

Concurrent observations of the ultraviolet nitric oxide  
and infrared O<sub>2</sub> nightglow emissions with Venus Express

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1 Abstract

2

3 Two prominent features of the Venus nightside airglow are the nitric oxide  $\delta$  and  $\gamma$  bands  
4 produced by radiative association of O and N atoms in the lower thermosphere and the O<sub>2</sub>  
5 infrared emission generated by three-body recombination of oxygen atoms in the upper  
6 mesosphere. The O<sub>2</sub> airglow has been observed from the ground, during the Cassini flyby and  
7 with VIRTIS on board Venus Express. It now appears that the global structure of the two  
8 emissions shows some similarities, but the statistical location of the region of strongest  
9 emission are not coincident. The SPICAV ultraviolet spectrograph has collected a large  
10 number of spectra of the Venus nitric oxide nightside airglow. VIRTIS spectral images have  
11 been obtained at the limb and in the nadir-viewing mode and have provided new information  
12 on the horizontal and vertical distribution of the emission. We present the first concurrent  
13 observations of the two emissions observed with Venus Express. We show that nadir  
14 observations generally indicate a low degree of correlation between the two emissions  
15 observed quasi-simultaneously at a common location. A statistical study of limb profiles  
16 indicates that the altitude and the brightness of the two airglow layers generally do not co-  
17 vary. We suggest that this lack of correlation is explained by the presence of strong horizontal  
18 winds in the mesosphere-thermosphere transition region. They carry the downflowing atoms  
19 over large distances in such a way that regions of enhanced NO emission generally do not  
20 coincide with zones of bright O<sub>2</sub> airglow.

21

21

## 22 I. INTRODUCTION

23 Planetary airglow is a powerful way to remotely probe the characteristics of  
24 atmospheres from an orbiting or flying-by spacecraft that complement in situ measurements.  
25 In particular, the study of airglow morphology, time variations and brightness may provide  
26 key observations concerning atmospheric composition, temperature structure, transport  
27 processes and their response to the solar photon and particle inputs. In the case of Venus, the  
28 Venera and Pioneer Venus missions have shed new light on some of these aspects more than  
29 twenty-five years ago. The presence of the delta and gamma bands of nitric oxide in the  
30 Venus nightglow was detected and identified by *Feldman et al.* [1979] using the ultraviolet  
31 spectrograph on board the International Ultraviolet Explorer (IUE). It was also observed by  
32 *Stewart and Barth* [1979] with the ultraviolet spectrometer on board the Pioneer Venus  
33 Orbiter (PV-OUVS). The emission is produced by radiative recombination through inverse  
34 predissociation of nitrogen  $N(^4S)$  and oxygen  $O(^3P)$  atoms and dominates the middle  
35 ultraviolet nightglow spectrum. In this process, excited NO molecules radiate the ultraviolet  $\delta$   
36 and  $\gamma$  bands between 180 and 310 nm:



38 giving rise to:



40 and



42 Emission from the  $C^2\Pi (v=0) \rightarrow A^2\Sigma (v=0)$  transition at 1.224  $\mu\text{m}$  which populates the  $A^2\Sigma$   
43 ( $v=0$ ) level was recently observed with VIRTIS-M in the Venus nightglow [*Garcia Muñoz et*

44 *al.*, 2009]. The N and O atoms are mainly produced by dissociation of N<sub>2</sub> and CO<sub>2</sub> on the  
45 dayside by extreme ultraviolet (EUV) photons and photoelectrons.

46 *Stewart et al.* [1980] obtained images of the Venus nightside in the  $\delta$  (0,1) band at 198  
47 nm every 24 hours with PV-OUVS when Pioneer-Venus was near apoapsis. They showed  
48 that the emission was highly variable in brightness and morphology over consecutive 24-  
49 hours periods. The location of the brightest spots ranged from 2130 to 0300 LT and 39° S to  
50 60°N [*Bougher et al.*, 1990]. This variability appeared to be caused by instabilities in the  
51 large-scale circulation, possibly as a result of wind shears near the terminator or time-varying  
52 wave drag from gravity waves. *Stewart et al.* [1980] built up a global map of the UV  
53 nightglow showing that the emission is concentrated in a bright spot located near 0200 local  
54 solar time, south of the equator. *Bougher et al.* [1990] estimated the emission rate of this  
55 enhanced emission to be ~1.9 kilo-Rayleighs (kR) whereas the average hemispheric nightside  
56 intensity is 0.48 kR. These observations confirmed the general picture where production of O  
57 and N atoms by solar EUV on the dayside is followed by global circulation to the nightside,  
58 downward transport and radiative recombination. The shift of the statistical bright spot toward  
59 dawn was interpreted as a signature of a residual super-rotation into the lower thermosphere.  
60 A determination of the altitude of the emission peak by *Gérard et al.* [1981] concluded that  
61 the emission peak is located close to 115 km. Using a one dimensional model, they derived an  
62 eddy diffusion coefficient  $K \sim 8 \times 10^{12} / n^{1/2} \text{ cm}^2 \text{ s}^{-1}$ , where  $n$  is the total number density. The  
63 general picture of production of O and N atoms followed by transport to the nightside by the  
64 subsolar to antisolar circulation, downward turbulent mixing and radiative recombination  
65 appeared quantitatively consistent with the PV-OUVS observations. This concept was  
66 numerically validated by three-dimensional simulations using the Venus Thermospheric  
67 General Circulation Model (VTGCM) [*Bougher et al.*, 1990]. The statistical location of the  
68 bright spot was reasonably well predicted by the three-dimensional model and implied zonal

69 winds of about 50-75 m s<sup>-1</sup> in the 115-150 km region. The observed shift toward dawn of the  
70 statistical location of the airglow maximum was reproduced by the VTGCM.

71 Limb observations of the spectrum of nightglow emission in the  $\delta$  and  $\gamma$  bands of NO  
72 with the SPICAV ultraviolet spectrometer on board Venus Express [*Titov et al.*, 2006;  
73 *Svedhem et al.*, 2007] have been recently reported by *Gérard et al.* [2008a]. The mean altitude  
74 of the emission layer was found to be located at 113 km, with variations between 95 and 132  
75 km. The mean limb brightness of the total NO emission at the limb was 32 kR, but it is highly  
76 variable with limb intensities as large as 440 kR and values below 5 kR at northern mid-  
77 latitudes. It was found that the mean altitude of the emission peak statistically drops with  
78 increasing latitude between 6° and 72° N. From model fits to observed profiles, they  
79 determined that the downward flux of N atoms at 130 km typically varies between  $1 \times 10^8$  to  
80  $4 \times 10^9$  atoms cm<sup>-2</sup> s<sup>-1</sup>. The eddy diffusion coefficient K deduced from comparisons to the  
81 observed limb profiles was significantly less than that determined from the observations made  
82 with the Pioneer Venus UV spectrometer at low latitudes during periods of high solar activity.

