

WASP-36b: A NEW TRANSITING PLANET AROUND A METAL-POOR G-DWARF, AND AN ANALYSIS OF CORRELATED NOISE IN TRANSIT LIGHT CURVES

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ABSTRACT

We report the discovery, from WASP and CORALIE, of a transiting exoplanet in a 1.54-d orbit. The host star, WASP-36, is a magnitude 12.7, metal-poor G2 dwarf ($T_{\text{eff}} = 5881 \pm 137$ K), with $[\text{Fe}/\text{H}] = -0.31 \pm 0.12$. We determine the planet to have mass and radius respectively 2.27 ± 0.07 and 1.27 ± 0.03 times that of Jupiter.

We have eight partial or complete transit light curves, from four different observatories, which allows us to investigate the extent to which red noise in follow-up light curves affects the fitted system parameters. We find that the solutions obtained by analysing each of these light curves independently are consistent with our global fit to all the data, despite the apparent presence of correlated noise in at least two of the light curves.

Subject headings: planetary systems – planets and satellites: detection – planets and satellites: fundamental parameters – stars: individual (WASP-36) – techniques: photometric

1. INTRODUCTION

Of the 171 confirmed transiting planetary systems⁸, the majority have been discovered from the ground, from surveys such as WASP (Pollacco et al. 2006) and HATnet (Bakos et al. 2004). Although the *Kepler* space mission is discovering an increasing number of planets and even more candidate planets (e.g. Borucki et al. 2010; Borucki et al. 2011), the ground-based discoveries have the advantage that the host stars are generally brighter. This allows radial velocity measurements to measure the planetary mass, and is conducive to further characterisation observations, such as measuring occultations in the infrared to probe atmospheric temperature and structure.

Many of the current questions in exoplanet science are being addressed by analysing the statistical properties of the growing ensemble of well characterised transiting planetary systems. Here we report the discovery of a transiting planet orbiting the $V \sim 12.7$ star WASP-36 (= 2MASS J08461929-0801370) in the constellation Hydra.

2. OBSERVATIONS

2.1. WASP photometry

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⁸ <http://www.exoplanet.eu>, 2011 October 13

WASP-36 was observed in 2009 and 2010 by WASP-South, which is located at SAAO, near Sutherland in South Africa, and by SuperWASP at the Observatorio del Roque de los Muchachos on La Palma, Spain. The instruments consist of eight Canon 200mm f/1.8 lenses, each equipped with an Andor 2048 × 2048 e2v CCD camera, on a single robotic mount. Further details of the instrument, survey and data reduction procedures are described in Pollacco et al. (2006) and details of the candidate selection procedure can be found in Collier Cameron et al. (2007) and Pollacco et al. (2008). A total of 13781 measurements of WASP-36 were made between 2009 January 14 and 2010 April 21.

WASP-South 2009 data revealed the presence of a transit-like signal with a period of ~ 1.5 days and a depth of ~ 15 mmag. The WASP light curve is shown folded on the best-fitting orbital period in Figure 1.

2.2. Spectroscopy

Spectroscopic observations of WASP-36 were made with the CORALIE spectrograph of the 1.2-m Euler-Swiss telescope. A total of nineteen spectra were taken between 2010 March 11 and 2011 January 11, and processed using the standard CORALIE data reduction pipeline (Baranne et al. 1996). The resulting radial velocity data are given in Table 1, and plotted in Figure 2. In order to rule out non-planetary causes for the radial velocity variation, such as a blended eclipsing binary system, we examined the bisector spans (e.g. Queloz et al. 2001), which exhibit no correlation with radial velocity (Figure. 2).

2.3. Follow-up photometry

We have a total of eight high-precision follow-up light curves of the transit of WASP-36b, summarised in Table 2. In each case differential aperture photometry was performed, using the IRAF/DAOPHOT package for TRAPPIST and FTN data, and the ULTRACAM pipeline

TABLE 1
RADIAL VELOCITY (RV) AND LINE BISECTOR SPAN (BS)
MEASUREMENTS OF WASP-36

BJD-2 450 000	RV (km s ⁻¹)	σ_{RV} (km s ⁻¹)	BS (km s ⁻¹)
55266.6926	-12.843	0.027	0.079
55293.5807	-13.591	0.021	-0.017
55304.6409	-13.212	0.057	0.025
55305.5493	-13.423	0.025	-0.046
55306.6464	-12.847	0.025	0.014
55315.5600	-12.981	0.023	0.013
55316.5339	-13.624	0.028	-0.005
55317.5642	-12.977	0.033	-0.046
55320.4833	-12.821	0.030	-0.034
55359.4568	-13.527	0.040	0.002
55547.8347	-12.842	0.024	0.007
55561.8393	-12.866	0.026	0.003
55562.8651	-13.306	0.025	-0.018
55563.8184	-13.402	0.033	-0.070
55564.7341	-12.920	0.024	-0.011
55565.8211	-13.467	0.029	0.056
55567.8088	-12.870	0.023	0.001
55570.7984	-12.931	0.023	-0.028
55572.7427	-12.944	0.023	-0.068

(Dhillon et al. 2007; Barros et al. 2011) for the LT data, with aperture radii optimised to give the lowest RMS.

3. DETERMINATION OF SYSTEM PARAMETERS

3.1. Stellar parameters

The individual CORALIE spectra of WASP-36 were co-added to produce a single spectrum with a typical S/N of around 50:1. The standard CORALIE pipeline reduction products were used in the analysis. In order to improve the line profile fitting for equivalent width measurements, the spectrum was smoothed using a Gaussian of width $\sigma = 0.05 \text{ \AA}$. For $v \sin i$ determinations the unsmoothed spectrum was used.

