

From direct to inverse analysis in flexible multibody dynamics

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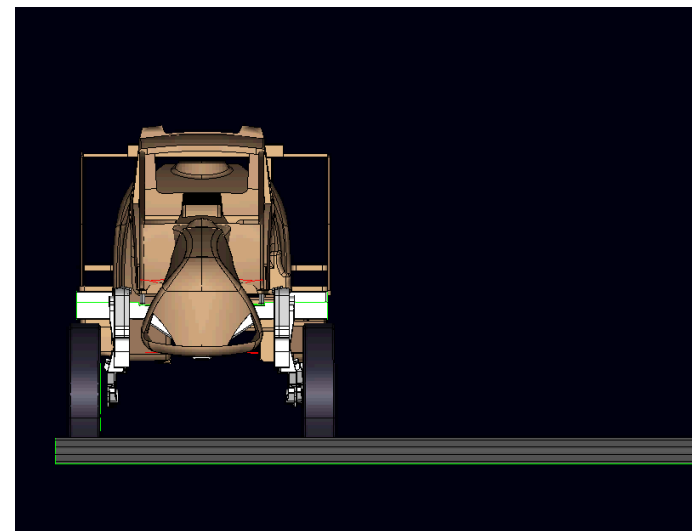
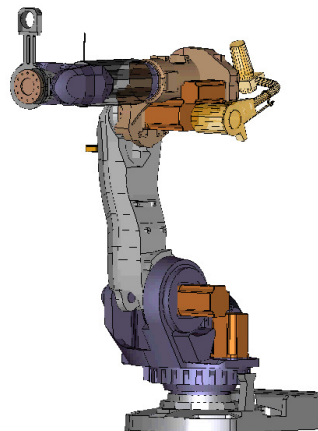
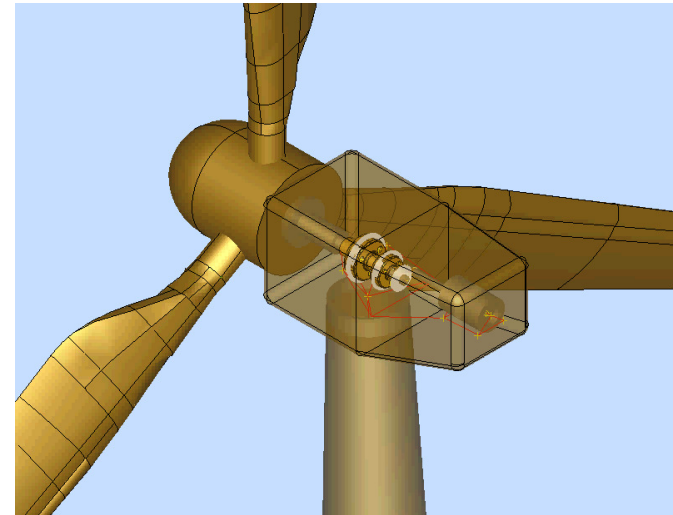
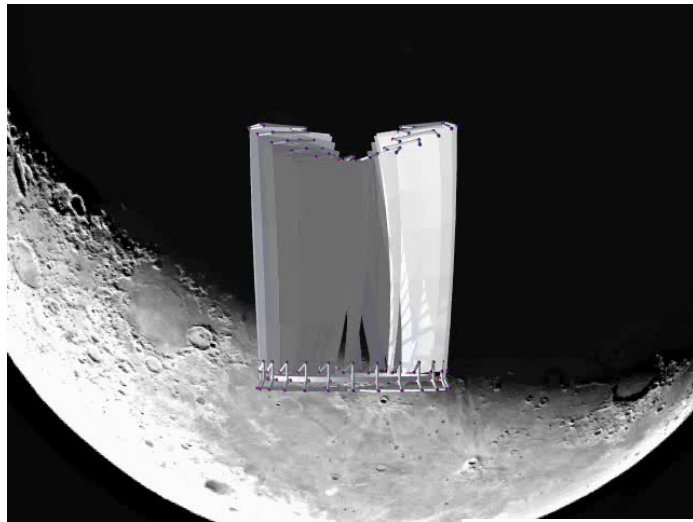
High-fidelity system-level simulation

1. Overview and added-value in mechanical applications
2. Potential and challenges in design optimization
3. Emerging formalisms for inverse analysis
(Lie group approach)

High-fidelity system-level simulation

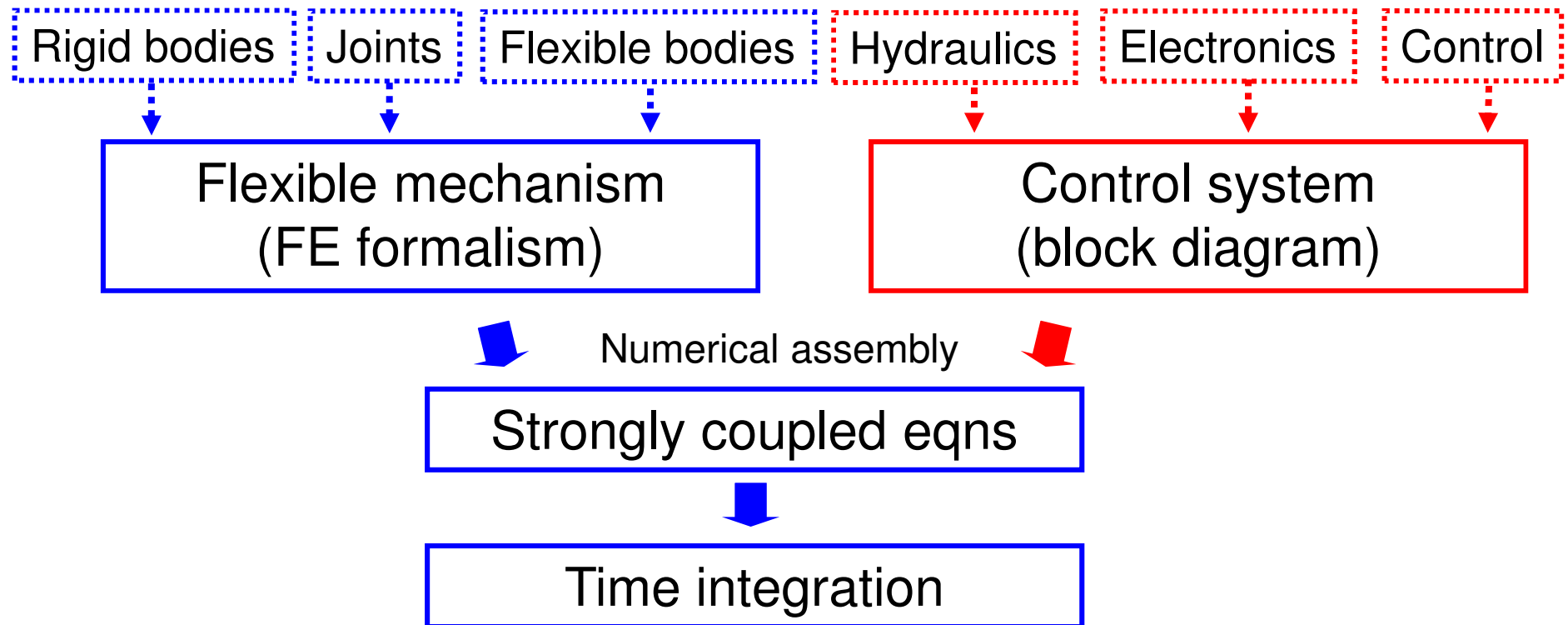
1. Overview and added-value in mechanical applications
 - The FE method in flexible multibody dynamics
 - Simulation of vehicle driveline
 - Simulation of compliant deployable structures
2. Potential and challenges in design optimization
3. Emerging formalisms for inverse analysis
(Lie group approach)

Examples of application fields



Mecano models (Courtesy: LMS Samtech)

Integrated simulation approach



FE method in flexible multibody dynamics

Technical system

Model

Equations of motion

Time integration

Results



Dynamic performance?

Mechanical loads & stresses?

FE method in flexible multibody dynamics

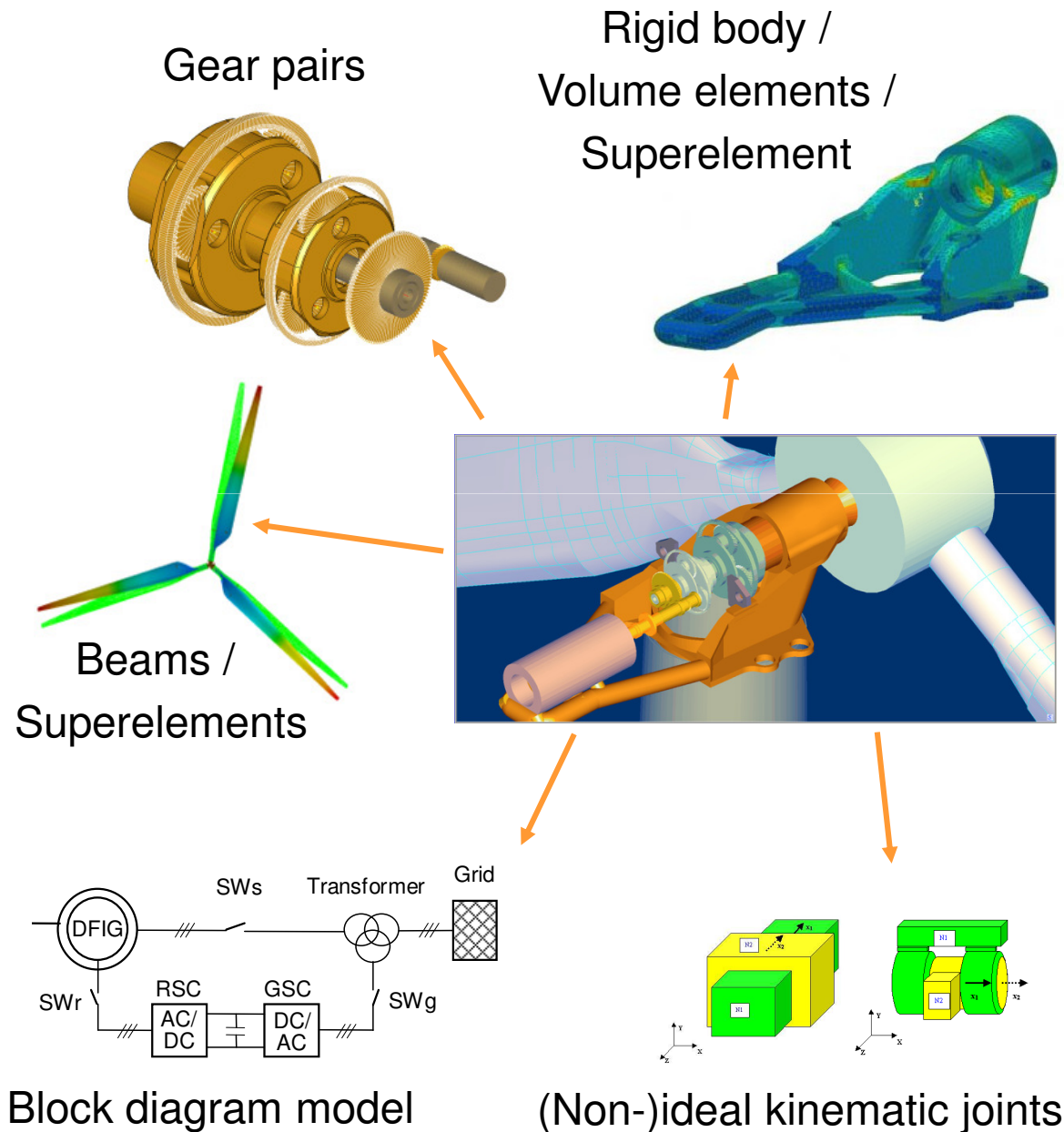
Technical system

Model

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Results



FE method in flexible multibody dynamics

(Gérardin & Cardona 2001)

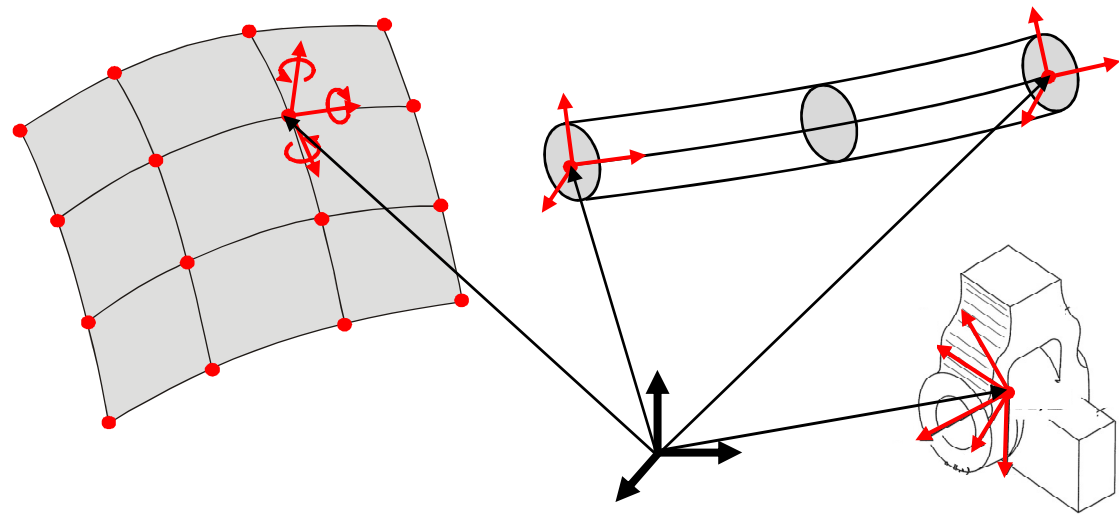
Technical system

Model

Equations of motion

Time integration

Results



If rotations are parameterized:

$$\mathbf{M}(\mathbf{q})\ddot{\mathbf{q}} = \mathbf{g}(\mathbf{q}, \dot{\mathbf{q}}, t) - \Phi_{\mathbf{q}}^T \boldsymbol{\lambda} + \mathbf{L}\mathbf{y}$$

$$\mathbf{0} = \Phi(\mathbf{q}, t)$$

$$\dot{\mathbf{x}} = \mathbf{f}(\mathbf{q}, \dot{\mathbf{q}}, \ddot{\mathbf{q}}, \boldsymbol{\lambda}, \mathbf{x}, \mathbf{y}, t)$$

$$\mathbf{y} = \mathbf{h}(\mathbf{q}, \dot{\mathbf{q}}, \ddot{\mathbf{q}}, \boldsymbol{\lambda}, \mathbf{x}, \mathbf{y}, t)$$

FE method in flexible multibody dynamics

Technical system

Model

Numerical solution of coupled
1st and 2nd order DAEs with index 3

Equations of motion

Generalized- α method
(Chung & Hulbert, 1993)