83 The oxygen airglow (0-0) emission at 1.27  $\mu\text{m}$  is the most intense non-thermal  
84 component in the Venus atmosphere. It belongs to the  $a^1\Delta_g - X^3\Sigma_g$  Atmospheric Infrared  
85 system and corresponds to an electric dipole forbidden transition with a radiative lifetime of  
86 about 70 min. [*Miller et al.*, 2001]. It was first discovered in ground-based observations of  
87 Venus by *Connes et al.* [1979] and subsequently imaged with ground-based telescopes [*Allen*  
88 *et al.*, 1992; *Crisp et al.*, 1996; *Lellouch et al.*, 1997; *Ohtsuki et al.*, 2008; *Bailey et al.*, 2008].  
89 It was measured from space during the Cassini flyby with a local maximum brightness of  
90 about 4 MR [*Drossart et al.*, 1993]. The oxygen IR nightglow appeared patchy, highly  
91 variable with asymmetries, often exhibiting multiple local maxima, with variations on  
92 timescales as short as 1 hour. A much weaker O<sub>2</sub> airglow was also detected in the Herzberg II  
93 visible wavelengths [*Krasnopolsky et al.*, 1986; *Bougher and Borucki*, 1994]. The 1.27  $\mu\text{m}$

94 emission is produced by recombination of oxygen atoms created by photodissociation of CO<sub>2</sub>  
95 and CO at thermospheric altitudes on the dayside. As previously explained, O atoms are  
96 transported to the nightside by the global thermospheric circulation. Three-body  
97 recombination of O atoms in the upper nightside mesosphere (95–110 km) leads to O<sub>2</sub>  
98 formation in excited states, followed by airglow emissions as the molecules relax to the X <sup>3</sup>Σ<sub>g</sub>  
99 ground state:

100



102 followed by



104

105 where O<sub>2</sub>\* indicates one of the excited states of the O<sub>2</sub> molecule and M is any neutral  
106 constituent. A fraction of the O<sub>2</sub> molecules, estimated to be about 7%, is formed directly in  
107 the a<sup>1</sup>Δ<sub>g</sub> metastable state. A substantial fraction of the upper excited states cascades into the  
108 <sup>1</sup>Δ state, so that the net efficiency of the production of this state in the three-body  
109 recombination may be close to 100% [Crisp *et al.*, 1996]. Below the emission peak, O<sub>2</sub> <sup>1</sup>Δ<sub>g</sub>  
110 molecules may be deactivated by collisions with CO<sub>2</sub>, causing non-radiative transitions to the  
111 O<sub>2</sub> ground state. The altitude of the peak of the 1.27 μm emission is thus controlled by the  
112 competition between vertical transport, recombination and quenching. The emission rate is  
113 related to the downward flux of oxygen atoms.

114 The VIRTIS-M O<sub>2</sub> airglow limb observations have been presented by *Drossart et al.*  
115 [2007a], *Gérard et al.* [2008b, 2009] and *Piccioni et al.* [2009a]. *Drossart et al.* [2007a]  
116 determined that the O<sub>2</sub> peak emission is located near 96 km, which is consistent with three-  
117 body recombination of oxygen atoms. *Gérard et al.* [2008b] found that limb profiles observed  
118 at northern mid-latitudes exhibit large intensity variations over short time periods. *Gérard et*

119 *al.* [2009] described further emission limb profiles extracted from the images. They  
120 determined the vertical distribution of O<sub>2</sub> (<sup>1</sup>Δ<sub>g</sub>) atoms using an Abel inversion of the radiance  
121 limb profiles. Assuming photochemical equilibrium for O<sub>2</sub> (<sup>1</sup>Δ), they used these density  
122 profiles combined with the CO<sub>2</sub> vertical distribution to determine the atomic oxygen density.  
123 *Piccioni et al.* [2009a] analyzed limb measurements from 42 orbits. They found that the peak  
124 altitude of the O<sub>2</sub> (<sup>1</sup>Δ<sub>g</sub>) volume emission rate is typically located between 95 and 100 km, with  
125 a mean value of 97.4 ± 2.5 km. The vertical profile is broader near the equator, with a full  
126 width at half maximum of 11 km, a factor 2 larger than at middle latitudes. They reported that  
127 a secondary peak is frequently observed between 103 and 105 km.

128 VIRTIS night-side observations from Venus Express have complemented ground-  
129 based observations at much higher spatial resolution. In addition, limb observations from an  
130 orbit around Venus have given unprecedented access to the vertical distribution of the airglow  
131 layer and provided key constraints on the models. *Drossart et al.* [2007a] confirmed that the  
132 O<sub>2</sub> nightglow exhibits a large spatial and temporal variability. Observations by VIRTIS in the  
133 nadir mode have been used to construct extensive maps of the Venus atmosphere in the O<sub>2</sub>  
134 emission band [*Gérard et al.*, 2008b, *Piccioni et al.*, 2009a]. In nadir viewing geometry, the  
135 contamination of the O<sub>2</sub> airglow by the thermal emission of the deeper atmosphere has to be  
136 subtracted to obtain clean O<sub>2</sub> airglow images. The mean value, integrated over the nightside  
137 of the southern hemisphere, is typically about 0.8 MR, which is in agreement with the early  
138 ground-based observations giving a mean brightness of 1.2 MR for the night side. *Hueso et*  
139 *al.* [2008] found that the airglow is highly inhomogeneous with the regions of highest  
140 intensity generally located at low latitudes near the midnight meridian. They showed that  
141 zonal velocity derived from the motion of airglow features is dominated by an intense  
142 prograde jet from dawn to midnight extending up to 22 hours in local time, with lower  
143 velocities and reversed sign from dusk. The brightest small-scale (~100 km) features appeared

144 correlated with locations of apparent convergence which may be a signature of compression  
145 and downwelling. *Piccioni et al.* [2009a] described the characteristics of the horizontal  
146 distribution of the airglow and showed that regions of high O<sub>2</sub> airglow intensity are associated  
147 with downwelling causing an increase of the infrared brightness temperature. A similar  
148 conclusion was reached by *Bailey et al.* [2008] who derived rotational temperatures in excess  
149 of the VIRA values in regions of enhanced O<sub>2</sub> <sup>1</sup>Δ emission rate. He associated these regions  
150 with conditions of larger downflow velocities where local temperature is increased by  
151 compressional heating. Similar and larger temperature enhancements were observed with  
152 SPICAV from measurements of UV CO<sub>2</sub> absorption measurements during stellar occultations  
153 by *Bertaux et al.* [2007].

154         The non-homogeneous, time-dependent distribution of the O<sub>2</sub> <sup>1</sup>Δ<sub>g</sub> nightglow indicates  
155 that the local downward flow of oxygen may differ substantially from the mean value, in  
156 response to variations in the efficiency of the global day-to-night transport, the focusing effect  
157 of the night-side subsidence, changing zonal wind speeds, eddy transport efficiency, and  
158 gravity wave breaking.

159         The NO and O<sub>2</sub> nightglows do not occur at the same altitude and thus provide  
160 information about different vertical levels: 95–105 km for O<sub>2</sub> airglow (*Piccioni et al.*, 2009a)  
161 and ~115 km for NO [*Gérard et al.*, 2008a]. In this study, we address the question of the co-  
162 variation of the two airglow layers. Earlier studies have established that the statistical location  
163 of the nightglow bright spots is not coincident. This result was unexpected since the two  
164 emissions are produced by recombination of atoms created on the dayside by  
165 photodissociation and transported to the nightside by the subsolar to antisolar global  
166 circulation. In this study, we take advantage of the unique opportunity offered by two  
167 instruments of the Venus Express mission (SPICAV and VIRTIS) to observe almost  
168 simultaneously the two emissions and to investigate if their characteristics co-vary in time and



169 space. The observations reported here were obtained both in nadir-viewing geometry, where  
170 horizontal variations of the emission rate can be mapped, and at the limb, where vertical  
171 variations are best investigated. We describe both types of observations and draw conclusions  
172 on the level of covariance we have observed and transport processes occurring in the Venus  
173 mesosphere-thermosphere transition region.

174

## 175 II. PARALLEL NADIR OBSERVATIONS

176

177 The Venus Express spacecraft moves along a quasi-polar eccentric orbit with a 24-hour  
178 period. The apocenter is located at 66,000 km, while the altitude of the pericenter (at 80° N)  
179 has varied between 250 km and 185 km. The orbit is fixed in the inertial space and therefore  
180 precesses at the rate of 1.6°/day. The precession of the orbital plane leads to a wide variety of  
181 configurations on the nightside as well as on the dayside. Several observation modes ('science  
182 cases') may be selected including nadir observations, star pointing for stellar occultations by  
183 Venus' atmosphere, fixed point tracking and limb observations [Titov *et al.*, 2006].