The spectral analysis was performed using the methods given in Gillon et al. (2009). The H_α line was used to determine the effective temperature, T_{eff} , while the Na I D and Mg I b lines were used as surface gravity, $\log g$, diagnostics. The parameters obtained from the analysis are listed in Table 3. The elemental abundances were determined from equivalent width measurements of several clean and unblended lines. A value for microturbulence, ξ_t , was determined from Fe I using the method of Magain (1984). The quoted error estimates account for the uncertainties in T_{eff} , $\log g$ and ξ_t , as well as for the scatter due to measurement and atomic data uncertainties.

The projected stellar rotation velocity, $v \sin i$, was determined by fitting the profiles of several unblended Fe I lines. A value for macroturbulence, v_{mac} , of $3.6 \pm 0.3 \text{ km s}^{-1}$ was assumed, based on the tabulation by Gray (2008), and we used an instrumental FWHM of $0.11 \pm 0.01 \text{ \AA}$, determined from the telluric lines around 630 nm. A best fitting value of $v \sin i = 3.2 \pm 1.3 \text{ km s}^{-1}$ was obtained. The measured $v \sin i$ is sensitive to the adopted value of v_{mac} . The recent work of Bruntt et al. (2010) indicates a lower value of macroturbulence, $v_{\text{mac}} = 2.6 \pm 0.3 \text{ km s}^{-1}$. Using this value, $v \sin i$ rises to $4.1 \pm 1.1 \text{ km s}^{-1}$.

3.2. Neighbouring objects

The Two Micron All Sky Survey catalogue (Skrutskie et al. 2006) reveals the presence of four fainter stars close on the sky to WASP-36. There is no evidence from analysis of catalogue proper motions that any of these stars are physically associated with WASP-36. The stars are separated from WASP-36 by $4''$, $9''$, $13''$ and $17''$, meaning that they fall well within the WASP photometric aperture, which has a radius of $48''$ (3.5 pixels), but outside of the $1''$ CORALIE fibre.

In the absence of reliable optical catalogue magnitudes for all of these objects, it was necessary to measure their fluxes to quantify the effects of blending in the photometry. The fluxes were measured from images taken during the two transits observed with the 1.2-m Euler-Swiss Telescope (see Table 2). The fluxes relative to that of WASP-36 are as follows: 0.012 (object at $4''$ separation from WASP-36), 0.00771 ($9''$), 0.00558 ($13''$) and 0.00827 ($17''$). Using these flux ratios, we corrected the WASP photometry to account for all four objects, and the high precision photometry to account for the object at $4''$, which is within the photometric apertures used. The magnitude of this correction is minimal, and had no significant ($\ll 1\sigma$) effect on the values of our best-fitting system parameters.

3.3. Planetary system parameters

CORALIE radial velocity data were combined with all our photometry and analysed simultaneously using the Markov Chain Monte Carlo (MCMC) method to determine the system parameters. Linear functions of time were fitted to each light curve at each step of the MCMC, to remove systematic trends. Our implementation of this technique is described in detail in Collier Cameron et al. (2007) and Pollacco et al. (2008). The MCMC proposal parameters we use are: the epoch of mid-transit, t_0 ; the orbital period, P ; the transit duration, t_{14} ; the fractional flux deficit that would be observed during transit in the absence of stellar limb-darkening, ΔF ; the transit impact parameter, b ; the stellar reflex velocity semi-amplitude, K_1 ; the stellar effective temperature, T_{eff} ; the stellar metallicity, $[\text{Fe}/\text{H}]$; and $\sqrt{e} \cos \omega$ and $\sqrt{e} \sin \omega$, where e is the orbital eccentricity, and ω is the argument of periastron (Anderson et al. 2011). The stellar mass was determined as part of the MCMC analysis, using an empirical fit to $[\text{Fe}/\text{H}]$, T_{eff} , and the stellar density, ρ_* (Enoch et al. 2010; Torres et al. 2010).

An initial MCMC fit for an eccentric orbit found $e = 0.012_{-0.008}^{+0.014}$ ($\omega = 43_{-107}^{+63}$ degrees), with a 3σ upper-limit to the eccentricity of 0.064, but we found this eccentricity is not significant. Following the F -test approach of Lucy & Sweeney (1971), we find that there is a 66 per cent probability that the apparent eccentricity could have arisen if the underlying orbit were actually circular. We therefore present here the model with a circular orbit, noting that the values of the other model parameters, and their associated uncertainties, are almost identical to those of the eccentric solution.

We tried fitting for a linear trend in the RVs with the inclusion of an additional parameter in our MCMC fit. Such a trend (such as that found in the RVs of WASP-34, Smalley et al. 2011) would be indicative of a third body in the system. The best-fitting radial acceleration is consistent with zero, indicating there is no evidence for an additional body in the system based on our RVs,

TABLE 2
OBSERVING LOG FOR FOLLOW-UP PHOTOMETRY

Light curve	Date	Telescope / instrument	Band	N_{obs}	$t_{\text{exp/s}}$	full / partial
(i)	2010 December 13	Euler ^a / EulerCam	Gunn <i>r</i>	94	120	partial
(ii)	2010 December 13	TRAPPIST ^b / TRAPPISTCAM	<i>clear</i>	756	10	partial
(iii)	2010 December 25	FTN ^c / Spectral camera	Pan Starrs <i>z</i>	176	60	full
(iv)	2011 January 02	TRAPPIST / TRAPPISTCAM	<i>I + z</i>	296	25	full
(v)	2011 January 05	TRAPPIST / TRAPPISTCAM	<i>clear</i>	179	18	partial
(vi)	2011 January 08	TRAPPIST / TRAPPISTCAM	<i>clear</i>	269	18	partial
(vii)	2011 January 15/16	LT ^d / RISE ^e	<i>V + R</i>	1290	9	full
(viii)	2011 January 21	Euler / EulerCam	Gunn <i>r</i>	167	60	full

