{\sigma_{111} n^{111}}{\sqrt{3}}\right)^2}$$

1.5 1.4

- Applicable when surface normal is in-between solid angle formed by $\hat{n}_1, \hat{n}_2 \& \hat{n}_3$
- 48 solid angles are identified in

$$\theta \in [0, 360] \, \text{ and } \phi \, \in \, [0, 180]$$



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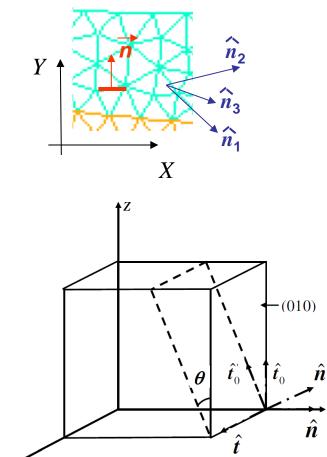


- Thickness effect
 - 2D-plane-stress model
 - Reality is 3D
 - Anisotropy
 - Weakest plane is not always the section
 - Find weakest plane passing through the interface edge
 - Iterate on θ
 - Compute new edge referential

$$\left\{ \begin{array}{c} \hat{n}' \\ \hat{t}'_0 \\ \hat{t}' \end{array} \right\} = \left[\begin{array}{c} \cos(\theta) & \sin(\theta) & 0 \\ -\sin(\theta) & \cos(\theta) & 0 \\ 0 & 0 & 1 \end{array} \right] \left\{ \begin{array}{c} \hat{n} \\ \hat{t}_0 \\ \hat{t} \end{array} \right\}$$

Compute normal and tangential stress
 in new referential

$$\mathbf{S}_{\text{nor}} = (\sigma \ \hat{n}) \cdot \hat{n'}, \tau = (\sigma \ \hat{n}) \cdot \hat{t'}$$
$$\tau_0 = (\sigma \ \hat{n}) \cdot \hat{t'}_0, \tau_{\text{resultant}} = \sqrt{(\tau)^2 + (\tau_0)^2}$$



Rotation of interface element along the thickness of MEMS



• Thickness effect (2)

- Find weakest plane passing through the interface edge (2)
 - Compute effective stress in the new referential

$$\mathbf{S}_{eff} = \begin{cases} \sqrt{(\mathbf{S}_{nor})^2 + (\boldsymbol{\beta})^{-2} (\boldsymbol{\tau}_{resultant})^2}, & \text{if } \mathbf{S}_{nor} \ge 0\\ \frac{1}{\boldsymbol{\beta}} \langle \langle |\boldsymbol{\tau}_{resultant}| - \boldsymbol{\mu}_c |\mathbf{S}_{nor}| \rangle \rangle, & \text{if } \mathbf{S}_{nor} < 0 \end{cases}$$

[Camacho & Ortiz, IJSS 1996]

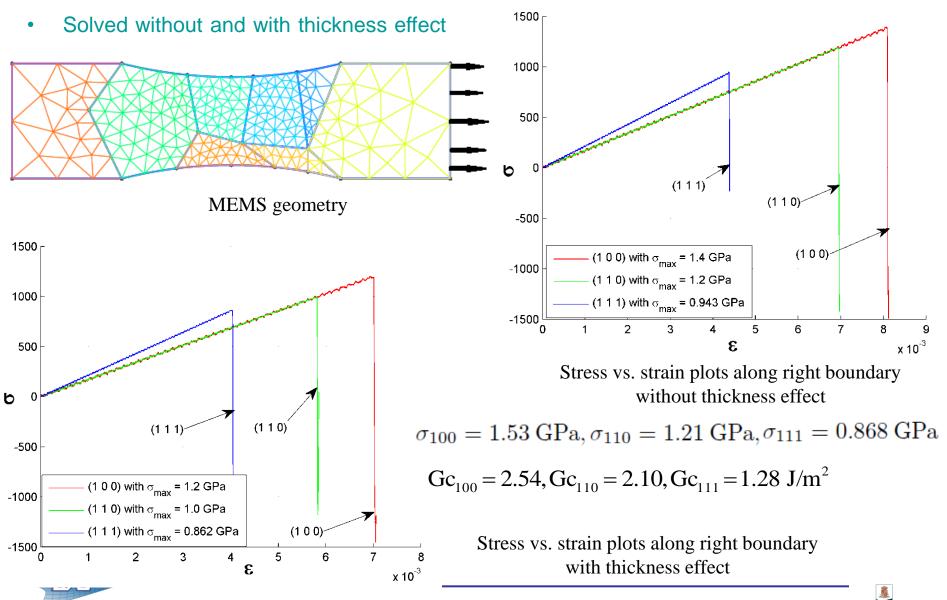
- Compare this value to effective fracture strength along \hat{n} ,
- Extrapolated as previously