184 The SPICAV instrument and its performances were described by Bertaux *et al.* [2007].  
185 The ultraviolet spectrometer covers a spectral range extending from 118 nm to 320 nm  
186 including the NO  $\delta$  (C <sup>2</sup> $\Pi$  - X <sup>2</sup> $\Pi$ ) and  $\gamma$  (A <sup>2</sup> $\Sigma$  - X <sup>2</sup> $\Pi$ ) emission bands, the only spectral  
187 features with Lyman- $\alpha$  observed in the Venus nightglow [Gérard *et al.*, 2008a]. The detector  
188 is a 407x288 pixel CCD and the angular field of view of one pixel is equal to 0.7x0.7 arcmin.  
189 For reasons of telemetry limitations and because of the time needed to read all the lines of the  
190 CCD, only 5 adjacent zones of the CCD detector are usually read out. In these nadir  
191 observations, the width of each spatial bin is 32 pixel lines, corresponding to a field of view  
192 of 3.7°. These lines are seen through the large (500  $\mu$ m) slit, providing a spectral resolution of  
193 about 12 nm. The planetary area intercepted by the field of view depends on the location on

194 its orbits. The spacecraft altitude ranged between 7350 and 9050 km during the nadir  
195 observations reported here. The SPICAV CCD is read out every second, but the actual  
196 integration period of each spectrum is 640 ms. The non-uniform dark current and offset  
197 values are carefully subtracted in each individual spectrum, using similar observations  
198 performed with a null amplification. The absolute calibration obtained by observing well-  
199 known hot stars spectra is then applied to obtain nitric oxide emission rates in kR [*Bertaux et*  
200 *al.*, 2007].

201 Spectral images have been regularly obtained in nadir geometry with VIRTIS mostly  
202 from segments of the orbit near apocenter. The VIRTIS [*Drossart et al.*, 2007b, *Piccioni et*  
203 *al.*, 2009b] pixel size of 0.25 mrad gives a spatial resolution of 15 km on Venus from  
204 apocenter. For this study, we use the VIRTIS M-mode which provides spectral cubes between  
205 0.25 and 5  $\mu\text{m}$  at a spectral resolution  $R \sim 200$ . Each spectral channel is  $\sim 9.5$  nm wide in the  
206 region of the O<sub>2</sub> Atmospheric Infrared system emission. A spatial scan, covering a 64 mrad x  
207 64 mrad field of view is obtained using a scanning mirror. However, even from apocenter,  
208 only a fraction of the Venus disk is observed during a mirror scan of the instrument and a  
209 spacecraft re-pointing is needed to collect a more extended coverage. For each VIRTIS  
210 image, the thermal contribution from the lower atmosphere is subtracted from the total signal  
211 using the VIRTIS fluxes measured in the three adjacent channels centered on 1.27  $\mu\text{m}$ . The  
212 count rate is expressed in radiative flux units and MR using the measured instrumental  
213 calibration and the O<sub>2</sub> (<sup>1</sup> $\Delta_g$ ) relative line intensity for a temperature of 200 K. Airglow  
214 radiation emitted downward and subsequently backscattered by the underlying clouds is  
215 accounted for using the correction factor derived by *Crisp et al.* [1996].

216 The SPICAV and VIRTIS databases have been searched to identify periods when the  
217 fields of view of the nadir-viewing observations of both instruments overlapped over a  
218 significant time span. Table 1 lists the orbit numbers, times, and locations of these

219 occurrences. As an example, Figure 1 illustrates the spatial coverage in the Venus atmosphere  
220 of the SPICAV slit (in grey) and the VIRTIS images (in black) during orbit 243. As can be  
221 seen, a common region was observed by both instruments northward of  $2^\circ$  N and southward  
222 of  $8^\circ$  S. The shape of the VIRTIS image coverage is defined by the combination of the  
223 decreasing spacecraft altitude during the 688 s of the VIRTIS exposure and a reorientation of  
224 the spacecraft close to the equator. Similarly, the footprint of the SPICAV slit moved at a  
225 nearly constant longitude, with a small deviation at low latitudes. Other cases of parallel  
226 observations with the two instruments present a similar pattern of spatial coverage. Once the  
227 regions of observation overlap have been identified, the brightness information is extracted  
228 from the VIRTIS nadir images along the track of the SPICAV slit. Since a VIRTIS image is  
229 constructed by combining adjacent pixel lines corresponding to successive positions of the  
230 mirror, the observations are not exactly coincident in time. The time difference is usually on  
231 the order of a few minutes. In the particular case of orbit 243 which covers a wide range of  
232 latitudes, the maximum time delay between measurements of the intensity with the two  
233 instruments at any given location varies between 4 and 16 minutes.

234 Figure 2a shows the latitudinal distribution of the NO and O<sub>2</sub> (<sup>1</sup>Δ<sub>g</sub>) nadir emission  
235 rates for this orbit measured at 0035 LT between 0752 and 0806 UT on Dec. 20, 2006. The  
236 NO brightness is sampled once per second and a smoothing function over 10 s has been  
237 applied. The O<sub>2</sub> signal has been extracted by averaging the intensity in the processed VIRTIS  
238 image over an area corresponding to the projection of one of the SPICAV spatial bins, as was  
239 shown in Figure 1. The region of parallel observations extends from close to the equator up  
240 to nearly  $50^\circ$  N. In this case, two different regions are identified. The first one, southward of  
241  $\sim 25^\circ$  N presents three successive peaks in the NO intensity with brightness ranging from 1 to  
242 2.7 kR. The O<sub>2</sub> airglow latitudinal variation shows structural similarities with the NO  
243 variations up to about  $14^\circ$  N, with two peaks nearly coincident with the NO maxima. Further

244 north, no clear increase of the O<sub>2</sub> brightness corresponds to NO maximum at 17°N.  
245 Northward of the location, no clear correlation is observed between the two signals. In  
246 particular, no feature is observed in the NO intensity at 33°N where the O<sub>2</sub> airglow increase  
247 by over a factor two. In the low latitude region of the two common intensity peaks, the  
248 O<sub>2</sub>/NO airglow intensity ratio is on the order of 500. A similar case is illustrated by Figure 2b  
249 for orbit 592 where the two emissions show a different latitudinal trend equatorward of 24°N,  
250 followed by a nearly coincident maximum near 30° reaching 5.6 kR in the NO bands and 1.2  
251 MR at 1.27μm. The two distributions show little correlation poleward of ~35°N. In this  
252 example, the average O<sub>2</sub>/NO intensity ratio is again close to 500. An example of uncorrelated  
253 structures of the two airglow emissions is illustrated in Figure 2c which shows the latitudinal  
254 distribution measured during orbit 342 (March 29, 2007) in the pre-midnight sector (2314  
255 LT). In this case, no correlation is observed between 10° S and 11° S. The O<sub>2</sub> airglow  
256 exhibits a peak reaching 1 MR on the equator and a decrease on either side of this maximum.  
257 A secondary peak is observed at 5°N. The NO airglow presents an equatorial dip with larger  
258 intensities up to 3.2 kR at 9°N. In this example, the brightness ratio of the two emissions is on  
259 the order of 800. Figure 2d is another example showing no correlation between the NO and  
260 the O<sub>2</sub> emissions observed during orbit 459. It extends from 20°S to 38°N, close to local  
261 midnight (2343 LT). Following correction for the thermal emission component, the O<sub>2</sub> (<sup>1</sup>Δ)  
262 emission rate is very weak southward of 10°N and hardly distinguished from the noise level.  
263 It continuously increases toward middle northern latitudes and nearly reaches the 1.2 MR  
264 level at the end of the observation sequence. Interestingly, the NO airglow shows a bright  
265 maximum of 2.2 kR near 2° N, a region where the O<sub>2</sub> emission is very weak. Inversely, the O<sub>2</sub>  
266 airglow increases poleward of 30° N, which corresponds to a region where the NO airglow  
267 drops with increasing latitude.

268 Concurrent sequences of the two airglow emissions have been collected during 14  
269 Venus Express orbits. None of them shows latitudinal distribution of the two emissions  
270 which are correlated over the full observation sequence. Instead, the two features are either  
271 totally uncorrelated as in Figures 2c and 2d or they exhibit some correlation between the  
272 locations of the intensity peaks over a limited latitudinal extent as in Figures 2a and 2b. Table  
273 1 summarizes the dates, times and locations of the parallel airglow observations and the  
274 correlation coefficients derived from each orbital sequence. From the examples in Figure 2  
275 and Table 1, we conclude that no large-scale correlation is generally observed in the  
276 latitudinal distribution of the vertical emission rate of the NO and O<sub>2</sub> airglow. Some of the  
277 brightness enhancements are co-located over a restricted region, such as on orbit 343, but the  
278 two latitudinal distributions may also show quite a different morphology over regions  
279 exceeding 50 degrees of latitude. We now examine concurrent limb observations of the two  
280 emissions to verify if the same conclusion holds and increase the sample size of parallel  
281 SPICAV-VIRTIS observations.