^a1.2-m Euler-Swiss Telescope, La Silla, Chile

^bTransiting Planets and Planetesimals Small Telescope, La Silla, Chile (Gillon et al. 2011; <http://www.astro.ulg.ac.be/Sci/Trappist>)

^cFaulkes Telescope North, Haleakala Observatory, Maui, Hawaii, USA

^dLiverpool Telescope, Observatorio del Roque de los Muchachos, La Palma, Spain

^eRapid Imaging Search for Exoplanets camera (Steele et al. 2008; Gibson et al. 2008)

TABLE 3
STELLAR PARAMETERS AND ABUNDANCES FROM ANALYSIS OF
CORALIE SPECTRA.

Parameter	Value	Parameter	Value
RA (J2000.0)	08h46m19.30s	[Fe/H]	-0.31 ± 0.12
Dec (J2000.0)	$-08^{\circ} 01' 36.7''$	[Na/H]	-0.39 ± 0.09
T_{eff}	5800 ± 150 K	[Mg/H]	-0.15 ± 0.10
$\log g(\text{cgs})$	4.5 ± 0.2	[Si/H]	-0.17 ± 0.11
ξ_t	0.9 ± 0.2 km s ⁻¹	[Ca/H]	-0.25 ± 0.14
$v \sin i$	3.2 ± 1.3 km s ⁻¹	[Sc/H]	-0.05 ± 0.09
$\log A(\text{Li})$	1.60 ± 0.15	[Ti/H]	-0.22 ± 0.12
Sp. Type	G2	[V/H]	-0.24 ± 0.15
Distance	450 ± 120 pc	[Cr/H]	-0.24 ± 0.23
Age	1 - 5 Gy	[Mn/H]	-0.54 ± 0.13
Mass	0.97 ± 0.09 M _⊙	[Co/H]	-0.30 ± 0.18
Radius	0.92 ± 0.23 R _⊙	[Ni/H]	-0.36 ± 0.12

Additional identifiers for WASP-36:

USNO-B1.0 0819-0221838

2MASS J08461929-0801370

ISWASP J084619.30-080136.7

Note: The spectral type was estimated from T_{eff} using the table of Gray (2008). The mass and radius were estimated using the Torres et al. (2010) calibration.

which span 10 months.

The system parameters derived from our best-fitting circular model are presented in Table 4. The corresponding transit and RV models are superimposed on our data in Figures 1 and 2.

3.4. System age

The lithium abundance of WASP-36 implies an age of 2 to 5 Gy according to Sestito & Randich (2005). The measured $v \sin i$ of WASP-36 gives an upper limit to the rotational period, $P_{\text{rot}} \simeq 15_{-5}^{+10}$ days. This corresponds to an upper limit on the age of $\sim 2_{-1}^{+3}$ Gy using the gyrochronological relation of Barnes (2007).

We searched the WASP photometry for periodic variations indicative of star spots and stellar rotation, but no significant variation was detected. We place an upper limit of 1.5 mmag at the 95 per cent confidence level on the amplitude of any sinusoidal variation. This null result is consistent with the low levels of stellar activity expected from a main-sequence G2 star. A lack of stellar activity is also indicated by the absence of calcium II H+K emission in the spectra. The uncertainties on the CaII emission index, $\log R'_{\text{HK}}$ are too large to allow meaningful constraints to be placed on the system age

by using an activity - rotation relation such as that of Mamajek & Hillenbrand (2008).

In Figure 3 we plot WASP-36 alongside the stellar evolution tracks of Marigo et al. (2008). From this we infer an age of 6_{-3}^{+5} Gy.

There is no evidence of any discrepancy between the ages derived from lithium abundance, gyrochronology and isochrone fitting. This suggests that the star has undergone little or no tidal spin-up, despite the presence of a massive planet in a close orbit.

3.5. Transit timing

We measured the times of mid-transit for each of the eight follow-up light curves, by analysing each light curve separately, without any other photometry (see Section 4.1). The times are displayed in Table 3.3, along with the differences, $O - C$ between these times and those predicted assuming a fixed epoch and period (Table 4). No significant departure from a fixed ephemeris is observed.

4. DETAILED ANALYSIS OF FOLLOW-UP LIGHT CURVES

Because we have several follow-up light curves of WASP-36 from different telescopes / instruments, whereas many planet discovery papers rely on only a single such light curve, we take the opportunity here to examine in detail the potential effects on the system parameters of using only a single light curve.

For survey photometry with low SNR, the durations of ingress and egress are ill-defined, leading to considerable uncertainty in the transit impact parameter and hence to large uncertainties in the stellar density and planetary radius. So-called ‘follow-up’ transit light curves are generally included in the analysis of new ground-based transiting planet discoveries, and are of significantly higher photometric precision than the light curves produced by survey instruments such as WASP. Such follow-up light curves are typically the result of observations with a 1 - 2-m class telescope, and are of critical importance to measuring precisely basic system parameters.