Time integration

Results

FE method in flexible multibody dynamics

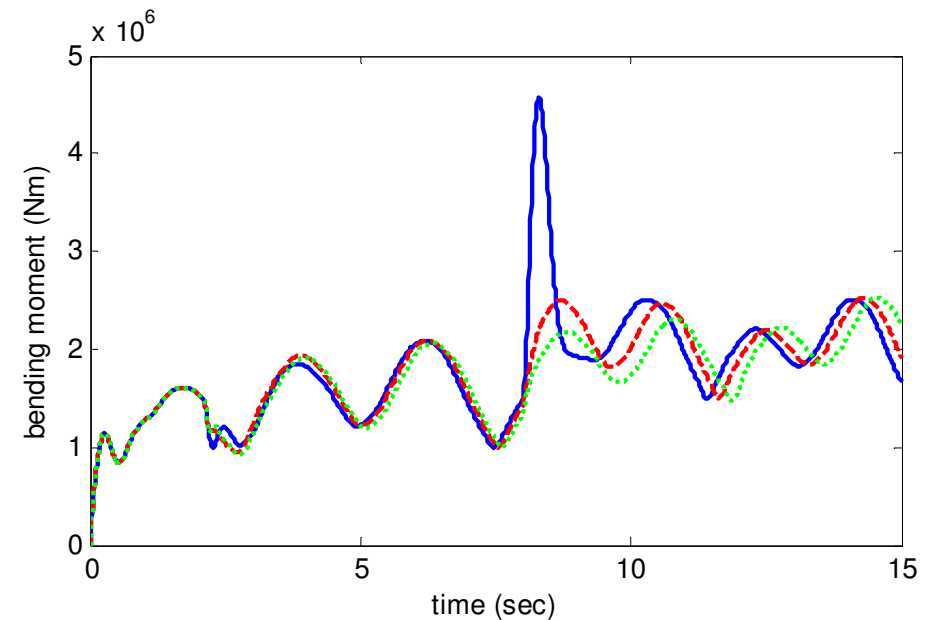
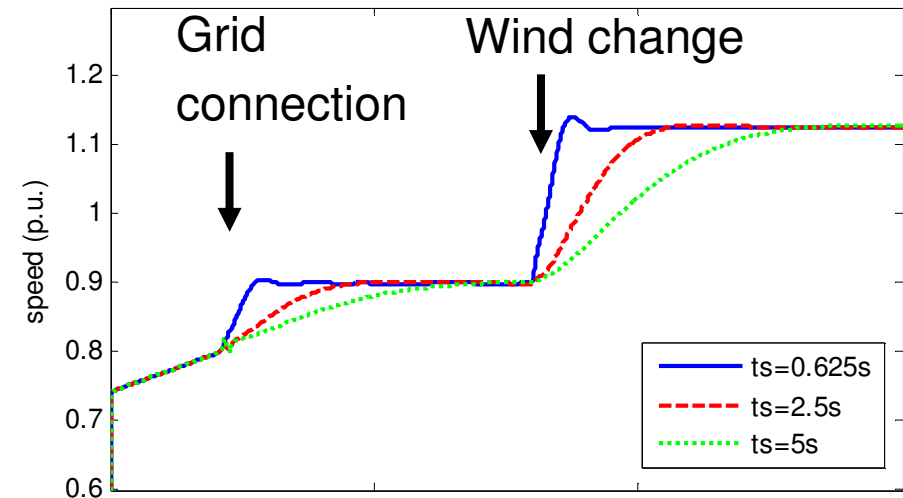
Technical system

Model

Equations of motion

Time integration

Results



Generalized- α time integrator

Stiff & large ODEs in structural dynamics (Chung & Hulbert 1993)

- Second-order accuracy
- Unconditional stability (A-stability)
- Controllable numerical damping (tuning parameter ρ_∞)
- Newmark and HHT are special cases
- Equivalent to a multistep method (Erlicher et al 2002)

Direct integration of index-3 DAEs

- Linear stability (Cardona & G rardin 1989)
- 2nd order accuracy (Arnold & B. 2007)
- Mechatronics (B. & Golinval 2008)

Index reduction methods

(Lunk & Simeon 2006, Jay & Negrut 2007, Arnold 2009, Arnold et al 2011)

Generalized- α time integrator

Newmark implicit formulae:

$$\mathbf{q}_{n+1} = \mathbf{q}_n + h\dot{\mathbf{q}}_n + h^2(0.5 - \beta)\mathbf{a}_n + h^2\beta\mathbf{a}_{n+1}$$

$$\dot{\mathbf{q}}_{n+1} = \dot{\mathbf{q}}_n + h(1 - \gamma)\mathbf{a}_n + h\gamma\mathbf{a}_{n+1}$$

$$\mathbf{x}_{n+1} = \mathbf{x}_n + h(1 - \theta)\mathbf{w}_n + h\theta\mathbf{w}_{n+1}$$

Generalized- α method (Chung & Hulbert, 1993)

$$(1 - \alpha_m)\mathbf{a}_{n+1} + \alpha_m\mathbf{a}_n = (1 - \alpha_f)\ddot{\mathbf{q}}_{n+1} + \alpha_f\ddot{\mathbf{q}}_n$$

$$(1 - \delta_m)\mathbf{w}_{n+1} + \delta_m\mathbf{w}_n = (1 - \delta_f)\dot{\mathbf{x}}_{n+1} + \delta_f\dot{\mathbf{x}}_n$$

Equations of motion at time t_{n+1}

$$\mathbf{M}(\mathbf{q})\ddot{\mathbf{q}} = \mathbf{g}(\mathbf{q}, \dot{\mathbf{q}}, t) - \Phi_{\mathbf{q}}^T \boldsymbol{\lambda} + \mathbf{L}\mathbf{y}$$

$$\mathbf{0} = \Phi(\mathbf{q}, t)$$

$$\dot{\mathbf{x}} = \mathbf{f}(\mathbf{q}, \dot{\mathbf{q}}, \ddot{\mathbf{q}}, \boldsymbol{\lambda}, \mathbf{x}, \mathbf{y}, t)$$

$$\mathbf{y} = \mathbf{h}(\mathbf{q}, \dot{\mathbf{q}}, \ddot{\mathbf{q}}, \boldsymbol{\lambda}, \mathbf{x}, \mathbf{y}, t)$$

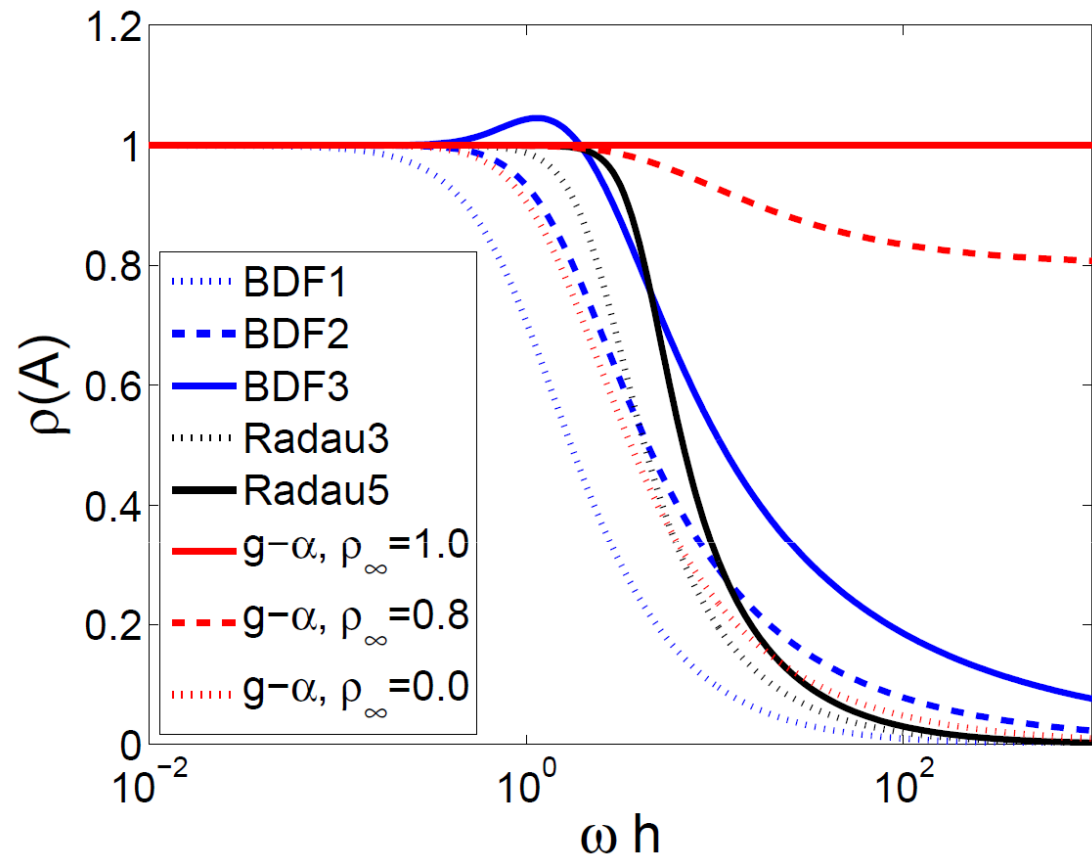
Comparison with other DAE solvers

Test equation:

$$\ddot{q} + \omega^2 q = 0$$

Numerical solution:

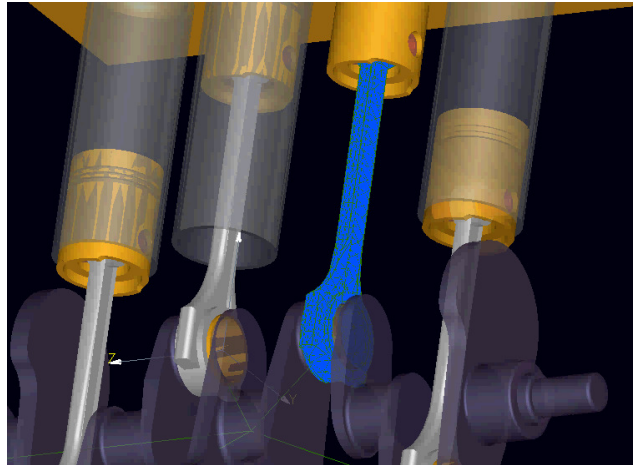
$$\mathbf{x}_n = (\mathbf{A}(\omega h))^n \mathbf{x}_0$$



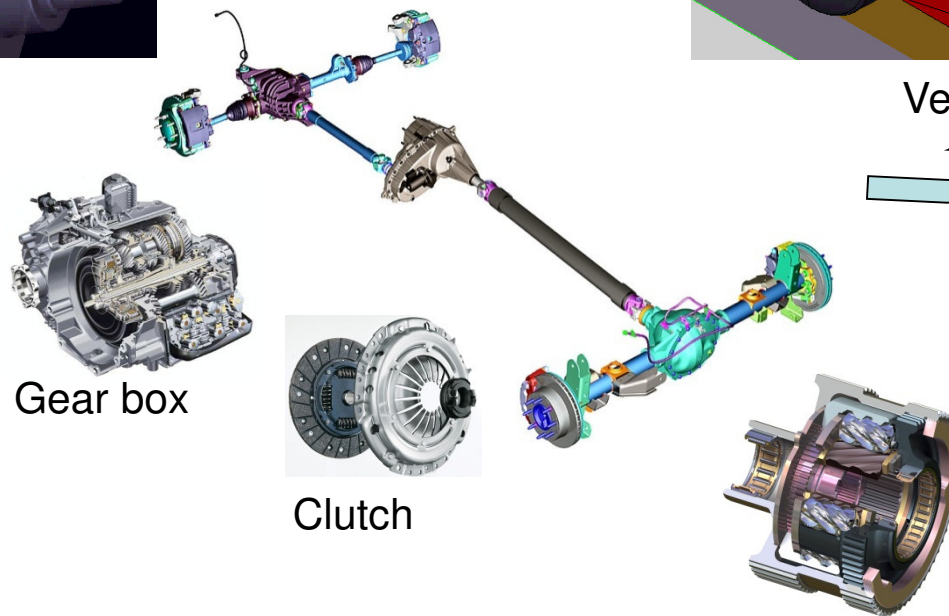
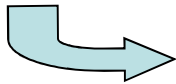
- Order of unconditionally stable BDF ≤ 2
- Less numerical dissipation with generalized- α method

Modelling of vehicle drivelines

(Courtesy: LMS Samtech)



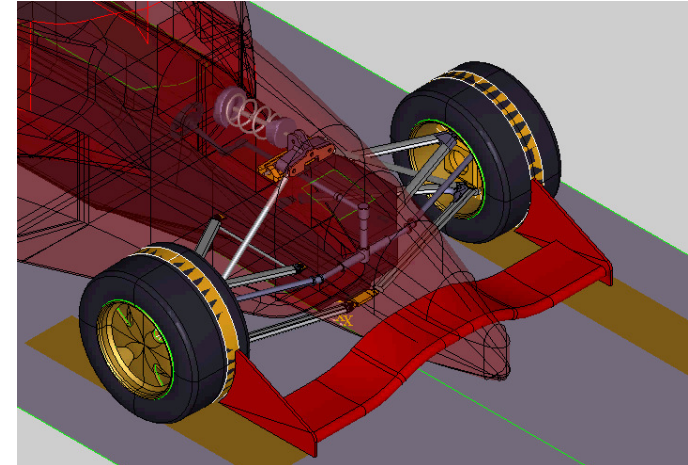
Motor



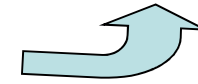
Gear box

Clutch

Differential



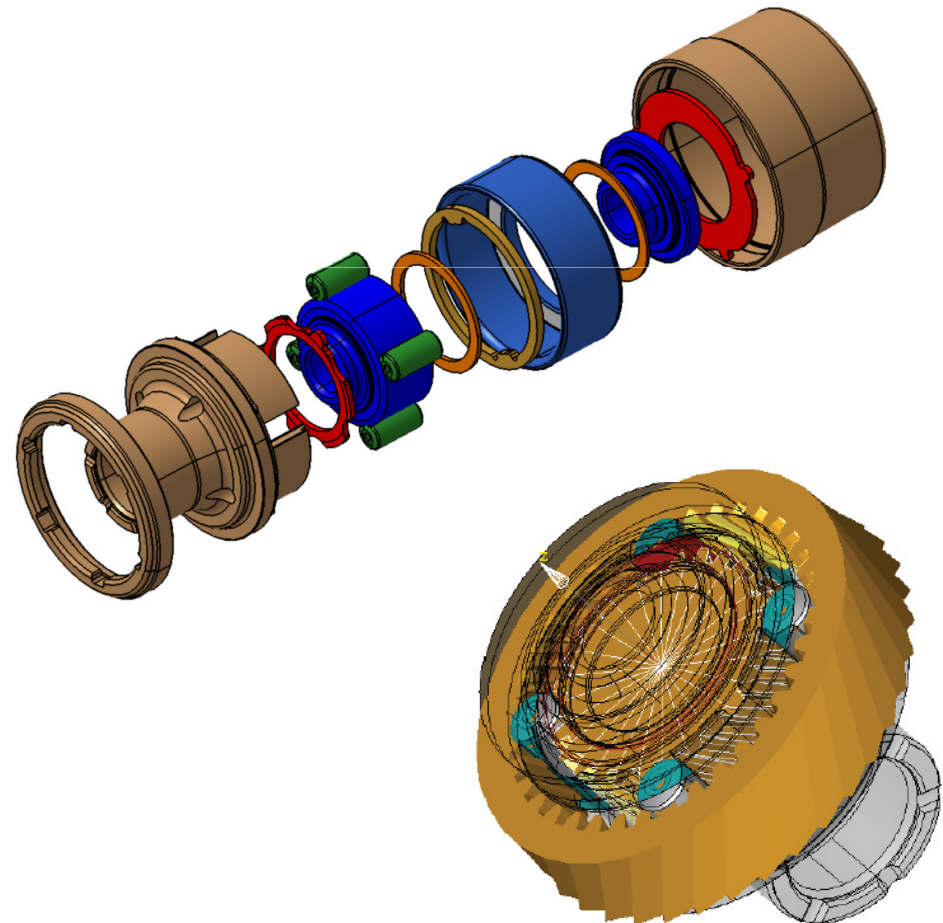
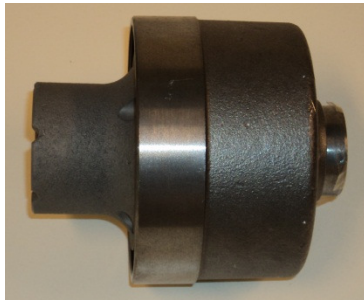
Vehicle dynamics