### 282 III. PARALLEL LIMB OBSERVATIONS

283 The observations used for this study were collected in the grazing (tangential) limb  
284 mode, where the line of sight is at some angle to the orbital plane and moves in such a way to  
285 maximize the time spent in the atmosphere [*Titov et al., 2006; Bertaux et al., 2007*]. The  
286 SPICAV line of sight scans a range of altitudes, generally between 70 km and 400 km and  
287 each second a full UV spectrum is obtained. In this mode, the line of sight crosses the dark  
288 limb several times during the ascending portion of the VEX orbit. Therefore, SPICAV  
289 supplied several sets of two (one for ingress, one for egress) altitude scans of five altitude  
290 profiles at each orbit [*Gérard et al., 2008*]. During these limb observations, the bin parameter  
291 varies between 2 and 32. The apparent altitude of the emission peak and its brightness depend  
292 on the value of the field of view projected on the limb. As was discussed by *Gérard et al.*

293 [2008a], the SPICAV field of view projected on the limb intercepts a vertical region whose  
294 size depends on the spacecraft-limb distance, the orientation of the slit and the bin parameter.  
295 It varies from 3 to 27 km, with a mean value of 14.9 km. This effect is accounted for by  
296 smoothing by the field of view of an emission layer having a vertical Chapman profile  
297 integrated along the line of sight. Each Chapman profile is constrained to show the same  
298 topside scale height as the observation. SPICAV data points have been corrected for this  
299 smoothing effect by setting the peak intensity and altitude to the values they would have had  
300 if the limb profiles had been observed with a negligibly small field of view. The 0.25 mrad  
301 pixel size of the VIRTIS-M detector projected on Venus limb corresponds to a spatial  
302 resolution of 1.9 km for a spacecraft distance of 7500 km, a typical value for a VIRTIS  
303 observation at 40° N. Analysis of the spectral cubes at the limb has indicated that the  
304 contribution of thermal radiation from the lower atmosphere is very small in the vicinity of  
305 1.27  $\mu\text{m}$  for altitudes of the tangent point above  $\sim 85$  km and corrections are negligible above  
306 90 km [Piccioni *et al.*, 2009a].

307         The possible correlation of the altitude and peak intensity of the NO and O<sub>2</sub> airglows  
308 at the limb has been investigated on a statistical basis. Periods when both SPICAV and  
309 VIRTIS-M were observing the same limb region have been identified and concurrent limb  
310 profiles of the two airglows layers have been extracted. The methodology consists in first  
311 determining the limb profile of the NO airglow distribution during ingress and egress of the  
312 SPICAV line of sight in the lower thermosphere and upper mesosphere. Figure 3 illustrates  
313 the example of a profile measured for a line of sight ingress. In this sketch, the emission  
314 peaks of the two airglow layers are separated by a vertical distance taken to be equal to 15 km  
315 on the basis of the NO and O<sub>2</sub> limb statistics. Once the NO emission profile is constructed, the  
316 corresponding VIRTIS limb image is scanned and pixels with a tangent point altitude located  
317 at an altitude  $\Delta h$  below each successive SPICAV data points are extracted from the VIRTIS

318 cube in the 1.27  $\mu\text{m}$  channel. In this way, an  $\text{O}_2$  limb profile is extracted from the VIRTIS–M  
319 cube so that the effects of horizontal inhomogeneity of the airglow is minimized. The time  
320 separating the acquisition of a SPICAV data point and the underlying  $\text{O}_2$  intensity is, at most,  
321 of a few minutes. In this way, quasi-simultaneous profiles of the two emissions are obtained  
322 for every ingress or egress when the two instruments were simultaneously operating. Figure 4  
323 shows a example obtained for orbit 323 (March 10, 2007) between 0035 and 0047 UT at 2258  
324 LT for an airglow layer separation of  $\Delta h = 15$  km. The NO limb profile shows a peak of 21  
325 kR at 111 km. The  $\text{O}_2$  profile reaches 69 MR at 99.8 km with a fast intensity drop above the  
326 peak. Statistical error bars are indicated on the VIRTIS and SPICAV data points. A total of  
327 249 such parallel limb profiles have been obtained between January 17, 2007 and January 9,  
328 2008. For each profile, the altitude and brightness of the emission peaks are determined and  
329 included in the database. The results of this statistical study are summarized in Figure 5 and 6.

330 The altitudes of the simultaneously observed NO and  $\text{O}_2$  emission profiles are shown  
331 in Figure 5. The scatter plot indicates that the intensity of the two airglow layers is not  
332 correlated, as confirmed by the very low value of the correlation coefficient  $R = -0.13$ . The  
333 dashed line indicates equal altitudes for the NO and  $\text{O}_2$  emission peaks. This plot also  
334 confirms that the  $\text{O}_2$  emission peak is in most circumstances located below the NO layer.  
335 However, the distance separating the two emission peaks varies from nearly zero to as much  
336 as 28 km, with an average of 15 km. A detailed analysis of the four data points where the  $\text{O}_2$   
337 peak is above the NO peak has been looked at in detail. These are specific cases when the NO  
338 airglow intensity shows a considerable latitudinal gradient as evidenced by the different peak  
339 altitudes obtained with the different SPICAV spatial bins. Since the method we use to extract  
340 the  $\text{O}_2$  limb profiles from the VIRTIS images is such that the NO and  $\text{O}_2$  profiles do not  
341 exactly correspond to the same observed volume, it is possible that the actual  $\text{O}_2$  emission  
342 peaks are not really located above the NO emission. The star indicates that the average

343 altitude of the NO airglow peak for this dataset is 113 km and 96.4 km for  $O_2(^1\Delta)$ . The large  
344 scatter in the distance between the two airglows layers probably reflects the widely changing  
345 dynamical regime prevailing in the transition region between the upper mesosphere and the  
346 lower thermosphere. A similar result is obtained when comparing the brightness of the  
347 emission peaks in parallel observations shown in Figure 6. For clarity, the NO and  $O_2$   
348 brightness has been plotted on a logarithmic scale since they vary by over a wide range of  
349 values. As it was found for the peak altitude, this plot indicates that the limb brightness of the  
350 two emissions is not significantly correlated ( $R=0.29$ ). The  $O_2/NO$  intensity ratio varies by  
351 nearly three orders of magnitude from 55 to  $\sim 40000$ . The black square indicates that the  
352 average limb brightness of the NO and  $O_2$  is 45 kR and 35 MR respectively in this dataset, in  
353 good agreement with the statistical values of 32 kR for NO [Gérard *et al.*, 2008a] and 29 MR  
354 for  $O_2(^1\Delta)$  [Saglam *et al.*, 2009]. The sensitivity of the results to the value of the distance  $\Delta h$   
355 separating the two emission peaks has been tested by varying it by  $\pm 5$  km. No significant  
356 difference was found in the results. It is thus concluded that neither the altitude of the  
357 emissions peaks nor the peak intensities at a given location in the Venus nightside atmosphere  
358 are correlated. This result is in full agreement with the independent conclusion derived from  
359 the nadir observations reported in section 2. We now examine possible explanations for these  
360 differences in the next section.