Any light curve may suffer from correlated noise, such as from observational systematics or from astrophysical sources such as stellar activity. To assess the levels of correlated noise in our follow-up light curves, we plot (Figure 4) the rms of the binned residuals to the fit of each light curve as a function of bin width, along with the white-noise expectation. For six of our light curves,

TABLE 4
 SYSTEM PARAMETERS

Parameter	Symbol	Unit	Value
Orbital period	P	d	1.5373653 ± 0.0000027
Epoch of mid-transit	T_c	HJD, UTC	$2455569.83731 \pm 0.000093$
Transit duration	T_{14}	d	0.07566 ± 0.00042
Ingress / egress duration	$T_{12} = T_{34}$	d	0.01531 ± 0.00052 -0.00055
Planet-to-star area ratio	$\Delta F = R_p^2/R_*^2$	-	0.01913 ± 0.00020
Transit impact parameter	b	-	0.662 ± 0.012 -0.014
Orbital inclination angle	i	°	83.65 ± 0.22 -0.20
Stellar orbital velocity semi-amplitude	K_*	km s ⁻¹	0.3914 ± 0.0083
System velocity	γ	km s ⁻¹	-13.2169 ± 0.0024
Orbital eccentricity (adopted)	e	-	0
Orbital eccentricity (3- σ upper-limit)	-	-	0.0647
Stellar mass	M_*	M_\odot	1.020 ± 0.032
Stellar radius	R_*	R_\odot	0.943 ± 0.019
log (stellar surface gravity)	log g_*	(cgs)	4.498 ± 0.012
Stellar density	ρ_*	ρ_\odot	1.218 ± 0.053 -0.048
Stellar effective temperature	$T_{*,\text{eff}}$	K	5881 ± 135 -138
Metallicity	[Fe/H]	dex	-0.31 ± 0.12
Planet mass	M_p	M_{Jup}	2.270 ± 0.068
Planet radius	R_p	R_{Jup}	1.269 ± 0.030
Planet density	ρ_p	ρ_{J}	1.112 ± 0.074 -0.064
Orbital major semi-axis	a	AU	0.02624 ± 0.00027
Planet equilibrium temperature (uniform heat redistribution)	$T_{p,A=0}$	K	1700 ± 42 -44

 TABLE 5
 TRANSIT TIMES

Light curve	E	T_C (HJD, UTC)	σ_{T_C} (min)	$O - C$ (min)
(i)	-17	2455543.70584	5.96	5.39
(ii)	-17	2455543.70379	1.16	2.43
(iii)	-9	2455556.00222	0.63	1.72
(iv)	-4	2455563.68807	0.34	0.32
(v)	-2	2455566.76703	8.86	6.41
(vi)	0	2455569.83686	0.86	-0.65
(vii)	5	2455577.52412	0.21	-0.02
(viii)	9	2455583.67344	0.27	-0.23

the rms of the binned residuals follows closely the white-noise expectation, indicating that little or no correlated noise is present in the data. Light curves (ii and vii) show deviation from the white-noise model, however, suggesting the presence of noise correlated on timescales of ~ 1 and ~ 10 minutes, respectively.

4.1. Method

After modelling all available data in a combined MCMC analysis (see Section 3.3), our ‘global solution’, we also ran several MCMCs each with just a single follow-up light curve in addition to the radial velocities and WASP photometry. Additionally, we re-ran each of these MCMCs applying a Gaussian prior to the stellar radius to impose a density typical of a main-sequence star (the ‘main-sequence constraint’). Such a constraint is usually applied when analysing a new planet which has poor quality follow-up photometry (such as a single light curve which covers only part of the transit), and there is no evidence that the star is evolved or otherwise non-main-sequence in nature. We also performed analyses where the only photometry included was a single follow-up light curve, i.e. the WASP photometry was excluded from the analysis. The purpose of this is to determine whether

the measured depth of transit is biased by inclusion of the WASP photometry. For these runs only, the orbital period was fixed to the value determined as part of our global solution, since this parameter is very poorly constrained by a single transit light curve and a few RVs. The epoch of mid-transit was treated as normal, and allowed to float freely. Finally, we performed an analysis excluding all follow-up photometry; the only photometry analysed was the WASP data.

4.2. Results

We produced correlation plots between several parameters, but choose to present here only plots showing impact parameter against planet radius and stellar radius versus stellar mass (Figures 5 and 6, respectively). Such plots, whilst representative of the ensemble correlation plots, are particularly instructive since b and R_p are two of the major quantities we wish to measure, are largely constrained by follow-up light curves rather than by survey photometry or by radial velocities, and can be significantly correlated with each other, indicating a strongly degenerate solution. The stellar density is measured directly from the transit light curve, and the stellar mass and radius, whilst interesting in themselves, are key in determining the values of several other system parameters of interest.

Several conclusions can be drawn from study of Figures 5 and 6, and similar plots, namely:

(1) each analysis including only a single follow-up light curve gives results that are consistent with our global solution, albeit with larger uncertainties. To measure the dispersion in the best-fitting parameter values obtained from each single follow-up light curve analysis, we calculated the weighted standard deviation. The standard deviations of b , R_p , R_* and M_* are 0.05, 0.08 R_{Jup} , 0.05 R_\odot and 0.006 M_\odot , respectively.