Modelling the components in their environment

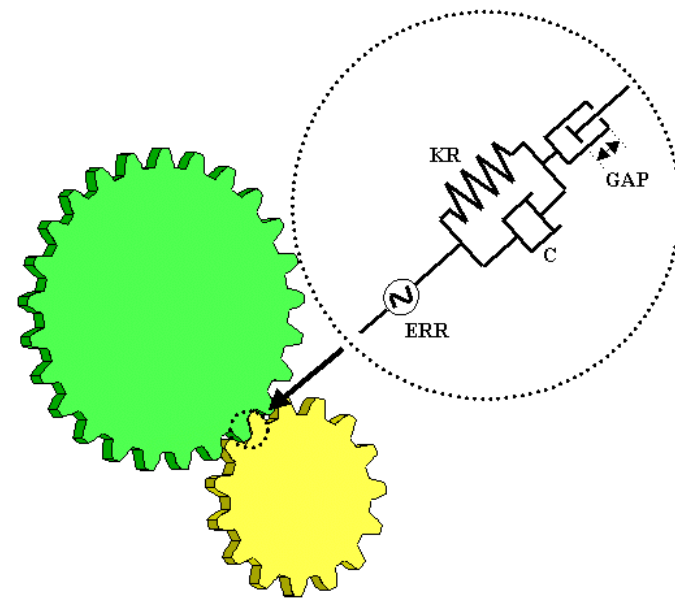
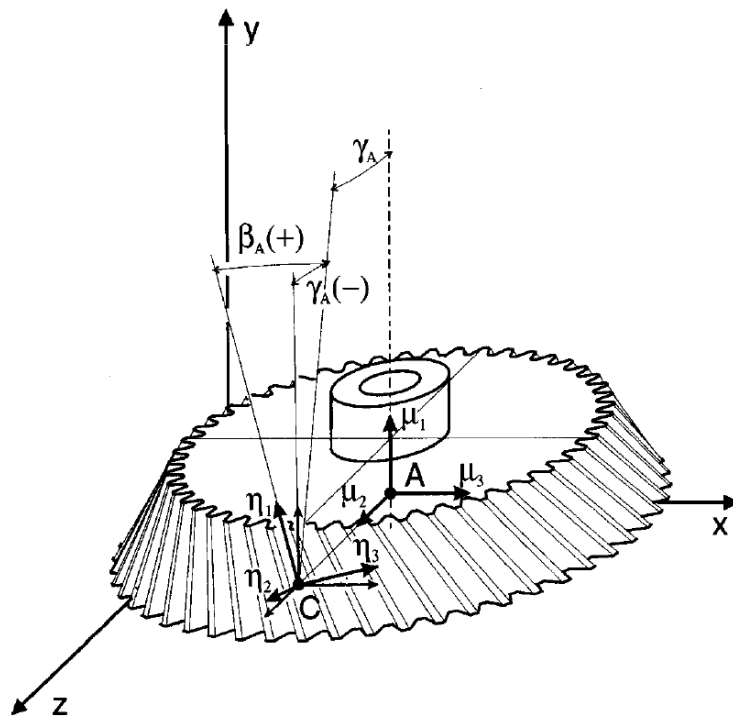
Torsen limited slip differential

- Variable torque distribution between the output shafts
- Locking by friction between gear pairs & thrust washers
- 4 working modes



Gear pair element

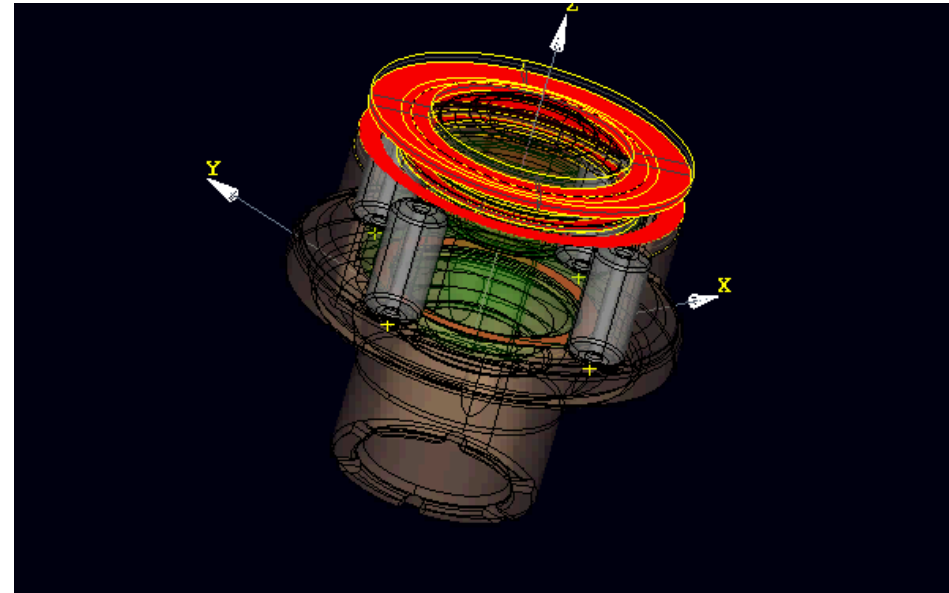
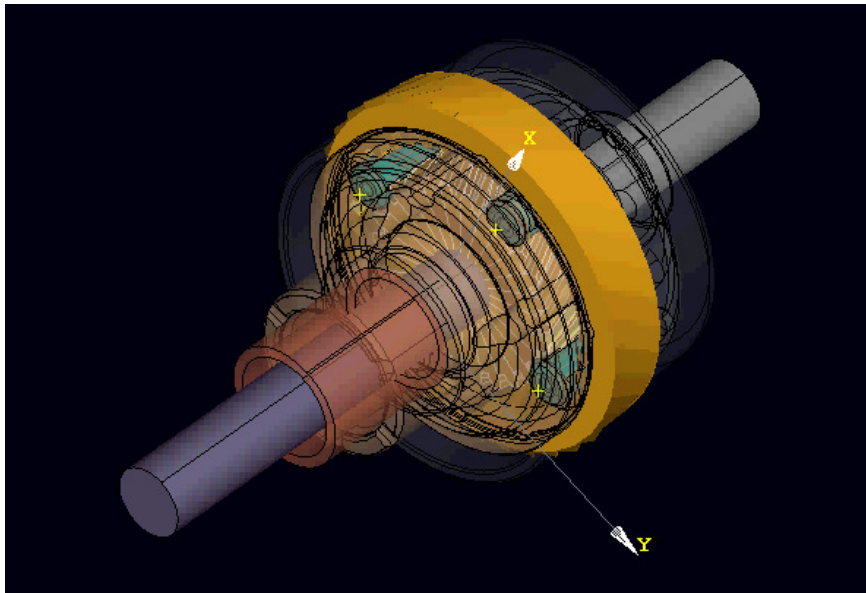
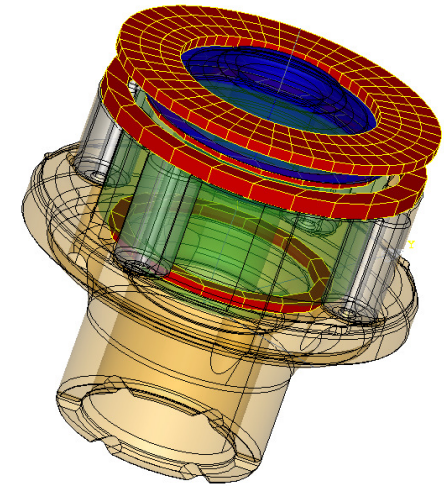
- ❑ Connection between two wheels modelled as rigid bodies
- ❑ Local flexibility : spring (KR) and damper (C)
- ❑ Time fluctuation of mesh stiffness (ISO 6336)
- ❑ Backlash (GAP), load transmission error (ERR), misalignment



Case 1: flexible washers

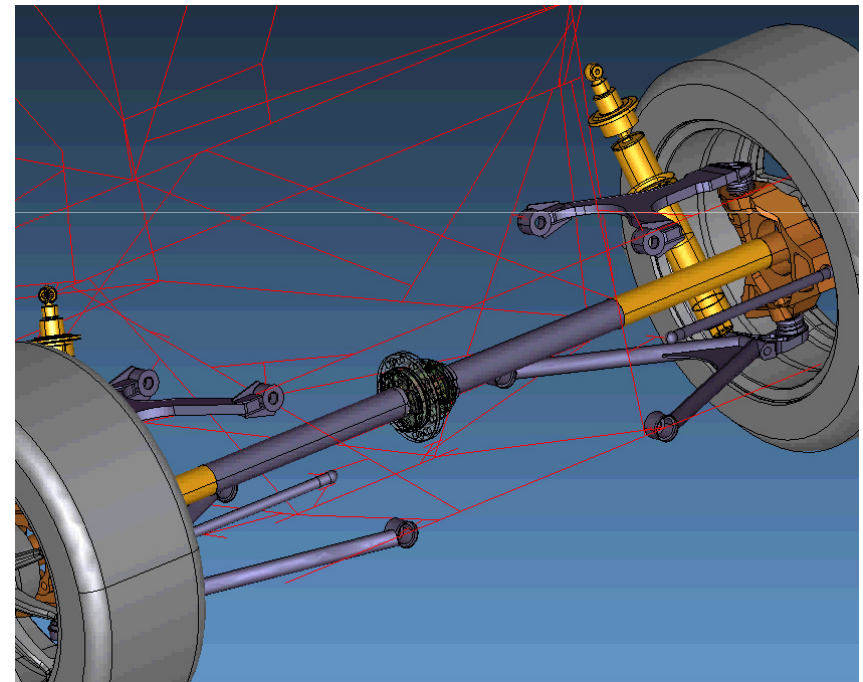
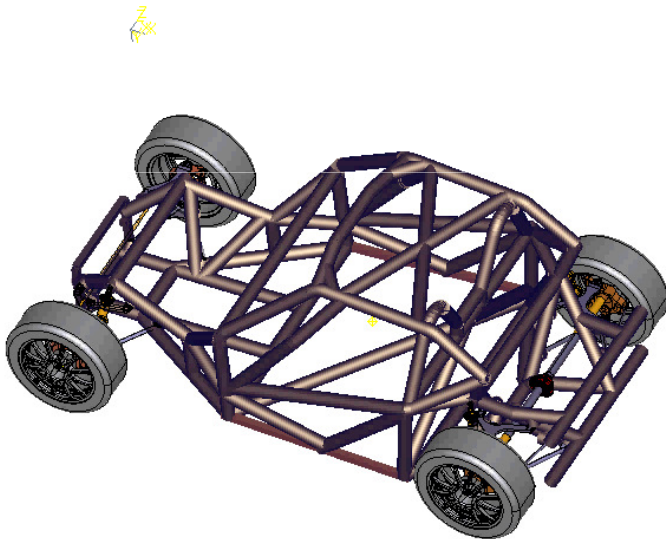
- Rigid/flexible contact conditions
- 8000 dofs

TDR1 numerical : 3.90
TDR1 experimental (Torsen) : 4.02



Case 2: rigid washers

Rigid/rigid contact condition: continuous impact model with a coefficient of restitution (Lankarani & Nikravesh 1994)

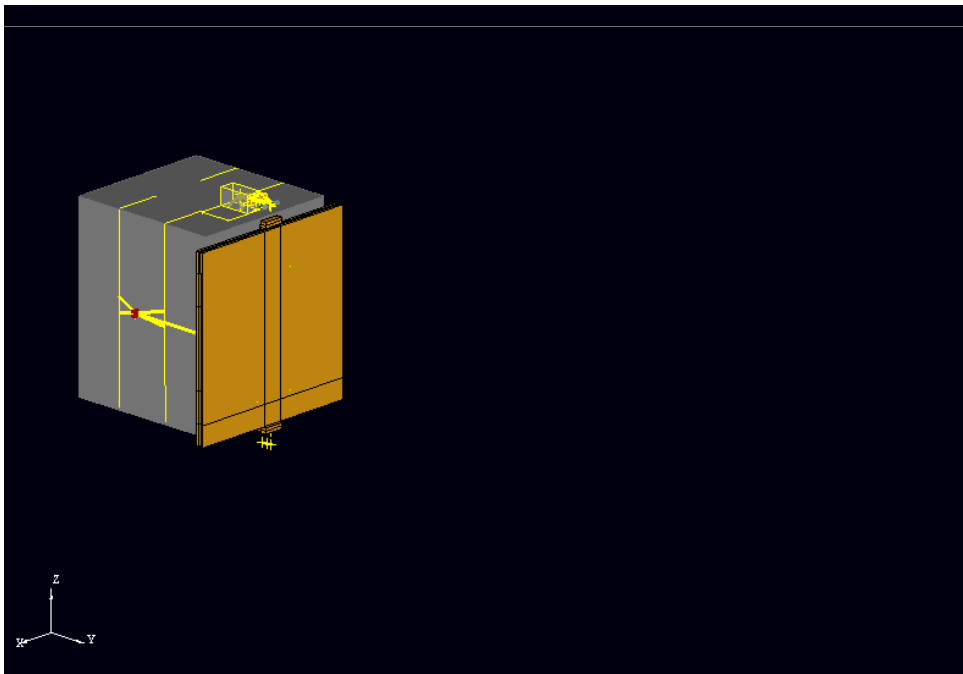


Alternative: nonsmooth description and time-stepping algorithm

Tape-spring hinge

MAEVA hinge (METRAVIB & CNES)