#### 361 IV. DISCUSSION

362

363 The observations collected with the SPICAV spectrograph and the VIRTIS-M spectral  
364 imager provide evidence that the molecular oxygen and the nitric airglow emissions are only  
365 weakly correlated. This conclusion is based on three different sets of observations. First, the  
366 statistical location of the NO and the  $O_2(^1\Delta)$  regions of enhanced emission are not coincident.  
367 This was demonstrated by the difference between the  $O_2$  spot centered on the equator at



368 midnight [*Gérard et al.*, 2008b; *Piccioni et al.*, 2009a] and the NO maximum which is shifted  
369 by about two hours toward dawn and southward of the equator [*Stewart et al.*, 1980; *Bougher*  
370 *et al.*, 1990]. Second, nadir quasi-simultaneous observations of the two emissions reported in  
371 section 2 of this study demonstrate that the distribution of the intensity along latitudinal cuts  
372 exhibits significant differences even though the two emissions may show similarities over a  
373 limited range. Third, limb observations indicate that neither the brightness nor the altitude of  
374 the emission peak co-vary in the northern hemisphere. The first aspect requires additional  
375 further studies and modeling. The shift of the statistical region of bright NO emission from  
376 the antisolar point [*Stewart et al.*, *Bougher et al.*, 1990, *Bougher and Borucki*, 1994] is an  
377 indication that the dawnward superotation observed at the cloud level persists in the upper  
378 hemisphere in such a way that the subsidence region of the global thermospheric circulation is  
379 statistically displaced by ~2 hours. This result needs further SPICAV observations to confirm  
380 that the shift is still observed in a period of low solar activity conditions. If the difference  
381 between the locations of the NO and O<sub>2</sub> bright regions is still observed during the Venus  
382 Express era, the picture of the vertical wind structure in the upper mesosphere-lower  
383 thermosphere transition region has to be revised accordingly.

384 The decoupling between the characteristics of the two emissions at a given location and  
385 time raises a different question. At first glance, the results reported in this study contradict  
386 the concept of a global subsolar-to-antisolar circulation carrying the O and N atoms from their  
387 dayside source region to the nightside location where they recombine to produce the NO and  
388 O<sub>2</sub> (<sup>1</sup>Δ) airglow emissions. In this view, the region of subsidence of the two species would be  
389 nearly coincident. We note however that, at any given time, global airglow images indicate  
390 that the location of the regions of bright emissions may considerably vary as was shown for  
391 NO [*Stewart et al.*, 1980; *Bougher et al.*, 1990] and O<sub>2</sub> (<sup>1</sup>Δ) [*Hueso et al.*, 2008; *Piccioni et*  
392 *al.*, 2009a]. This important feature shows that the location of the subsidence of the SS-AS

393 circulation is very variable, presumably as a consequence of an intrinsic variability of the  
394 circulation (possibly caused by the drag due to gravity waves) and horizontal winds. As  
395 mentioned before, apparent wind velocities locally as large as about  $100 \text{ ms}^{-1}$  have been  
396 deduced by *Hueso et al.* [2008] from the displacement of  $\text{O}_2$  airglow features between  
397 successive VIRTIS-M images.

398 We now take a close look at the role of horizontal winds in the de-correlation between the  
399 two airglow emissions. The two airglow layers are separated in altitude typically by  $\sim 16 \text{ km}$ .  
400 Although this distance is limited, the time required for vertical down transport of atoms from  
401 the NO layer to the  $\text{O}_2$  emission peak is fairly long, leaving the possibility that a region richer  
402 in oxygen at 113 km has moved over considerable distances by the time it reaches the 97 km  
403 level. This situation is illustrated by the sketch in Figure 7 where an air parcel of enhanced  
404 NO emission (in white-blue) is carried over an horizontal distance  $\Delta L$  by the time it has  
405 reached the altitude of maximum  $\text{O}_2$  recombination  $\Delta h$  km below the NO layer. The effect of  
406 this horizontal transport in a situation of non-uniform downward flux above the NO airglow  
407 layer is potentially able to explain the observed and calculated wind and vertical transport  
408 velocities. We first need to estimate the downward velocity of an air parcel to travelling the  
409 distance  $\Delta h$  separating the two airglow layers. The vertical transport velocity  $w$  has not been  
410 directly measured but may be estimated using models and circumstantial evidence. We  
411 examine three possible ways to derive typical vertical velocity values. First, in the one-  
412 dimensional chemical transport model used by *Cox et al.* [2008] and *Gérard et al.* [2008a],  
413 vertical transport below the homopause is parameterized by an eddy diffusion coefficient  $K$ .  
414 We adopt a standard value of  $\Delta Z = 16 \text{ km}$  for the distance separating the average altitudes of  
415 the NO and  $\text{O}_2$  emission peaks. Based on numerical simulations using the expression  $K =$   
416  $A/n^{1/2}$ , we find a velocity on the order of  $3 \text{ cm s}^{-1}$ . A second approach to estimate  $w$  is based  
417 on results from the three-dimensional model by *Bougher et al.* [1990] where eddy diffusion is

418 relatively small. In a recent version of this model, the vertical advection velocity between 115  
419 and 95 km is on the order of 15-20 cm s<sup>-1</sup> near the antisolar point (Bougher, private  
420 communication, 2009). Finally, an estimate of the downflow velocity was also given by  
421 *Bailey et al.* [2008] for a nightside region with a 20 K temperature enhancements observed  
422 over a two-day period. He obtained a vertical flow velocity of ~20 cm s<sup>-1</sup> for such a region of  
423 airglow brightening associated with enhanced vertical transport. The corresponding transit  
424 time for atoms flowing from the NO to the O<sub>2</sub> airglow peak ranges between 22 hours and 6  
425 Earth days.

426 The horizontal wind velocity  $v$  at the level of the O<sub>2</sub> airglow layer is very variable as was  
427 recently summarized by *Lellouch et al.* [2008]. Winds velocities, derived from CO millimeter  
428 observations, are typically on the order of 30-50 m s<sup>-1</sup> at ~93 km and 90-120 m s<sup>-1</sup> near 102  
429 km. *Hueso et al.* [2008] derived values of the effective wind velocity at 97 km of a few tens  
430 of m/s from the motions of bright spots of O<sub>2</sub> airglow. Horizontal winds calculated with the 3-  
431 D model values are on the same order in the region separating the two emission layers. A  
432 crude estimate of the horizontal distance  $\Delta L$  crossed by moving air parcel is then given by  $\Delta L$   
433  $= v/w \Delta Z$ . Adopting a value of  $v = 100 \text{ m s}^{-1}$  as an upper limit, we find values on the order of  
434 ~53000 km if vertical transport is parameterized by eddy diffusion or ~8000 km if the  
435 estimate for  $w$  from Bailey et al. adopted. Using  $v = 10 \text{ m s}^{-1}$  as a wind velocity probably  
436 closer to the average, we estimate a typical horizontal transport range of ~ 800 to 5300 to km.  
437 Actually, the situation is further complicated by the occurrence of chemical reactions and  
438 collisional quenching which limit the chemical lifetimes of the ground-state O atoms and the  
439 excited O<sub>2</sub>(<sup>1</sup>Δ) molecules respectively. In any case, it appears that the downward flow is  
440 much slower than the horizontal transport and atoms may travel considerable horizontal  
441 distances during their transit from 113 to 97 km. It is therefore a direct explanation of the  
442 lack of correlation between the two emissions observed concurrently by SPICAV and

443 VIRTIS-M. The very fact that the two emissions exhibit significant differences in their  
444 horizontal distribution strongly argues for the presence of strong horizontal winds in the  
445 thermosphere-mesosphere transition region. Other factors such as the different altitude where  
446 the N and the O atoms are formed on the dayside may play an additional role in the lack of  
447 correlation since the streamlines followed during their transport to the nightside are slightly  
448 different.

449         On the basis of these arguments, it may now appear more difficult to understand the  
450 co-variation occasionally observed in latitudinal cuts such as in Figure 2a and 2c. We  
451 speculate that the two emissions may co-vary at least under two particular conditions. First,  
452 when the distance between the two airglow layers are less than the average separation of 16  
453 km. The existence of such conditions is testified by Figure 6 which shows the existence of  
454 cases when this distance may be reduced to only a few kilometers. Under such circumstances,  
455 the vertical downflow time and the horizontal travelled distance may also be reduced  
456 accordingly. Second, maps of the horizontal displacements of bright O<sub>2</sub> features have  
457 indicated that regions with much smaller horizontal velocities, sometimes quasi-null values,  
458 have been observed near 97 km. It is possible that these horizontal stagnation regions persist  
459 over a sufficiently long period of time, allowing a more vertical downflow of the oxygen  
460 atoms. Finally, local enhancements in the brightness of the O<sub>2</sub> airglow suggest the presence of  
461 strong downflows, possibly associated with an increased vigor of local downward turbulent  
462 transport. These conditions may well explain that the two emissions can exhibit some degree  
463 of co-variation, which is otherwise absent in areas of strong horizontal winds and/or weaker  
464 vertical velocity.