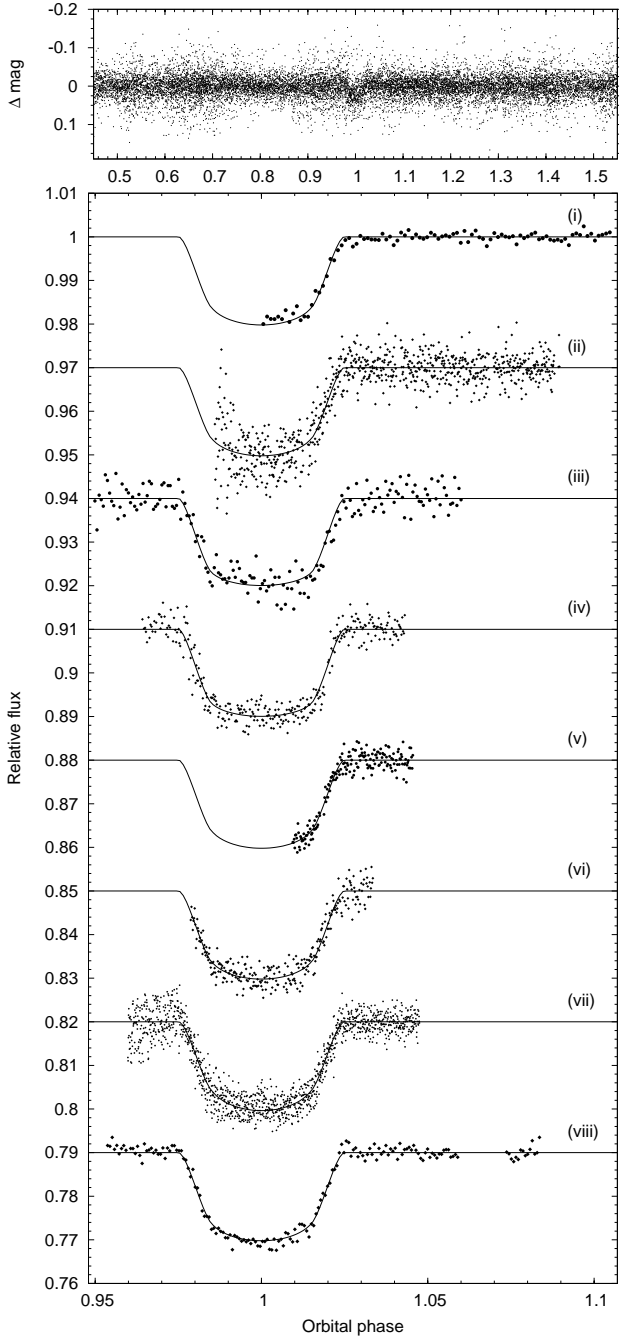


FIG. 1.— Photometry. *Upper panel:* WASP-36 b discovery light curve folded on the orbital period of $P = 1.5373653 d$. For display purposes, points with an error greater than three times the median uncertainty are not shown. *Lower panel:* High-precision transit photometry, over-plotted with our best-fitting model (solid lines). Each individual dataset is offset in flux for clarity, and is labelled with a numeral corresponding to that in the first column of Table 2.

(2) The largest uncertainties are obtained for follow-up light curves that cover the smallest fraction of the transit (light curves i and v), as expected.

(3) The analyses which exclude the WASP photometry give larger uncertainties, but these are only significantly so when the follow-up photometry is poor. This indicates that the WASP photometry only makes a significant contribution to constraining the shape and depth of

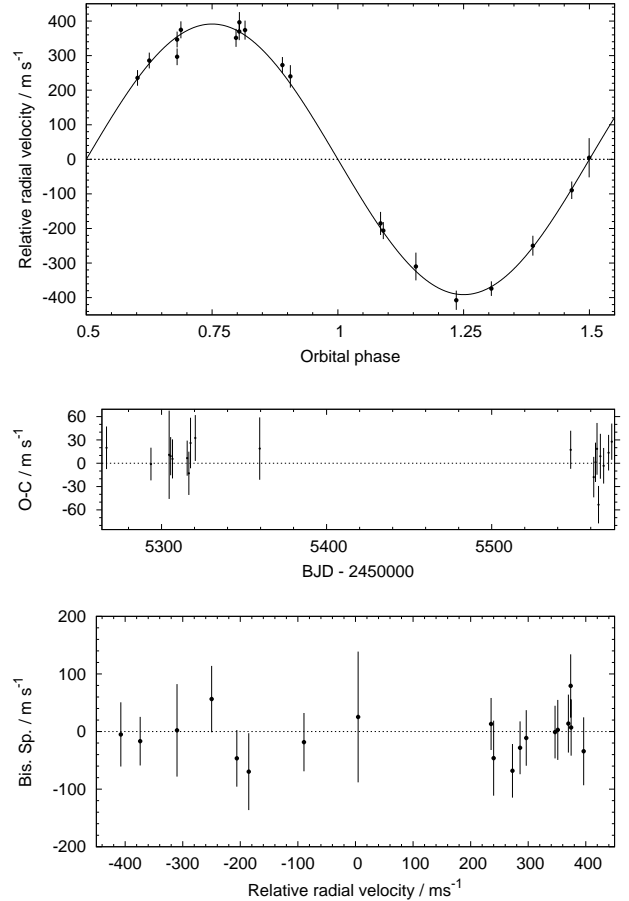


FIG. 2.— Radial velocities. *Upper panel:* Phase-folded radial velocity measurements (Table 1). The centre-of-mass velocity, $\gamma = -13.2169 \text{ km s}^{-1}$, has been subtracted. The best-fitting MCMC solution is over-plotted as a solid line. *Middle panel:* Residuals from the radial velocity fit as a function of time. *Lower panel:* Bisector span measurements as a function of radial velocity. The uncertainties in the bisectors are taken to be twice the uncertainty in the radial velocities.