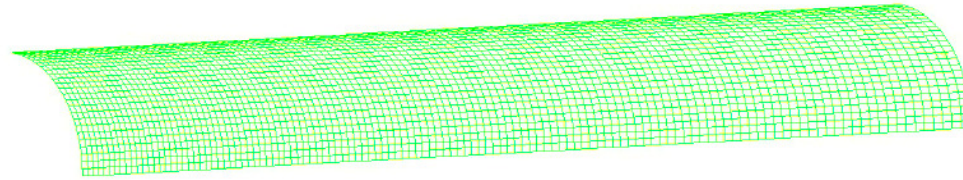
- Guiding, driving and locking functions
- No contact between sliding surfaces



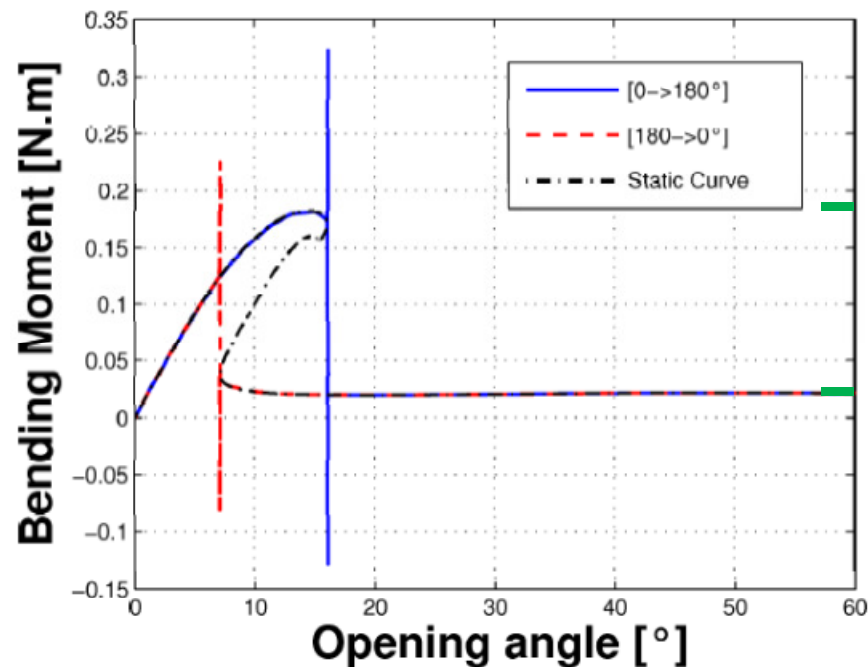
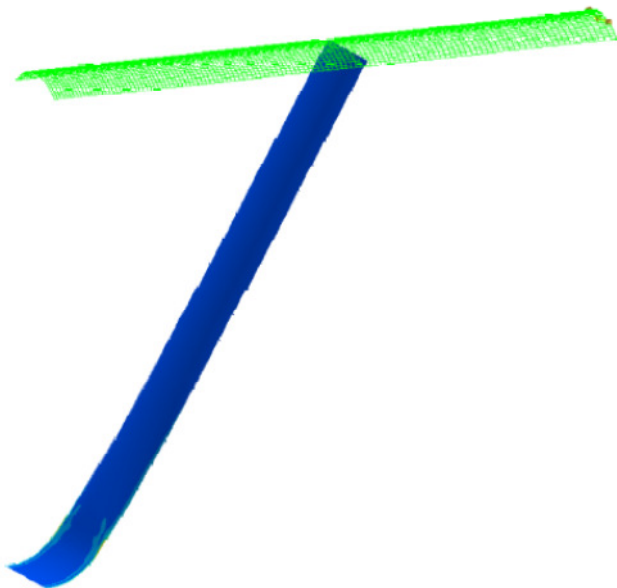
First model: ideal hinge

- No 3D behaviour
- No self-locking

Static analysis of a single tape-spring



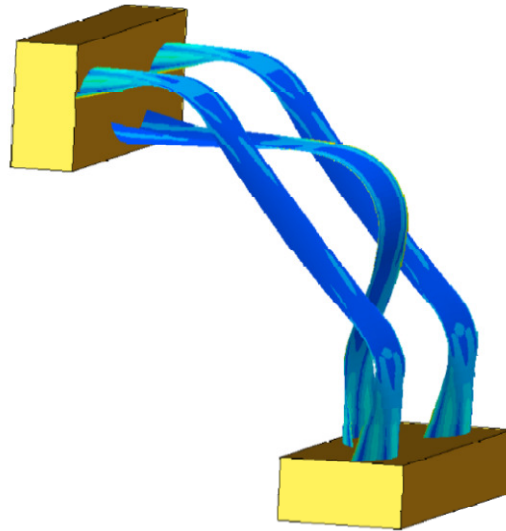
- Fine mesh with second order Mindlin shell elements
- Symmetry is exploited
- Continuation vs. pseudo-dynamic methods



Holding torque

Driving torque

Static behaviour of a full hinge



Numerical results

Driving torque : 0.152 Nm

Holding torque : 6.67 Nm

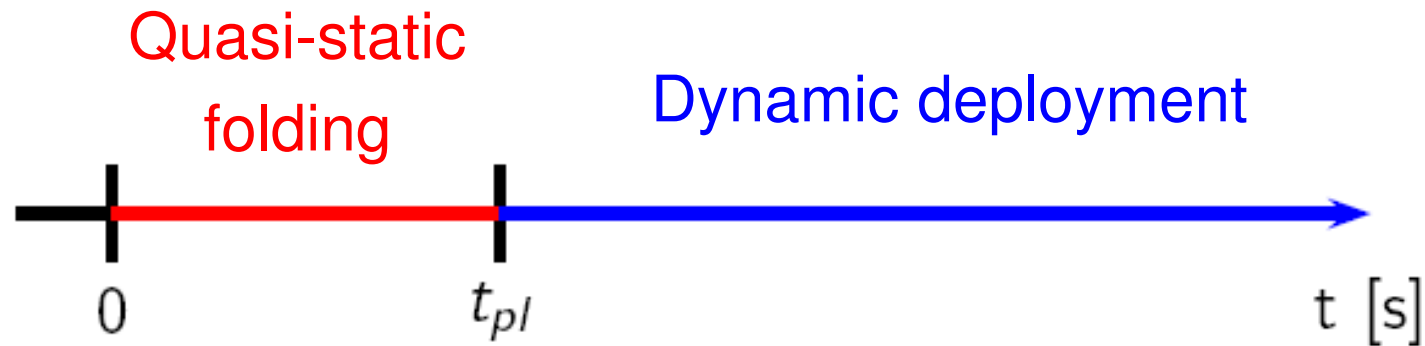
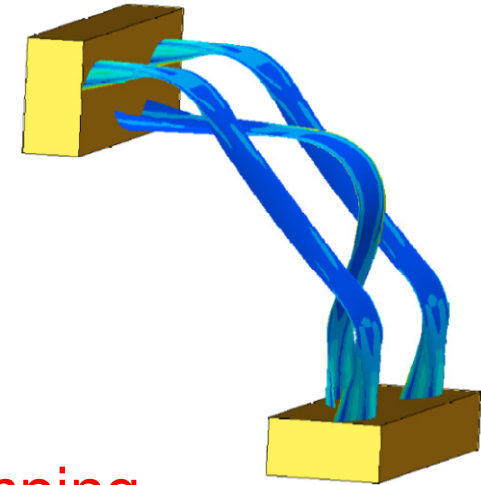
Experimental tests (Metravib):

Driving torque > 0.15 Nm

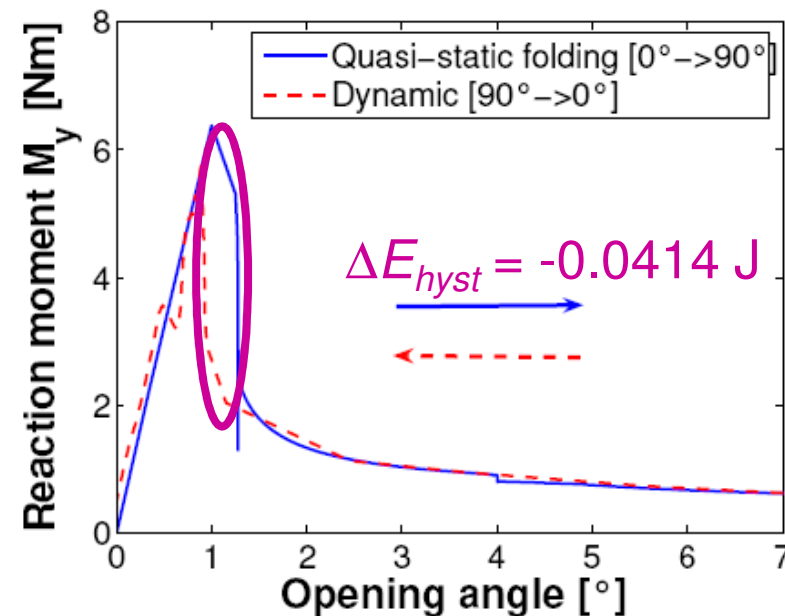
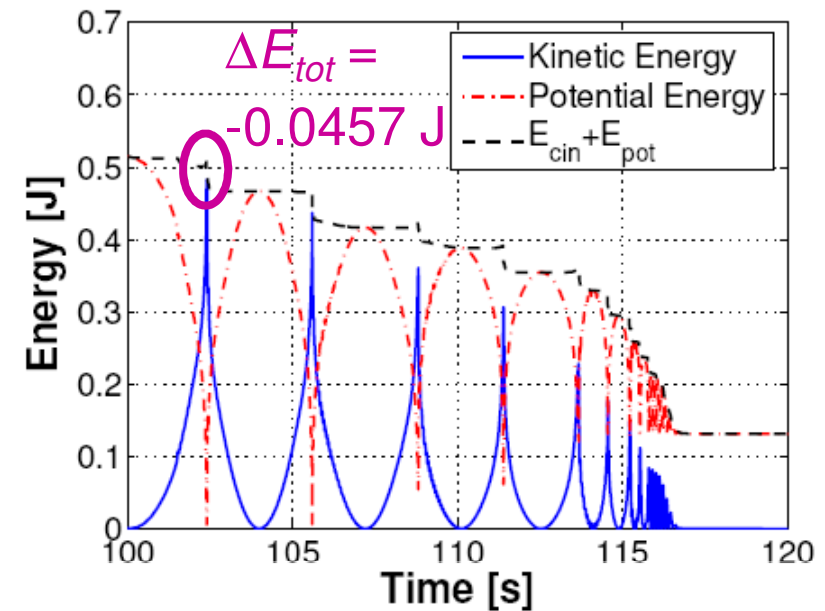
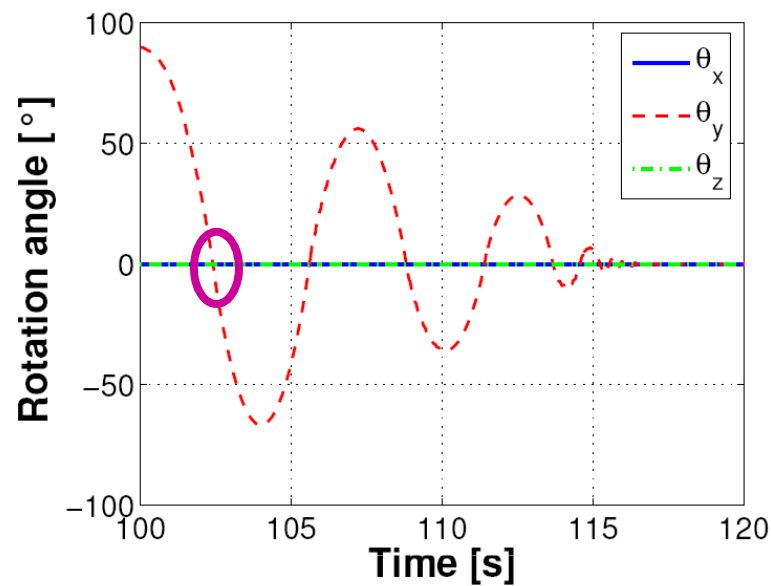
Holding torque > 4.5 Nm

Dynamic behaviour of a full hinge

- Inertia of the rigid appendix (solar panel)
- No structural damping but **numerical damping**

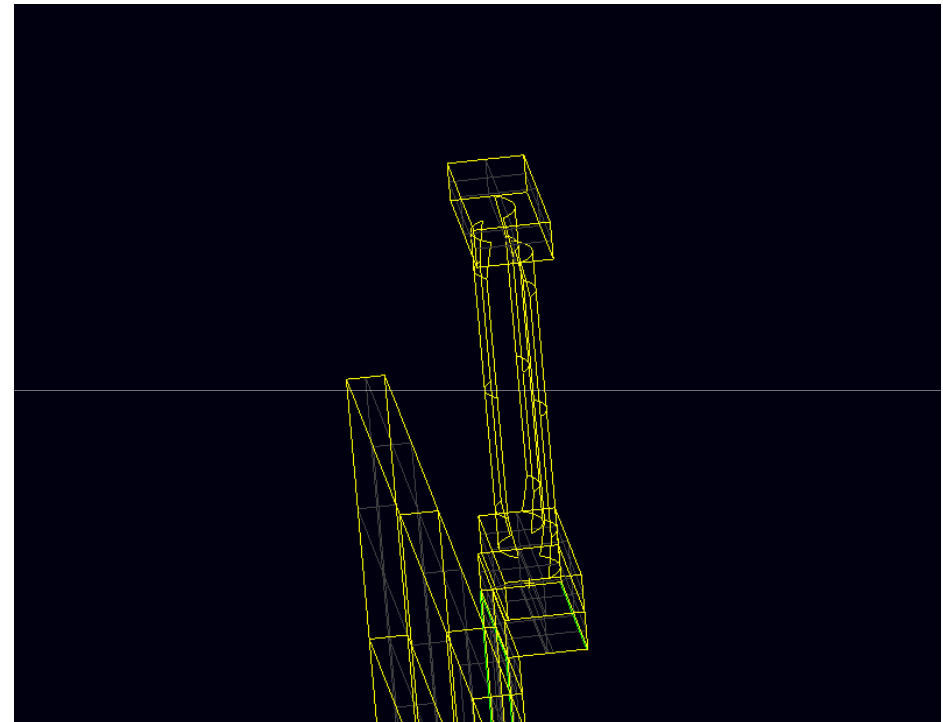
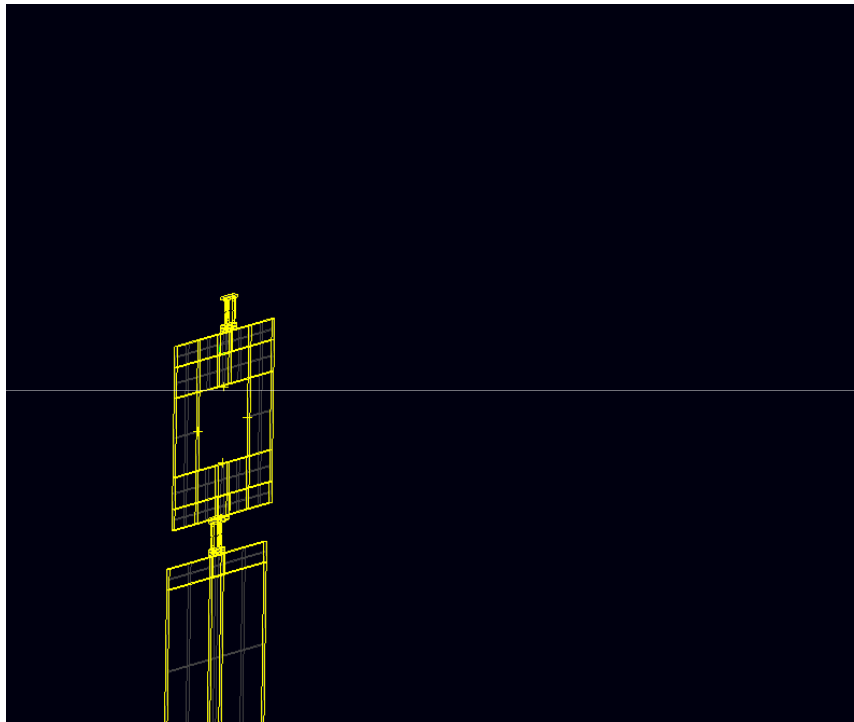


Full hinge - Torsional mode blocked



Full hinge - Torsional mode free

(Hoffait et al 2010)



- Self-locking is caused by the hysteresis phenomenon
- The global dynamic response is acceptable even though the physical dissipation is not modelled!