465 V. CONCLUSION

466 The  $O_2(^1\Delta)$  and the NO  $\delta$  and  $\gamma$  bands night airglows of Venus have been observed  
467 concurrently for the first time by two instruments on board the Venus Express spacecraft.  
468 These observations include both nadir and limb viewing geometries. During nadir  
469 observations, the  $O_2(^1\Delta)$  emission intensity has been extracted from VIRTIS images  
470 observing the same locations as the footprint of the SPICAV slit in the Venus atmosphere on  
471 the same orbit. Occasional positive correlations between the latitudinal distributions of the  
472 two emissions have been observed over a limited latitudinal range. However, the overall co-  
473 variability of the two airglows is low, as shown by some of the results illustrated in this study  
474 and by the globally low correlation coefficients obtained between the two sets of  
475 observational sequences. A similar conclusion is reached from the statistical comparison of  
476 the altitude and peak brightness of the two emissions simultaneously observed in the grazing  
477 limb geometry. The very low correlation coefficients of both the altitude and the intensity of  
478 the two airglow layers indicate that the transport of O and N atoms considerably deviates  
479 from a steady state vertical flow, even in the region close to the anti-solar point. We suggest  
480 that the airglow de-correlation frequently observed is a consequence of the transport of the  
481 downward moving air mass by strong horizontal winds in the transition region that  
482 dynamically decouples the airglow in the two layers in a given vertical column. Our simple  
483 calculation of horizontal transport indicates that the atoms may travel considerable distances  
484 during the transit time of their vertical transport by vertical advection and eddy diffusion.  
485 Occasional correlations of the latitudinal intensity variations may be associated with  
486 situations when the distance between airglow layers is small or when the ratio of the  
487 horizontal and vertical transport velocity components is reduced. Simulations with a two- or  
488 three-dimensional chemical-transport model are needed to assess this scenario and quantify  
489 this effect.

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593 TABLE 1: time, location, maximum brightness and correlation coefficient R of the nitric  
 594 oxide and O<sub>2</sub> (<sup>1</sup>Δ) airglow concurrent observations by SPICAV and VIRTIS-M at nadir

595

<b>Orbit</b>	<b>Start time (UT)</b>	<b>Stop time (UT)</b>	<b>Local time</b>	<b>Min latitude</b>	<b>Max latitude</b>	<b>Imax NO (kR)</b>	<b>Imax O<sub>2</sub> (MR)</b>	<b>R</b>
243	2006.12.20 07:52:46	2006.12.20 08:06:29	0035	1.8°	48.9°	2.7	1.1	0.55
341	2007.03.28 05:49:01	2007.03.28 05:58:43	2307	-9,3°	-5.1°	1.0	0,6	-0.84
342	2007.03.29 05:47:05	2007.03.29 06:06:39	2314	-10.2°	10.9°	3.1	1.1	-0.41
343	2007.03.30 05:45:17	2007.03.30 05:57:22	2318	-10.9°	-0.4°	1.0	2.2	0.74
345	2007.04.01 05:42:15	2007.04.01 05:56:04	2330	-11.7°	-2.3°	2.1	0.9	0.23
346	2007.04.02 05:40:55	2007.04.02 06:05:27	2336	-12.0°	14.1°	3.7	3.5	0.32
453	2007.07.18 07:10:30	2007.07.18 07:35:42	2306	-16.4°	38.1°	1.1	0.6	-0.13
459	2007.07.24 06:50:06	2007.07.24 07:18:12	2343	-18.8°	37.5°	2.2	0.9	-0.08
567	2007.11.09 00:55:00	2007.11.09 01:23:40	2312	-10.9°	31.6°	3.5	1.1	0.86
571	2007.11.13 00:57:30	2007.11.13 01:21:04	0010	16.4°	19.0°	2.5	1.1	-0.59
592	2007.12.04 02:13:35	2007.12.04 02:31:45	0156	13.4°	48.3°	6.2	1.4	0.25
901	2008.10.08 08:11:21	2008.10.08 08:30:11	2301	-13.1°	20.2°	1.7	0.8	0.44
905	2008.10.12 08:19:58	2008.10.12 08:38:48	2328	-12.8°	20.2°	2.5	0.7	0.49
907	2008.10.14 08:24:18	2008.10.14 08:43:08	2339	-12.7°	20.4°	3.9	0.8	0.26

596

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598

599

600 Figure captions

601

602 Figure 1: spatial coverage of SPICAV (in grey) and VIRTIS-M (in black) nadir observations  
603 of the Venus nitric oxide and O<sub>2</sub> (<sup>1</sup>Δ) airglow during Venus Express orbit 243.

604

605 Figure 2: examples of concurrent observations of the latitudinal distribution of nightglow  
606 intensities at nadir by VIRTIS-M and SPICAV as a function of latitude: (a) orbit 243, (b)  
607 orbit 592, (c) orbit 342, (d) orbit 459. Note the different brightness scales used for the NO and  
608 the O<sub>2</sub> airglow emission rates. The brightness at 1.27 μm has been corrected for backscattered  
609 emission and both emission rates have been corrected for the view angle.

610

611 Figure 3: sketch illustrating the methodology to generate the data point for the study of the O<sub>2</sub>  
612 and NO concurrent limb observations shown in Figures 5 and 6. The projection of the  
613 SPICAV field of view in the atmosphere is represented by the yellow dotted line. The green  
614 dotted line shows the parallel line shifted downward by a distance Δh along which the values  
615 of the O<sub>2</sub> intensity is extracted to trace emission limb profiles such as illustrated in Figure 4.

616

617 Figure 4: example of concurrent limb profiles of the nitric oxide and O<sub>2</sub> infrared airglow  
618 measured during Venus Express orbit 323.

619

620 Figure 5: altitude distribution of the emission peaks of the NO and O<sub>2</sub> (<sup>1</sup>Δ) airglow layers  
621 measured quasi-simultaneously at the limb. The observations corresponding to each point  
622 have been collected as illustrated in Figures 3 and 4 (see text). The dashed line indicates  
623 equal altitudes for the two airglow emission peaks. The correlation coefficient is -0.05,

624 indicating the lack of co-variation of the altitude of the two emissions. The full square  
625 indicates the mean value of the peak altitude of the two emissions.

626

627 Figure 6: brightness distribution of the emission peaks of the NO and O<sub>2</sub> (<sup>1</sup>Δ) airglow layers  
628 measured quasi-simultaneously at the limb. The observations corresponding to each point  
629 have been determined as illustrated in Figures 3 and 4 (see text). The correlation coefficient is  
630 0.29, indicating the lack of co-variation of the brightness of the two emissions. The full  
631 square indicates the mean value of the peak intensity of the two emissions.

632

633 Figure 7: sketch illustrating the role of horizontal wind in the mesosphere-thermosphere  
634 transition region as a source of spatial de-correlation between a bright spot of NO airglow and  
635 a region of enhanced O<sub>2</sub> (<sup>1</sup>Δ) nightglow. The region initially enriched in oxygen atoms has  
636 traveled a horizontal distance ΔL by the time the blob of O-rich gas reaches the altitude of the  
637 O<sub>2</sub> nightglow layer located Δz km below the O<sub>2</sub> emission peak.