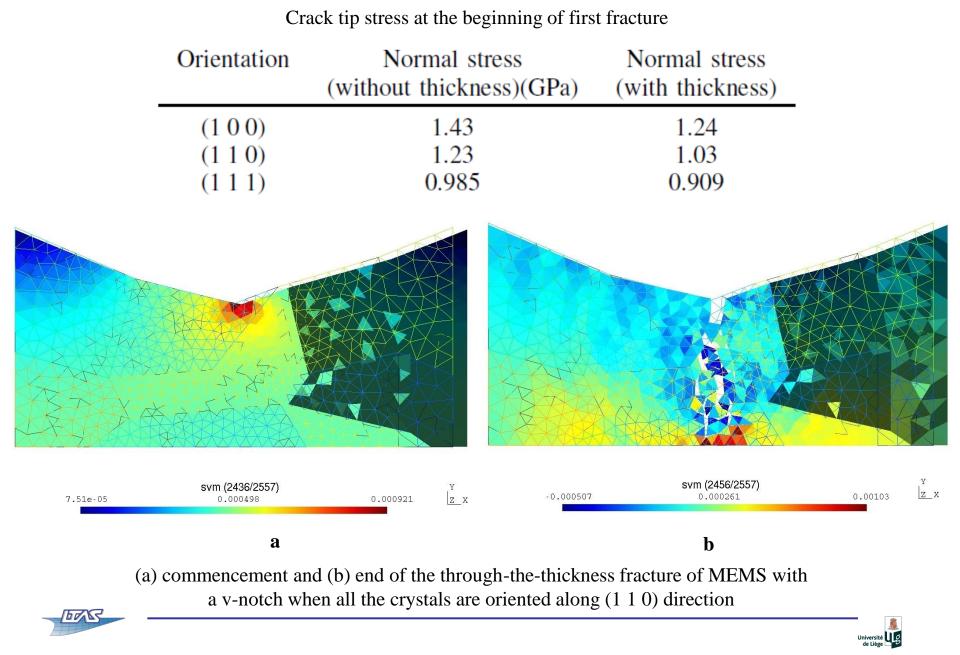
Preliminary results

MEMS modelled by 9 crystal lattices with 534 elements

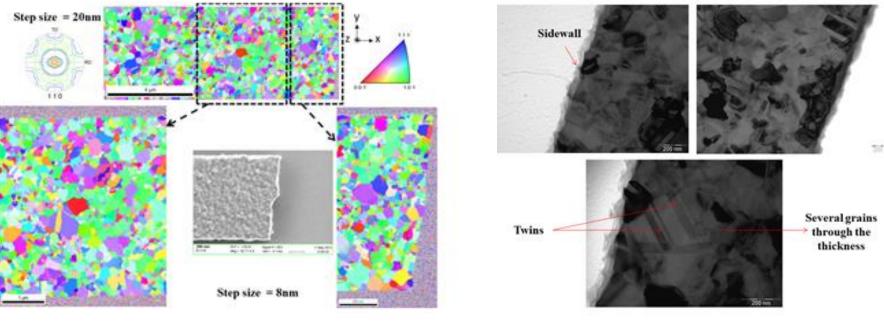


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Preliminary results (2)



- On-chip tensile microstructure fabricated to test MEMS for fracture
 - Extraction of Young's modulus and fracture strain by SEM and TEM
 - Automated crystallographic orientation mapping on transmission electron microscope (ACOM-TEM) technique to determine local orientations of grains



a

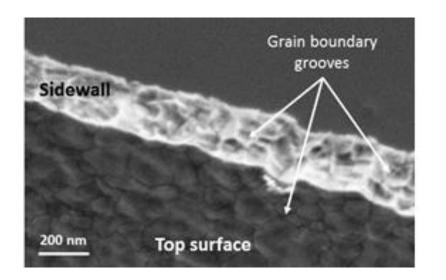
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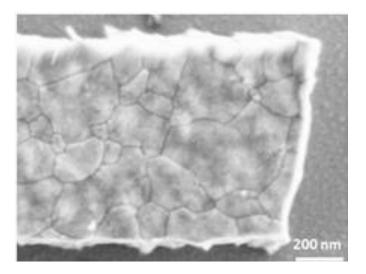
(a) Top view of the out-of-plane orientation map of 240 nm-thick polysilicon sample and (b) bright field TEM image of polysilicon sample