Summary

Fully integrated approach in flexible multibody dynamics

- Nonlinear finite element method
- Block diagram language
- Monolithic generalized- α time integration

Added-value in applications:

- Motion, vibration & control analysis
- Stress computation with accurate dynamic loadings
- Analysis of compliant systems

Can we use these simulation tools for inverse analysis?

The unknowns may be

- the externally applied loads
- the mechanical design

High-fidelity system-level simulation

1. Overview and added-value in mechanical applications
2. Potential and challenges in design optimization
 - Inverse dynamics
 - Structural optimization
 - Sensitivity analysis
3. Emerging formalisms for inverse analysis
(Lie group approach)

Inverse dynamics of flexible MBS



Flexibility \Rightarrow underactuated system

Forward integration for differentially flat or minimum-phase systems
(Blajer & Kolodziejczyk 2004, Seifried 2010)

Stable inversion for systems in nonlinear I/O normal form (Seifried & Eberhard 2009)

$$\begin{aligned} M(q)\ddot{q} + g(q, \dot{q}, t) + \Phi_q^T \lambda &= A u \\ \Phi(q) &= 0 \\ y(q) &= y_d(t) \end{aligned}$$

Optimal control for flexible MBS in DAE form

Inverse dynamics of flexible MBS

$$\min_{\mathbf{q}(t), \boldsymbol{\lambda}(t), \mathbf{u}(t)} G(\mathbf{q}(t_f), \dot{\mathbf{q}}(t_f)) + \int_{t_0}^{t_f} F(\mathbf{q}, \dot{\mathbf{q}}, \boldsymbol{\lambda}, \mathbf{u}) d\tau$$

$$\text{s.t. } \mathbf{M}(\mathbf{q})\ddot{\mathbf{q}} + \mathbf{g}(\mathbf{q}, \dot{\mathbf{q}}, t) + \boldsymbol{\Phi}_{\mathbf{q}}^T \boldsymbol{\lambda} = \mathbf{A}\mathbf{u}$$

$$\boldsymbol{\Phi}(\mathbf{q}) = \mathbf{0}$$

$$\mathbf{y}(\mathbf{q}) = \mathbf{y}_d(t) + \text{other constraints}$$

Direct collocation method \Rightarrow large but sparse NLP problem

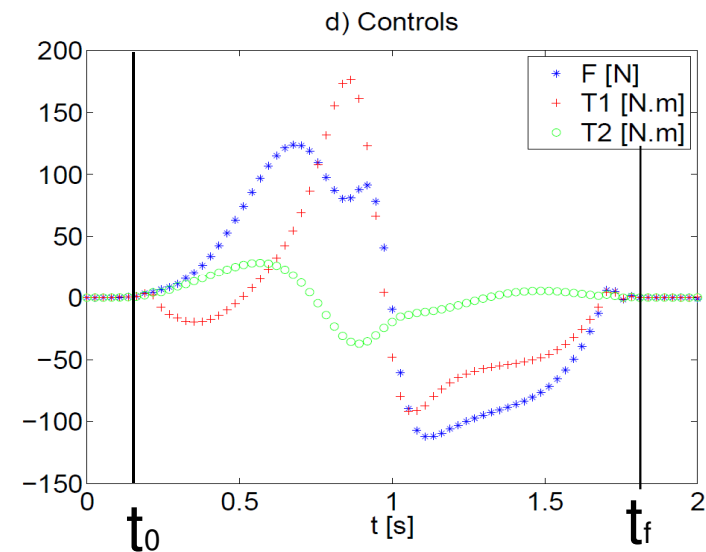
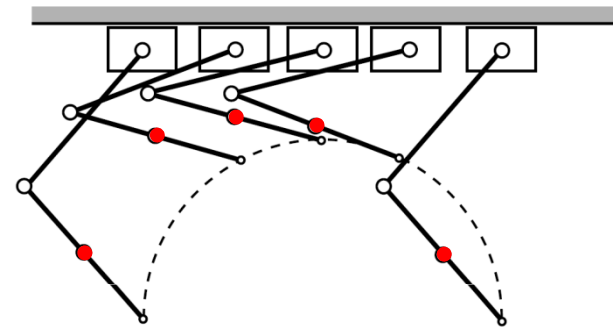
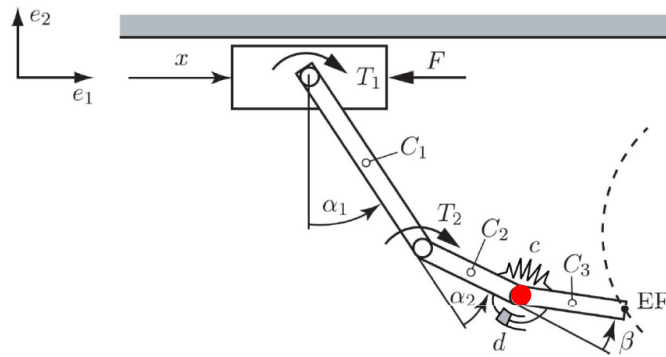
$$\mathbf{x} = (\mathbf{q}_1, \dot{\mathbf{q}}_1, \ddot{\mathbf{q}}_1, \mathbf{a}_1, \boldsymbol{\lambda}_1, \mathbf{u}_1, \dots, \mathbf{q}_N, \dot{\mathbf{q}}_N, \ddot{\mathbf{q}}_N, \mathbf{a}_N, \boldsymbol{\lambda}_N, \mathbf{u}_N)$$

$$\min_{\mathbf{x}} G(\mathbf{q}_N, \dot{\mathbf{q}}_N) + \sum_{n=2}^N h F(\mathbf{q}_n, \dot{\mathbf{q}}_n, \boldsymbol{\lambda}_n, \mathbf{u}_n)$$

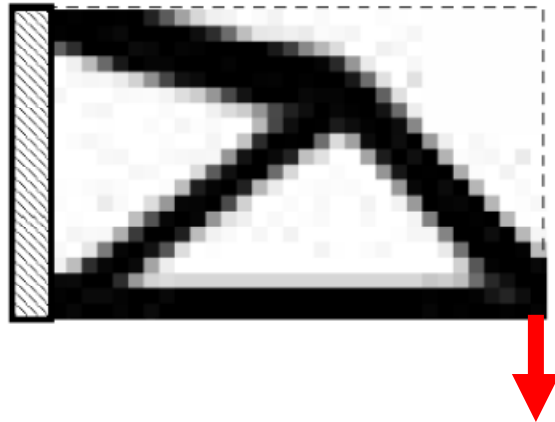
$$\text{s.t. } \left\{ \begin{array}{l} \text{equations of motion} \\ \text{integration formulae} \end{array} \right\} \quad \text{at each step}$$

Inverse dynamics of flexible MBS

Manipulator with one passive joint (non-minimum phase)
(Bastos, Seifried & B. 2011)



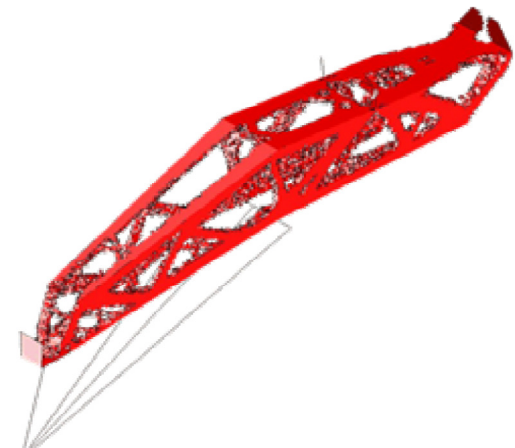
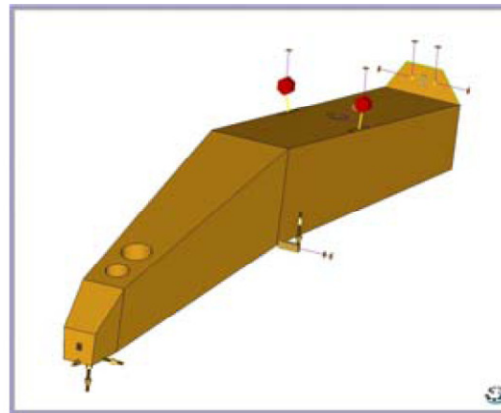
Structural optimization



Topology optimization

(Bendsøe & Kikuchi 1988, Sigmund, 2001)

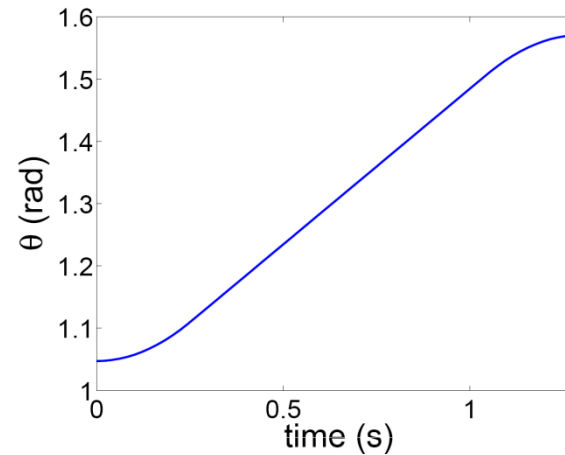
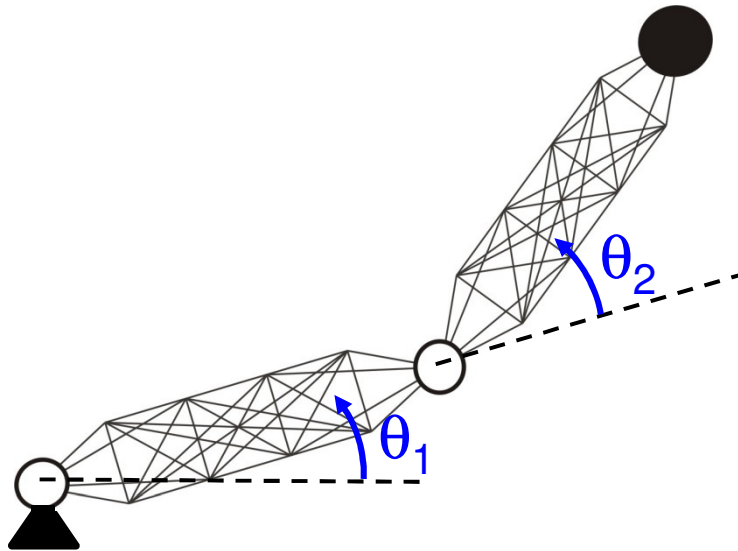
Complex structures can be optimized w.r.t. static loads:



Optimization of articulated systems with dynamic load cases?

Topology optimization of a planar robot arm

(B. et al 2007)