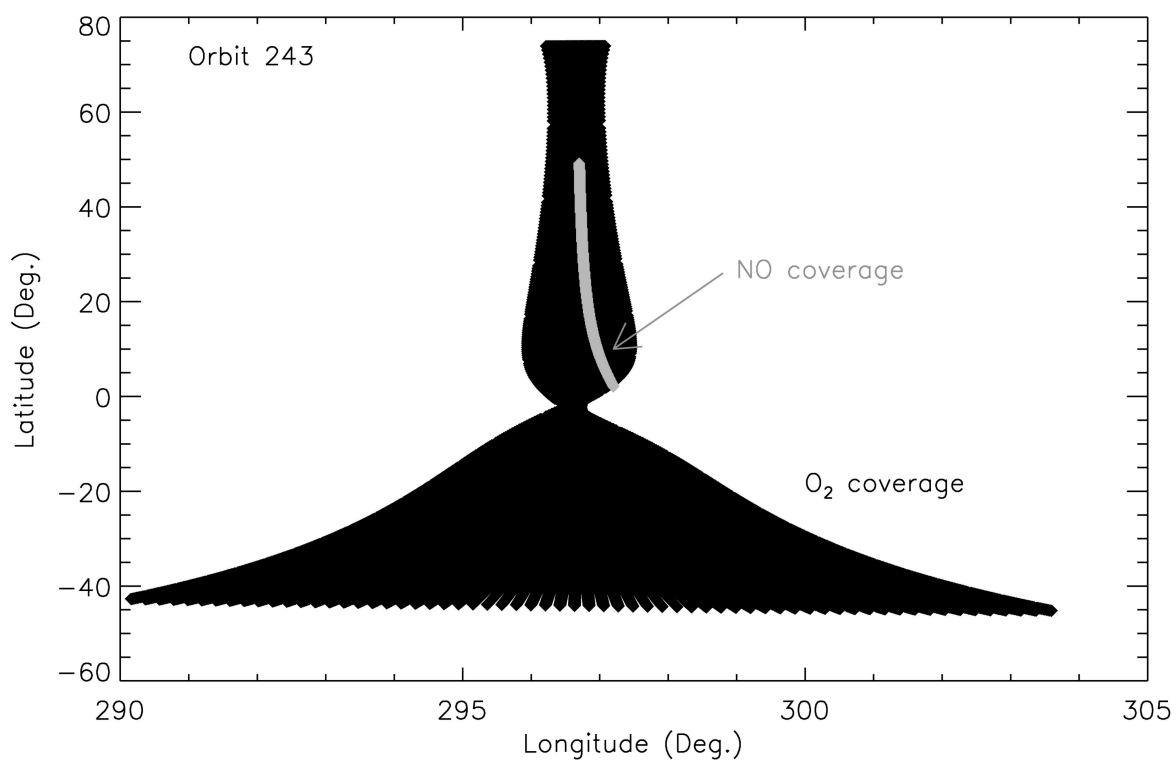


Figure 1

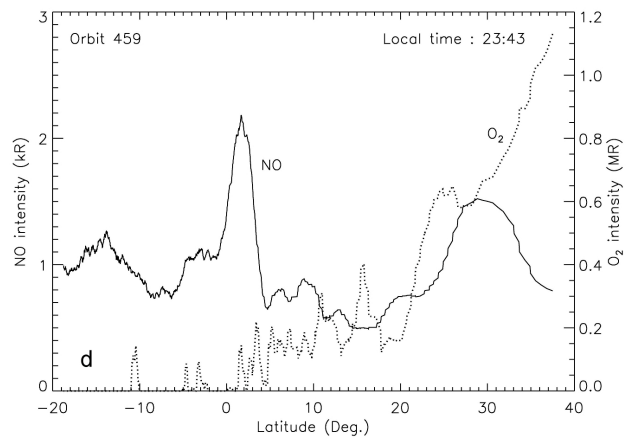
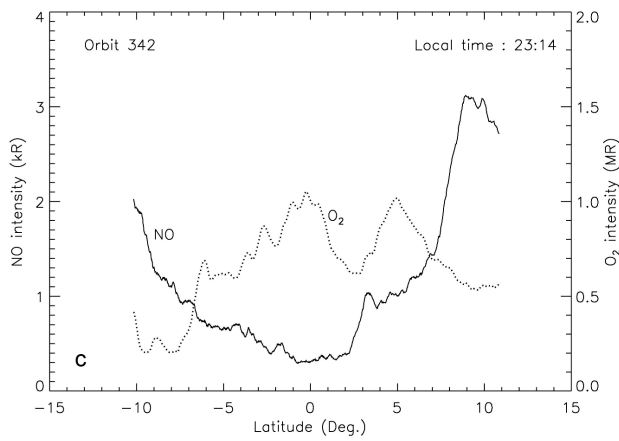
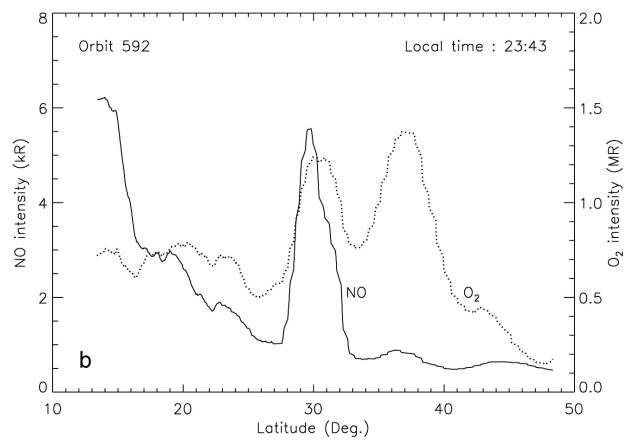
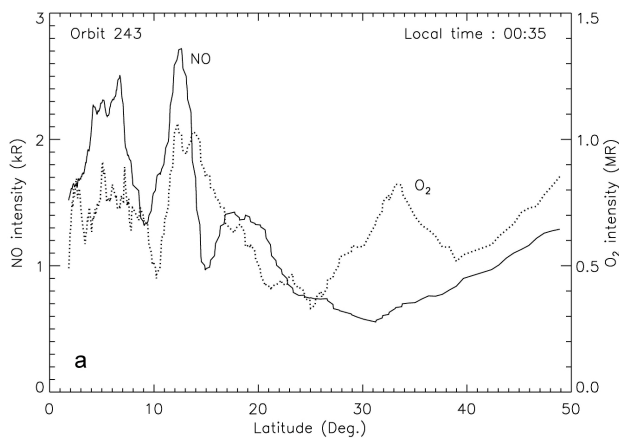