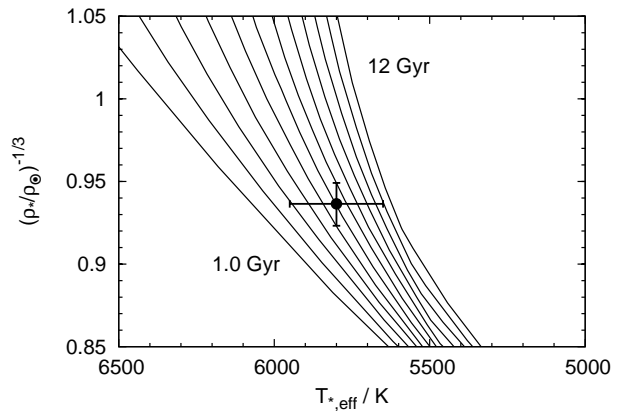


FIG. 3.— Modified Hertzsprung-Russell diagram. WASP-36 is plotted alongside isochrones from the evolutionary models of Marigo et al. (2008). The isochrones are for $Z = 0.0093$, and are regularly spaced at intervals of 1 Gy, from 1 to 12 Gy.

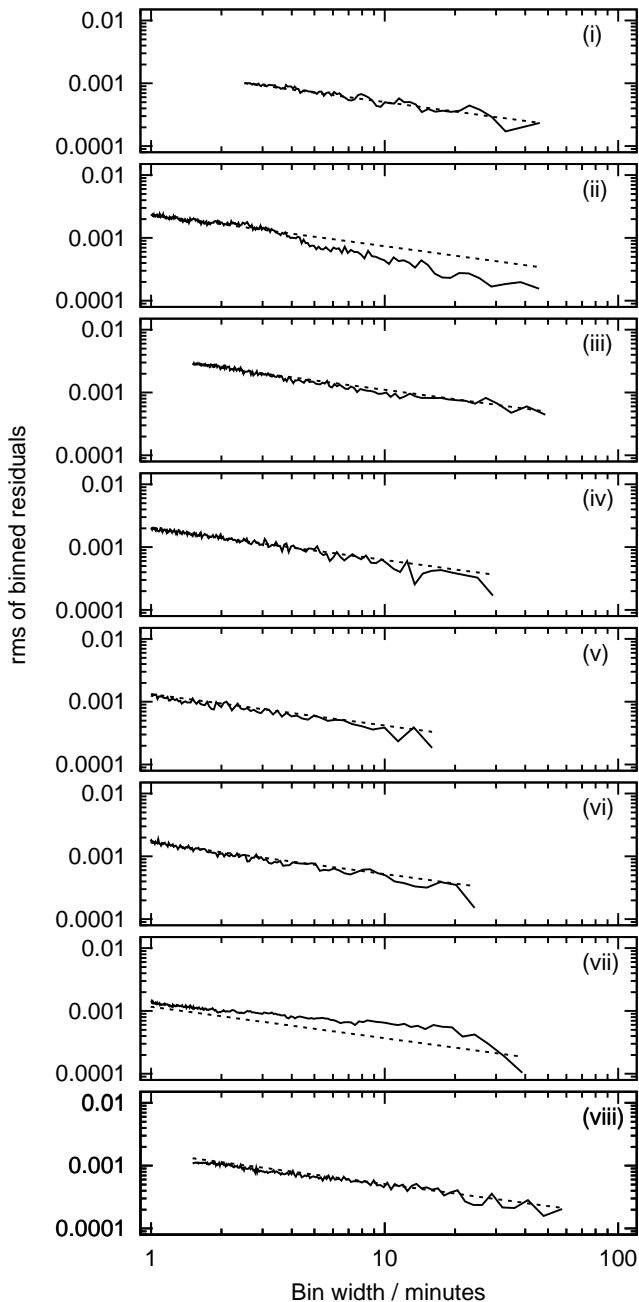


FIG. 4.— Correlated noise in follow-up light curves – rms of binned residuals versus bin width, for light curves, 1 - 8 (solid lines). The white-noise expectation, where the rms decreases in proportion to the square root of the bin size, is indicated by the straight, dotted lines.

the transit when the follow-up light curve is incomplete.

(4) Even a partial transit light curve improves the precision of the measured system parameters enormously compared to those derived solely from the WASP photometry and the RVs.

(5) The imposition of a main-sequence constraint does not significantly alter the parameters or uncertainties for high-precision light curves that are complete, thus indicating that WASP-36 is a main-sequence star. When the

follow-up light curve does not well constrain the range of possible models, however, limiting the star to the main-sequence can significantly reduce the large degeneracy in the possible solutions. This is best illustrated by light curve (iv), where the effects of the constraint are to decrease the stellar density we find and confine the solution to a smaller area of parameter-space, close to the global solution, while largely resolving the degeneracy between b and R_p .

In summary, if only one of the follow-up light curves had been available, we would have reached a solution compatible with the current best-fitting model, although the uncertainties on the model parameters may have been much greater, if the light curve was not of the highest precision. Obtaining additional light curves is clearly of benefit if one only has a light curve that partially covers transit. It is also useful to have multiple high-precision light curves for systems where stellar activity may bias the observed transit depth by varying amounts at different epochs, as may be the case for WASP-10b (Christian et al. 2009; Johnson et al. 2009; Dittmann et al. 2010; Maciejewski et al. 2011b,a).

5. DISCUSSION AND CONCLUSION

WASP-36 is a metal-poor, Solar mass star which is host to a transiting planet in a 1.54 d orbit. We find the planet to have a mass of $2.27 M_{\text{Jup}}$, and a radius $1.27 R_{\text{Jup}}$, meaning it is slightly denser than Jupiter. There is an observed correlation between planetary radius and insolation (e.g. Enoch et al. 2011), with the more bloated planets generally receiving a greater flux from their star. WASP-36b is somewhat larger than predicted by the models of Bodenheimer et al. (2003), which predict radii between 1.08 (for a planet with a core at 1500 K) and 1.20 (for a core-less planet at 2000 K).