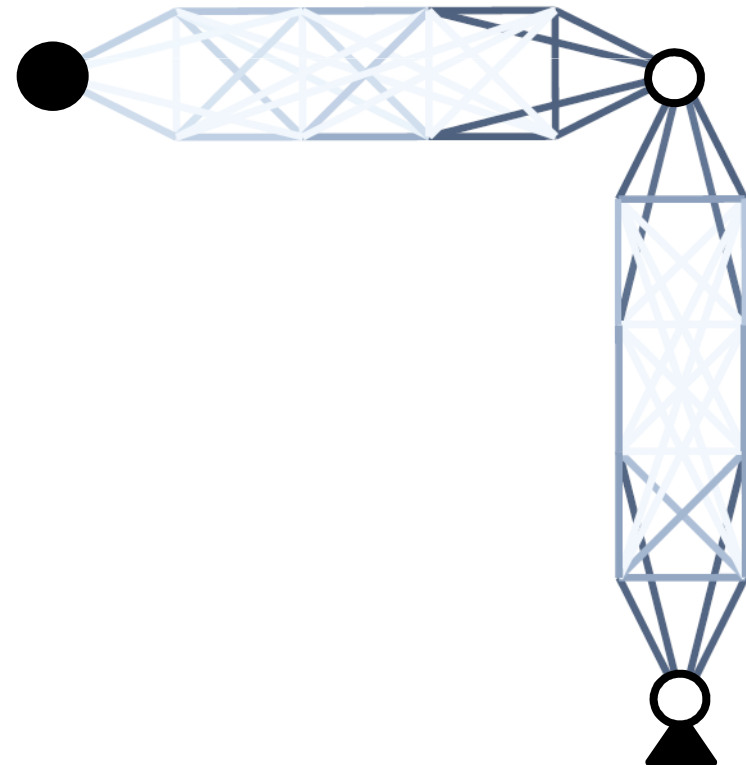
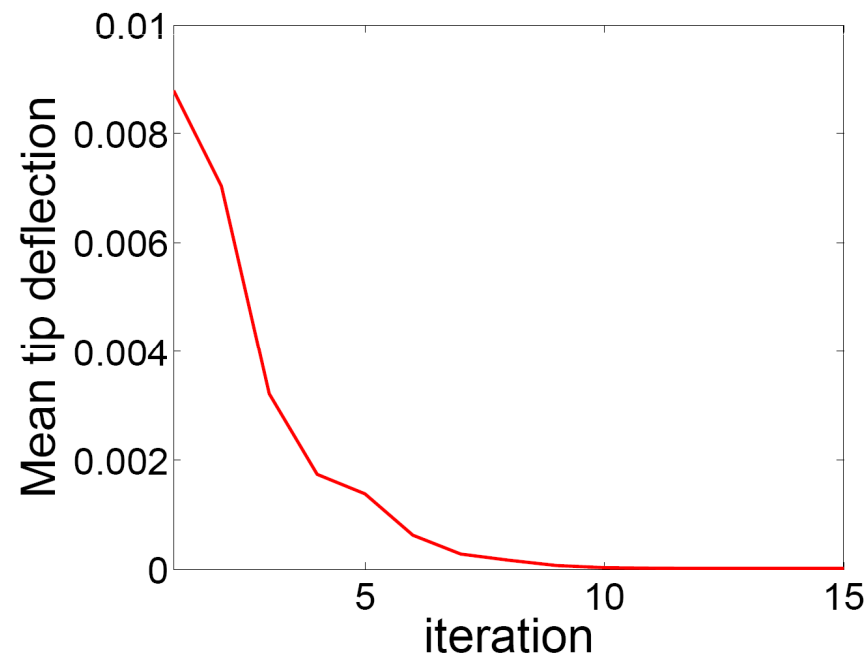
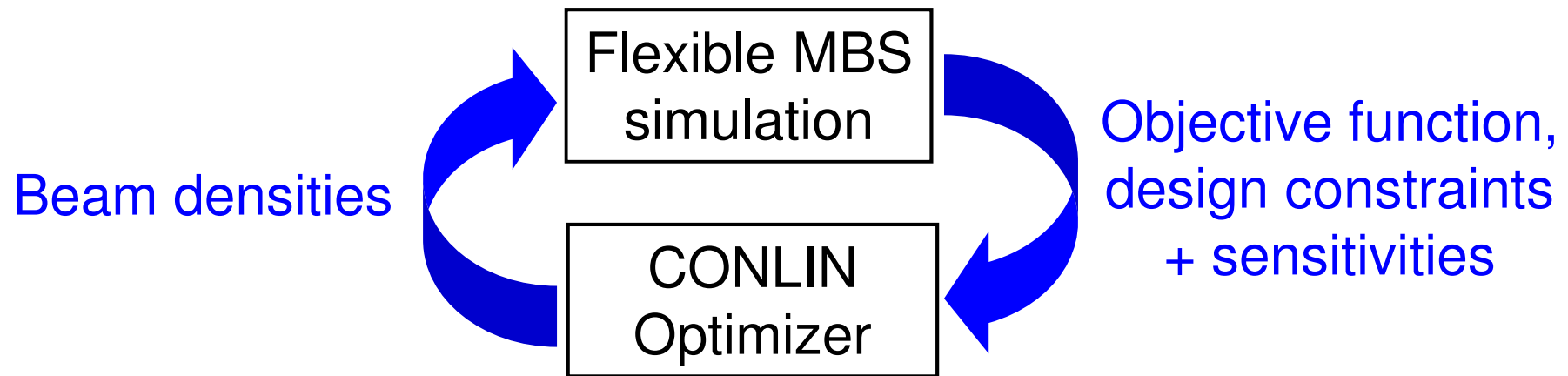
Point-to-point
joint trajectory

One topology variable per beam (SIMP penalization)

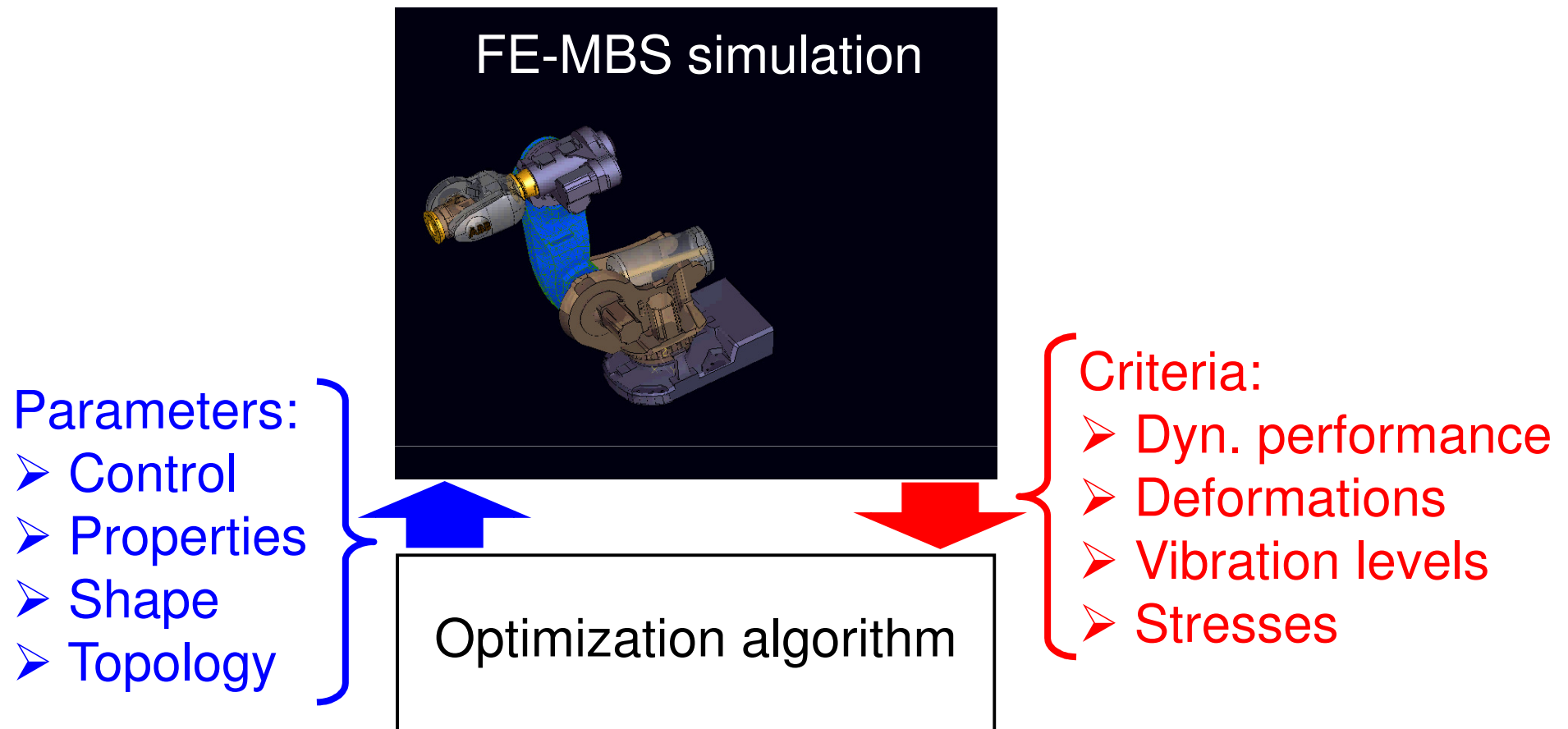
$$\min \quad \frac{1}{t_f} \int_0^{t_f} \|\mathbf{r} - \mathbf{r}_{rigid}\|^2 dt$$

$$\text{subject to } V_{(i)} \leq 0.4 V_{full,(i)}$$

Topology optimization of a planar robot arm



Optimization of full-scale 3D MBS?



Gradient-based & sparse methods (SQP, IP, CONLIN, MMA, etc)

- The problem should be carefully formulated
- An efficient evaluation of the **sensitivities** is needed

Methods for sensitivity analysis

High cost of finite differences for large scale problems

- n_p additional simulations for fwd/bwd differences (order 1)
- $2 n_p$ additional simulations for central differences (order 2)

Automatic differentiation

- High reliability but suboptimal code, so that a manual post-processing of the code is often required

Semi-analytical methods (direct differentiation / adjoint variable)

- Optimized but manual implementation
- Tend to amplify the intricacy of a simulation code
- Feasible for flexible MBS?

Large rotations in 3D: the whole story

$$\mathbf{q} = (\mathbf{x}_1, \boldsymbol{\alpha}_1, \dots, \mathbf{x}_{k_n}, \boldsymbol{\alpha}_{k_n})$$

Rotational equilibrium of a free body: $\mathbf{M}(\boldsymbol{\alpha}) \ddot{\boldsymbol{\alpha}} + \mathbf{g}(\boldsymbol{\alpha}, \dot{\boldsymbol{\alpha}}) = \mathbf{0}$

$$\mathbf{M}(\boldsymbol{\alpha}) = \mathbf{T}^T(\boldsymbol{\alpha}) \mathbf{J} \mathbf{T}(\boldsymbol{\alpha})$$

$$\mathbf{g}(\boldsymbol{\alpha}, \dot{\boldsymbol{\alpha}}) = \mathbf{T}^T(\boldsymbol{\alpha}) (\mathbf{J} \dot{\mathbf{T}}(\boldsymbol{\alpha}, \dot{\boldsymbol{\alpha}}) + (\mathbf{T}(\boldsymbol{\alpha}) \dot{\boldsymbol{\alpha}}) \times \mathbf{J} \mathbf{T}(\boldsymbol{\alpha}) \dot{\boldsymbol{\alpha}})$$

Updated Lagrangian strategy (Cardona & G  radin 1989)

$$\mathbf{R}(t_{n+1}) = \mathbf{R}(t_n) \mathbf{R}_{inc}(t_{n+1})$$

- Only the incremental rotation is parameterized
- Geometrically exact and singularity-free approach
- Equivalent to a reparameterization at each time step

Successful for simulation codes but **challenging for SA!**

(B. & Eberhard 2008)

Summary

Dynamic response optimization

- The FEM in flexible multibody dynamics can be exploited for inverse dynamics & structural optimization
- This leads to large scale optimization problems involving transient analyses
- More efficient transient/sensitivity analyses are needed for the optimization of full-scale 3D systems

High-fidelity system-level simulation

1. Overview and added-value in mechanical applications
2. Potential and challenges in design optimization
3. Emerging formalisms for inverse analysis
 - Lie group approach
 - Sensitivity analysis

Lie group formulation

- The configuration of a MBS is described as an element of a matrix Lie group.
- The equations of motion are formulated on the Lie group
- Numerical solution is computed on the Lie group

Properties:

- parameterization-free (geometric) approach
- simpler formulations and numerical procedures

Lie group description of a MBS

Example: $\mathbf{R}(t) \in SO(3)$

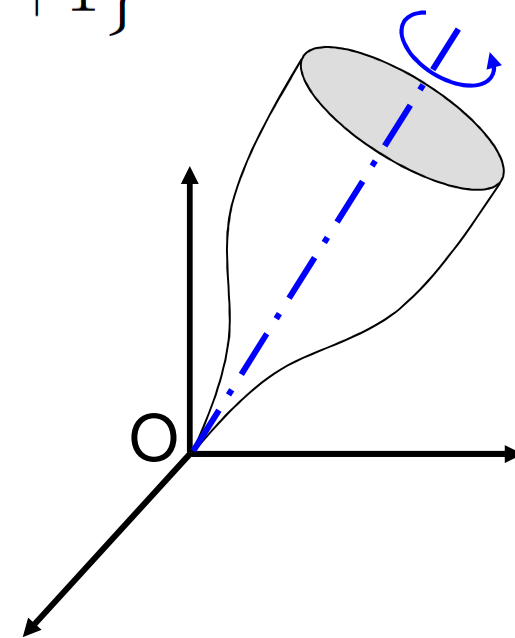
$$SO(3) = \{ \mathbf{R} \in \mathbb{R}^{3 \times 3} | \mathbf{R}^T \mathbf{R} = \mathbf{I}_3, \det \mathbf{R} = +1 \}$$

$$\dot{\mathbf{R}} = \mathbf{R} \tilde{\boldsymbol{\Omega}}$$

$$\boldsymbol{\Omega} = [\Omega_1 \ \Omega_2 \ \Omega_3]^T \in \mathbb{R}^3$$

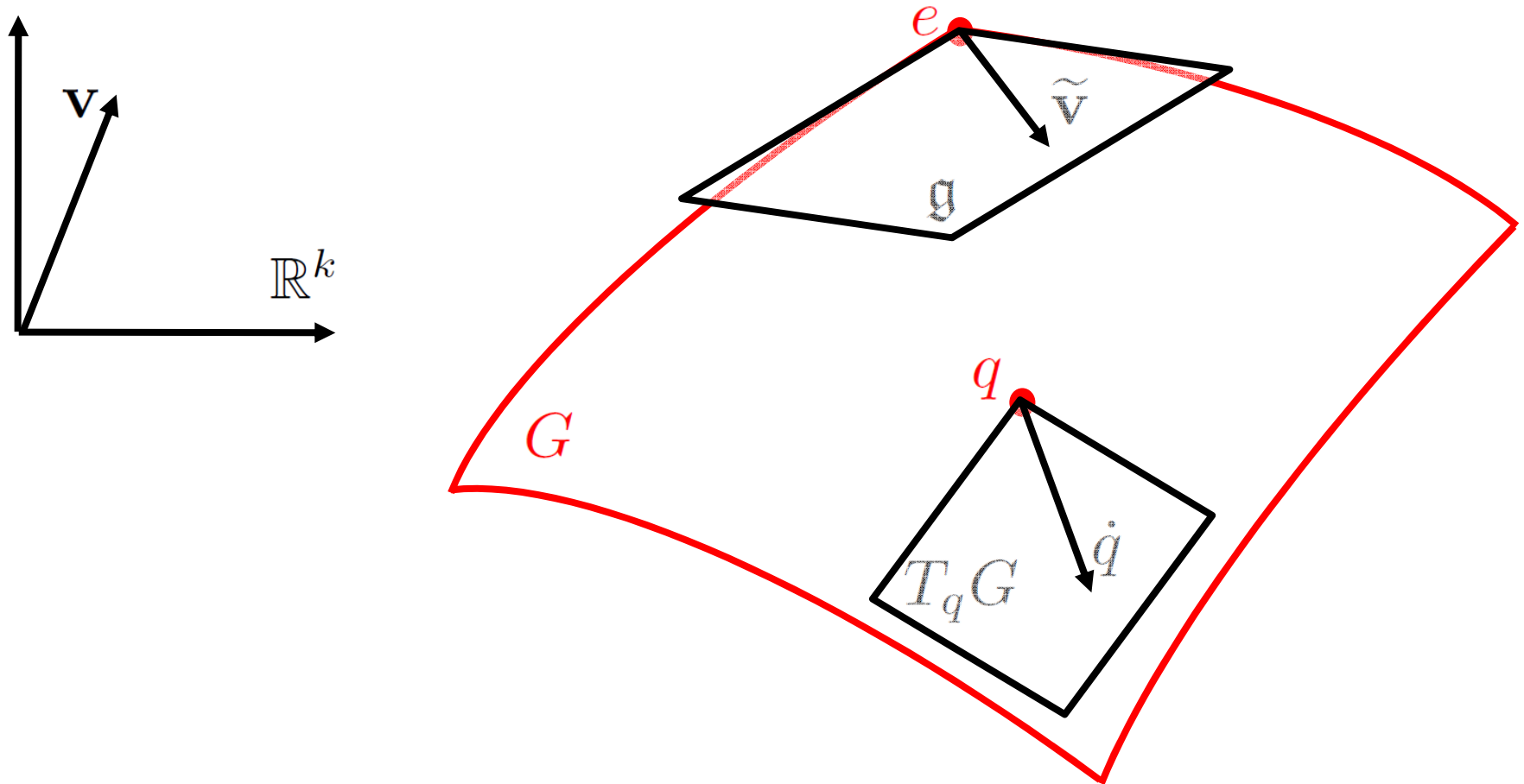
$$\tilde{\boldsymbol{\Omega}} = \begin{bmatrix} 0 & -\Omega_3 & \Omega_2 \\ \Omega_3 & 0 & -\Omega_1 \\ -\Omega_2 & \Omega_1 & 0 \end{bmatrix}$$

$$\in \mathfrak{so}(3) = \{ \tilde{\boldsymbol{\Omega}} : \tilde{\boldsymbol{\Omega}} + \tilde{\boldsymbol{\Omega}}^T = \mathbf{0} \}$$



A Lie group is not a linear space!