Experimental observations continued ...



a





(a) SEM image of the side wall of 240 nm-thick polysilicon sample and (b) SEM image of the top view of fracture zone of polysilicon sample

- SEM observation shows the presence of one or two grains along the thickness of sample
- Average local preferential orientation (1 1 0) in the out-of-plane direction and in-plane orientations are random
- Fracture initiated due to the flaws on sidewalls created during sample preparation





- Maximum stresses along the loading edge and crack tip are close to effective fracture strength
 - Validate the correctness of the computation of effective fracture strength
- Maximum fracture stress at crack tip is slightly lower with thickness effect as compared with without thickness effect
 - Maximum stress at fracture is either $S_{
 m nor}$ or au
 - First fracture is detected when $S_{eff} \ge \sigma_{(111)}$
 - Irrespective of the orientation of crystal lattices, there will be at least one interface plane orientated in the direction (1 1 1)
 - Verifies experimental observation that, independent of the orientation of crystal lattices, crack propagates in the direction (1 1 1)
 - $\sigma_{\max}(100) \ge \sigma_{\max}(110) \ge \sigma_{\max}(111)$, as (111) is weakest plane





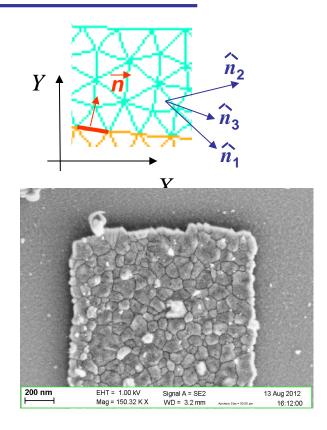
- Experimentally observed fracture strain 0.96% (+/- 0.07%) and fracture stress 1.41 (+/- 0.1) GPa
 - fracture stress in between the fracture strengths along (1 0 0) and (1 1 0) cleavage planes, as these planes influence in-plane fracture behaviour
- Numerically observed fracture strain 0.7% (+/- 0.1%) and fracture stress 1.1 (+/- 0.1) GPa
 - Fracture stress is slightly lower than experimentally observed value
 - Effective fracture strength is computed by weighted average values of fracture strengths along the (1 0 0), (1 1 0), and (1 1 1) orientations
 - Experimental sample has random in-plane orientations with higher influence of (1 0 0) and (1 1 0) orientations
- Transgranular crack path





Future work

- Inter-granular strength
 - Characterize strength
 - In terms of mis-orientations
- Compare with experiments
 - Grains orientations by automated crystal oriented mapping (ACOM)
 - Analysis of the competition between intergranular versus trans-granular crack path with respect to grain orientation

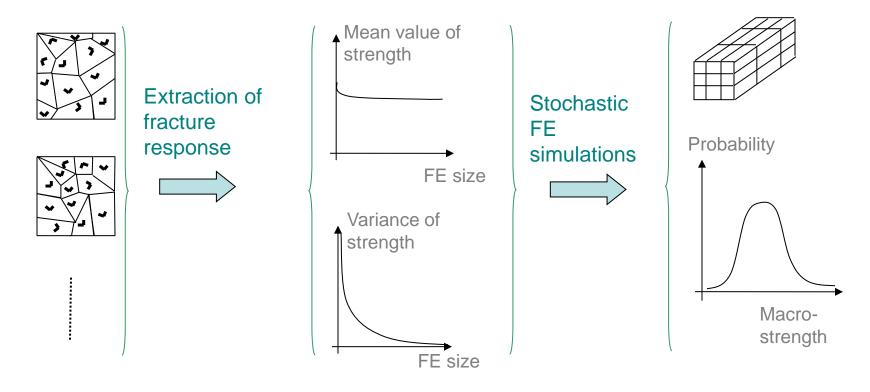


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Future work

- Robust-design
 - Statistical fracture strength at meso-scale from micro-scale simulations involving different grain sizes and grain orientations
 - Stochastic numerical method considering statistical distribution of fracture strength







Thank you