The close orbit and large radius of the planet make it a good target for measuring the planetary thermal emission, via infra-red secondary eclipse (occultation) measurements with, for example, *Spitzer*. The expected signal-to-noise ratios of the occultations in *Spitzer* channels 1 ($3.6 \mu\text{m}$) and 2 ($4.5 \mu\text{m}$) are around 10 and 9 respectively.

One of the striking properties of the WASP-36 is the low stellar metallicity ($[\text{Fe}/\text{H}] = -0.31 \pm 0.12$). Giant planets are known to be rare around such low-metallicity stars (e.g. Santos et al. 2004; Fischer & Valenti 2005), although several other low-metallicity systems are known, including the transiting systems WASP-21 ($[\text{Fe}/\text{H}] = -0.46 \pm 0.11$, Bouchy et al. 2010), WASP-37 ($[\text{Fe}/\text{H}] = -0.40 \pm 0.12$, Simpson et al. 2011) and HAT-P-12 ($[\text{Fe}/\text{H}] = -0.29 \pm 0.05$, Hartman et al. 2009).

Such systems will be critical in probing our understanding of the planet-metallicity correlation; proposed explanations for the correlation include insufficient material for proto-planetary cores to attain the critical mass needed for runaway accretion, and the suggestion that the high density of molecular hydrogen in the inner galactic disk is responsible for the effect (Haywood 2009). WASP-36b may also play a key role in determining whether stellar metallicity is the key parameter influencing whether or not a hot Jupiter's atmosphere exhibits a thermal inversion. Insolation was initially propounded as this parameter (Fortney et al. 2008), although this has now been largely disproved. More recently stellar

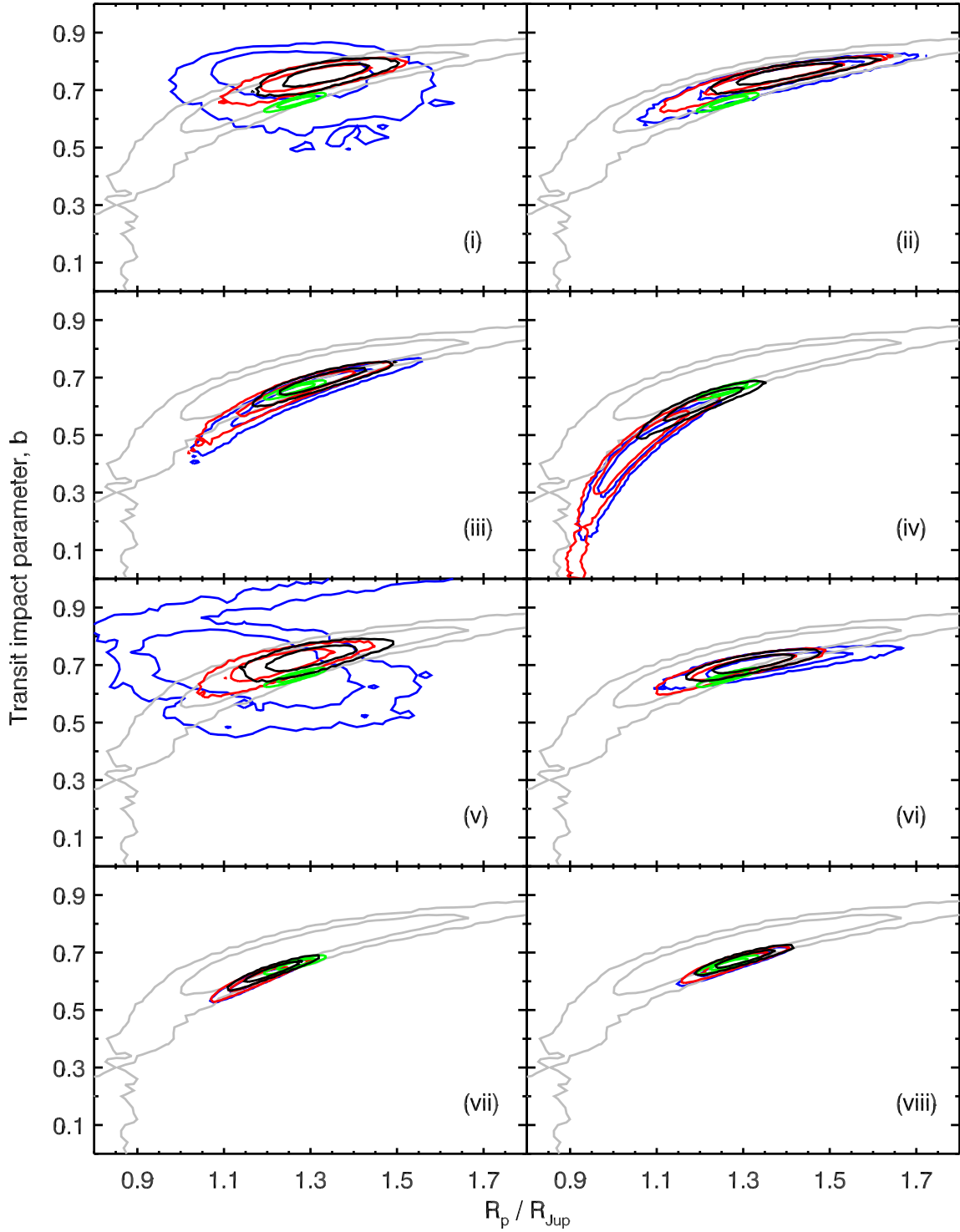


FIG. 5.— Analysis of follow-up light curves I. The MCMC posterior probability distributions for R_p and b for each of the follow-up light curves. The numbering of each panel corresponds to the light curve numbering in Table 2 and the $1\text{-}\sigma$ and $2\text{-}\sigma$ contours are shown. In each case red corresponds to analysis of a single follow-up light curve plus the WASP photometry, black to the single light curve plus the WASP photometry with the main-sequence constraint imposed, and blue to that of a single light curve with *no* WASP photometry. The green contours indicate our global solution, and the grey contours the no follow-up photometry solution, and are therefore identical in each panel.