Figure 2

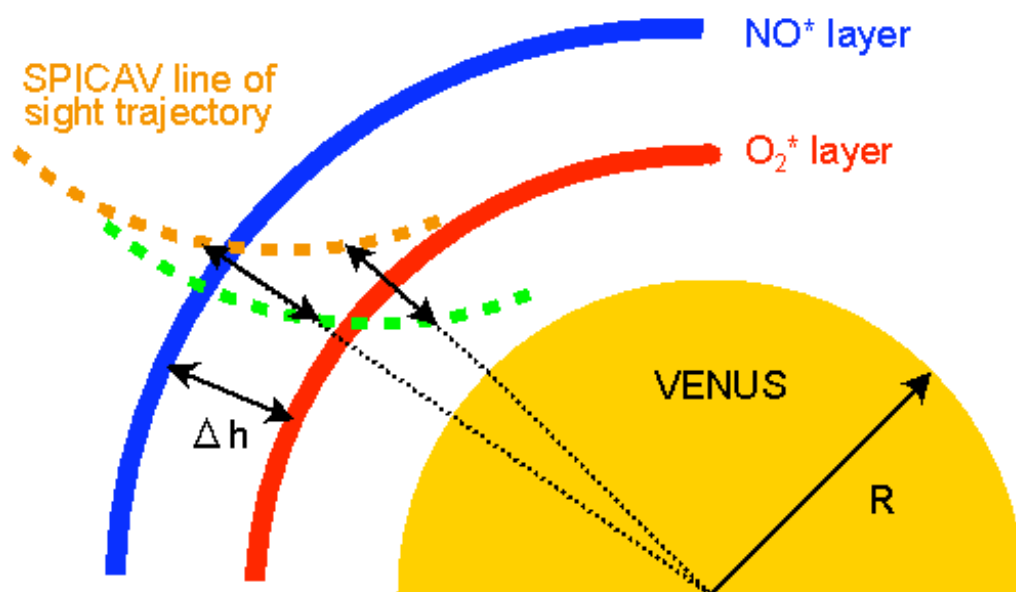


Figure 3



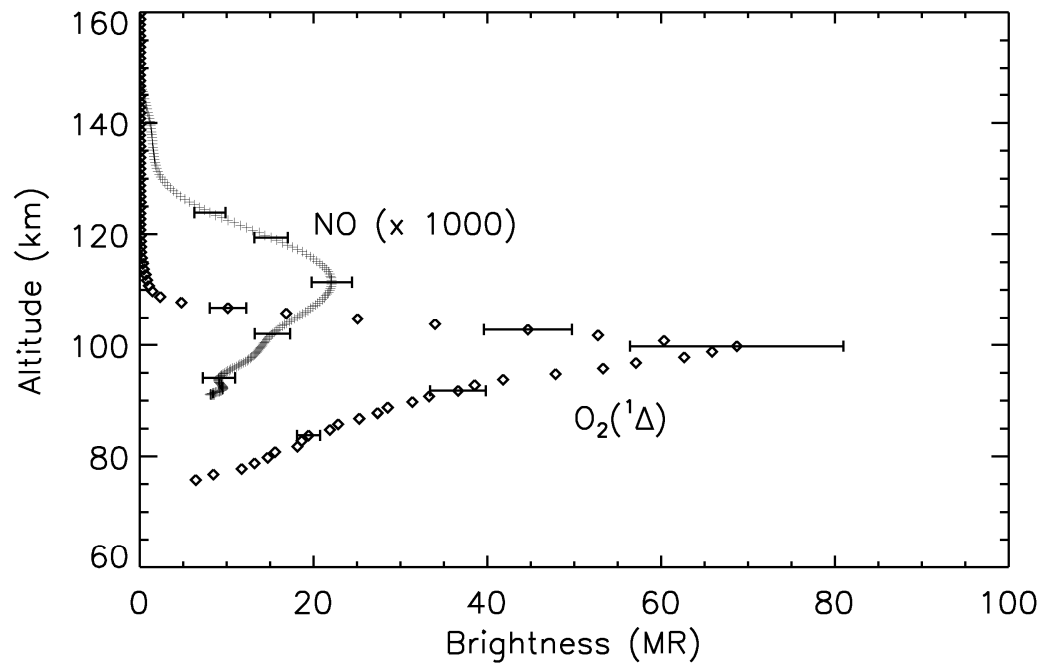


Figure 4

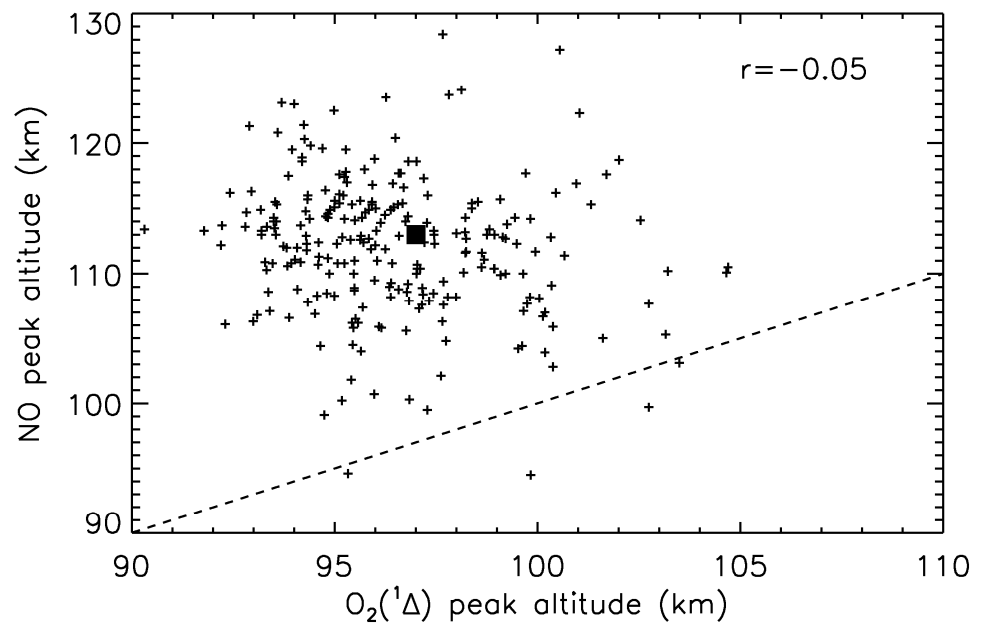


Figure 5

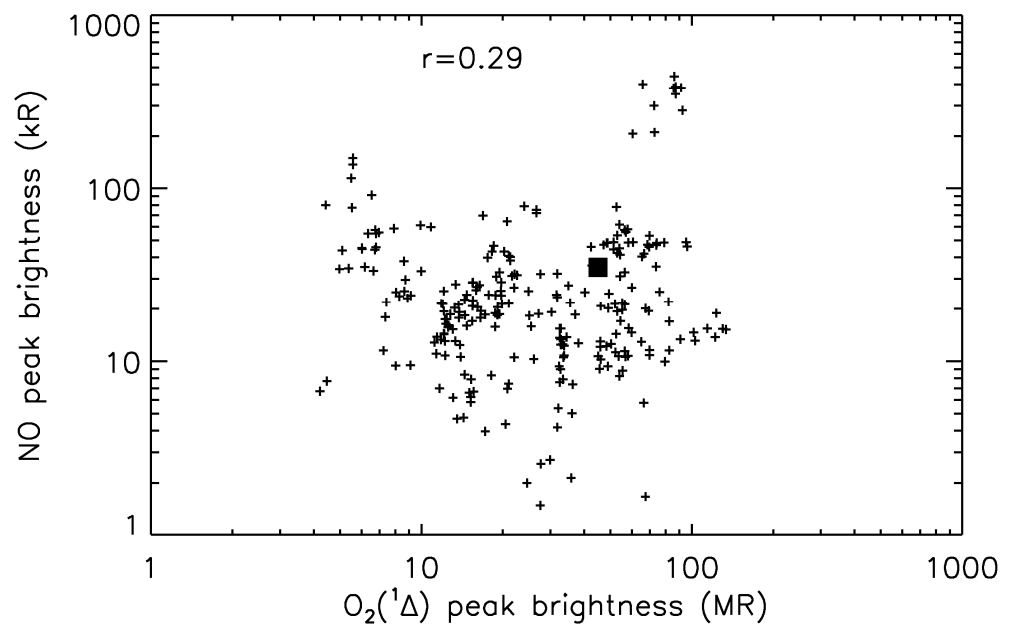


Figure 6

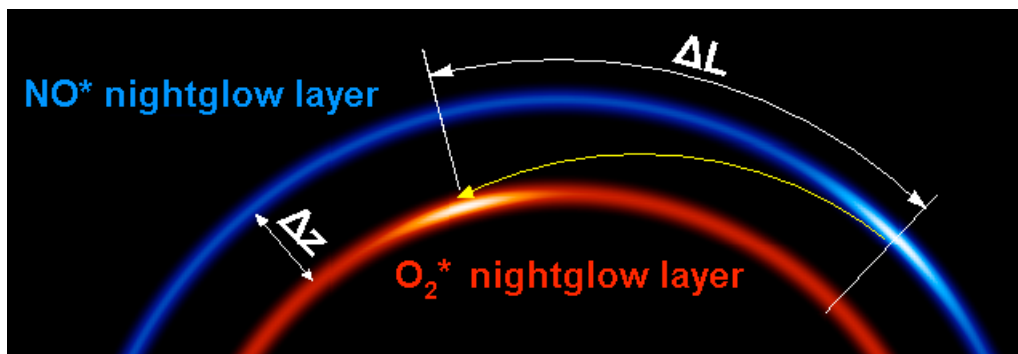


Figure 7