Lie group, tangent vectors and Lie algebra



Kinematic compatibility equation (left translation map): $\dot{q} = q\tilde{v}$

Lie group description of a nodal variable

$$\mathbb{R}^3 \times \text{SO}(3): \quad q = (\mathbf{x}, \mathbf{R})$$

➤ Composition: $(\mathbf{x}_1, \mathbf{R}_1) \circ (\mathbf{x}_2, \mathbf{R}_2) = (\mathbf{x}_1 + \mathbf{x}_2, \mathbf{R}_1 \mathbf{R}_2)$

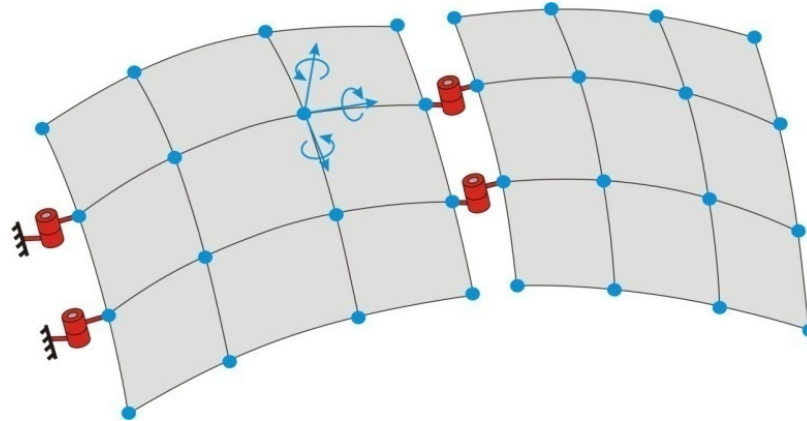
➤ Velocity vector: $\mathbf{v} = \begin{bmatrix} \mathbf{u} \\ \boldsymbol{\Omega} \end{bmatrix}$ with $\dot{\mathbf{x}} = \mathbf{u}$

$$\text{SE}(3): \quad q = \begin{bmatrix} \mathbf{R} & \mathbf{x} \\ \mathbf{0}_{1 \times 3} & 1 \end{bmatrix}$$

➤ Composition: product of 4x4 homogenous transf. matrices

➤ Velocity vector: $\mathbf{v} = \begin{bmatrix} \mathbf{U} \\ \boldsymbol{\Omega} \end{bmatrix}$ with $\dot{\mathbf{x}} = \mathbf{R}\mathbf{U}$

Configuration of a multibody system



- $q \in G$ is a collection of nodal variables, so that,

$$G = \mathbb{R}^3 \times SO(3) \times \dots \times \mathbb{R}^3 \times SO(3)$$

$$\text{or } G = SE(3) \times \dots \times SE(3)$$

- m kinematic constraints $\Phi(q)$

Equations of motion on a Lie group

Hamilton principle \Rightarrow Index-3 DAE on a Lie group

$$\begin{aligned}\dot{q} &= q\tilde{\mathbf{v}} \\ \mathbf{M}\dot{\mathbf{v}} - \hat{\mathbf{v}}^T \mathbf{M}\mathbf{v} + \mathbf{g}(q, t) + \mathbf{B}^T(q)\boldsymbol{\lambda} &= \mathbf{0} \\ \Phi(q) &= \mathbf{0}\end{aligned}$$

- The configuration is described by the matrix q
- The velocity is described by a vector \mathbf{v} ,
related with the matrix $\tilde{\mathbf{v}}$
- The mass matrix is constant, inertia forces are quadratic
- Parameterization-free formulation!

Overview of Lie group integration methods

Local (incremental) parameterization of the equations of motion

- Cardona & Géradin (1989): HHT method for flexible MBS
- Munthe-Kaas (1995, 1998): RK method for ODEs
- Bottasso & Borri (1998): RK and EC methods for flexible MBS

Integration formulae on a Lie group using the exponential map

- Simo (1988, 1991): Newmark and EC scheme for NL structures
- Crouch & Grossman (1993): RK and multistep methods for ODEs
- B. et al (2010, 2011): Generalized- α method for flexible MBS

Lie group generalized- α method

Solution of DAEs on a Lie group (B. & Cardona 2010)

$$\begin{aligned}\mathbf{M}\dot{\mathbf{v}}_{n+1} - \hat{\mathbf{v}}_{n+1}^T \mathbf{M} \mathbf{v}_{n+1} &= -\mathbf{g}(q_{n+1}, t_{n+1}) - \mathbf{B}(q_{n+1})^T \boldsymbol{\lambda}_{n+1} \\ \boldsymbol{\Phi}(q_{n+1}) &= \mathbf{0}\end{aligned}$$

$$q_{n+1} = q_n \exp(\widetilde{\Delta \mathbf{x}_{n+1}})$$

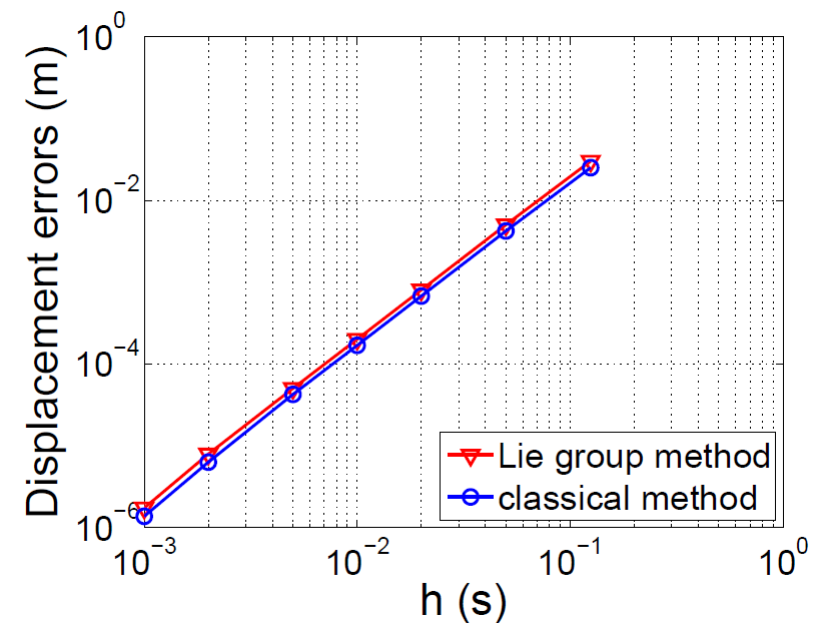
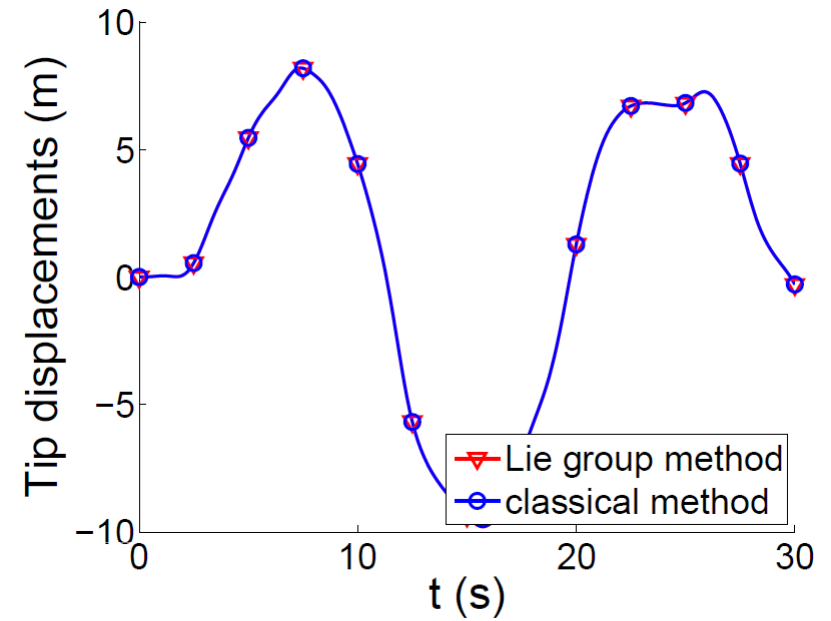
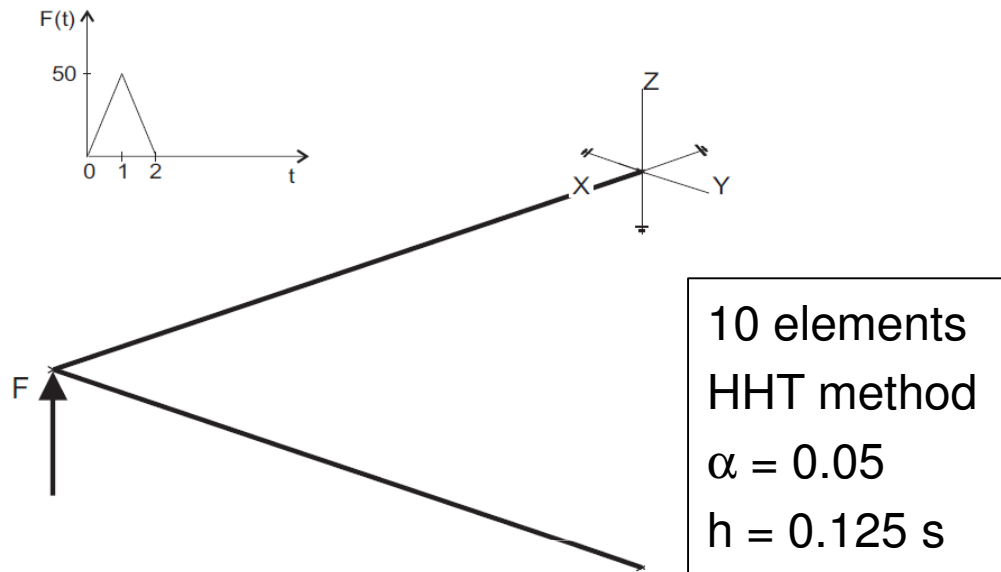
$$\Delta \mathbf{x}_{n+1} = h \mathbf{v}_n + (0.5 - \beta) h^2 \mathbf{a}_n + \beta h^2 \mathbf{a}_{n+1}$$

$$\mathbf{v}_{n+1} = \mathbf{v}_n + (1 - \gamma) h \mathbf{a}_n + \gamma h \mathbf{a}_{n+1}$$

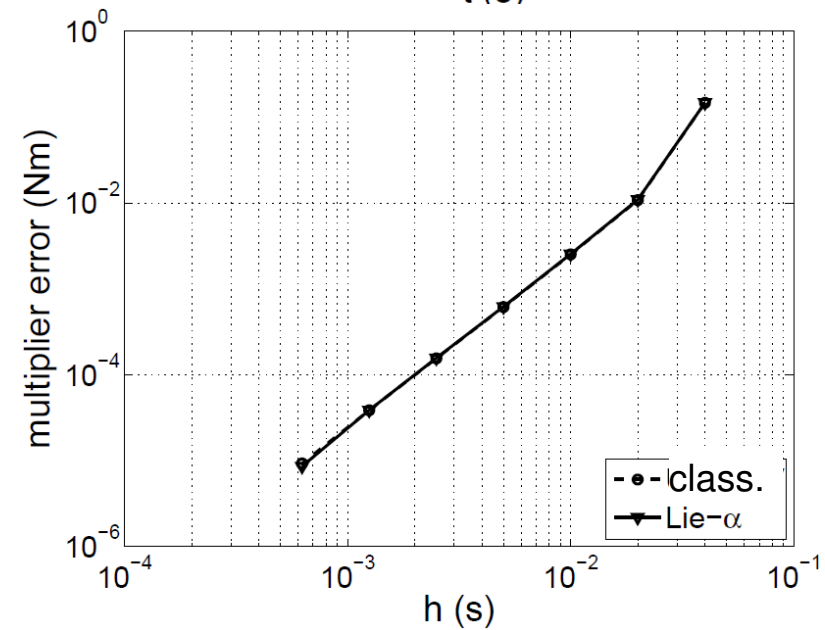
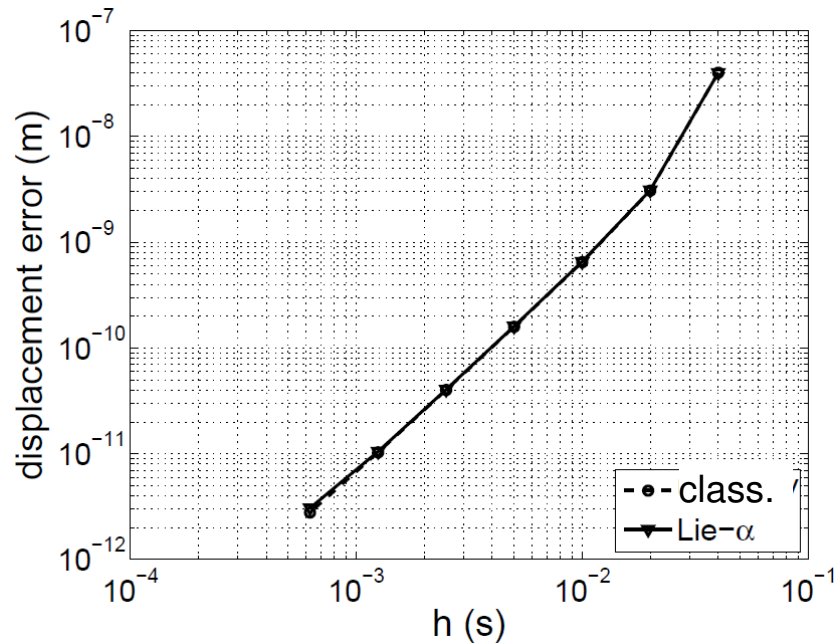
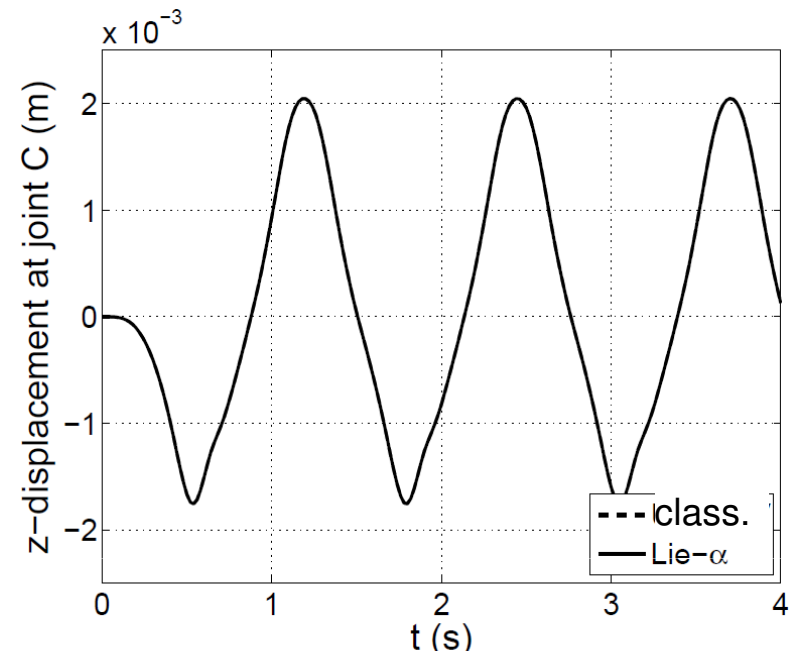
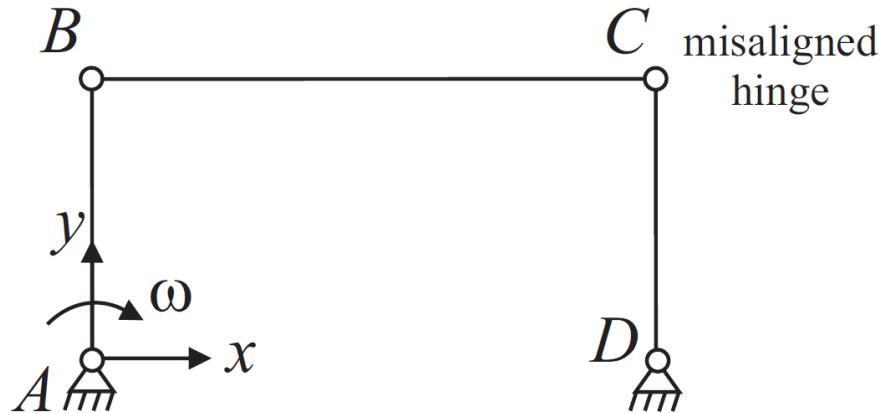
$$(1 - \alpha_m) \mathbf{a}_{n+1} + \alpha_m \mathbf{a}_n = (1 - \alpha_f) \dot{\mathbf{v}}_{n+1} + \alpha_f \dot{\mathbf{v}}_n$$