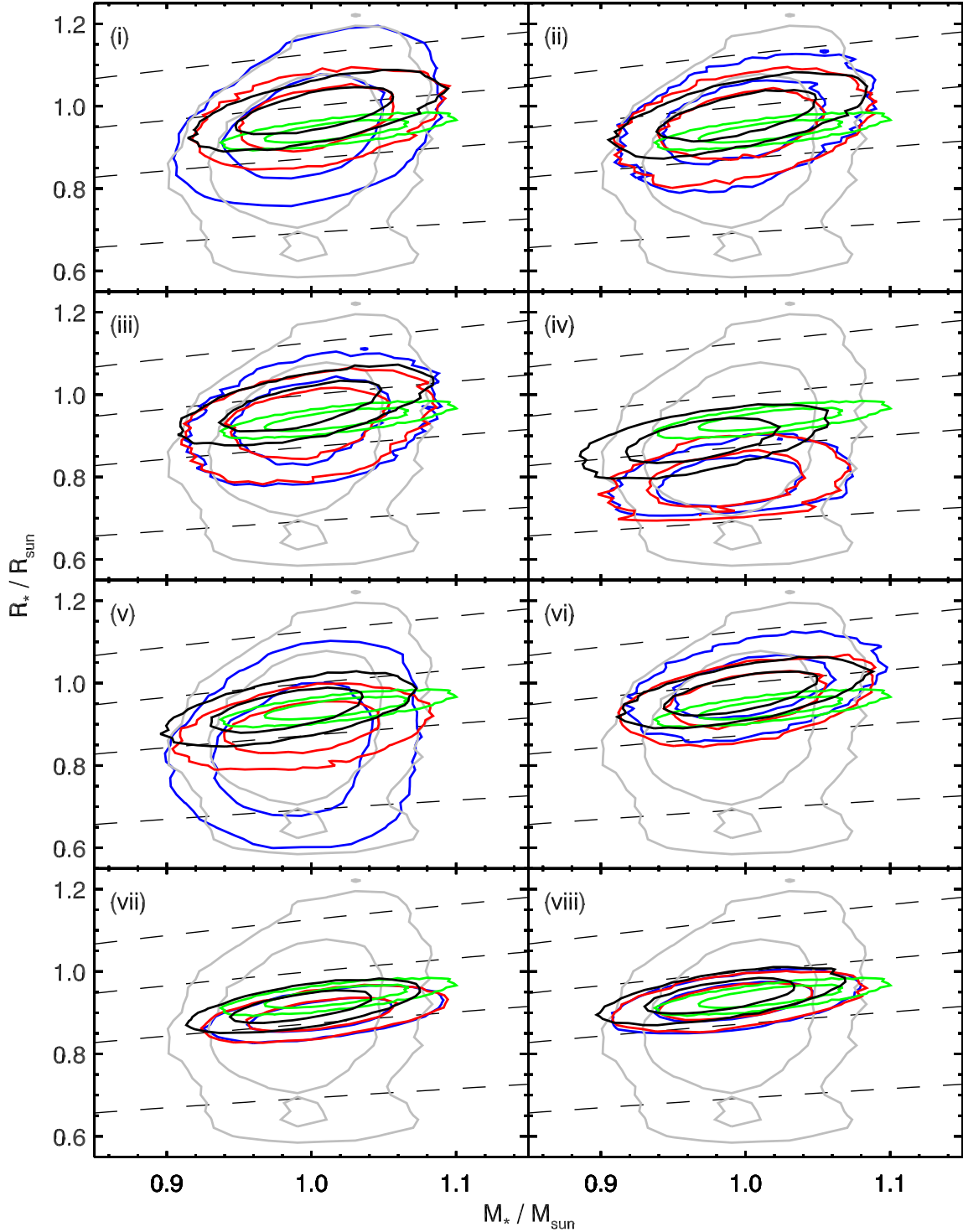


FIG. 6.— Analysis of follow-up light curves II. The MCMC posterior probability distributions for M_* and R_* for each of the follow-up light curves. The numbering of each panel corresponds to the light curve numbering in Table 2 and the $1\text{-}\sigma$ and $2\text{-}\sigma$ contours are shown. In each case red corresponds to analysis of a single follow-up light curve plus the WASP photometry, black to the single light curve plus the WASP photometry with the main-sequence constraint imposed, and blue to that of a single light curve with *no* WASP photometry. The green contours indicate our global solution, and the grey contours the no follow-up photometry solution, and are therefore identical in each panel. Also in each panel are dashed lines which are contours of constant stellar density, corresponding, from top to bottom, to 0.7, 1.0, 1.5, and 3.0 times solar density.

activity (Knutson et al. 2010) and metallicity have been advanced instead; work aiming to resolve this issue is ongoing.

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