- Inspired by Newmark / generalized- α methods
- Analytical form of the exponential map
- Newton iterations for vector unknowns (not matrices)
- Second-order convergence (B., Arnold, Cardona 2011)
- Reduced-index formulation (Arnold et al 2011)

Rightangle flexible beam



Flexible four bar mechanism



Sensitivity analysis on a Lie group

Let us consider a single design parameter p

$$\begin{aligned}\dot{q} &= q\tilde{\mathbf{v}} \\ \mathbf{M}(p)\dot{\mathbf{v}} - \hat{\mathbf{v}}^T \mathbf{M}(p)\mathbf{v} + \mathbf{g}(q, p, t) + \mathbf{B}^T(q, p)\boldsymbol{\lambda} &= \mathbf{0} \\ \Phi(q, p) &= 0\end{aligned}$$

and a single criterion function

$$\Psi(p) = G(q(t_f), \mathbf{v}(t_f), p) + \int_{t_0}^{t_f} F(q, \mathbf{v}, \boldsymbol{\lambda}, p) dt$$

Sensitivity in the Lie algebra: $q' = q\tilde{\mathbf{w}}$

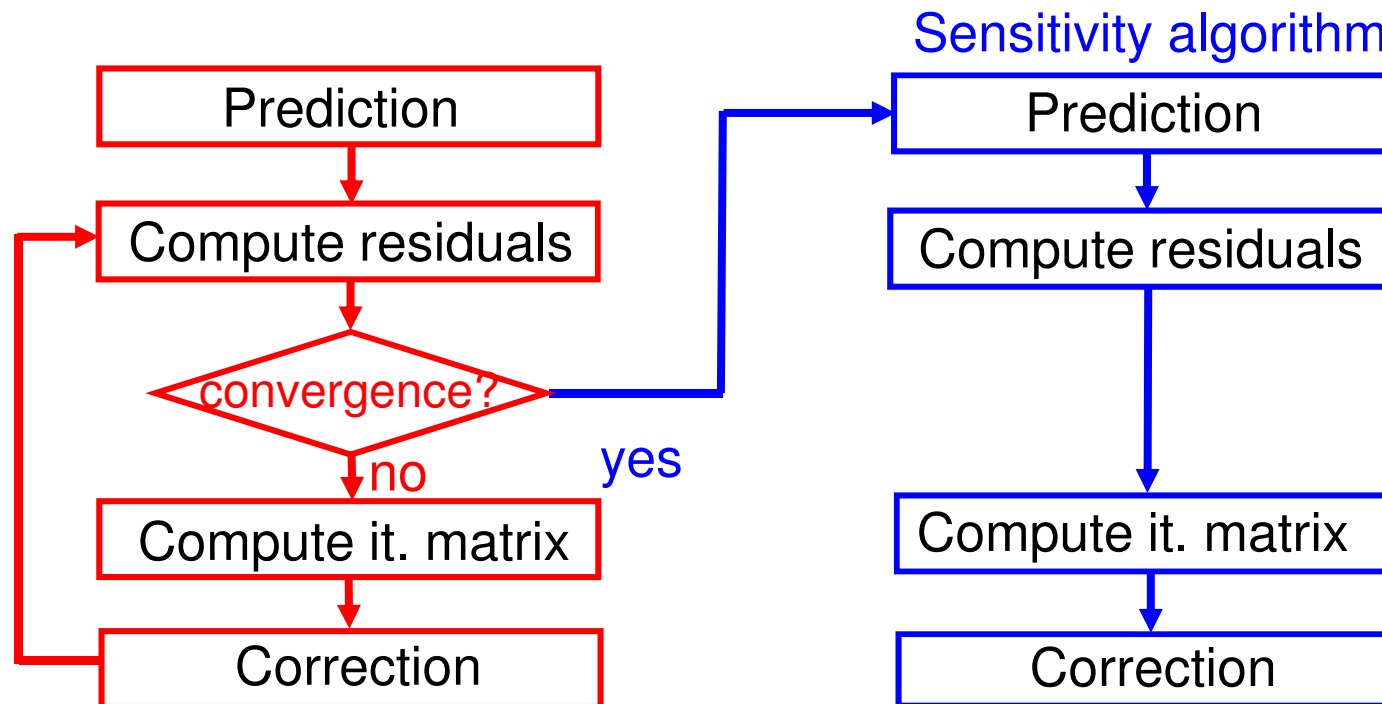
Extension to several parameters and criteria is straightforward

Direct differentiation on a Lie group

$$\begin{aligned}\dot{\mathbf{w}} &= \mathbf{v}' - \tilde{\mathbf{v}}\mathbf{w} \\ \mathbf{M}\dot{\mathbf{v}}' + \mathbf{C}_t\mathbf{v}' + \mathbf{K}_t\mathbf{w} + \mathbf{B}^T\boldsymbol{\lambda}' &= -\mathbf{res}' \\ \mathbf{B}\mathbf{w} &= -\boldsymbol{\Phi}'\end{aligned}$$

With : $\mathbf{res}' = (\partial\mathbf{M}/\partial p)\dot{\mathbf{v}} - \hat{\mathbf{v}}(\partial\mathbf{M}/\partial p)\mathbf{v} + (\partial\mathbf{g}/\partial p) + (\partial\mathbf{B}/\partial p)^T\boldsymbol{\lambda}$
 $\boldsymbol{\Phi}' = \partial\boldsymbol{\Phi}/\partial p$

For each design variable, one linear DAE for \mathbf{w} , \mathbf{v}' and $\boldsymbol{\lambda}'$



Adjoint variable method

$$\delta\Psi = (G_p + \rho^T \chi_p + \pi^T \zeta_p) \delta p + \int_{t_0}^{t_f} (F_p + \mu^T \mathbf{r}_p + \nu^T \Phi_p) \delta p \, dt$$

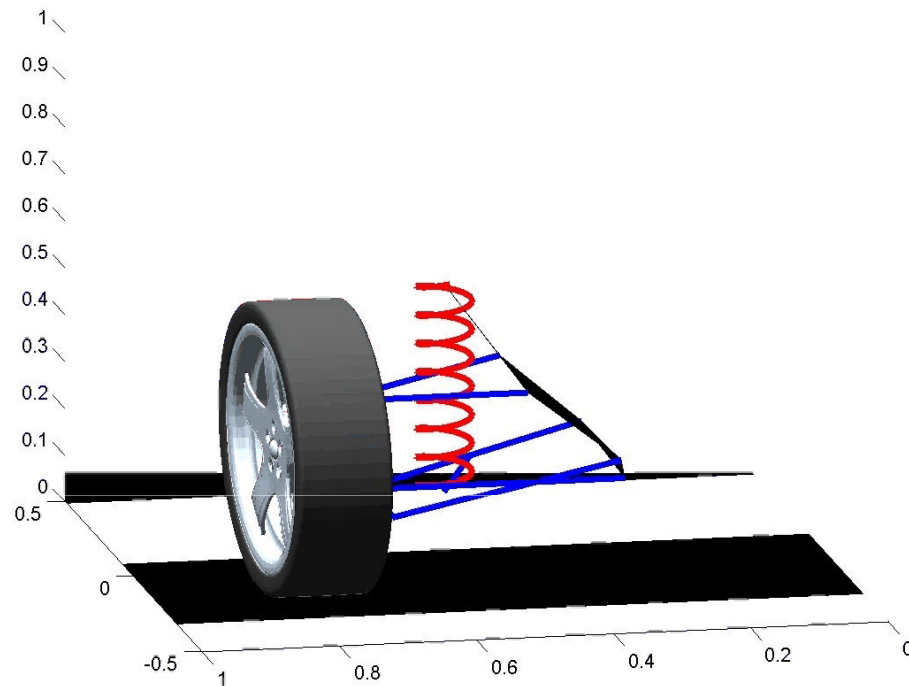
provided that the adjoint variables satisfy

$$\begin{aligned} \mathbf{M}\ddot{\boldsymbol{\mu}} - (\mathbf{M}\hat{\mathbf{v}} + \mathbf{C}_t)^T \dot{\boldsymbol{\mu}} + (\mathbf{K}_t + \mathbf{C}_t^T \hat{\mathbf{v}} - \dot{\mathbf{C}}_t)^T \boldsymbol{\mu} + \mathbf{B}^T \boldsymbol{\nu} &= -F_{q*}^T \\ \mathbf{B} \boldsymbol{\mu} &= -F_{\lambda}^T \end{aligned}$$

$$\begin{aligned} \text{With : } \mathbf{r}_M^T \boldsymbol{\mu}(t_f) &= -(G_C)_{t_f}^T \\ \mathbf{r}_M^T \dot{\boldsymbol{\mu}}(t_f) &= (F_C + \boldsymbol{\mu}^T \mathbf{r}_C + G_K)_{t_f}^T \\ \chi_C^T \boldsymbol{\rho} &= (\boldsymbol{\mu}^T \mathbf{r}_M)_{t_0}^T \\ \zeta_K^T \boldsymbol{\pi} &= (F_C + \boldsymbol{\mu}^T \mathbf{r}_c - \dot{\boldsymbol{\mu}}^T \mathbf{r}_M - \boldsymbol{\rho}^T \chi_K)_{t_0}^T \end{aligned}$$

For each active criterion function, one linear DAE for $\boldsymbol{\mu}$ and $\boldsymbol{\nu}$, which can be solved backward in time

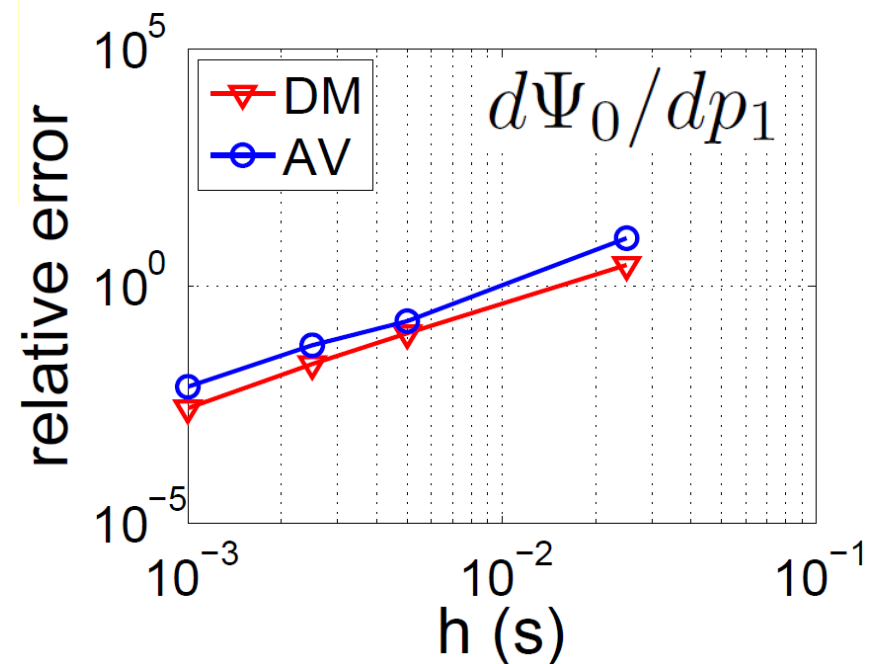
Numerical example



p_1 = damping coefficient

P_2 = stiffness coefficient

$$\Psi_0 = \int_{t_0}^{t_f} \dot{v}_{z,chassis}^2(t) dt$$



Conclusion

The FE method in flexible multibody dynamics has a high potential in mechanical applications for:

- simulation (virtual prototyping)
- dynamic response optimization

However, gradient-based methods require

- a careful formulation of the optimization problem
- efficient transient and sensitivity analysis

Lie group methods may improve the efficiency of 3D models

- parameterization-free formulations and time integration
- simplified algorithms but similar levels of accuracy
- well-suited for sensitivity analysis

Thank you for your attention!

From direct to inverse analysis in flexible multibody dynamics

Olivier Brüls

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